A MATHEMATICAL INVESTIGATION OF THE EFFECT OF TUBE SPACING, EXCESS AIR, AND BRIDGEWALL AND STACK TEMPERATURES UPON PIPE STILL DESIGN

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ABSTRACT

A MATHEMATICAL INVESTIGATION OF THE EFFECT OF TUBE SPACING, EXCESS AIR, AND BRIDGEWALL AND STACK TEMPERATURES UPON PIPE STILL DESIGN

by William Vere D'A. Saunders

A mathematical model is presented for use in the design of petroleum furnaces. This model includes the requirements of heat load and pressure drop. A program for the solution of the design equations has been prepared, and solutions were obtained for specific furnace requirements by varying tube spacings, per cent excess air, and bridgewall and exit stack temperatures.

In order to predict pressure drop in the presence of twophase flow an equation similar to the Fanning equation was developed. It is difficult to judge the reliability of this equation because data was not available.

Furnace designs obtained were correlated in terms of the distribution of heat loads to radiant and convection sections, and were plotted as functions of the bridgewall and exit stack temperatures. Radiant heat transfer rates were related to the fraction of the total heat input absorbed in the radiant section. This correlation included the effect of tube spacing.

The principal significance of excess air is its effect on the fraction of total heat input absorbed in the radiant section. At a fixed bridgewall temperature and normal conditions of operation, this fraction can be increased by a factor of approximately 1.5 with a 50 per cent decrease in the excess air.

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A THESIS

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TABLE OF CONTENTS

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Abstract
Acknowledgement ii
List of Figures
List of Tables
List of Appendices
Introduction
Tube Still Heaters
History and Development
Development and Selection of Equations
Radiation
Convection
Heat Balances
Continuous Balance Vaporization - Per Cent Vaporized 20
Physical Properties of the Crude 21
Calculation of the Density of the Crude
at any Temperature
Thermal Conductivity of Petroleum Liquids 22
Pressure Drop Calculations
Statement of the Problem
Restrictions and Limitations
Results of Calculations 27
Discussion of Results 33
Conclusions and Recommendations 40

Page

LIST OF FIGURES

Figure		Page
1	Box Type Furnace	7
2	Types of Pipestill Heaters	8
3	Radiation Between a Plane and One Tube Row Parallel to the Plane	13
4	Variation in Q_r/Q_c With Exit Stack Temperature	35
5	Variation in Q_r/Q_c With Exit Stack Temperature	36
6	Variation in Q_r/Q_c With Exit Stack Temperature	37
7	Per Cent Excess Air Versus Fraction of Total Heat Absorbed in Radiant Section	3 8
8	Fraction of Total Heat Absorbed in Radiant Section vs $Q_t / \alpha A_{cp} $	39
9	Calculations Flow Around One Tube Increment	48
1 0a	Flow Diagram-Convection Section	54
10B	Flow Diagram-Convection Section	55
10C	Flow Diagram-Radiant Section	56
10D	Flow Diagram-Radiant Section	57
11	True Boiling Point Curve and Gravity Profile for Crude	67
12	Flash Curve	68

LIST OF TABLES

Table		Page
I.	Mean Length of Radiant Beams in Various Gas Shapes	11
II.	Service Requirements of Furnace	2 8
III.	Results of Calculations	29
IV.	Results of Calculations	32
٧.	Polynomials for the Evaluation of Ψ	105
VI.	Molar Heat Capacities of Flue Gas Components	107

LIST OF APPENDICES

1

Pa	ge
Bibliography	2
Iomenclature	4
Program Abstract	6
Description of Program 4	7
Sample Calculations	8
Fixed Point Orders 6	9
Floating Point Orders	9
Furnace Design Program	2
Polynomials	5

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INTRODUCTION

The design of modern petroleum furnaces is tending toward larger and more efficient units, sufficiently flexible to adapt to wide variations in physical characteristics of the charging stock, and more responsive to accurate control of the finished product.

It is the current trend in design practice to increase the ratio of heat receiving surfaces to refractory surfaces and to increase the radiant heat transmission rates. These increases affect the size and duty of the convection section and, unless a suitable distribution of heat is obtained throughout the furnace, uneven heating and inefficient units result. There is also a tendency to increase the fraction of vapor discharged from the still to the fractionating tower because increases in efficiency are obtained under these conditions.

It is the purpose of this investigation to study the effect of the tube spacing, percent excess air, and bridgewall and exit stack temperatures and to propose a suitable means of allowing for their effects in the design of tube still heaters.

The successful design of a tube still heater must include the requirements of heat load to the furnace and pressure drop of oil through the tubes. This involves the calculation of (1) flash vaporization data to determine the sensible and latent heat required as well as the fractions of liquid and vapor present in the tubes; (2) suitable heat transfer rates in the radiant and convection sections; (3) the tube length and diameter best suited for the type of still and discharge conditions required.

The mechanisms of heat transfer in combustion chambers are complicated by the phenomenon of radiant heat transmission; as a result, the design of tube still heaters has developed on an empirical basis. Based on the fundamental principle of

- 1 -

radiation postulated in the Stefan-Boltzmann equation, $E_{b} = \delta T^{4}$, and with the aid of certain generalizations, an equation was developed by Lobo and Evans (1),

$$q = \delta \left[T_1^4 - T_2^4 \right] \alpha A_{cp} \psi$$

which is extensively used in the design of heaters.

Calculations for pressure drop were made using the Fanning equation, with a correction for added frictional loss caused by continuous vaporization of crude throughout the still tubes.

By arbitrarily dividing the tubes into a number of zones, the problem of obtaining heat balances and pressure drops over these sections can be solved but necessitates an elaborate trial-and-error calculation. These calculations not only involve the simultaneous solution of heat transfer equations, but also the evaluation of sixty third degree and seven fourth degree polynomials. This has been accomplished with the aid of a digital computer. A program for the computer was prepared whereby changes in tube spacing, per cent excess air, and bridgewall and exit stack temperatures could be made in a design.

The results of these calculations have shown that the distribution of total heat loads to both sections of petroleum furnaces is an important factor in determining their design. Various distribution ratios can be obtained by changing the bridgewall and exit stack temperatures and the per cent excess air. However, the choice of a specific ratio can only be made after the economic permissible heat transfer rates have been determined.

TUBE STILL HEATERS

One of the most important commercial applications of radiant energy transmission is encountered in petroleum refinery furnaces. These furnaces, or tube still heaters, are extensively used in atmospheric or vacuum crude distillation, high temperature gas processing, and thermal cracking, as well as in various heating, treating, and vaporization services.

History and Development

"The early stills used by the oil-refining industry were of the simplest kind. Holding but a few barrels, they were set directly over a coal or tar-fired furnace. The ascending vapors were condensed in a coil submerged in water with no attempt at fractionating further than the gravity indication of the overhead condensate." (2) The stills were generally potshaped and, owing to their construction, were often called shell stills.

In the early years of the petroleum industry, progress was slow. The only attempts made to improve the design of the shell still were increases in its size, leading later to the "cheese-box" still. "The cheese-box still with its eventual capacities up to 1,000 barrels replaced the shell stills at a number of refineries, starting in the late sixties. These cylindrical stills usually had what was termed a 'vapor chest' connected to the still by vertical pipes. The still had a dome-shaped top and a double curved steel plate bottom. The still was supported by a series of arches." (3)

The increased heat efficiency and capacity of these stills reduced the costs per barrel of throughput in comparison with the earlier shell stills. However, the small effective heating surface and the large volume of charged stock caused them to be very inefficient and to have low rates of heat input.

With the advent of the fractionation art and the introduction of cracking operations, it became necessary to construct heaters that could withstand the high temperatures and pressures of the

- 3 -

cracking process. These requirements resulted in a continuous operation permitting the greater use of heat exchange and the steady improvement of design and operating efficiency.

An early attempt at continuous operation involved construction of a battery of shell stills connected in series, with the first still emptying into the second and then in series on to the number of shell units in operation.

"Comparatively successful application of tubular heaters on a small scale for dehydration and refining of emulsified oils led to the gradual adoption of tubular heaters for general refining purposes and eventual substitution for shell heater for large-scale refining operations." (4) The earlier tubular heaters were similar in design to the shell stills, but the stills were displaced by a bank of tubes. This improvement, in some instances, doubled the heat transfer rates from 3000 Btu/ft^2 -hr with the shell type batch operation to 5000 or 6000 Btu/ft²-hr with the tubular heater.

Increases in heat load to the furnace led to localized heating of the tubes and created zones with excessively high transfer rates of 15,000 to 20,000 Btu/ft²-hr, called "hot spots." Coke formation and tube failures occurred at these points. As these furnaces were designed to obtain the major portion of heat transfer by convection, hot spots were attributed to radiation from flames in the fire box. Furnaces were then built with tube banks shielded from the flame by perforated or solid walls to protect the tubes from its radiation. This resulted in a more uniform heat distribution within the furnace, with higher overall heat transfer rates.

It was then discovered that, even with the shield, the tubes first exposed to the combustion products became easily overheated. Unless the gas temperature was reduced below a certain minimum, hot spots would still occur. This reduction of gas temperature was accomplished either by diluting the fuel with excess air and then with recirculated flue gas, or by installing tubes in the combustion chamber to cool the hot gases by absorbing radiant energy from the flame, the gases, and the refractory surfaces of the chamber.

- 4 -

The basic construction of most tube-still heaters is similar. Such units almost always consist of a radiant section, where the major portion of the heat supplied to the process stream is by radiation, and a convection section, where the major portion of the heat is supplied by convection from the gaseous combustion products.

Heaters vary widely in shape and size, and are designed to meet various requirements for such variables as charging stock, heat distribution, thermal efficiency and time-temperature effect. Because thermal decomposition is a rate process, the degree of decomposition is a function of both temperature and time and is described as the time-temperature effect.

Petroleum heaters have been divided (5) into three main groups according to the amount of decomposition obtained.

(1) Heaters used only for heating, with little or no decomposition;

(2) Heaters where, in addition to heating, substantially all of the decomposition desired for the refining process is obtained;

(3) Heaters where only partial decomposition is obtained in the heater, the remainder in the reaction chambers or soaking drums which are usually not heated externally.

Heaters of the first group are employed in operations where no chemical change is desired in the charged stock, as for nondestructive distillation processes. These heater are designed to obtain a minimum time-temperature effect with a maximum temperature.

Heaters of the second group are used primarily in cracking operations where the decomposition of the stock takes place within the heating coils. These heaters are designed to give maximum time-temperature effect at the highest operating temperature allowable.

Heaters of the third group, used for thermally sensitive residual cracking stocks, must be designed for a time-temperature effect that will permit the highest outlet temperature without excessive decomposition from soaking within the heating coils.

- 5 -

A number of typical furnace arrangements are shown in Figures (1) and (2) to illustrate diagrammatically the arrangement of tubes and the direction of fluid flow in the basic types of tube still heaters.

Figure 1 shows a typical box type furnace fired from the end walls. Radiant tubes cover the side walls, roof, and bridgewall (partition between radiant and convection section) surfaces. The tendency in modern furnace design is to fill the radiant section with cold tube surfaces; to accomplish this tubes may also be placed on the floor surfaces.

In cracking operations, oil is preheated in the upper and lower rows of the convection bank then passed through the radiant tubes. After reaching an elevated temperature conducive to the cracking process, the oil is passed through a large number of convection section tubes wherein it is maintained at a high temperature for a sufficient time to accomplish the desired degree of cracking.



FIGURE 1



FIGURE 2. Suppos of lipestill leaters.

DEVELOPMENT AND SELECTION OF EQUATIONS

Evaluation of the rate of heat transmission to the cold surfaces in tube still heaters is accomplished by considering the extent to which each mechanism of heat transfer influences the overall rate.

Transfer of heat energy liberated by the chemical union of molecules in the flame takes place first in the radiant section to the surrounding tubes primarily by radiation though some convection occurs. As the gaseous products of combustion progress through the convection section, heat transmission is principally caused by the mechanism of forced convection accompanied by small amounts of gas radiation.

Thermal energy transferred within the furnace enclosure must be equated to the change in enthalpy of the entering and exit streams. To accomplish this, the following equations must be obtained: heat transfer by radiation, heat transfer by convection, and heat balances about the oil.

Radiation

Radiation in a combustion chamber originates from three distinct sources (6):

- (1) the chemical union of molecules in the flame,
- (2) the hot products of combustion, and

(3) the luminosity or soot content of the flame. The magnitude of radiation emitted from the first source is dependent on the composition of the fuel, the maximum temperature attained, and the absorbing characteristics of the flame for its own radiation. However, in muffle furnaces where the flame is shielded from the surfaces of the combustion chamber, heat from this source is transferred to the combustion products by conduction and convection.

Radiation of greatest magnitude originates from the combustion products and is dependent on composition, temperature, and shape

- 9 -

and size of the gas mass. Of the gases comprising the combustion products, carbon dioxide, carbon monoxide, the hydrocarbons, and water vapor are the only ones with emission bands of sufficient energy to merit consideration. Gases with simple symmetrical molecules, such as N_2 , H_2 , and O_2 , which also comprise the total gas mass, show no absorption bands in the region of importance in radiant heat transmission. Moreover, carbon monoxide and hydrocarbons are present in such small amounts as to be negligible compared with water vapor and carbon dioxidé. Finally, the third source of radiation from the flame, its soot content, is dependent on the degree of combustion and the design of the combustion chamber.

Using data obtained from investigations on the infrared spectra of carbon dioxide and water vapor, Hottel (7) has presented charts for use in calculating the quantity of heat transmitted from these gases. He has also shown that the energy emitted from a gas mass to a unit area of bounding surface is a function of the gas and surface temperatures, the absorptivity of the surface, and the product PL, where P is the partial pressure in atmospheres of the radiating constituents, and L is the average length of a blanket of flue gas in all directions for each of the points of the bouding surface of the furnace. Values of L for furnaces of various shapes were determined by Hottel and Table I presents a digest of these values for furnace calculations.

Incident radiation is not completely absorbed by its ultimate heat receiving surfaces immediately but is reflected and absorbed in an infinite series of interchanges between source and surface. Consequently, radiant interchange between the surfaces of an enclosure must involve consideration of the view the surfaces have of each other as well as their emitting and absorbing characteristics.

The absorptivities of bodies are generally dependent on the wavelength of incident radiation and also on the factors affecting their emissivities. Absolute values of the emissive power of bodies are not readily obtainable, however, the ratio of the actual emissive power to the black body emissive power,

- 10 -

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TABLE I

Mean Length of Radiant Beams in Various Gas Shapes

Dimensional Ratio (length, width, height in any order)

Rectangular Furnaces

 1-1-1 to 1-1-3 1-2-1 to 1-2-4
 1-1-4 to 1-1-∞
 1-2-5 to 1-2-8
 1-3-3 to 1-∞-∞

2/3^{3.} V Furnace Volume

.

 L_{B}

1 x smallest dimension

1.3 x smallest dimension

1.8 x smallest dimension

defined as the emissivity, has been determined for many materials and data are presented in most textbooks of heat transfer. In systems such as furnaces composed of walls and pipes, it becomes difficult to evaluate the manner in which radiant energy falls on these surfaces. The next flux between source and surface occurs by a complex process involving multiple reflection from all surfaces forming the enclosure. The new concept necessary here is F, and has been defined by Hottel (8) as the direct interchange factor, dependent on the angle factors between the refractory surfaces and the surfaces surrounding it, together with the emissivities of the source and sink surfaces.

Since heat receiving surfaces or heat sinks in most industrial furnaces are composed of a multiplicity of tubes disposed over walls, roof, and floor of the combustion chamber, it is necessary to evaluate the effective heat transfer area. The development of Hottel, almost exclusively used in design work, assumes that the heat source is a radiating plane parallel to the tube row. The effectiveness factor, α , is the factor by which the surface of a plane replacing the tube row with assumed emissivity of 1.0, must be multiplied to obtain the equivalent cold plane surface. For a detailed development of α , reference should be made to Hottel (9). Figure 3 presents values of α for radiation to single rows of tubes with refractory behind them.

In view of the complexity of the problem, numerous investigators have correlated furnace performance by means of empirical and semitheoretical equations. The most acceptable of these, is the semitheoretical equation proposed by Lobo and Evans (1). Using an equation of the Stefan-Boltzman type in correlating data from 85 tests on 19 different petroleum furances, they developed the following equation:

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$$q = \delta (T_{g}^{4} - T_{s}^{4}) \alpha A_{cp} \psi + h_{c} A^{\dagger} r (T_{g} - T_{r}) + h_{c} A_{c} (T_{g} - T_{s})$$
(1)

- 12 -



- 13 -

= Pactor of C mparison with Two Parellel Flaces

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In this equation the direct interchange factor, F, has been replaced by an overall exchange factor, ψ . ψ includes, in addition to direct interchange, the contributions due to multiple reflection at all surfaces as well as such contributions by reradiation from zones at which the net radiant-heat transfer at the wall surface is zero.

Since both the external losses from the furnace and the net heat transferred to the refractory by convection, given by the term $\mathbf{h_c A'_r}(\mathbf{T_g - T_r})$, are usually small, the two may be assumed equal without appreciably affecting the results. Equation (1) may be rewritten

 $q = \delta \left[T_g^4 - T_s^4\right] \alpha A_{cp} \psi + h_c A_c (T_g - T_s)$ (2) By making the following assumptions, Lobo and Evans further simplified their equation:

The convection coefficient lies normally between 2 and
 Btu/hr sq ft - ^oF;

2. In most furnaces A_c equals (2 αA_{cp}) approximately;

3. The overall exchange factor ψ has a value of about 0.57. Therefore, the terms h_c and A_c in Equation (2) can be expressed in terms of and ψ , thus:

$$\frac{{}^{h}{}_{c}{}^{A}{}_{c}}{\alpha A_{cp}}\psi = \frac{(2)(2)}{(0.57)} = 7.0$$

or

$$h_c A_c (T_g - T_s) = 7(\alpha A_{cp} \psi) [T_g - T_s]$$

Making the substitution in Equation (2),

$$\frac{\mathbf{q}}{\alpha \mathbf{A}_{\mathbf{c}\mathbf{p}}\mathbf{\psi}} = \delta[\mathbf{T}_{g}^{\mathbf{\mu}} - \mathbf{T}_{s}^{\mathbf{\mu}}] + 7 (\mathbf{T}_{g} - \mathbf{T}_{s})$$
(3)

In the combustion chamber T_g , the mean temperature of the hot gases in the furnace and the temperature of the exit gases will undoubtedly differ. However, the assumption of complete mixing in the furnace and that T_g could be replaced by the exit gas temperature was justified by satisfactory results. The exact evaluation of ψ is tedious and complicated, however, their development included a plot of ψ versus the ratio $\frac{A_r}{CP}$ with the flame emissivity as a parameter. Results indicate that the ψ plot represents an accurate and simple method of simultaneously allowing for the effect of flame emissivity and the amount of refractory surface present.

Hottels' charts giving the values of the radiant heat transfer flux due to CO_2 and water are most conveniently used in calculating the emissivity of the flame. The radiant flux of H_2O and CO_2 are additive, although a small correction must be included to allow for the influence of one type of molecule with radiation from the other. The flame emissivity is given by the equation:

$$\epsilon_{g} = \frac{(q_{c} + q_{w})_{T_{g}} - (q_{c} + q_{w})_{T_{s}}}{(q_{b})_{T_{g}} - (q_{b})_{T_{s}}}$$
(4)

Radiant heat transmission in the convection section is particularly significant to the uppermost tubes or "shield tubes" in the convection bank where the temperature of the gases is still high. In spite of the small beam length, Monrad (10) has shown that radiation may account for 5 to 30 percent of the total heat transfer in the convection section.

Evaluation of the radiant heat coefficient was accomplished by adapting the method of Lobo and Evans and the simplified charts provided by Hottel. The mean length of the radiant beam for exchange between tubes, given in Table I, is based on the center to center spacing $(\not e - \not e)$, and the outside diameter of the tubes (D_0) .

$$L_{B} = 0.4 [(\not e - \not e) - 0.567] D_{O}$$
 (5)

$$q_{rc} = \epsilon_{s} \left[\left(q_{c} + q_{w} \right)_{T_{g}} - \left(q_{c} + q_{w} \right)_{T_{s}} \right] A \left(\frac{100 - \%}{100} \right)$$
(6)

since

$$h_r = \frac{q_{rc}}{A\Delta T}$$

$$h_{r} = \epsilon_{s} \frac{\left[\left(q_{c} + q_{\mu}\right)_{T_{g}} - \left(q_{c} + q_{\mu}\right)_{T_{s}}\right]}{T_{g} - T_{s}} \frac{\left(100 - \%\right)}{100}$$
(7)

where the emissivity of the surface ϵ_s is assumed to have a value of 0.95. (11)

Convection

Heat transfer by convection varies widely with gas velocity and size of gas passage, somewhat with temperature of the gas, and very little with gas composition. Although there are numerous relationships available for obtaining convection coefficients, little work has been done to develop a satisfactory relationship for furnace design work. The empirical equation by Monrad (10) is the only comprehensive formulation of convection transfer rates. For direct convection from the gases he proposes the relationship:

$$h_{c} = \frac{1.6 \ G^{2/3} \ T^{0.3}}{D' \ D'}$$
(8)

The equation applies to any conventional arrangement of the tubes in the convection section. However, the coefficient h_c is the pure convection coefficient and it does not include radiation from the hot gases or from the walls. Monrad has made a study of these factors. The first of these is designated as h_{rg} or the coefficient of heat transfer from the gas by radiation. A formula for the evaluation of this coefficient was presented in the previous section on radiation.

In his calculations, Monrad also included a correction for the increased thickness of the gas layer at the top of the tube bank. His assumption was that this radiation could be approximated by that of a plane equal in area to the top tube bank at PwL = PcL = 1.0, between the temperature of the gas above the bank and the temperature of the tube. He reasoned that since radiation had already been calculated for PwL and PcL based on the center to center spacing and the tube diameter (h_{rg}) , the added gas radiation would be equivalent to that at the top gas temperature between PcL = PwL = 1.0 and PcL and PwL based on the center to center spacing and the diameter; consequently, the correction:

 $h'_{rg} = h(at PL = 1) - h_{rg}$

This however, appears to be too severe a correction as PL values for the gas layer above the bank could hardly approach a value of unity. Water vapor, the radiating constituent of highest concentration in the gas mass rarely exceeds a partial pressure of 0.20 atms. Consequently, the average beam length would necessarily have to be greater than 50 feet (for lower concentrations of water vapor) to cause values of PL = 1.0. For this reason, it was concluded that sufficiently accurate results would be obtained without including this correction.

The area of the walls surrounding the tubes comprise a fairly large fraction of the tube area. These walls pick up heat from the gases by convection and radiation, and reradiate to the tubes by black body radiation according to Stefans' Law. With the assumption that factors such as reabsorption of heat and the differences in heat transfer coefficients to the wall and tubes are negligible, the following equations were presented:

$$[h_{c} + h_{rg}][T_{g} - T_{w}]A_{w} = heat to walls$$

$$= [h_{rb}][T_{w} - T_{t}]A_{t} = heat from the walls = heat$$

$$to tubes from walls$$

Therefore the per cent increase in heat absorption by tubes above that received directly:

$$\frac{\mathbf{h}_{rb} \left[\mathbf{T}_{w} - \mathbf{T}_{t} \right] \mathbf{A}_{w} \times 100}{\left[\mathbf{h}_{c} + \mathbf{h}_{rg} \right] \left[\mathbf{T}_{g} - \mathbf{T}_{t} \right] \mathbf{A}_{t}} = \frac{\mathbf{h}_{rb} \left[\mathbf{T}_{w} - \mathbf{T}_{t} \right] \mathbf{A}_{w} \times 100}{\mathbf{A}_{t} \left[\left[\mathbf{h}_{c} + \mathbf{h}_{rg} \right] \left[\mathbf{T}_{g} - \mathbf{T}_{w} \right] + \left[\mathbf{h}_{c} + \mathbf{h}_{rg} \right] \left[\mathbf{T}_{w} - \mathbf{T}_{t} \right] \right]}$$
(9)

$$h_{rb} \text{ may be approximated by:}
h_{rb} = 0.00688 \epsilon_{s} \left[\frac{T}{100}\right]^{3}$$

$$\epsilon_{s} = 0.95$$
(10)

T = Temperature of the tube surfaces The complete coefficient of heat transfer in the convection section was computed from the preceding items as follows:

$$h_{\rm u} = \frac{(100 + \% \text{ wall effect})}{100} \quad (h_{\rm c} + h_{\rm rg}) \tag{11}$$

In considering radiation from the walls to the individual rows of tubes in the bank, the correction for the wall effect becomes less pronounced as the gases cool on their way to the stack. The wall effect then, will be of most significance to the tubes in the shield section. If a correction is made for the shield section (suppermost rows surrounded by gas at a temperature of 1300°F or above) only, it will be necessary to evaluate the fraction of the wall surface that will "see" the tubes in the section. An approximation was made by assuming that the area of the walls surrounding the shield tubes would be the only surfaces that could see that section.

Corrections for the wall effect and gas radiation were made only for tubes in the shield section, whereas the effect of radiation was neglected in calculations for the film coefficient for the rest of the tubes in the convection bank.

In general, the coefficient of heat transfer on the gas side is controlling; although the resistance of the liquid film may be assumed negligible, approximate values were estimated using a modified Dittus-Boelter (11) equation.

$$\frac{h_{1}D_{1}}{K} = 0.027 \left(\frac{DG}{\mu}\right)^{0.8} \left(\frac{C\mu}{K}\right)^{1/3} \left(\frac{\mu}{\mu_{W}}\right)^{0.14}$$
(12)

Based on the outside pipe diameter, the overall coefficient of heat transfer (U) is given by:

$$U = \frac{1}{\frac{1}{h} + \frac{D_{o}(D_{o}-D_{i})}{K_{m}(D_{o}+D_{i})} + \frac{D_{o}}{D_{i}h_{i}}}$$
(13)

and the heat flux to the oil evaluated using the following equation:

$$q_{c} = U (\pi D_{o} \Delta L) (\Delta t)$$
(14)

Heat Balances

The energy balances made throughout these calculations involving the physical properties of the crude oil, can only be fair approximations due to the complex nature of the hydrocarbon mixture. Calculations based on the heat content of the charged stock will involve some error since the difficulty of obtaining the accurate specific heat of the oil, its vapor and its latent heat of vaporization is quite large. The difficulty in estimating external heat losses from the furnace is another source of error.

The liquid heat content of Mid-Continent source crude oils was calculated using the equation obtained by Weir and Eaton (13).

 $(H - Ho)_{L} = (15d - 27) + (0.811 - 0.465d)t + 0.000290t^{2}$ (15) H = total heat continent above 32°F. Ho = heat continent at 32°F. t = temperature °F.

d = sp. gr. of material at 60°F.

The heat content of liquids do not vary appreciably with pressure, therefore, Equation (15) was used to calculate the liquid heat content over the entire range of pressures encountered in the pipe still. Like the liquid data, Weir and Eaton found it possible to incorporate the vapor heat content versus temperature relationship in a single recommended equation.

$$(H - H_0)_v = (215 - 87d) + (0.415 - 0.104d)t$$

$$+ (0.000310 - 0.000078d)t^2$$
(16)

It is necessary to include a correction for the variations with pressure. Using the equation of state proposed by Linde to predict PV data, and the constants as obtained by Baklke and Kay (14) an equation for the total heat content of the vapor was developed.

$$PV = AT - P \frac{(C + EP)}{T^3} + P(D + FP)$$
 (17)

$$V = \frac{AT}{P} - \frac{C + EP}{T^3} + D + FP$$
(17b)

(b)

$$T = {}^{\bullet}R$$

$$V = ft^{3}/lb$$

$$P = \#/in^{2}$$
A, C, D, E, and F = constants
$$d_{H} = C_{p}dT - [T(\frac{\partial V}{\partial T})_{p} - V] dP$$
(a)

Differentiating Equation (17)

$$P\left(\frac{\partial V}{\partial T}\right)_{P} = A + \frac{3 CP}{T^{4}} + \frac{3 EP^{2}}{T^{4}} \quad \text{or}$$
$$T\left(\frac{\partial V}{\partial T}\right)_{P} = \frac{AT}{P} + \frac{3C}{T^{3}} + \frac{3EP}{T^{3}}$$

Substituting (b) and (17b) into (a) and simplifying

$$d_{H} = C_{p}dT - \left[\frac{4C}{T^{3}} + \frac{4EP}{T^{3}} - D - FP\right] dP$$

integrating

H - Ho =
$$\int_{32}^{T} C_{p} dT + \frac{1}{9331.7} \left[DP + \frac{FP^{2}}{2} - \frac{4CP}{T^{3}} - \frac{2EP^{2}}{T^{3}} \right]$$
 (18)

and

$$\int_{32}^{T} C_{p} dT = (215 - 8Td) + (0.415 - 0.104d)t + (0.000031 - 0.000078d)t^{2}$$

The constants, as obtained by Bahlke and Kay are: A = 157 $C = 7234 \times 10^7$ D = 20 $E = 102 \times 10^7$ F = 0.52

Energy balances on the gas side were obtained from heat capacity equations for constituent in the combustion product. These equations were summed according to the total moles of each constituent and evaluated as a single equation. These equations are presented in the appendix.

Continuous Equilibrium Vaporization - Per cent Vaporized

In most cases of distillation of such complex mixtures as crude oil, continuous equilibrium vaporization is used. It is then necessary to know the relationship between equilibrium vaporization temperature and per cent vaporized for any given pressure if intelligent design calculations are to be made.

Piromoor and Beiswenger (15) have established a widely used correlation which enables the flash curve of the crude (flash zone temperature versus per cent oil vaporized) to be estimated from the true boiling point curve of the crude. These correlations were later modified by Maxwell (16) and are based upon the empirical facts that:

 The True Boiling Point curves (TBP) of many commonly encountered crudes and fractions are nearly straight lines between their 70% and 10% vaporized points.

- 2. There exists a fairly close relationship between the slope of the TBP curve (straight position) and the flash vaporization curve.
- 3. There exists a relationship between the 50% distillation, TBP temperature and the 50% distillation point of the flash curve.

An example of the use of these relationships is outlined in reference 17. To correct flash curves to other pressure, the flash curve is displaced parallel to itself at a higher or lower temperature (depending on whether the pressure is higher or lower than atmospheric) as determined by the vapor pressure versus temperature chart (Cox Chart). The vapor pressure versus temperature relationship with an atmospheric temperature corresponding to that at the 50% distilled point on the flash curve is chosen.

Physical Properties of the Crude

The physical properties of an oil are found to vary gradually throughout the range of compounds that constitute the oil. The properties such as specific gravity and viscosity are found to be different for each drop or fraction of the material distilled. The rate at which these properties change may be plotted as mid per cent curves; i.e., a plot of the desirable property versus percentage distilled. A mid per cent yeild curve was used to determine the specific gravity at 60°F of the crude at its different stages of vaporization. The viscosity was obtained using the relationship obtained by Nelson (18) for the high temperature viscosities of hydrocarbons.

Calculation of the Density of the Crude at any Temperature (19)

It is assumed that the thermal expansion of any sample may be represented by an equation of the form:

 $D_{+} = D_{m} + A(t-T) + B(t-T)^{2}$

 $\mathbf{D}_{\!\!+}$ = density at any temperature t.

 D_m = density at a standard temperature

and A and B are based on the change at 25°C.

Since the specific gravities of most materials are given at $15.56^{\circ}C$ (60°F), this temperature is chosen as standard. Then,

 $SG_{60} = SG + [\alpha_{T} + 2\beta(t-25)][t-15.56] + \beta[t-15.56]^2$ (19) converted to degrees Fahrenheit

$$SG_{60} = SG + \left[\frac{t-60}{1.8}\right] \alpha_{T} + \frac{\beta_{T}}{1.8} [3t-217]$$
 (19a)

Values of $\boldsymbol{\alpha}$ and $\boldsymbol{\beta}$ were obtained from charts as functions of the specific gravity of the material.

Thermal Conductivity of Petroleum Liquids (20)

The thermal conductivity of the crude is given by the following equation:

$$K = \frac{0.813}{12(S.G)} [1 - 0.0003 (t-32)]$$
(20)

Pressure Drop Calculations

The problem of calculating pressure drop in tube still heaters cannot be solved by the conventional Fanning equation; since vaporization of the crude with increasing temperatures, results in the presence of two phases, making the equation inapplicable.

Pressure drops encountered in two-phase systems are higher than those resulting from single phase flow for a number of reasons. The energy change of the phase transition, the frictional energy remaining in the system due to the internal shear at the boundaries of each phase, and the reduced crosssectional area of flow for one fluid produced by the presence of a second fluid all affect the pressure drop. Consideration must also be given to the Hydraulic Energy (PV) of the fluid mass which not only changes with temperature and pressure, but also with the composition of the liquid and vapor phases.

The complex conditions of multiphase flow and the number of variables involved has been the subject to intensive survey, and relationships (21, 22, 23) have been proposed which correlate two-phase flow data. These investigations have been restricted to isothermal conditions and no attempts have been made to propose a correlation for non-isothermal systems in which there is a continuous change of phase. Therefore it was assumed that each section of pipe would behave isothermally at the average volume fractions and physical properties of the crude. The data of Reid, Reynolds, et al (22) was chosen because their investigations were conducted on pipes of similar diameter to those encountered in tube stills.

By a proper definition of terms, pressure drop calculations for two phase flow could be predicted by a Fanning type equation:

$$\Delta P_{\rm TP} = \frac{2 \ \rho LG^2}{\lambda \ g_{\rm c} D_{\rm i}} \tag{21}$$

 λ = pseudo density of the two phase mixture.

ø = pressure coefficient analogous to the friction factor
 of the Fanning equation.

In order to calculate λ , the following assumptions were made:

- The volume occupied by the liquid plus the volume occupied by the vapor, at any instant, must equal the total volume of the pipe.
- 2. The linear velocities of vapor and liquid phases are equal.
- 3. Flow is sufficiently turbulent to cause complete mixing of both phases.

From (2)
$$\frac{W}{\lambda D_{i}^{2}} = \frac{W_{L}}{P_{L}D_{L}^{2}} = \frac{W_{v}}{P_{v}D_{v}^{2}}$$

From (1) $D_{L}^{2} + D_{v}^{2} = D_{i}^{2}$

$$\frac{\mathbf{W}_{\mathrm{L}}}{\mathbf{\mathcal{P}}_{\mathrm{L}}^{\mathrm{D}}\mathbf{L}^{2}} = \frac{\mathbf{W}_{\mathrm{v}}}{\mathbf{\mathcal{P}}_{\mathrm{v}}^{\mathrm{D}}\mathbf{v}^{2}}$$

Substituting for D_v



$$\lambda = \frac{W_{L} P_{L} P_{V}}{P_{L} W_{V} + P_{V} W_{L}} + \frac{W_{V} P_{L} P_{V}}{P_{L} W_{V} + P_{V} W_{L}}$$

The liquid volume fraction (LVF) = $\frac{W_L P_v}{W_L P_v + W_v P_L}$

The vapor volume fraction (VVF) = $\frac{W_V f_L}{W_L f_V + W_V f_L}$

$$\lambda = \mathcal{C}_{L} (LVF) + \mathcal{C}_{V} (VVF) = LVF (\mathcal{C}_{L} - \mathcal{C}_{V}) + \mathcal{C}_{V}$$
(22)

Since ϕ is an empirical constant it must be obtained from a correlation of two phase data.

Reid, Reynolds, et. al. (23) have shown that for liquid volume fractions above 10 per cent, one follows the relationship:

$$\Delta P_{\rm TP} = \Delta P_{\rm L} (\rm LVF)^{-1}$$
 (23)

In this equation the following assumptions are made: (1) single phase friction factor correlating charts are applicable to two phase flow problems; (2) the superficial average densities and velocities can be employed.

 ΔP_L is the liquid phase pressure drop as calculated from single phase correlations if the liquid were flowing alone in the pipe at the same rate as the two phase flow.

If the assumption is made that $\lambda = LVF(P_L)$, a comparison of the two Equations (21 and 23) indicates that ϕ may be approximated by the friction factor of the liquid phase. For Reynolds numbers, Re>2.5 x 10⁶ Perry (24) recommends the following equation for friction factors:

 $f = 0.0014 + 0.09 \left(\frac{\mu}{DG}\right)^{27}$ Hence

 $\phi = 0.0014 + 0.09 \left(\frac{\mu}{M_T}\right)^{0.27}$ Where $0.1 \le LVF \le 1.0$ (24a)

At liquid volume fractions below 0.1 the vapor can be assumed to behave independently of the liquid. Therefore values of ϕ were approximated from the single phase friction factor correlating charts at vapor Reynolds numbers greater than 2 x 10⁵. These values varied from 0.0014 to 0.005. It was also observed from two phase data that ϕ decreases as the LVF decreases, and increases as the Reynolds' number of the liquid phase decreases.

In view of this, it was concluded that the following equation could predict sufficiently accurate values of ϕ at low liquid volume fractions

 $\phi = 0.001 \mu + LVF \left(\frac{\mu}{DG}\right)^{0.27}$ Where LVF < 0.1 (24b)

- 24 -

STATEMENT OF THE PROBLEM

Restrictions and Limitations

This investigation was initiated to study the effect of the variables of per cent excess air, center to center spacing, bridgewall and exit stack temperatures on the design of Petroleum Furnaces.

Based on a semi-theoretical equation proposed by Lobo and Evans for the calculation of heat transfer rates in the radiant section of tube still furnaces, and on an empirical equation developed by Monrad, a series of design equations were written. These equations were solved for different values of the variables and their effect on heat transfer rates and furnace dimensions studied.

The Equation of Lobo and Evans was developed for applications only to already designed or completed furnaces. However, as it was used for the actual design in the investigation, certain assumptions were necessary. These are:

1. View factors to individual tubes can be calculated using the empirical relationship of Lobo and Evans.

2. The mean beam path to the individual tubes can be approximated by the average beam length of the furnace.

The furnaces considered in this investigation were limited to the conventional box-type with the width equal to the height, single rows of tubes in the radiant section and with equal center to center spacing in the convection and radiant sections. Four tubes per row were placed in the convection bank and the length of the furnace arbitrarily fixed and equal to the length of tubes.

Calculations for the inside film coefficient of the tubes were based on the liquid phase only. The errors caused by this approximation are very slight as in most cases the outside film coefficient controls. Errors in the tube temperature calculated from this coefficient do not appreciably affect the heat transfer rates as the difference between the fourth power of the gas and

- 25 -

the surface temperatures are extremely large.

The reliability of the final design is limited by the accuracy of the physical properties data, and the validity of the assumptions used in obtaining the design equations. Perhaps the most severe of these restrictions is that encountered in the approximation of friction factors used in pressure drop calculations.
RESULTS OF CALCULATIONS

The results of the calculations are summarized in Tables III and IV. Table II gives the service requirements used in the investigation, as well as the furnace characteristics. The results presented in Table III were obtained without considering pressure drop requirements. Table IV shows the results obtained when the diameters are calculated to satisfy the requirements.

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TABLE

Service Requirements of Furnaces

Per Cent Excess Air		50	50	50	50	50	25	I
Tube Diam.	Ft.	14.0	0. ⁴ 1	0.45	0.45	0.45	0.45	ı
Tube Length	Ft.	ର	8	25	25	25	25	I
Tube Spacing Tube	Diam.	1.5	2.0	1.5	2.0	1.5	1.5	1.5
Exit Press.	Psia	I	ı	I	ı	ı	50	22
Inlet Press.	Psia	7.7גנ	7.7LL	240.0	240.0	240.0	240.0	240.0
Exit Temp.	۰ بر	670	670	750	750	300	800	800
Inlet Temp.	e ₽	014	014	400	400	400	1400	1400
Rate	Bbls/hr	133	133	ı	I	I	ı	ı
Flow	1b/hr	39 , 262.4	39,262.4	80,000.0	80,000.08	80,000.0	80,000.08	80,000.08
Furnace Group Svmbol	Unit	Ч	CJ	m	4	5	9	7

-	t																		
G	ad xlo		3.85	- - - - - - - - - - - - - - - - - - -	K0.1	4.08	5.43	6. 6	4.08	4.74	10•7			4.6	5.9	4.37	6.27	4.7	5.93
	⁴ 0		1.837	1.581	000.00	2.391	1.579	1.161	2.808	2.154	1.266			1.671	1.247	2.139	1.342	2.461	1.761
	$^{\rm R}_{\rm A}$		0.3907	0.3689	0062.0	0.3963	0.3484	0.3011	0.3895	0.3605	0.2949			0.3744	0.3316	0.3812	0.3199	0.3749	0.3362
uer Je	NC NC		11	27 14 14	S	32	9	1 17	54	5 8	36			8	76	89	74	36	1 17
lmun to f	NR R		26	2 2 2 2	t -1	2 8	ର	16	2 8	57	16			18	77	ଷ	77	ଷ	16
	Size		8.123	9##•/.	0.410	9.477	7.466	5.416	9.477	7.446	5.416			8.123	6.318	9.026	6.318	9-026	7.221
-	Press.		63 • 59	64•29	00.00	72.85	74.43	75.83	79.88	80.78	82.49			49.68	44.59	70.19	65.52	80.17	+ 77•77 +
	Temp.		00°TL9	18.6/9		671.89	678.34	667.22	671.53	670.12	666.95			674.52	674.26	668.14	669.32	669.84	672.28
ත්	x10 ⁻⁶		4.7980	026/.• †	1.1400	5.5997	5.5791	5.5256	6.4296	6.4003	6.3222			4.7980	4.7980	5.5497	5.5837	6.3966	6.44.78
ď	т. х10-6		2.4702	2.8138	2. JYEY	2.1199	2.8566	3.2630	1.8905	2.2697	3.1170			2.6777	3.1682	2.2458	3.0129	1190.2	2.6032
ď	x10 ⁻⁶		4.6606	4.4495 0011 0	v	5.0685	4.5095	3.7871	5.3092	4.8382	3.9474			₶╻┧╱५५	3.9513	4.8026	4.0437	5.0714	4.5847
Lb. Flue	Gas Per Hour		13870.7	13870.7		14914.5	14859.7	14717.3	15861.6	15789.5	15596.7			13870.7	13370.7	14781.4	14781.4	15780.3	15906.5
Bridge	Wall Temp.	Ч	1416.04	1494 • 67	11.0/OT	1395.58	1554.33	1646.57	1417.78	1495.73	1671.58	٩	' 	1463.61	1575.14	1426.84	1586.82	1454.11	1558.48
	Stack Temp.	Furnace	830	830 020	000	930	930	930	1030	1030	1030	Furnace		830	830	930	930	1030	1030

.

TABLE III

4 2A _{cp}	6.02 5.63 5.56	5.0 5.0 66 66 7.6 7.6 7.6 7.6 7.6 7.6 7.6 7.6 7	5.56
୶ୗ୶ୖ	1.44 1.892 2.460	2.022 1.508 2.636 1.926 3.803 2.529	3.39
$^{ m R}_{ m A}$	0.3412 0.3662 0.3858	0.3369 0.3481 0.4036 0.4288 0.4288 0.3865	0.3899
R S C S C S C S S S S S S S S S S S S S	8885	25 40 86 25 40 86	36
Numl to 1 NR R	45 0 4 10 0 7 10 0	846 35 34 35 36	38
S1ze	14.038 14.038 15.516	19.703 14.77 20.668 15.76 22.66	18.72
ditions Press.	128.56 143.96 155.85	90.43 92.77 128.62 129.52 147.57 152.47	163.35
Exit Cor Temp	753.20 753.26 755.72	755.93 757.31 751.21 756.43 754.88	761.05
ALO-6	1.5003 1.6266 1.7545	1.5003 1.5003 1.6266 1.6266 1.7545	2.0729
ور ×10 ⁻⁶	8.4268 7.1333 6.0147	6.8074 8.2264 7.0836 4.3152 5.8139	4.8174
е _в × 10 ⁻⁶	1.2135 1.3495 1.4798	1.3765 1.2404 1.4806 1.3640 1.6409	1.6329
Lb. Flue Gas Per Hour	42648.8 42648.8 42648.8	42648.8 42648.8 42648.8 42648.8 42648.8 42648.8	46931.4
Bridge Wall Temp.	ce 3 1496.59 1494.34 1506.17	ce 4 1375.82 1481.73 1381.23 1490.65 1379.5 1491.65	1430.67
Stack Temp.	Furna 850 950 1050	Furnau 850 950 1050 1050	0 <u>5</u> 11

TABLE III (continued)

40 x10-4		3. 94 6.98	3. 86 6.8	3. 9 6.96		3.68 4.15	3.36 4.45	3.28 4.39
gf g		2.43 1.36	2.99 1.58	4.06 1.88		2.63 2.6	4.14 2.99	: :
$^{ m R}_{ m A}$		0.4345 0.3539	0.4230 0.3461	0.4404 0.3548		0.438 0.475	0.51 ⁴ 8 0.48	0.5073 0.4674
NC Ser		56 7 2	5 29 20	5 8		5 2 56	36 #4	24 32
Numl of Tul R		378	96 38	3 3 33 33		5 2 58	% S	54 24
Size		22.166 13.3	23.643 14.78	23.643 14.78		22.166 19.21	24.38 19.21	25. 86 20. 69
ditions Press.		55.81 108.68	83.25 125.59	104.89 144.71		86.82 94.34	96.63 122.32	104.30 132.95
Exit Con Temp.		808.07 306.64	80 2.60 805.80	805.89 800.50		800.99 808.61	805.65 803 .1 6	805.45 803.17
er x10-6		1.5003 1.5003	1.8176 1.8382	1.9605 1.5827		1.2397 1.2397	1.3443 1.3443	1.5951 1.6087
ac x10-6		6.9533 1.0075	5.9169 9.2273	4.7119 8.1738		5.9672 6.6423	4.6252 6.0090	3.645 5.1182
an Role		1.6368 1.3737	1.7662 1.4614	1.9138 1.5399		1.7477 1.7283	1.9168 1.7981	2. 0178 1.3609
Lb. Flue Gas Per Hour		42648.8 42648.8	47655.9 48196.6	47655.9 48196.6		361 23. 8 361 23. 8	361 23. 8 36123.8	39731.8 40069.2
Br1dge Wall Temp.	5	1386.76 1618.03	1357 . 10 1570.55	1372.12 1595.98	9	1404.46 1465.63	1378.28 1502.57	1355.50 1472.61
Stack Temp.	Furnace	850 850	950 950	1050 1050	Furnace	850 850	950 950	1050 1050

TABLE III (continued)

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		$^{ m R}_{ m A}$		0.487	0.47	0.433	0.428	0.344	0.464	0.291	0.345	0.209	0.208
	Tube Ma.	Ft.		0.4128	0.4115	0.4100	0.4157	0014.0	0.4281	0.4184	0.4219	1204.0	1614.0
nber Je	thes	NC N		52	56	5	56	72	6 1	5	52	80	72
μ Ν	́ н́	$^{\rm N}_{ m R}$		64	54	42	64	36	₽	5	76	32	43
		Size		22.166	18.403	14.22	22.166	12.864	22.166	18.003	27.794	10.35	15.89
	ditions	Press.		19.13	60.86	6.47	43.32	69.15	55.0	61.98	0.04	37.13	54.0
	Exit Cor	Temp.		17.108	805.93	809.07	802.38	800.15	804.35	800.23	41.867	801.44	803.8
e = 850 ° F	، م ^ل)01x		1.2397	1.2397	1.2397	1.500	1.500	1.500	2.1508	2.1497	2.1 497	2.75988
emperatur	، م	7-01x		5.8314	6.60916	8.0838	6.8156	1.00313	5.6239	9.87721	7.6223	1.38224	1.2835
t Stack T	с <mark>н</mark>	,		1.7304	1.68906	1.56244	1.6351	1.31014	1.78375	1.28784	1.5427	9.3180	1.06017
Ext	Lb. Flue Gas	Per Hour		36123.8	36123.8	36123.8	42648.7	42648.7	42648.7	64669.5	64635.6	64635.6	83354.9
	Bridge Wall	Temp.		1392.7	1462.66	1593 . 9	1376.4	1614.8	1286.5	1386.4	1266.7	1593.4	1391.8
	Per Cent Excess	Air	Furnace 7	25	25	25	5	ß	20	100	100	100	125

TABLE IV

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- 32 -

DISCUSSION OF RESULTS

The design of tube still heaters should be made by considering distributions of the heat load between both sections of the furnace. When a distribution ratio has been chosen, the design of a petroleum furnace can then be established on a basis of the permissible average radiant heat transfer rate.

The results illustrated in Figures IV and V show the heat distribution ratio to be a function of bridgewall temperature; exit stack temperature; and furnace capacity. These curves include the results obtained using two tube spacings, indicating that the heat distribution ratio is independent of this variable.

Figure VI is presented in order to permit visualizing the effect of excess air on the distribution ratio. At a specific bridgewall temperature a greater percentage of heat can be distributed to the radiant section by decreasing the percentage excess air. Using the bridgewall temperature as a parameter, this effect was correlated in terms of the fraction of the total heat input absorbed in the radiant section, and plotted versus the per cent excess air. These results are shown in Figure VII.

In the preceeding discussion, emphasis is placed upon available distributions of the total heat load between both sections of the furnace without considering the transfer rates per unit area of cold surfaces. It is quite difficult to generalize regarding allowable rates of heat transfer as this would naturally depend upon the rate at which the oil removes heat from the tubes and the maximum temperature to which the tube may be heated without causing corrosion, distortion of the tube, thermal cracking of the crude and coke deposition inside the tubes. It should not be concluded however, that a choice of the maximum allowable rate results in the best furnace design, as this maximum may only be attainable in one section of the furnace. The choice, should be based upon rates in both sections which will yield the lowest total tube surface area.

- 33 -

The results of these calculations indicate that this situation is accompanied by high bridgewall temperatures and small quantities of excess air.

In the interests of economy, it is the current trend to fill the radiant section with cold tube surfaces. Figure VIII gives an indication of the relationship between the cold plane area in the radiant section and the fraction of the total heat input absorbed in this section. These curves show that this fraction, although independent of the tube spacings used, varies with excess air and furnace capacity. These curves are also significant in that they give an indication of the optimum size of the radiant section. For example, if the maximum allowable rate of heat transfer was found to be 20,000 Btu/ft² for Furnace 1 (illustrated by curve 1 in Figure VIII), then a design could not be made such that the fraction of the total heat absorbed in the radiant section is less than 0.32. Also, if the permissible rate were close to the maximum, for example 18,000 Btu/hr ft² then the resulting small radiant section would be obtained at a loss of economy. This, however, can be avoided by increasing the per cent excess air and thus increasing the fraction absorbed for a specific radiant rate.



Fridge-Wall Temperature *F



- 36 -

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Brid e-Wall Temperature

- 37 -





Percentage Excess Air



- 39 -

CONCLUSIONS AND RECOMMENDATIONS

Accurate furnace designs can be obtained using mathematical models similar to that employed in these calculations. Although these models are useful in determining the effect of various phenomenon in furnace characteristics, they cannot be used to directly determine optimum designs. In order to obtain the best design for a furnace, models must include conomic considerations as well as those factors included in this model.

The factors most significant in determining the ultimate design of furnaces are the distribution of total heat load to the radiant and convection section, and the average heat transfer rates.

The heat distribution ratio of a furnace is dependent upon the per cent excess air and bridgewall and exit stack temperatures. Increases in any of these characteristics will result in a larger distribution ratio.

Average heat transfer rates in the radiant section are almost entirely dependent upon the fraction of the total heat input absorbed in this section and increases as the fraction decreases. This fraction is dependent on the per cent excess air and bridgewall temperatures and will increase as either variable decreases.

The total tube area required for a furnace with a specific duty is usually lowest when the bridgewall temperature is high and the per cent excess air is low.

The number of calculations made throughout this investigation was severely restricted by the size of the computer available and the type of programming employed. Should a similar investigation be attempted, it should be conducted on a larger computer using a faster method of interpretive programming. It is also recommended that a more accurate method of predicting pressure drop in the presence of two phase flow be obtained. APPENDIX ES

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NOMENCLATURE

		0
A _c	=	total outside tube area, ft ² .
A _R	=	effective refractory area, ft ² .
A' _R	=	actual refractory area, ft ² .
A _T	=	total wall area, ft ² .
A cp	=	area of plane replacing tubes, ft ² .
C	=	heat capacity, Btu/lb °F.
D	=	pipe diameter, ft.
D'	=	pipe diameter, ins.
DL	=	equivalent diameter: of liquid phase.
Dv	=	equivalent diameter of vapor phase.
ď	=	fraction distilled.
Е _р	=	black body emissive power, Btu/hr ft ² .
F	=	geometric factor.
Ψ	=	overall exchange factor.
f	=	Fanning friction factor.
G	=	superficial mass velocity, lb/hr ft ² .
H	=	enthalpy, Btu/lb.
Δн	=	heat absorbed by crude per node, Btu/lb.
h _c	=	pure convection coefficient, Btu/hr ft ² °F.
h _r	=	radiant coefficient, Btu/hr ft ² °F.
ĸ	=	thermal conductivity of crude, Btu/hr ft ² [•] F/ft.
ĸ	=	thermal conductivity of tube, Btu/hr ft ² $^{\circ}$ F/ft.
\mathtt{L}_{B}	=	mean length of radiant beam, ft.
L	=	length of radiant section, ft.
ΔL	=	length of tube node, ft.
Р	=	pressure, lb/in ²
P _{CO2}	=	partial pressure of CO2, atms.
P _{H2} 0	Ξ	partial pressure of H_2^0 , atms.
q _	=	heat transferred to the oil, Btu/hr.
S.G	Ħ	specific gravity.
т	=	temperature, [•] R.
t	=	temperature, °F.

•

u	=	overall heat transfer coefficient, Btu/hr ft 2 °F.
W	=	mass flow rate, 1b/hr.
Z	=	height of furnace, ft.
α	Ξ	factor by which A must be reduced to obtain effective
		cold surface, αA_{cp}^{or} (effective tube area).
€ _G	=	gas emissivity.
ø	=	pressure drop coefficient.
μ	=	viscosity, lb/ft hr.
λ	=	pseudo density of two phase mixture.
¢-¢	=	center to center spacing of the tubes.
δ	=	Stefan-Boetzmann constant, 0.173 x 10 ⁻⁸ Btu/hr ft ² •R ⁴
ρ	H	density, lb/ft ³ .

Subscripts

.

Ъ	=	black body.
с	=	convection.
g	=	gas •
i	=	inside.
m	=	arithmetic average.
0	=	outside.
r	=	radiation.
ß	=	surface.
W	=	wall.
1	=	inlet.
2	=	exit.

PROGRAM ABSTRACT

TITLE: Pipe Still Heater Design. AUTHOR: William V. Saunders.

DESCRIPTION

The program calculates the dimensions, number of tubes and their diameter in the radiant and convection sections of the furnace. The method used includes the correlations of Lobo and Evans for the evaluation of radiant heat transfer rates and those of Monrad for calculating convection coefficients.

COMPUTER

MISTIC, 1024 cathode-ray tube memory locations, perforated tape input and output.

PROGRAM LANGUAGE

Fixed point and floating point coding.

RUNNING TIME

Six to ten hours depending on the accuracy of initial estimates of the guessed quantities.

COMMENTS

The engineer can have the calculation stop after any of the two sections: radiant section, convection section. Program has been successfuly used over forty times in designing furnaces.

AVAILABILITY

A manual for the description of the codes used in this program is available in the Computer Laboratory library at Michigan State University. This program is available from the Computer Library in the Chemical Engineering Department.

DESCRIPTION OF PROGRAM

Description

To handle the lengthy calculations involved in designing the tube-still furnace, a procedure was developed for use with a small sized digital computer. The machine routine is such that a choice of tube length, flow rate of crude stock, centerto-center spacing of the tubes, percentage excess air, and bridgewall and exit stack temperatures may be varied in considering the different designs.

The furnace may be calculated in any increment of tube length desired; however, as the computer routine used was an extremely slow one, it was necessary to shorten the calculations as much as possible in order to have expediency of calculation time. Such being the case, an increment of four tube lengths was chosen in the convection section and two tube lengths in the radiant section.

The flow of calculations around a tube increment is shown in Figure 9. At the inlet of the tube, the temperature, pressure and liquid volume fraction, (t, T, P, LVF) is known from the previous tube, or if the first calculation from the inlet conditions to the furnace. The inlet conditions (t_2 and P_2) are guessed and their arithmetic averages computed. Based on these average values, new values of t_2 and P_2 are calculated.

The calculated values of the outlet temperature and pressure are compared with the assumed values. If the difference between calculated and assumed values are not within tolerance, new values for the exit conditions are chosen and the procedure repeated until the outlet values are within tolerance. These values are then used for the next tube.

When the gas temperature above the tube reaches a certain maximum, which is set as the highest allowable bridgewall temperature, calculations for the convection section are stopped and calculations for the radiant section commences. When the exit temperature and pressure from a tube increment compares

- 47 -



FICRE 9. Calculations Flow Around one Care Encrement

with the desired discharge conditions, the dimensions of the furnace are calculated from the number of tubes in this section and compared with the assumed dimensions. If the discharge conditions do not check within tolerance a new diameter is assumed and all calculations repeated. Calculations within the radiant section are repeated when the assumed dimensions do not compare with those calculated.

Finally a heat balance is made about the furnace and the fuel rate modified until the heat liberated equals the heat absorbed plus the heat lost from the furnace.

Method of Fitting Equations to Data

The numerous charts necessary to the solution of this problem had to be programmed so that they could be interpreted by the computer. The most convenient and accurate method of doing this is by fitting the curve to polynomials.

Polynomials were obtained by reading a set of points from each curve and by fitting these points to a polynomial of the form

$$F(x) = 1/2 \sum_{s=0}^{n-1} A_s x^s$$

The criterion of excellence for each polynomial was that the sum of the squares of the deviations

$$M = \sum_{i=0}^{n-1} [F(X_i) \cdot f(X_i)]^2 A_w(X_i)$$

 $f(X_1) = points from the curve$

Should be a minimum with respect to arbitrary variations of the coefficients $\mathbf{A}_{_{\!\!\mathbf{G}}}$.

Two library routines, K_3 and L7S, were available for the evaluation of these polynomials. The constants, as obtained for the curves used in these calculations are presented in Table 5.

MACHINE REQUIREMENTS

The automatic computer used for the furance calculations was MISTIC, a binary, fractional, single address computer with a word length of 40 bits, a memory of 1024 words, and which puts two instructions in a word.

MISTIC like most digital computers is composed of five units: input, memory or storage, control, arithmetic, and output.

Input

This unit includes as an input medium a 5-level perforated Teletype paper tape by which the problem is communicated to the computer, and as an input device a photo-electric reader which is able to pull the tape past a light. The light shines through the holes activating a photo-sensitive surface which converts the spots to electrical impulses, equivalent to the number represented by the coded character on tape. These impulses are sent directly to an assembly register.

Memory

The numbers received by the assembly register are sent under the control of the computer to memory. The memory is an electronic device composed of 40 vacuum tubes. On the face, or gird, of each tube 1024 spots can be stored. Each spot is assigned a certain address or location, and corresponds to a binary digit from every word. Each tube represents a different binary digit, that is, a power of two (referred to as a bit) to comprise a total of 40 bits, or one word.

Information stored in memory usually consists of a series of instructions (a program), directing the machine to execute certain operations and a set of data on which these operations are performed.

Arithmetic

The instructions and data held by memory are used both by the arithmetic and control sections. The arithmetic section is that section of the computer in which mathematical and logical operations are performed. These operations include: addition, subtraction, multiplication, division, and the comparison of two quantities and are associated with three ⁴O bit registers and an adder. These are the A register (accumulator), the Q register (quotient), the R register and the adder. In an addition, A holds one of the terms and the result, Q holds the quotient

- 50 -

of a division and has no additive properties. The adder is a register on which the number in A is added to the number in R.

Control

The control section of the computer directs every operation executed by the computer. MISTIC stores the program provided by the operator in memory, and before any operation can be executed it must be summoned from memory by the control section. This section is comprised of an instruction register (IR) which holds the instruction currently being executed by control and a control counter which holds the address of the next instruction pair to be sent to IR.

Output

When the result of a calculation has been obtained which must be transferred to the operator, it is sent to the output section. The output section is similar to input in that electrical impuses are converted to a suitable code and translated by an electro-magnetic devise and presented to the reader on tape.Teletypewriters are available and are used for printing on paper the symbols represented by the punched tape.

Subroutines

MISTIC operates in binary, and in view of this, special attention must be given to the location of the decimal or binary points of numbers before every arithmetic operation. Its arithmetic unit necessitates the binary point to be fixed so that any number X used in computation must be in the range $-1\le X\le 1$. It is necessary then, that each number at every stage of a calculation be scaled within the capacity of the Machine. Many problems encountered in engineering calculations involve numbers of various magnitudes making scaling an additional inconvenient complexity. For these calculations floating point routines may be used. These routines represent numbers as $x = a \times 10^b$ and store a and b. Thus they can represent numbers in the range of $10^{-63} \le x < 10^{63}$. The floating point routine used in the preparation of this program was a standard library routine designated as A1.

Running Time and Accuracy

The running time for a complete computation cannot be predicted exactly as this would depend on the accuracy of the initial estimates of the guessed quantities and the number of increments chosen. The running time necessary for an increment of tube length to converge was approximately three minutes and that for a complete computation averaged to about five hours. A conservative estimate of the time required to perform a complete computation by hand, would be 110 hours. The accuracy for each arithmetic operation is set by that of the Al Routine which provides an answer to at most nine decimal digits. However, the elaborate trial and error computation involved necessitated a choice of limits-of-convergence which resulted in an accuracy within 0.8 per cent of the correct answer for the exit temperature and 1.5 per cent for the exit pressure for each tube increment.

Error Stops

During the necessarily lengthy calculation periods, the operator was kept informed of progress and possible errors by special features included in the program. At the end of each calculation cycle, print outs showed; (1) dimensions of the furnace, (2) difference between the desired and calculated pressure drops and also tube diameter, and (3) number of tubes in the radiant section and the fuel rate. Obvious errors in any of these values could easily be detected however, the less obvious mistakes could only be detected in the final answers.

Diagrammatic Flow Chart

The backbone of the automatic computation system consists of two separate but complementary computer programs for the solution of the design problem. The first program treats the convection section as an independent unit and makes calculations around tube increments until the gas temperature above an increment reaches a certain maximum. This maximum is set by the bridgewall temperature desired in the radiant section. The second program utilizes this temperature and the exit conditions from the convection section to solve the problem of furnace dimensions, number of tubes required, tube diameter and also the heat balance about the furnace.

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The calculation procedure is diagrammed in Figure 10. With the exception of the preliminary calculations (box 1), each step in the diagram corresponds to a series of calculations listed in the smaple problem given on the following pages. In order to shorten machine running time as much as possible certain quantities were evaluated at the beginning of the program to avoid repetition during any cycle for which these values do not change. These are the preliminary calculations and include the evaluation of $1.6G^{1/3}/D'_{0}$ from the convection coefficient, Equation (8); the mean beam path of the gases between banks of tubes, Equation (5) and the ratio A_v/A_t .



FIGURE 10A. Flow Diagram. Convection Section.



FIGURE 107. Flow Diagram. Convection Section.



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FICHE 10C. Flow Diagram. Radiant Section.



FIGURE 10D. Flow Diagram. Nadiant Section.

SAMPLE CALCULATIONS

The following computations serve to illustrate the flow of calculations around a tube node in the convection and radiant sections. It should be clearly understood that this illustration represents only a small fraction of the computations necessary to complete a solution of the furnace design problem.

Furnace data:

Inlet temperature = $400 \, {}^{\circ}$ F. Inlet pressure = $240 \, {}^{\circ}$ Jb/in². Stack temperature = $870 \, {}^{\circ}$ F. Tube spacing = 1.5 tube diameters. Fuel - CH₄ fired with 50% excess air.

Crude data:

Flow rate = 80,000 lb/hr.

API gravity = 35.7

True boiling point and mid-gravity curves are presented in Figure 11. From the true boiling point curve, a flash curve is drawn at 238 lb/in^2 and is shown in Figure 12.

Calculations

Tube increment = 100 ft.

(2) Assume:

Exit temperature = $409.517 \, {}^{\circ}F.$ Exit pressure = $236.29 \, {}^{\circ}F.$ Tube diameter (inside) = $0.41 \, \text{ft.}$ $\therefore t_m = 404.76 \, {}^{\circ}F; P_m = 238.15 \, \text{lb/in}^2.$

(3) Calculate per cent vaporized: $d = \frac{t - t(b)}{100m}$ (equation for the flash curve) t(b) = t(p) - 50m

 $t_{(p)}$ is the vapor pressure-temperature equation for a hydrocarbon with an atmospheric boiling point corresponding to the 50% point of the flash curve.

 $t_{(p)} = 426.58 + 3.2P - 0.0113 P^2 + 15.4 \times 10^{-6P3}$ Solving for $t_{(p)}$ at $p = 238.15 lb/in^2$, $t_{(p)} = 759.74 °F$ $\therefore t_{(b)} = 379.74 °F$. Since $t_{(b)} < 404.76$, vaporization has occurred. $d = (404.76 - 379.74) \div (100) (7.6) = 0.033$ This result can also be obtained by reading the fraction distilled from the flash curve Figure 12. (6) Calculate the liquid volume fraction (L.V.F.):

$$L \cdot V \cdot F \cdot = \frac{W_{L}}{W_{L} + \frac{W_{V} \rho_{L}}{\rho_{y}}}$$

$$\begin{bmatrix} \frac{W_{L}}{\rho_{L}} \end{bmatrix}$$

$$\begin{bmatrix} \frac{W_{L}}{\rho_{c}} \end{bmatrix}$$

$$\begin{bmatrix} \frac{W_{C}}{\rho_{c}} \end{bmatrix}$$

$$\begin{bmatrix} \frac{W_{C}}{\rho_{c}} \end{bmatrix}$$

$$W_{c} = 80,000 \text{ lb/hr}$$

$$(SG)_{c} = 0.8474$$

The specific gravity of the liquid can be obtained from the midgravity curve (Figure 11), or from the following polynomial at d = 0.033.

$$(SG)_{L} = 0.599 + 0.965d - 1.487 d^{2} + 1.017 d^{3} = 0.63$$

$$\therefore W_{L} = 57,450 lb/hr \qquad W_{V} = 22,550 lb/hr$$

$$\rho_{V} = \frac{1}{V} = 1728 / \left\{ \frac{157T}{P} - \frac{1}{T^{3}} \left[7234 + 102P lx 10^{7} + 0.52P + 20 \right] \right\}$$

$$T = 864.76 ^{\circ}R \quad P = 238.15 lb/in^{2} \therefore \rho_{V} = 2.6 lb/ft^{3}$$

$$(SG)_{L} = (SG)_{L}^{\circ} + (\frac{60-t}{1.8}) \left\{ \alpha_{T} + \frac{\beta_{T}}{1.8} \left[3t - 217 \right] \right\}$$

$$\alpha_{T} = 0.00122 + 0.00276 (SG)_{L}^{\circ} - 0.0079 (SG)_{L}^{\circ^{2}} + 0.0046 (SG)_{L}^{\circ^{3}}$$

$$\beta_{T} = 1.92307 \times 10^{-6} (SG)_{L}^{\circ} - 1.5538 \times 10^{-6}$$

$$(SG)_{L}^{\circ} = 0.63$$

$$\therefore \rho_{L} = (SG)_{L} 62.4 = 29.95$$

$$LVF = \frac{W_{L}\rho_{L}}{W_{O} + W_{O}}$$

$$LVF = \frac{U U}{W_L \rho_V + W_V \rho_L}$$

... LVF = 0.018

Calculate **∆**H:

$$H_{L} = 15 (SG)_{L} - 26 + [0.811 - 0.465 (SG)_{L}]t = 0.00029t^{2}.$$

$$H_{V} = [215 - 87 (SG)_{L}] + [0.415 - 0.104 (SG)_{L}]t + [3.1 - 0.78 (SG)_{L}]t^{2} \times 10^{-4} + \frac{1}{9331.7} \left\{$$

$$[20.47P + 0.258P^{2} - \frac{P}{T^{3}} (28.970 + 203.4P)x10^{7}] \}$$

$$\Delta H = [W_{L}H_{L} + W_{V}H_{V}]_{2} - [W_{L}H_{L} + W_{V}H_{V}]_{1}$$

$$t_{1} = 400^{\circ}F; P_{1} = 240 \text{ lb/in}^{2}; t_{2} = 409.517^{\circ}F; P_{2} = 236.29^{\circ}F.$$

$$\Delta H = 5.35 \times 10^{5} \text{ Btu/hr.}$$

(7) Calculate T₀. (Gas temperature above node)

$$\Delta H = \int_{T_{1}}^{T_{2}} C_{p}dT$$

$$c_{p} = 1.9x10^{4} + 1.7T - 9.12x10^{-5}T^{2} \text{ Btu/}^{\circ}R$$

$$5.35x10^{5} = 1.9x10^{4}[T_{2} - 1330] + 0.85[T_{2}^{2} - (1330)^{2}]$$

$$- 3.04x10^{-5}[T_{2}^{3} - (1330)^{3}]$$

- 60 -

Newton's method is used to obtain the nth approximation of T_2 , $T_n = T_{n-1} - \frac{f(T)}{f'(T)}$ where $f(T) = \int_{T_1}^{T_2} C_p dT - \Delta H$

Following this procedure, T_2 is solved for by trial and error. $T_2 = 923$ °F. (9) Calculate q(= UA Δ t)

$$A = \pi D_{o} \Delta L; \quad \text{ft}^{2}; \quad \Delta t = [\text{tg - tm}]$$
$$1/U = \frac{1}{hu} + \frac{D_{o}(D_{o}-D_{i})}{(D_{o}+D_{i})K_{m}} + \frac{D_{o}}{D_{i}h_{i}}$$

Km, the thermal conductivity of the tube, is assumed independent of temperature. Km = 26 Btu/ft² $^{\circ}F/ft$. . hu = $\frac{1.6 \ G^{2/3}Tg^{0.3}}{D_{o}^{1/3}}$

G = (gas flow rate)/(minimum cross sectional area) $= 8.34x10⁴/LD_o($$\empsilon - $$\empsilon - $$\mathbf{l}$] area)$ $D'_o = D_o/12; Tg = (923 + 870)/2 + 460 = 1356.5 °R.$... hu = 7.8 Btu/hr ft² °F. $<math display="block">\frac{h_i D_i}{K_L} = 0.023 \left(\frac{D}{\mu}\right)^{0.8}_L \left(\frac{C_p \mu}{K}\right)^{1/3}_L$
- 61 -

$$K_{\rm L} = \frac{0.813}{12(SG)} [1 - 0.003(t-32)] = 0.095$$
 Btu/hr ft² °F/ft
 $C_{\rm p} = \frac{\Delta H}{\Delta t} = 0.75$ Btu/lb °F.
(kinematic viscosity) = 5.995 - 0.023t + 3.93x10⁻⁴t² - 2.32x10⁻⁶t³

 μ_{K} (kinematic viscosity) = 5.995 - 0.023t + 3.93x10 t⁻ - 2.32x10 t⁻ $\mu = \mu_{K}$ (SG)(2.42) = 1.0T lb/ft hr. Solving:

$$h_{i} = 806.9 \text{ Btu/hr ft}^{2} \, ^{\circ}\text{F}$$

$$u = 7.66 \quad \text{Btu/hr ft}^{2} \, ^{\circ}\text{F}.$$

$$q = 7.66 \, (\pi)(0.466)(100)[896.5 - 404.76]$$

$$q = 5.5 \text{xlo}^{5} \text{ Btu/hr}.$$

Since q and ΔH are approximately equal, the assumed t_2 is correct. However, if $q \neq \Delta H$, within tolerance, another value of t_2 would be calculated using the following relationship:

$$t_2^1 = t_2 + (q - \Delta H)(t_2 - t_1)/(\Delta H).$$

and all calculations repeated.

(16) Calculate
$$\Delta P$$
.

$$\Delta P = \frac{2 \not o \Delta L G^{2}}{\lambda q_{c} D_{1}(144)}$$

$$\lambda = LVF (\rho_{L} - \rho_{V}) + \rho_{V} = 0.18(29.95-2.6) + 2.6 = 7.52 \text{ lb/ft}^{3}$$

$$\phi = 0.0014 + 0.09 \left(\frac{\mu}{DG}\right)_{L} = 0.0049$$

$$\Delta L = 100 + equivalent length of 4 return bends = 220 \text{ ft.}$$

 $\Delta L = 100 + equivalent length of 4 return bends = 220 ft. Solving$

 $\Delta P = 4.2 \text{ lb/in}^2$

Since $P_2 + \Delta P \doteq P_1$ the assumed pressure is correct.

For $P_2 + \Delta P \neq P_1$ repeat calculations using $P_2^{l} = P_1 + \Delta P$.

(10) Shield Section

For tube nodes in the shield section

$$h_{u} = \frac{(100 + \% \text{ wall effect})}{100} (h_{c} + \text{hrg}) \qquad \text{Btu/hr ft}^{2} \bullet F.$$

$$h_{c} = \frac{1.6 \text{ G}^{2/3} \text{Tg}^{0.3}}{\text{D}_{o}^{1/3}}$$

hrg =
$$\epsilon_{s} \left[\frac{(q_{c}+q_{w})_{Tg} - (q_{c}+q_{w})_{Ts}}{(Tg-Ts)} \right] \frac{(100-\%)}{100}$$

% wall effect = $\frac{\text{hrb x Aw x 100}}{[\text{hc + hrg = hrb]At}}$ hrb = 0.00688 $\epsilon_{s} \left[\frac{T}{100}\right]^{3}$ and $\epsilon_{s} = 0.95$

This necessitates a trial and error solution since the tube surface temperature, t, is unknown. A solution is obtained by approximating h_u and calculating t_s .

$$t_{s} = t_{m} + \Delta t_{s}$$
$$\Delta t_{s} = \left[\frac{D_{o}/\text{Dihi} + D_{o}(D_{o}-D_{i})/(D_{o}+D_{i})\text{Km}}{1/h_{u} + D_{o}/\text{Dihi} + D_{o}(D_{o}-D_{i})/(D_{o}+D_{i})\text{Km}}\right] \Delta t$$

$$\Delta t = tg - tm$$

From t_s , h_u is calculated and compared with the assumed value. If their difference is small, the calculated h_u is used as the correct value. If their difference is not negligible, another value of h_u is assumed and the process repeated.

When a tube node has converged, the calculated values of t_2 and P_2 are used as inlet conditions to the adjacent node and calculations repeated until Tg is within the temperature range assigned to the bridgewall temperature. At this point, calculations for the radiant section begin.

The method of evaluating exit conditions from tube increments in the radiant section is similar to that employed in the convection section. However, q, previously the heat transferred by convection, must be replaced by q_r , the heat transferred to the oil by radiation. The evaluation of q_r is illustrated in the following example.

Length of tube = 25 ft. Length of node = 50 ft. Bridgewall temperature = 1391.84 °F. Inlet temperature = 797.83 °F Inlet pressure = 51.35 lb/in² Per cent excess air = 125%.

(1) Assume:

Dimension of radiant section = $16 \times 16 \times 25$ ft.

Exit temperature = 303.85 °F.
Exit pressure = 53.99 lb/in².
(2) Calculate (previously illustrated).
vaporized = 79.42 #.

$$\rho_{L} = 34.086 \ lb/ft^{3}.$$

 $\rho_{V} = 0.497 \ lb/ft^{3}.$
 $W_{V} = 6.187 \times 10^{4} \ lb/hr.$
 $W_{L} = 1.822 \times 10^{4} \ lb/hr.$
 $W_{F} = 0.00429$
 $\Delta H = 4.0839 \times 10^{5} \ Btu/hr.$
 $K_{L} = 0.0556 \ Btu/hr \ ft^{2} \ ^{\circ}F/ft.$
 $\mu_{L} = 0.48 \ lb/ft \ hr.$
 $C_{p} = 0.83 \ Btu/lb \ ^{\circ}F.$
(3) Calculate t_{g}
 $\Delta H = u \ \pi \ D_{0} \Delta L \ [ts-tm]$
 $1/u = 1/h1 + D_{0} (D_{0} - D_{1})/(D_{0} + D_{1}) Km$
 $\frac{h1D1}{K_{L}} = 0.023 \ (\frac{DG}{\mu})^{2}L \ (\frac{Cp\mu}{K})^{1/3}$
 $D_{0} = 0.46 \ ft.; D_{1} = 0.42 \ ft.$
Solving $t_{g} = 815 \ ^{\circ}F.$
(4) Calculate q_{r}
 $q_{r} = \alpha A_{cp} \Psi \left\{ 0.173 \times 10^{8} \ [Tg^{4} - Ts^{4}] + 7 \ [Tg-Ts] \right\}$
 $L_{B} = 2/3 \ [25 \times 16 \times 16]^{1/2} = 12.28 \ ft.$
 $P_{CO_{2}} = 0.0466 \ atms. P_{H_{2}O} = 0.0392 \ atms.$
 $P_{L}(CO_{2}) = 0.5493; \ PL(H_{2}O) = 1.099 \ atm. \ ft.$
A tube spacing of 1.5 tube diameters corresponds to
 $\alpha = 0.07$ (Figure 3)

ponds to $\alpha = 0.97$ (Figure 3) $\therefore \alpha A_{cp} = \alpha \quad L \notin D_{o} = 33.5 \text{ ft}^2$

ft.

$$\epsilon_{\rm G} = \left[\left(q_{\rm c} + \frac{q_{\rm w}}{{}_{\rm Tg}} - \left(q_{\rm c} + q_{\rm w} \right)_{\rm Ts} \right] \left[\frac{100 - \%}{100} \right] \quad \text{Btu/hr.}$$

To obtain q_c and q_w at their respective PL values, it is necessary to interpolate between two polynomials. The following interpolation formula is used.

$$f(x) = f(x_1) + \frac{P - P_1}{P_2 - P_1} [f(x_2) - f(x_1)]$$

 $P_1 < P < P_2$ and $f(x_2) > f(x_1)$.

 $P \neq \text{desired parameter (PL value), } f(x_1) \text{ and } f(x_2) \text{ are the polynomials corresponding to parameters } P_1 \text{ and } P_2 \text{ respectively.}$

Emission due to CO₂ molecules:

 $2q_{e} (at PL = 0.4) = 4823.29 - 12.153t + 0.0097198t^{2} - 7.351 \times 10^{5}t^{3}$ $2q_{e} (at PL = 0.6) = 7588 - 18.258t + 0.013484t^{2} - 1.1916 \times 10^{-4}t^{3}$ Emission due to H₂0 molecules: $2q_{w}(at PL=1.0) = 622.9 - 1.6958t + 0.0043397t^{2} + 1.848 \times 10^{-4}t^{3}$ $2q_{w}(at PL=1.25) = 2283.7 + 5.2497t - 0.000627t^{2} + 3.4113 \times 10^{-4}t^{3}$ Black body radiation: $2q_{b} = -5926 + 27.632t - 0.03172t^{2} + 2.550 \times 10^{-3}t^{3}$

Polynomials are also available for the percent correction. These polynomials are identified by the parameter $S = P_{cL} + P_{wL}$, and evaluated as functions of $R = CO_2/(H_2O + CO_2)$. Interpolation is again necessary.

$$P_{cL} + P_{wL} = 1.648$$

$$R = \frac{CO_2}{CO_2 + H_2O} = 0.334$$

$$2(\%), (at S=1) = 2.5018 + 49.2874R - 92.454R^2 + 102.165R^3 - 52.9670R^4$$

$$2(\%), (at S=2.0) = 4.232 + 51.983R - 101.885R^2 + 111.444R^3 - 54.630R^4$$
Solving for ϵ_{g} :

$$\epsilon_{g} = 0.394.$$

$$A_{R}/\alpha A_{cp} = \frac{2Z(Z + 2L)}{\alpha L(2Z + f)} - 1$$

If tubes are placed on the Bridge wall, f = 2/3Z. If not, f = 0. The first iteration through the radiant section is made with f = 0. After convergence to the correct discharge temperature, f is calculated from the number of tubes in this section. The magnitude of f determines whether tubes should be placed on the bridge wall.

$$A_{\rm R}/\alpha A_{\rm cp} = \frac{2 \times 16 [16 + 2 \times 25]}{0.97 \times 25 [2 \times 16 + 0]} - 1 = 1.72$$

24, (at $\epsilon_{\rm G} = 0.38$) = 0.7336 + 0.3505 $\left(\frac{A_{\rm R}}{\alpha A_{\rm cp}}\right)$ - 0.0495 $\left(\frac{A_{\rm R}}{\alpha A_{\rm cp}}\right)^2$ + 0.0024 $\left(\frac{A_{\rm R}}{\alpha A}\right)^3$

$$2\Psi, (at \epsilon_{G} = 0.40) = 0.7671 + 0.3792 \left(\frac{A_{R}}{\alpha A_{cp}}\right) - 0.06246 \left(\frac{A_{R}}{\alpha A_{cp}}\right)^{2} + 0.0037 \left(\frac{A_{R}}{\alpha A_{cp}}\right)^{3}$$

Solving

$$q_r = (33.5) (0.6152) \begin{cases} 0.173 \times 10^{-8} [(1851.4)^4 - (1275)^4] + \\ 7 [1851.4 - 1275] \end{cases} = 4.22 \times 10^5 Btu/hr.$$

Since $q_r \doteq \Delta H$, t_2 will not be recalculated. Calculations for ΔP are similar to those performed in the convection section. Calculations are repeated until the exit temperature, t_2 , from a tube node is equal to the desired discharge temperature from the still.

(12) Check furnace dimensions:

N = number of tubes calculated from assumed dimensions

$$N_a = \frac{2Z + f}{e e D_o} = \frac{(2)(16)}{(15)(0.46)} = 46$$
 tubes.

In this example, 48 tubes were counted in the radiant section. Since $N_a = 48$, the assumed dimension Z is approximately correct and will not be recalculated. If $N_a \neq 48$, N_a in the above equation would be replaced by 48 and Z calculated. This Z could be used during the next iteration through the radiant section.

Calculate f.

 $f = 48 (e'e') D_0 - 2Z.$

(a) If f < L/2, tubes are not placed in the bridge wall, consequently f is set equal to zero.

(b) If f > L/2, tubes are placed in the bridge wall and f is set equal to 2/3Z. Since f = 48 (1.5) (0.46) - 32 = - 3.2 < L/2, the initial assumption, f = 0, is correct. Calculations will not be repeated. When case (b) exists, the iteration is repeated using $Z = 3/8 [N_c \not e \not c D_o]$. N_c = number of tubes counted in the radiant section.

(15) Heat balance:

In this example

heat absorbed in the radiant section = 1.06017×10^7 Btu/hr. heat absorbed in the convection section = 1.2835×10^7 Btu/hr. total heat lost = 2.75988×10^7 Btu/hr.

total heat liberated = 5.10355×10^7 Btu/hr.

From the flow rate of the fuel and its net heating value $(584.2 \text{ Btu/lb. gas}) (83354.9 \text{ lb./hr.}) = 4.85 \times 10^7 \text{ Btu/hr.}$

Since the total heat liberated as calculated from the heat balance and the net heating value of the fuel are almost equal, it will not be necessary to repeat the calculations.

(18) Tube diameter:

Finally, P₂ is compared with the desired discharge pressure. If their difference is not within tolerance another tube diameter is assumed and all calculations repeated. The assumption is made using the relationship:

$$Di_n = Di_{n-1} \left[\frac{\Delta P_c}{\Delta P_a} \right]^{1/5}$$

 $\Delta P_c = \text{calculated pressure drop after n-lth iterations}$ $\Delta P_a = \text{desired pressure drop}$ $Di_n = \text{diameter assumed for the nth iteration}$ $Di_{n-l} = \text{diameter previously used.}$





FIGURE 12

FIXED POINT ORDERS

ORDER	OPERATION
L5 n	Transfer contents of location n to A.
50 n	Transfer contents of location n to Q.
Ll _i n	Add contents of n to A.
LO n	Subtract contents of n from A.
75 n	Multiply Q by contents of n.
40 n	Transfer contents of A to n.
S5	Transfer contents of Q to A.
22 n	Transfer control to right order at location n.
32 n	If contents of $A \ge 0$, execute 22.
26 n	Transfer control to left order at location n.
36 n	If contents of $A \ge 0$, execute 26.
F5 n	Transfer contents of n to A and increase the address digits at the right side by 1.

FLOATING POINT ORDERS:

Let F be the floating decimal number in the floating accumulator and let F(n) be the floating decimal number in location n.

ORDER	OPERATION
80 n	Replace F by $F - F_{(n)}$
81 n	Replace F by -F(n)
82 n	Transfer control to the right hand interpactive order in n if $F \ge 0$
83 n	Transfer control to the left hand interpactive order in n if F ≥ 0
84 n	Replace F by $F + F(n)$
85 n	Replace F by F(n)

ORDER	OPERATION
86 n	Replace F by F/F(n)
87 n	Replace F by F x F(n)
88 O	Replace F by one number read from the input tape punched as sign, any number of decimal digits, sign, and two decimal digits to represent the exponent.
89 n	Punch or print F as a sign, n decimal digits, sign, two decimal digits to represent the exponent and two spaces.
8K n	Replace F by n if $0 \le n < 200$
85 n	Replace F _(n) by F
8N n	Replace F by $/F/ - /F_{(n)}/$
8J n	Transfer control to the fixed point order at the left side of location n
8F n	Give a carriage return and line feed and arrange to print a block of numbers having n columns.

If the first function digit of a floating point order is 0, 1, ...7, it refers to one of a set of control registers, or b-registers in the floating decimal routine which are similarly numbered. These are used to count the number of passages through loops and for increasing the addresses of these orders on successive passages. These addresses are increased only if the function digit of the order corresponds to that of the control register currently used.

ORDER LIST WITH b # 8

The index Cb is used for counting purposes to determine the number of passages through a loop. The index gb is used for advancing the address of floating point orders.

ORDER	OPERATION
b0 n	Replace F by $F - F(n + gb)$
bl n	Replace F by -F(n + gb)

ORDER	OPERATION
b2 n	Replace gb, Cb by gb + 1, Cb + 1 Then transfer control to the right hand (if b2 n)
b3 n	or left hand floating point (if b3 n) order in n if $Cb + 1 < 0$. This transfer is used at the end of the loop.
b4 n	Replace F by F + F(n + gb)
b5 n	Replace F by F(n +gb)
b6 n	Replace F by $F/F(n + gb)$
b7 n	Replace F by F x F(n + gb)
bk n	Replace gb, Cb by O, -n. This floating point order is used for preparing L cycle around loop n times
bS n	Replace $F(n + gb)$ by F
bN n	Replace F by F - $F(n + gb)$
bL n	Replace gb, Cb by gb + n, Cb. This order is used when one wishes to change addresses by some increment other than +1 in a loop. If one places bL 1022 in a loop, the effect will be to decrease addresses by two on each passage.

FURNACE DESIGN PROGRAM

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Program Outline

LOCATION	<u>CON TEN TS</u>
3	Specifies Floating Accumulator
4	Specifies Al routine
5	Specifies A3 routine
6	Specifies SA3 (Natural Log) routine
7	Specifies SA2 (Exponential) routine
8	Specifies routine to calculate physical properties
9	Specifies routine which selects correct polynomials
10	Specifies location of data
11	Specifies location of generated data
12	Specifies location of data
13	
14	Specifies location of routine to calculate $\epsilon_{\mathbf{c}}$
15	Specifies location of routine to calculate ΔP
16-19	Temporary storages
20-157	Convection section routine
158-270	Radiant section routine
635-660	Continuation of radiant section routine
271-379	Subroutine to calculate physical properties
380-449	Subroutine to select correct polynomials
450-508	Data
509-546	Data (depending on type crude)
547-589	Data (generated)
590-634	Polynomials for ψ
665-678	Routine to calculate ϵ_{G}

LOCATION	CONTENTS
679-704	Routine to calculate ΔP
705-738	Storages for selected polynomials and their parameters
739-743	Answer storages
739	Counter for shield section
740	Number of rows in convection section
741	Heat absorbed in convection section
742	Heat absorbed in radiant section
743	Number of tubes in radiant section
744	Heat lost from furnace

P	F	RC	G	F	Y	LΜ

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003K		Dire	ective Specifying Location Routines
00F 00745F			Floating accumulator
00F 00830F			A1
00F 00773F			A3
00F 00800F			SA3
00F 00747F			SA2
00F 00271F			Calculate physical properties •
00F 00380F			Select correct polynomials
00F 00509F			Data
00F 00547F			Data
00F 00450F			Data
00F 00651F			
00F 00665F			Calculate Co
00F 00679F			Calculate ΔP
LOCATION	ORDE	R	NOTES
<u>0020K</u>			
0	22 50	L L	
1	26 S ОК Ц	54 5F	Read in and store constants for This routine is later overread
2	88 OS 5	F 190F	
3	03 2 8 3 9	999F	Transfer control to continue reading and storing program.
00555к			
0	22 50	L L	

LOCATION	ORDER		NOTES	
1	26 OK	S4 23F	Read in and store constants	
2	88 OS	F 100F		
3	03 OK	2L 59F		
4	88 OS	F SN		
5	O3 OK	4L 38F		
6.	88 OS	F SK		
7	03 8K	6L 1F		
8	8K 83	1F 20F	Transfer control to location 20	
	24	555N	Start executing program at location 5	55
<u>0020K</u>			Program for the Convection Section	
0	85 87	52SN SK	Preliminary calculations	
1	84 85	53SN SS		
2	0K 05	5F 1SK		
3	0S 02	1SS 2L	· ·	
4	ОК 85	2F 45SN	Calculate PL of CO_2 and H_2O for the constitute hash	on-
5	87 07	SS 21SK	vection bank	
6	0S 02	2855 LL		
7	8J 8J	8L 8L		

LOCATION	ORI	DER	NOTES
8	41 F5	739F 8L	Entrance to subroutine to select correct polynomials for $\epsilon_{\mathbf{c}}$
9	24 OK	S9 ЦF	Transfer control to subroutine
10	8K 8S	F 195K	
11	87 04	35K 165K	
12	03 8S	11L 744F	
13	8K 8S	F 36SS	
14	85 80	6sk 1sn	
15	87 87	SS 7SK	
16	85 85	5F 155K	
17	86 86	5F SN	
18	8S 87	5F 5F	
19	86 86	SS 36SN	
20	8J 86	S6 3SN	•
21	8J 87	S7 LISN	
22	8S 8к	15F 1F	
23	8 K 84	1F 14sk	
24	86 87	145K 6 5 K	
25	86 85	41SN 32SS	

LOCATION	ORDER		NOTES
26	85 84	LISS 5SS	Assume P2
27	86 85	2SN 8SS	Store P _m
2 8	85 84	1SS 2SS	Assume T ₂
29	86 85	2Sn 7SS	Store t _m
30	ОК 8К	Ц г Г	
31	8S 87	19SK 3SS	
32	04 02	165 k 31L	
33	8S 8J	17F 34L	Entrance to subroutine to calculate
34	22 F5	34L 34L	physical properties
35	26 85	\$8 2155	Transfer control to this subroutine Calculate temp. of flue gases above nth
36	84 85	17F 5F	tube (Ig)
37	81 85	5F 195 K	
38	OK 1K	3F F	
39	05 17	165 k 3SN	
40	OS 1L	8F 1Q23F	
41	03 OK	39L 3F	
42	8K 87	F 20SK	Assume T _a
43	04 02	8F 42L	

LOCATION	ORDER		NOTES
<u>1</u> 11	8S OK	5F 4F	
45	8ĸ 87	F 20SK	
46	04 02	16SK 45L	
47	8S 85	6F 24SN	
48	8n 82	6F 52L	Is this the correct temperature? Y es: Transfer control to 52L
49	85 86	6F 5F	No: Modily g
50	85 85	5F 20SK	
51	80 85	5F 205K	,
52	82 85	41 L 205K	Repeat calculations using new T _g Calculate h _c
53	84 86	3SS 2SN	
54	85 8J	2655 56	
55	87 8J	44SN S7	
56	87 85	15F 27SS	
57	85 80	26SS 10SN	
58	85 80	31SS 1009F	Is this a shield tube?
59	82 8k	64L 1F	No: Continue at 100L without including calculations for a shield tube
60	83 85	120F LSS	These orders are used in the radiant section
61	80 84	1002F 1008F	$P_1 - P_x \ge 0$?

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LOCATION	ORDER		NOTES
62	83 89	250F 4F	Yes: Continue calculations for another node No: Point out their difference and also the tube diameter. Then change the tube
63	85 89	SK 4F	diameter and repeat all calculations.
64	83 84	645F 1F	Go to 645 to change diameter. Calculations for shield tubes begin
65	84 85	739F 739F	Increment counter for shield tubes
66	85 84	SS SK	Calculate T _s
67	87 85	211 SK 5F	
68	85 80	SS SK	
69	87 86	SS 5F	
70	8S 85	5F SS	
71	86 86	SK 25SS	
72	84 85	5F 5F	
73	8 K 86	1F 23SK	Assume h
74	84 85	5F 6F	
75	85 80	31SS 7SS	
76	87 86	5F 6F	
77	84 85	7SS 30SS	
78	85 80	31SS 30SS	

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LOCATION	ORDER		NOTES
79	8S 8J	16F 80L	Entrance to subroutine to evaluate $\left[\left(\alpha + \alpha\right)\right] = \left[\left(\alpha + \alpha\right)\right] = \left[\left(\alpha + \alpha\right)\right]$
80	22 F5	80L 80L	$\left[\left(\frac{q_{c}}{c} + \frac{q_{w}}{T} \right)^{T} - \left(\frac{q_{c}}{c} + \frac{q_{w}}{T} \right)^{T} \right] \left[\frac{100 - \pi}{100} \right]$
81	26 85	SF 10F	Transfer control to subroutine Evaluate h _{rg}
82	86 87	16F 46SN	
83	8S 8K	13F 100F	Evaluate h rb
84	8S 85	5F 30SS	
85	84 86	10SN 5F	
86	8S 87	6F 6F	Evaluate percent wall correction
87	87 87	6F 46SN	
88	87 85	47SN 6F	
89	84 84	13F 27SS	
90	8S 85	7F 6F	
91	87 86	32SS 7F	
92	84 85	1SN 8F	
93	85 84	13F 27SS	
94	87 85	8F 345S	
95	80 85	23SK 5F	Is the h_u assumed equal to h_u calculated
96	85 8n	39SN 5F	

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LOCATION	ORI	DER	NOTES
97	83 85	99L 34SS	Yes: Transfer control to 99L No: Modify h _u
98	8S 83	23SK 66L	repeat calculations using new h
99	85 85	34SS 27SS	$27SS = h_{tube}$ if n^{th} tube is not a shield
100	85 84	SS SK	$27SS = (wall correction)(h_{th}) if n^{th}$
101	87 85	24sk 5f	tube is a shield tube -
102	85 80	SS SK	Calculations for U
103	87 86	SS 5F	
104	8S 85	5F SS	
105	86 86	SK 25SS	
106	84 85	5F 5F	
107	8ĸ 86	1F 27SS	
108	84 85	5F 5F	
1 Q 9	8ĸ 86	1F 5F	
110	8S 87	35SS SS	Calculate $q = uA\Delta t$
111	87 87	7SK 41SN	
112	8S 85	5F 31SS	
113	80 87	7SS 5F	
114	8S 80	5F 21SS	Does $q = \Delta H$?

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LOCATION	ORDER		NOTES
115	8 5 85	6F 1004F	
116	8N 83	6F 126L	Yes: continue at 126L
117	85 80	2SS 1SS	No: modify t ₂
118	87 86	6F 21SS	
119	84 85	2SS 2SS	
120	8ĸ 84	1F 40SS	If t_2 does not converge after the 7 th iteration use last calculated value of
121	8S 80	4055 1008F	t ₂ as the correct value
122	83 85	126L 22SN	No convergence after 7 th iteration, proc e ed 126L
123	80 83	739F 48F	Was this a shield tube? No: repeat calculations using modified t ₂
124	85 80	739F 1SN	Yes: reset counter
125	8S 83	739F 48F	Repeat calculations using modified t ₂
126	85 86	7SK 14SK	
127	85 85	3855 145K	
128	8S 8J	3955 3F	Reset counter for number of iterations t_0 t_2
129	JO F5	130L 139L	
130	26 85	SL 20SK	Transfer control to subroutine to calculate ΔP at this point the tube node has converged
131	8S 85	3 88 988	Make a summation of the ΔHn 's
132	8S 85	36SS 21SS	

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LOCATION	ORDER		NOTES
133	84 8 5	741F 741F	
134	8K 84	1F 740F	Make a summation of the number of rows (n's)
135	85 85	740F 6SN	Was this the last tube in the shield section?
136	80 83	739F 26L	No: repeat calculations for another tube
137	8K 8S	2F 3955	Tes: Floceed to the fadfant Section.
<u>003K</u>	1 4	1000	
0	41 26	149F	Reset counter for the number of iterations of t_2 to zero.
00158K			Program For the Radiant Section
0	85 80	20SK 10SN	Store T as the bridgewall temperature g
1	8S 85	31SS 1SS	Store t_1 as the cross-over temperature
2	8S 85	6SS 4SS	Store P_1 as the cross-over pressure
3	8S 8K	37SS F	Set counter for ΔH_R and $n_R = 0$
4	8S 8S	742F 743F	
5	85 85	6SS 1SS	
6	85 85	37SS 4SS	
7	85 87	38SS 1000F	Assume Z, and calculate L_B
8	87 8J	1000F S6	
9	86 8J	3SN S7	

LOCATION	ORDER		NOTES
10	87 86	2SN 3SN	•
11	8S 87	5F 21SK	Calculate PL of $\rm CO_2$ and $\rm H_2O$ in the radiant section
12	85 85	2855 5F	
13	87 85	225K 2955	
14	8K 87	2F 1000F	Calculate $\mathbf{A}_{\mathrm{R}}/\alpha \mathbf{A}_{\mathrm{CP}}$ for the radiant section
15	8S 84	5F 42SN	
16	87 87	38SS 1001F	
17	85 8k	6F 2F	
18	87 84	38SS 1000F	
19	87 86	5F 6F	
20	80 85	1SN 32SS	
21	85 87	3855 3955	Calculate aA _{cp} per node
22	87 87	1001F SS	
23	87 85	6SK 41SS	
24	85 84	SS SK	
25	87 85	24SK 5F	
26	85 80	SS SK	
27	87 86	SS 5F	

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LOCATION	ORDER		NOTES
28	8S 8J	19F 29L	
29	22 F5	29L 29L	Select correct polynomials for ϵ_{lpha}
30	26 85	S9 4SS	Assume P ₂
31	84 86	5SS 2SN	Calculate P _m
32	8S 85	8SS 1SS	Assume t ₂
33	84 86	2SS 2SN	Calculate t _m
34	8S 8J	7SS 35L	
35	22 F5	35L 35L	
36	26 8K	S8 1F	Calculate percent vaporized, H_{L} , H_{v} , etc. Calculate T _s from q = uA[t _s - t _m]
37	86 84	25SS 19F	
38	85 85	5F 21SS	
39	87 86	5F SS	
40	86 86	41SN 39SS	
41	86 84	3855 755	
42	85 8j	3055 43L	
43	22 F5	43L 43L	Calculate $[(q_{c} + q_{u})T_{B} - (q_{c} + q_{u})_{T}][100-\%]$
44	26 2K	SF 2F	Calculate $(q_{BB})_{T_B} - (q_{BB})_{T_s}$

LOCATION	ORDER		NOTES
45	окµғ 8к	F	
46	27 04	3088 548N	
47	03 2S	461. 5F	
48	23 80	451. [.] 5f	
49	8S 85	6F 10F	Evaluate ϵ_{G}
50	86 85	6F 10F	
51	OK 1K	9F F	Select correct polynomial for ψ
52	15 80	590F 10F	
53	83 85	59L 718F	
54	3K 15	4F 591F	
55	3S 1L	7 10F 1F	
56	32 1 L	541. 1F	
57	03 85	52L 46SN	
58	8S 83	33SS 72L	
59	8S 3K	719F ЦF	
60	15 3S	591F 714F	
61	1L 33	1F 60L	
62	2K OK	2F F	Interpolate between the polynomials with values of ϵ_{on} closest to ϵ_{G}

LOCATION	ORDER		NOTES
63	1K 8K	4F F	
64	87 04	32SS 710F	
65	OL 13	1F 64L	
66	2S 23	8f 63l	
67	80 85	8F 7F	
68	85 80	719F 718F	
69	8S 81	5F 718F	
70	86 87	5 F 7F	
71	84 85	8F 33SS	Store V
<u>00230K</u>			
0	2K 25	2F 30SS	Evaluate q _r
1	84 8J	10SN S6	

1	84 8J	10SN S6
2	87 8J	Losn S7
3	2S 22	5F L
4	80 87	5F 1018F
5	8S 85	6F 31SS
6	80 87	30SS 1008F
7	84 87	6F 33SS

LOCATION ORDER		DER	NOTES	
8	87 80	41SN 21SS	Is $q_r = \Delta H$?	
9	8S 85	6F 1004F		
10	8N 82	6F 1ЦL	Yes: proceed to 14L	
11	85 80	255 155	NO: modily t ₂	
12	87 86	6F 21SS		
13	84 85	2SS 2SS		
14	82 8J	190F 15L	Repeat calculations using new t ₂	
15	JO L5	16L 15L		
16	26 8K	SL 2F	Calculate ΔP At this point the tube node has co nverged.	
17	84 85	743F 743F	Make a summation of n _R 's	
18	85 84	742F 21SS	Make a summation of the ΔH_R 's	
19	8S 82	742F 80F	Transfer control to location 80	
20	85 80	1003F 1SS	is $t_x - t_2 > 0$	
21	82 8K	188F 2F	Yes: Repeat calculations for another node No: Stop calculating nodes and test for	
22	87 85	1000F 5F	convergence of furnace	
23	85 87	743F 6sk	,	
214	87 85	SS 9F		
25	80 87	5F 2SN	If $2f - dL < 0$ use $f = 0$ 2f - dL > 0 use $f = 2/3 Z$	

. 88

LOCATION	ORI	DER	NOTES
26	80 83	38SS 30L	
27	8K 8S	F 42SS	Use $f = 0$
28	85 86	9F 2SN	
29	8S 82	9F 33L	
30	85 87	9F 44sn	Use $f = 2/3 Z$
31	86 85	42SN 9F	
32	87 86	2SN 3SN	
33	85 85	42SS 1000F	Does Z assumed = Z calculated?
34	80 8 N	9F 42SN	
35	82 85	639F 741F	No: proceed to location 639 Yes: Is the assumed flue gas flow rate
36	84 84	742F 744F	Correcti
37	86 85	1004F 12F	
38	86 80	155K 15N	
39	85 8K	5F 1F	
40	83 83	635F 635F	Continue these calculations at 635
00635K			
0	85 8n	39SN 5F	
1	82 8F	6L 2F	Yes: Continue at 641 No: print out N _R last calculated and flue gas flow rate last used

LOCATION	ORI	DER	NOTES
2	85 89	743F 3F	
3	85 89	12F 6F	
24	83 85	3SJ 9F	Proceed to modify flue gas flow rate Modify Z
5	85 89	1000F 5F	Print out modified Z
00641к			~
0	82 85	161F LSS	Repeat calculations in radiant section using Does $P_2-P_x = 0$ within limits? new Z
1	80 89	1002F 3F	Print out difference
2	85 85	5 F 1008F	
3	8n 82	5F 8L	Yes: Furnace design complete, proceed to 8L
- '4	85 80	14SK 4SS	No: modify tube diameter
5	86 8J	1005F S6	v
6	86 8j	1006F S7	
7	87 85	SK SK	Using new diameter, repeat all calculations.
8	82 85	7SJ 1000F	First, reset all counters to zero at 7SJ. Furnace design complete. Print out answers.
9	89 OK	8F 6F	
10	05 89	739F 8F	
11	03 8J	10L 12L	
12	OF OF	F F	Stop all calculations

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LOCATION	ORDER		NOTES			
13	86 85	15sk 10f	(3SJ)	Modify gas flow rate		
14	0K 05	3F 16sk				
15	87 05	10F 165K				
16	02 85	14L 12F				
17	8S OK	155K 6F		Set all counters to zero		
18	8K OS	F 739F				
19	03 83	18L 20F		Repeat all calculations from beginning		
<u>0027 1K</u>				Subroutine for Physical Properties		
0	42 50	106L L				
1	26	S 4				
	OK	ЦF		Calculate percent vaporization		
2	8K 87	F SS				
3	04 02	265K 2L				
4	8S 85	5F 8sk				
5	87 85	58N 6F				
6	85 80	5F 6F				
7	80 85	7SS 5F		Was there any vaporization in this node?		
8	83 81	107L 5F		No: Go to 107L. and pick up previously calculated percent Yes: Is the amount of vapor negligible?		

LOCATION	ORDER		NOTES			
9	86 86	6F 2SN				
10	85 85	955 95K				
11	80 83	955 15L	Yes: Set percent vaporized = 0			
12	OK 8K	ЦF F	a function of the percent vaporized.			
13	87 04	9SS 30SK				
1)4	03 82	13L 16L				
15	8K 8S	F 9SS	Set percent vaporized = 0			
16	85 85	105 K 1055	Calculate the density of the crude at t_m			
17	85 87	1055 115K	.)			
18	84 86	12SK 6SN	· · ·			
19	85 85	5F 7SS				
20	87 80	3SN 7SN				
21	87 85	5F 5F				
22	ОК 8К	ЦF F				
23	87 04	1055 485N				
24	03 84	23L 5F				
25	8S 85	5F 8SN				
26	80 86	7SS 6SN				

LOCATION	ORDER		NOTES
27	8 7 84	5F 1055	
28	87 85	9SN 11SS	
29	85 80	1SN 9SS	Calculate W _L
30	87 87	135K 1055	
31	86 85	105K 1255	
32	85 80	135K 1255	Calculate $w_y = w - w_L$
33	85 85	1355 755	Calculate the vapor density.
34	84 85	10SN 5F	
35	87 86	11SN 8SS	4
36	85 85	8F 8SS	
37	87 84	13SN 12SN	
38	87 86	17SN 5F	
39	86 86	5F 5F	
40	85 85	7F 8SS	
41	87 84	14SN 15SN	
42	84 80	6F 7F	
43	8S 85	6F 16SN	
2424	86 85	6F 14SS	

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LOCATION	ORDER		NOTES		
45	85 87	13SS 11SS	Calculate the liquid volume fraction		
46	86 84	1455 1255			
47	85 85	6F 12SS			
48	86 85	6F 1588			
49	85 87	185N 1055	Calculate H_{L} and H_{v}		
50	80 85	195N 10F			
51	85 87	20SN 10SS			
52	80 85	21SN 9F			
53	85 85	22SN 8F			
54	1K OK	2F 3F			
55	8ĸ 17	F 1SS			
56	04 02	8f 55l			
57	1S 12	16SS 54L			
58	0K 05	3F 23SN			
59	87 85	1055 6F			
60	05 80	26SN 6F			
61	0S 02	8f 58l			
62	1K 0 K	2F 3F			

LOCATION	ORI	ER	NOTES
63	8K 17	F 1SS	
64	04 02	8f 63l	
65	1S 12	1 8 SS 62L	
66	0K 85	2F 30SN	
67	07 84	455 295N	
68	07 8S	4SS 6F	
69	85 07	32SN 4SS	
70	84 07	31SN 4SS	
71	87 85	175 N 7F	
72	05 84	1SS 10SN	
73	8S 85	8F 7F	
74	86 86	8F 8F	·
75	86 85	8F 9F	
76	85 80	6F 9F	
77	86 04	33SN 18SS	
78	0S 02	1855 66L	
79	85 80	7SS 37SN	Calculate the thermal conductivity of the crude.
80	87 85	34SN 6F	

NOTES

LOCATION	ORDER		NOTES		
81	8K 80	1F 6F			
82	87 86	35SN 36SN			
83	86 85	1055 2055			
84	85 80	1755 1655	Calculate AH		
85	87 85	1255 6F			
86	85 80	1955 1855			
87	87 84	1355 6F			
88	8S 85	21SS 2SS			
89	80 85	155 6F			
90	85 80	1755 1655			
91	86 85	6F 22SS			
92	ОК 8 К	ЦF F	Calculate the viscosity of the vapor		
93	87 04	7SS 34SK			
94	03 87	93L 11SS			
95	86 87	9SN 38SN			
96	8S 85	23SS 40SN	Calculate (R _e) _L		
97	87 86	1255 415N			
98	86 86	SK 23SS			
LOCATION	ORI	DER	NOTES		
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99	85 8J	2455 56	Calculate h _i		
100	87 8J	42SN S7			
101	87 85	43SN 5F			
102	85 87	22SS 23SS			
103	86 8J	2055 56			
104	86 8J	3SN S7			
105	87 85	5F 25SS			
106	8ĸ 82	1F []F	Transfer control to main routine		
107	85 85	3655 955			
108	8ĸ 82	1F 10L			

Program to Select Correct Polynomials for

<u>00380к</u>			This program will select, from tape, two polynomials with PL values closest to polynomial desired. It will also interpolate between these polynomials to obtain ϵ_{c}
0	42 50	26L L	
1	26 2K	54 2F	
2	зк 35	2F 2855	Store PL CO_2 , and then PL H_2O
3	8S 8J	706F 41	
4	22 F5	4L 4L	

LOCATION	ORD	ER	NOTES
5	26 OK	408F 4F	Transfer control to 408 Store polynomials in the locations designated for them
6	05 25	710F 720F	
7	05 2S	715F 724F	
8	2L 03	1F 6L	
9	85 3S	709F 736F	
10	2L 32	ЦF 2L	Repeat to select and evaluate q_{CO_2} . At T_B and T_S
11	85 84	2855 2955	Calculate parameter for percent correction Parameter = PL_{CO_2} + PL_{H_2O}
12	85 8j	706F 440F	T. C. to 440
13	85 80	2SN 706F	If 2 - $(PL_{CO_2} + PL_{H_2O}) < 0$ let $(PL_{CO_2}^{CO_2} + PL_{H_2O}) = 1.8$
14	83 85	16L 6SN	
15	8S 83	706F 16L	
16	85 86	2855 706F	
17	8S 8J	10F 18L	
18	22 F5	18L 18L	
19	26 OK	408F 5F	Select polynomial Evaluate
20	8K 87	F 10F	
21	04 02	715F 20L	
22	85 8k	5F 100F	

LOCATION	ORDE	ER	NOTES
23	8S 80	6F 5F	
2]4	86 85	6F 738F	
25	8 j 8k	442F 1F	Transfer to 442
26	8K 82	1F []F	
∞4 <u>08k</u>			
0	42 50	12L L	
1	26 8J	54 446F	Transfer control to 446
2	85 80	707F 706F	Is $PL_m - PL > 0$?
3	82 85	5l 708f	Yes: Proceed to 5L No: Store difference and proceed to 421
4	8 j 8k	421F 1F	
5	82 8J	1L 429F	Read in and check another polynomial Transfer to 429
6	8 j 85	448f 58sn	Transfer to 448 Was this the 1ast polynomial in the set?
7	80 82	7 05F 9L	Yes: Use the two polynomials in memory
8	8j 8k	436F 1F	No: Proceed to 430
9	83 85	6l 708f	Interpolate between the PL values available.
10	84 80	706F 707F	
11	86 85	708F 709F	
12	8K 82	1F []F	Transfer control to 385

LOCATION	ORD	ER	NOTES
ооццок			
0	F5 40	51. 51.	Prepare to operate on μ^{th} degree polynomials
1	26 001F	2954	
2	L5 LO4L	5L	Reset to operate on 3 rd degree polynomials
3	40 26	51. 2954	
4	00F 001F	2	
5	00F 004F	,	
6	81 40	40f 707f	Read a sexadecimal character from tape (representing a value of PL) and store in 707
7	22 26	7L 2954	Transfer control to the order following
8	81 40	40F 705F	Read in and store the next PL on tape
9	22 26	9L 2 9 34	Proceed to order following the last 8J order executed
00421К			
0	22L L5	6L	
1	42 41	2L 7L	
2	81 40	40F 7 10F	Read in and store a polynomial
3	F5 42	2L 2L	
4	F5 40	7L 7L	
5	LO 36	445F 2984	Proceed to last 8J order executed
6	26 00	2L 710F	

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LOCATION	ORI	DER	NOTES
7	00 00	F F	
8	L5 42	9L 10L	Read in and s tore polynomial associated with PL last checked
9	41 LO	7L 715F	
10	81 40	40F 715F	
11	F5 42	10L 10L	
12	F5 40	7L 7L	
13	LO 36	445F 2954	Proceed to order following 8J order last
14	26 26	10L 10L	executed
15	41 81	7L 40f	Read in and dump a polynomial
16	F5 40	7L 7L	
17	LO 36	445F 2984	Transfer to order following the last 8J
18	22 22	15L 15L	order executed
00665к		Subrouti	ne to Evaluate Polynomials for and to interpolate between them
0	42 50	13L L	
1	26 0K	SL 2F	
2	1K 2K	F 2F	
3	3К 4К	2F 4F	
4	8K 07	F 30SS	

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LOCATION	ORDER			
5	14 1L	720F 1F		
6	42 35	ЦL 5F		
7	32 80	3L 5F		
8	26 84	736F 5F		
9	2S 23	8F 3L		
10	84 05	8F 10F		
11	03 80	2L 10F		
12	87 85	738F 10F		
13	8K 82	1F []F		

00679К

171		
0	46 50	22L L
1	26 85	S 4 2455
2	8J 87	56 1014F
3	8J 8S	S7 5F
4	85 80	15SS 1015F
5	83 85	19F 1588
6	87 84	5 f 1016f
7	87 87	37SN 13SK

NOTES

Subroutine to Evaluate ΔP

LOCATION	ORI	DER	NOTES
8	87 OK	13SK 5F	
9	86 03	SK 9L	
10	86 85	1017F 5F	
11	85 80	1155 1455	
12	87 84	1555 1455	
13	8S 85	6F 1011F	
14	87 84	SK 38SS	
15	87 87	39SS 5F	
16	86 84	6F 5SS	Does P_2 assumed = P_2 calculated?
17	80 85	4SS 10F	
18	8n 82	39SN 22L	No: Modify P ₂ and repeat calculations
19	85 85	5SS 4SS	Yes: Replace P ₁ by P ₂
20	85 85	2SS 1SS	Replace t_1 by t_2
21	84 85	Losn 2SS	
22	8 2 85	[]F 5SS	Transfer control to program
23	80 85	10F 5SS	
24	84 86	4SS 2SN	
25	8S 82	855 158	

104

LOCATION ORDER <u>0019K</u> 0

85 83 1015F 685F NOTES

Per cent Corr	ection:				
2% = a + bR +	$dR^2 + dR^3$	³ + eR ⁴ ; R	$= \infty_2/(\infty_2 +$	H ₂ O)	
Parameter = PL	a	b	с	đ	e
0.01 0.25 0.5 0.75 1.0 1.5 2.0	0.3857 0.3558 4.5758 1.0423 2.5018 3.2912 4.2321	4.3900 21.9324 7.3524 51.7054 49.2875 55.1680 51.9834	20.1857 -16.2491 18.9710 -98.6457 -92.4543 -109.9346 -101.8847	-33.8957 5.5128 -19.1826 106.2914 102.1649 120.0117 111.4446	9.7907 -8.6094 -6.4339 -53.2739 -52.9669 -58.3493 -54.6296

Table 5. Polynomials for the Evaluation of $\boldsymbol{\psi}$

Radiation due to Carbon Dioxide:

 $2E = a - bt + ct^2 - dt^3$

rarameter	Pa	ran	net	.er
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= $PL_{\mathcal{O}_2}$	a	b	c x 10 ³	d x 10 ⁷	
0.001	51.7010	0.1513	0.2078	-0.2919	
0.002	54.0670	0.1 802	0.3150	0.3364	
0.003	211.8020	0.5723	0.6414	0.7937	
0.004	173.6900	0.6191	0.8155	1.0992	
0.005	30.0980	0.3734	0.7349	0.8113	
0.006.	35.6890	1.0696	1.2016	1.4618	
0.008	252.8100	0.9766	1.2920	1.4689	
0.01	651.9700	1.9534	2.0423	2.6610	
0.015	754.5600	2.2446	2.3347	2.7522	
0.02	1278.9000	3.50 58	3.2896	4.2098	
0.03	1503.5400	4.1913	3.9372	5.0236	
0.04	1103.5000	3.3869	3.5666	3.7936	
0.06	1346.5700	4.1166	4.2543	4.3270	
0.08	668.4400	2.5749	3.2179	1.4470	
0.10	1817.8000	5.3986	5.3289	5.2498	
0.15	3927.0900	10.0878	8.3592	9.5874	
0.20	3990.4400	10.2342	8.4211	8.1694	
0.30	5781.0000	14.3987	1.1211	11.8823	
0.40	4823.2900	12.1531	9.7198	7.3506	
0.6	7588.0000	18.2576	13.4843	11.9161	
1.00	6246.9900	16.0636	12.7290	8.7444	
2.00	8057.2900	19.8252	15.1587	9.2603	
4.00	6090.0000	16.0543	13.1966	1.7166	

Radiation Due to Water Vapor:

 $2E = a - bt + ct^2 - dt^3$

Paramet

= PL _{H2} O	a	b	c x 10 ³	d x 10 ⁷	
0.01	55,5228	0.1204	0,2066	0.0937	
0.015	-110,9520	-0.2520	0.0261	-0.4142	
0.02	25.6710	0.0940	0.3523	0.1163	
0.025	321.7550	0.7316	0.7960	0.7295	
0.03	265.0900	0.6268	0.8141	0.6336	
0.04	229.8430	0.5751	0.8985	0.5156	
0.05	430.1100	1.0786	1.3918	1.1563	
0.06	25.3160	0.2532	1.0233	0.1046	
0.08	288.6660	0.7717	1.3913	0.3541	
0.10	139.8600	0.3633	1.2072	-1.1428	
0.15	537.3800	1.3519	2.2078	0.6104	
0.20	1908.1100	4.3802	4.4417	1.7536	
0.25	8.4300	0.3724	2.2078	-3.4239	
0.30	1196.2760	3.1741	4.4304	-0.5380	
0.40	453.3800	1.1607	3.0303	-0.6590	
0.50	-90.2000	-0.0144	2.5594	-9.9 998	
0.60	631.0400	1.9334	4.1433	-9.2052	
0.80	2025.3000	4.6138	5.8862	-11.4342	
1.00	622.9000	1.6958	4.3397	-18.4798	
1.25	-2283.7000	5.2497	-0.6273	-34.1129	
1.50	1408.0000	0.0252	3.8746	-31.9589	
2.00	5548,0000	10.9464	8.7011	-32.9599	
3.00	1362,3000	-0.3637	-0.5046	-64.4918	

Black Body Radiation:

 $2E_{B} = -5926.00 + 27.6324t - 31.7241t^{2} + 255.0130t^{3}$

Overall Exchange Factor:

 $2 \psi = a + bR - cR^2 + dR^3; R = A_R / \alpha A_{cP}$

Parameter =	a	b	c x 10 ²	d x 10 ⁴
0.2	0.3875	0.3009	0.3655	0.1787
0.22	0.3934	0.3569	0.5188	0.2995
0.24	0.4655	0.3292	0.4177	0.2023
0.26	0.4981	0.3467	0.4726	0.2480
0.28	0.5460	0.3480	0.4765	0.2437
0.30	0.5827	0.3634	0.5380	0.2990
0.32	0.6170	0.3740	0.5750	0.32811
0.34	0.6583	0.3651	0.5726	0.30 38
0.36	0.6830	0.3818	0.6028	0.3486

Table 5 (cont.)					
Parameter =	a	b	c x 10 ²	d x 10 ⁴	
0.38 0.40 0.45 0.50 0.55 0.60 0.65 0.70	0.7336 0.7671 0.8699 0.9567 1.0559 1.1404 1.2232 1.2863	0.3505 0.3791 0.3668 0.3472 0.3379 0.3245 0.3068 0.2976	0.4950 0.6236 0.6148 0.5738 0.6076 0.6003 0.5993 0.6122	0.2447 0.3728 0.3657 0.3290 0.3765 0.3765 0.3929 0.4172	_

Table 6. Molar Heat Capacities of Flue Gas Components

= a + bT + cT ² ;	T = OK		
Compound	a	b x 10^{3}	сх 10 ⁶
Oxygen Carbon dioxide Water vapor Nitrogen	6.0954 6. 39 30 7.219 6.4492	3.2533 10.100 2.374 1.4125	-1.0171 -3.405 0.267 -0.0807

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