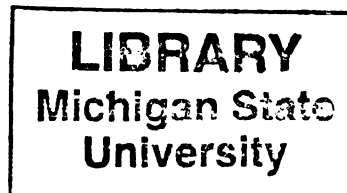




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**THE APPLICATION OF MODULAR MODELING METHODS IN FINITE
ELEMENT ANALYSIS**

By

Xin Li

A DISSERTATION

**Submitted to
Michigan State University
in partial fulfillment of the requirements
for the degree of**

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ABSTRACT

The Application of Modular Modeling Methods in Finite Element Analysis

By

Xin Li

Structural modeling analysis is important to engineering design and development. An important issue is the interconnectivity problem for multi degree of freedom subsystems. The traditional finite element method requires re-meshing subsystems for mesh compatibility at interconnections. As the number of subsystems increase, the computational burden of re-meshing increases rapidly, sometimes exponentially.

The Modular Modeling method is applied here to develop a new finite element analysis approach. Modular Modeling through the use of input / output ports on the boundaries of interconnected parts is introduced for this purpose. Using input / output ports and the Modular Modeling method to connect parts in system design does not require expensive remodeling and re-meshing. This approach uses the original FEA model and combines parts with the input / output ports among their boundaries. As the number of available alternate parts in a structural system increases, the modular FEA method minimizes the time and cost for the system integration.

A two-dimensional example of FEA elastic strain-stress problems is used to illustrate this technique. This example shows it is not necessary to re-mesh components to assemble a structural system. Computation error is evaluated to show that this modular finite element method produces accurate system models for engineering design.

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INTRODUCTION

During the design of complex systems, reliable predictions of detailed stress states for coupled components are often critical to correctly predict the performance of those systems. The finite element analysis (FEA) yields accurate, detailed stress predictions for individual system components with the resolution of mesh distribution. When more precision is needed, a higher precision mesh needs to be generated. In engineering practice, system components can be so complex that the finite element method is mandatory for the design phase. It is computationally intensive and often requires extensive re-modeling effort.

Coupling system components into assemblies is a time intensive task frequently requiring reformulation for finite element analysis models. For example, the Chrysler large car models used for mechanical geometry studies contain representations of over 5500 interconnected physical subsystems [Computers in Engineering: "Chrysler designs paperless cars", 1998]. To reduce the length and cost of the design process for a complex system, efficient modeling techniques that guarantee reliable predictions of detailed stress states are needed. These techniques must be not only computationally efficient, but also dramatically reduce or eliminate the necessity of re-modeling. According to [Computers in Engineering: "Chrysler designs paperless cars", 1998], Chrysler engineers solved design issues by developing a computer model instead of

physical prototypes, reducing the cycle time from 39 to 31 months and saving the company more than \$75 million.

Finite element analysis of system of components typically requires compatible node geometry at each component connection. When constructing the two-dimensional FEA models, it is a mandatory requirement to keep mesh nodally compatible [Zienkiewicz and Taylor, 1989]. At present, the methods available to ensure FEA nodal compatibility include the global /local analysis method, the coupling analysis method, and the interface element method, [M. Aminpour, 1995; V. S. Kothnur, 1999; J. B. Ransom, 1993]. These methods have drawbacks that prevent them from being widely accepted as standard design tools, [Z. Zhao, 1998]. The common limitation of these methods is that they can only be applied to mesh discretizations with a one-to-one nodal correspondence across boundaries between structural subsystems. Methods that do not require nodal compatibility on the component boundaries have much more modeling flexibility, and eliminate the need for re-meshing.

A new finite element analysis approach is introduced in this paper which employs the fixed input / output structure "modular modeling method" [B. Byam and C. Radcliffe, 1999]. In Modular Modeling, the mathematical model that describes each component remains the same independent of the system model in which it is used. Fixed input / output structure means the input and output ports are standardized, therefore the internal equations of mathematical subsystem

models of engineering systems have the same modularity as the engineering system. Fixed input / output structure modular modeling is a power-based, physically intuitive, top-down methodology systematic modeling methodology that eliminates equation reformulation from large model development across multiple energy domain [B. Byam and C. Radcliffe, 1999]

This paper demonstrates the modular modeling method in FEA. The work presented here developed an analytical framework with the Modular Modeling analysis method, a solution strategy, and a demonstration of this solution on a FEA example. Finally, the accuracy and convergence of a Modular FEA Model will be evaluated by comparison with various conventional FEA solutions to a 2-D plate problem (Fig, 1.1).

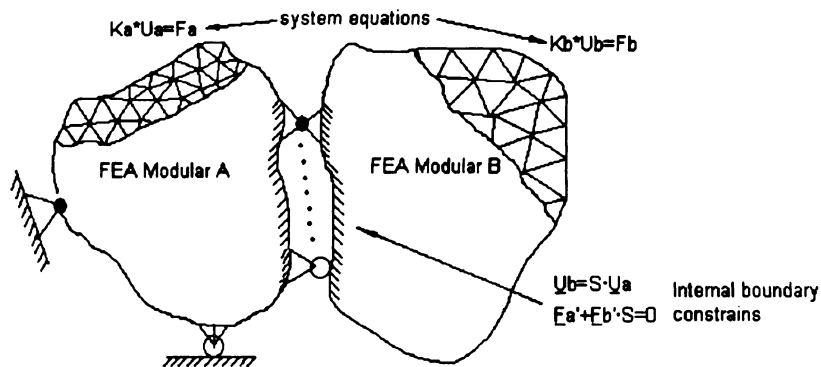


Figure 1.1. The modular modeling using input / output Ports (joint)

MULTI-COMPONENT MODULAR MODELING FOR FINITE ELEMENT ANALYSIS

2.1. FEA Modular Modeling and the related DOF analysis

Modular modeling is a new modeling method designed to eliminate the model reformulation and enhance model performance verification. It is defined by a standard input-output structure, in which connection model physical interaction between components via power transfer. Power Transfer between components requires the product of two signed variables such as force times velocity, pressure times flow rate, and voltage times current. The two physical variables form a component input-output pair. For each type of physical interaction there is a standard input-output definition for modular modeling components. [B. Byam and C. Radcliffe, 1999]

Modular modeling can be applied to FEA model components. In the discussion below, a basic 2-D elastic stress-strain FEA problem will demonstrate the basic approach for connecting different FEA component models together, even when their nodes are geometrically incompatible. The following figure shows the basic concept for modular modeling in FEA. In this approach, a single, multi-node input/output port is defined to connect two FEA models. The power, force and displacement transfers through this connection port have standard definitions [B. Byam and C. Radcliffe, 1999].

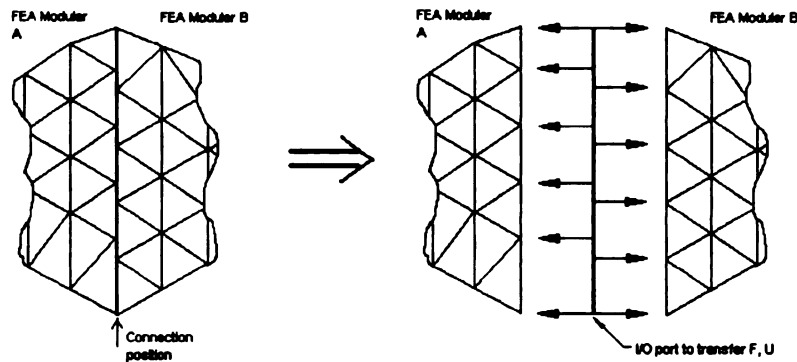


Figure 2.1. The modular modeling using input / output Ports (connectors)

The analysis of system degrees of freedom (DOF) is the foundation of all other analysis. We can not begin any static or dynamic analysis before we decide the system DOF. In FEA modular modeling, the connection of two independent FEA components produces a modular model with a unique solution if it is the zero DOF. In another words, modular connectors provide enough constraints for whole system to generate a non-singular system equation. For example, if we combine two independent FEA components, each with 6, unconstrained DOF, we need to find at least $2 \times 6 = 12$ constraints for displacement to generate a system model with stable, unique solutions for different force inputs. These constraints should provide both external boundary conditions and internal boundary conditions. In Figure 1.1, a constraint set would fix one component to the ground (external constraint) and the other component is fixed on the first one (internal constraint). This combination of constraints will remove all rigid body degrees of freedom and generate system equations with a unique solution.

Supplemental boundary constraint equations are critical to control the FEA modular model DOF. In the first step of the proposed analysis, we combine system equations and their degrees of freedom. For example, in 2-D stress-strain analysis, the three degree of freedom unconstrained components have a singular stiffness matrix, whose rank is $N-3$ where N is the dimension of the stiffness matrix. When we combine two three degree of freedom components together, the

combined stiffness matrix has a rank of $N+M-6$, where N and M are the dimensions of the two component stiffness matrices. External boundary conditions combined with connection constraints must increase the combined system rank to $N+M$ to generate a model with unique solutions to external force inputs. This idea can be easily understood if we consider this with a physical model. In two-dimensional space, a 3-DOF subject can be held fixed if 3 boundary conditions are provided. When we connect two 3-DOF subjects together, an additional 3 boundary conditions prevent relative motion between the two components. When the 3 boundary conditions on one component are combined with 3 internal constraints, the resulting stiffness matrix is non-singular and yields unique solutions to input forcing.

2.2. Visual Component Boundary Port — FEA Input/Output Connector

The modular modeling method generates non-singular equations by using the original component system equations and the modular input/output equation. Actually, this is one of the advantages of the modular modeling method. With this method, a great deal of effort can be saved because no remodeling and re-verification is needed [B. Byam and C. Radcliffe, 1999]. When applying this idea into FEA, the benefit is clear: an FEA model can be used directly in system re-mesh and re-verification of component models are not necessary.

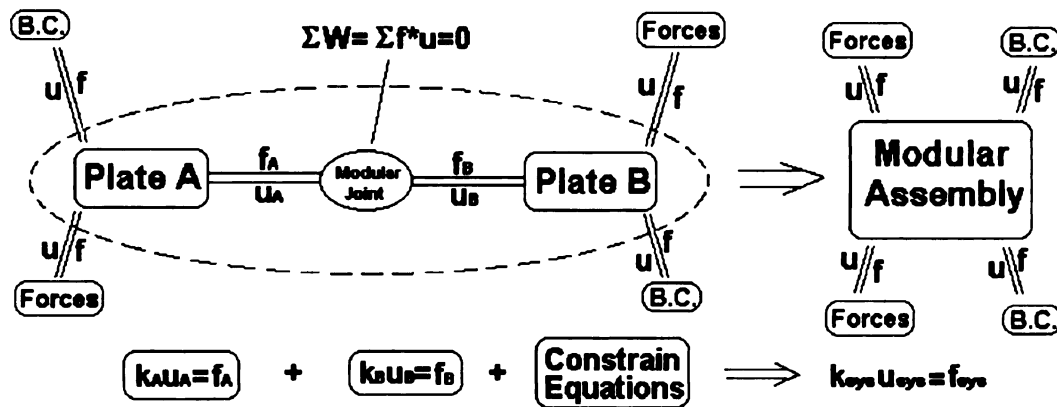


Figure 2.2. The modular modeling method for the FEA application

Definition of an I/O port on each component connection will provide proper constraint equations to make the combined system equation non-singular. The component system equations will be applied directly to generate the combined system equations. The Modular Modeling Method (Fig. 2.2) allows the development of standard integration and computer automation methods. Development of I/O port constraint equations begins with the basic properties of FEA model. Force and displacement transferred across a component connection port must satisfy basic energy and work conservation requirements [Yang, T.Y. 1986,]. The total work transferred through the I/O port must sum to zero.

$$\delta W_A + \delta W_B = 0 \Rightarrow \delta W_A = -\delta W_B \quad \dots\dots\dots(2.1)$$

in which :

$$\delta W_A = \sum f_{aV} \cdot u_{aV} = f_{2a1-1} \cdot u_{2a1-1} + f_{2a1} \cdot u_{2a1} + f_{2a2-1} \cdot u_{2a2-1} + f_{2a2} \cdot u_{2a2} + \dots + f_{aV} \cdot u_{aV}$$

$$\delta W_B = \sum f_{bM} \cdot u_{bM} = f_{2b1-1} \cdot u_{2b1-1} + f_{2b1} \cdot u_{2b1} + f_{2b2-1} \cdot u_{2b2-1} + f_{2b2} \cdot u_{2b2} + \dots + f_{bM} \cdot u_{bM}$$

In the upper equations, the indices a1, a2,... aN, and b1,b2,... bN refer to the nodal number along the connection forming the I/O port.

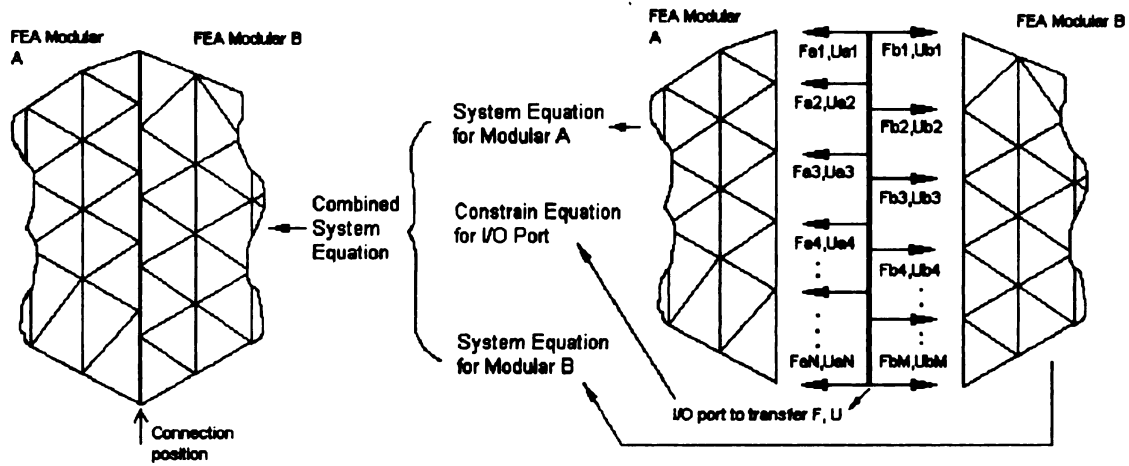


Figure 2.3. The modular modeling input / output Ports along the FEA boundary

Theoretical solution is provided here for a two-dimensional, straight-line boundary and three degree of freedom elastic stress-strain analyses. Future work will be needed for three-dimensional, curve-surface boundary, different type of elements and six degree of freedom finite element analysis.

2.3. FEA Modular Modeling Procedures and Formulations

New formulations for the modular modeling need to be developed for the two-dimensional FEA models. The procedures and formulations for one-dimensional FEA modular modeling have been demonstrated [B. Byam and C. Radcliffe, 1999]. The two dimensional problems are different and need updated procedures and formulations. Two-dimensional FEA modular modeling does not require mesh nodal compatibility along the connection boundary. To make FEA modular modeling, we must develop new procedures and formulations for making the FEA modular model work under the situation of mesh nodal non-compatibility along the connection boundary. Methods to connect modular FEA models without compatible nodal point geometry are given below.

In this paper, four steps develop a two-dimensional FEA modular modeling formulation, which is demonstrated in section §2.3.1, §2.3.2, §2.3.3 and §2.3.4. The combined system equation of FEA models is provided by step one. Step two applies the boundary displacement constraint equations to transform the system equations. Step three puts the system equations into standard linear algebraic equation form by moving the unknown force into the unknown variable vector. In step four, the boundary force constraint equation will supplement the system equation, and make the system stiffness matrix square. The system model equations are now complete, non-singular and have a unique solution given a set of independent inputs.

2.3.1. Sub-system Equation Combination

The first step is to get the combined system equation. Consider 2 FEA components: A and B. A has N nodes and B has M nodes and their stress-strain equation is following. To avoid confusion, The displacement and force on subject A and B are represented by u_A , u_B and f_A , f_B . They have same definitions and units.

$$\mathbf{K}_A \cdot \mathbf{U}_A = \mathbf{F}_A \quad \text{or}$$

$$\begin{bmatrix} \mathbf{k}_{11} & \mathbf{k}_{12} & \cdots & \mathbf{k}_{1N} \\ \mathbf{k}_{21} & \mathbf{k}_{22} & \cdots & \mathbf{k}_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ \mathbf{k}_{N1} & \mathbf{k}_{N2} & \cdots & \mathbf{k}_{NN} \end{bmatrix}_A \cdot \begin{bmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \\ \vdots \\ \mathbf{u}_N \end{bmatrix}_A = \begin{bmatrix} \mathbf{f}_1 \\ \mathbf{f}_2 \\ \vdots \\ \mathbf{f}_N \end{bmatrix}_A \quad \dots\dots\dots(2.2)$$

$$\mathbf{K}_B \cdot \mathbf{U}_B = \mathbf{F}_B \quad \text{or}$$

$$\begin{bmatrix} \mathbf{k}_{11} & \mathbf{k}_{12} & \cdots & \mathbf{k}_{1M} \\ \mathbf{k}_{21} & \mathbf{k}_{22} & \cdots & \mathbf{k}_{2M} \\ \vdots & \vdots & \ddots & \vdots \\ \mathbf{k}_{M1} & \mathbf{k}_{M2} & \cdots & \mathbf{k}_{MM} \end{bmatrix}_B \cdot \begin{bmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \\ \vdots \\ \mathbf{u}_M \end{bmatrix}_B = \begin{bmatrix} \mathbf{f}_1 \\ \mathbf{f}_2 \\ \vdots \\ \mathbf{f}_M \end{bmatrix}_B \quad \dots\dots\dots(2.3)$$

In both equations, K_A and K_B , are singular. For two dimensional problems the subsystem models allow “free body” motion with two orthogonal deflection and one rotation. So the ranks of K_A and K_B are $N-3$ and $M-3$. For every individual equation could be solved if boundary condition such as following provided

$$\mathbf{u}_A^\bullet = \begin{pmatrix} \mathbf{u}_a \\ \mathbf{u}_b \\ \mathbf{u}_c \end{pmatrix}_A = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}; \quad \text{and} \quad \mathbf{u}_B^\bullet = \begin{pmatrix} \mathbf{v}_a \\ \mathbf{v}_b \\ \mathbf{v}_c \end{pmatrix}_B = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad \dots\dots\dots(2.4)$$

Any 3 U value equal to zero could eliminate the rank of K by 3, and make it non-singular. Now we combine both components together without any re-mesh work. From this equation, we can see that it contains all of the old elements and nodes.

$$\mathbf{K}_{\text{sys}} \cdot \mathbf{u}_{\text{sys}} = \mathbf{f}_{\text{sys}} \dots\dots\dots (2.5)$$

in which

$$\mathbf{K}_{\text{sys}} = \begin{pmatrix} \mathbf{K}_A & \vdots & \mathbf{0} \\ \cdots & \cdot & \cdots \\ \mathbf{0} & \vdots & \mathbf{K}_B \end{pmatrix}; \quad (\mathbf{u}_{\text{sys}}) = \begin{pmatrix} \mathbf{u}_A \\ \cdots \\ \mathbf{u}_B \end{pmatrix}; \quad (\mathbf{f}_{\text{sys}}) = \begin{pmatrix} \mathbf{f}_A \\ \cdots \\ \mathbf{f}_B \end{pmatrix}$$

Obviously, the rank of $\mathbf{K}_{\text{sys}}$ could be available from following equation:

$$\text{Rank}(\mathbf{K}_{\text{sys}}) = \text{Rank}(\mathbf{K}_A) + \text{Rank}(\mathbf{K}_B) = (N-3) + (M-3) = N+M-6 \dots\dots\dots (2.6)$$

From equation 2.6 we know that the combined system needs at least 6 internal and external boundary condition to become solvable.

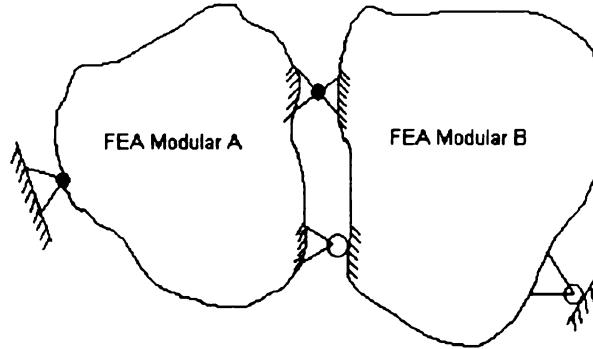


Figure 2.4. At least 6 constraint are needed for 2-D FEA modular modeling

2.3.2. Applying Boundary Displacements Constraints

The boundary displacement constraint needs to be defined at first, like most regular FEA procedures. This procedure will make the system equation (2.5) a full-rank linear equation. Therefore, this is the crucial step to make the system solvable. According to the previous analysis, the internal (boundary between 2 components) displacement constraint equation must be more than rank of 3.

Assume the following equation:

$$(\mathbf{u}_A)_{n \times n} = \mathbf{S}_{m \times n} \cdot (\mathbf{u}_B)_{m \times m} \quad \dots\dots\dots (2.7)$$

is the boundary displacements constraint equation, we need to make sure that:

$$\text{Rank}(\mathbf{S}_{m \times n}) \geq 3 \quad \dots\dots\dots (2.8)$$

The simulation result with the FEA model is a linear approximation for real situation [Zienkiewicz, O. C. and Taylor, R. L., 1989]. Suppose we are using linear elements, the boundary displacements constraint analysis should be based on every nodal value on both components boundaries linear approximation. The finite element method provides a displacement value within every element only according to the nodal value in itself. Therefore the boundary constraint equation could be found by the local boundary elements interpolation. Figure 2.4 demonstrates the mean of the local boundary elements nodal displacement relationship.

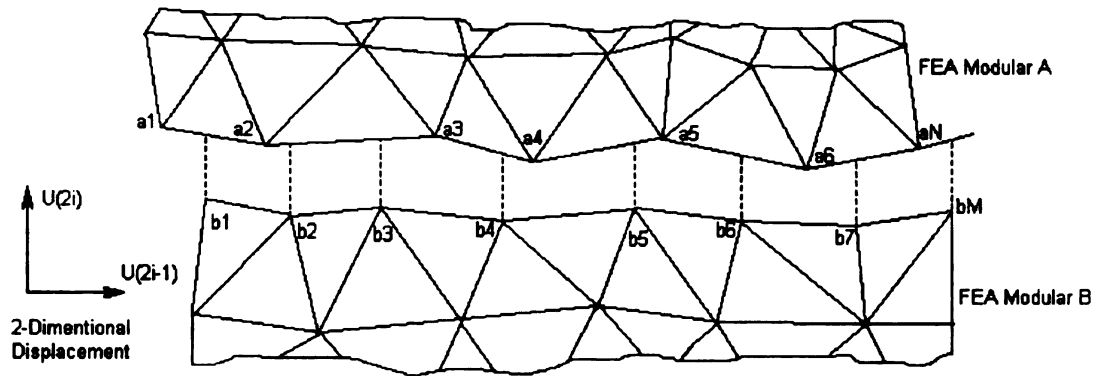


Figure 2.5. Nodal displacements along the boundary elements

The boundary displacement in Figure 2.4 at node “i” was defined as “u_{2i-1}” and “u_{2i}”. If the boundary nodes on modular A and B is a₁, a₂,... a_N and b₁, b₂,... b_M, then the boundary displacement constraint equation could be written as.

$$\begin{pmatrix} u_{b1} & u_{b2} & \cdots & u_{bM} \end{pmatrix}_B^T = S \cdot \begin{pmatrix} u_{a1} & u_{a2} & \cdots & u_{aN} \end{pmatrix}_A^T \dots\dots\dots(2.9)$$

specially, the 2 - D equation situation :

$$\begin{pmatrix} u_{2b1-1} & u_{2b1} & u_{2b2-1} & u_{2b2} & \cdots & u_{2bM-1} & u_{2bM} \end{pmatrix}_B^T = S_{2M \times 2N} \cdot \begin{pmatrix} u_{2a1-1} & u_{2a1} & u_{2a2-1} & u_{2a2} & \cdots & u_{2aN-1} & u_{2aN} \end{pmatrix}_A$$

In most cases the rank of equation (2.9) will be more than 3, otherwise at least one component has only one node on the boundary. So, if we can develop boundary displacement constraint relationships, equation (2.9) will have enough constraints as required. For both subjects (Fig. 2.5), the nodal displacement value along the boundary should satisfy the boundary interpolation [Langhe, K.De, 1995].

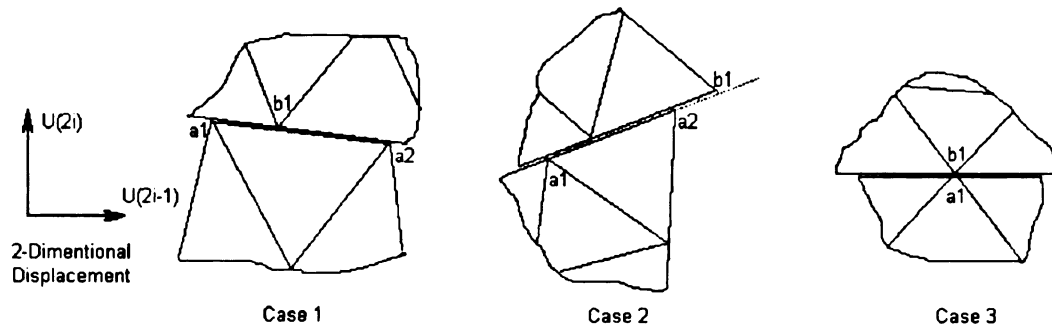


Figure 2.6. Nodal compatible situation along the boundary elements

Figure 2.5 shows the position of node “b1” has only 3 possibilities relative to the corresponding elements. Case 1 is located on the edge of corresponding boundary element. Case 2 is located on the straight extended line of the edge of corresponding boundary element. Case 3 is coincided with one node of the corresponding boundary element.

Independent of which component has more boundary nodes than another, both sides of the port can be used to develop boundary displacement constraint equations. In other words, equation 2.9 has 2 equivalent forms and they are equivalent constraints. Here we present results for nodes on module B along the boundary.

CASE 1. From the definition of finite element, any internal point value can be represented by the linear transformation of nodal values [Segerlind, L. J., 1984]. For the boundary nodes in another component, we also use similar linear interpolation approach to figure out their values [Langhe, K.De., 1995 and Kothnur, V.S. , Yu Xie, 1999]. In this paper, the discussion was focused in the basic theoretical solution for two-dimensional, straight-line boundary and three degree of freedom elastic stress-strain problem. So we can apply the similar approach from reference [Langhe, K.De., 1995].

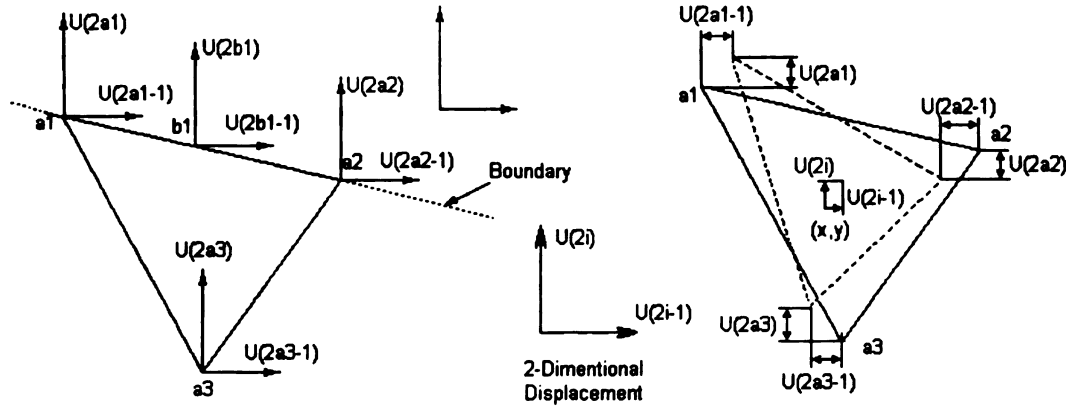


Figure 2.7. Nodal displacements in the boundary elements

As it shows in figure 2.6, we can get the shape function of an arbitrary point in the element which has the coordinates of (x, y) [Yang, T.Y. 1986]:

$$\begin{aligned} u_{2i}(x, y) &= u_{2a1} + c_1 x + c_2 y \\ u_{2i-1}(x, y) &= u_{2a1-1} + c_3 x + c_4 y \end{aligned} \quad \dots\dots\dots (2.10)$$

The c_1, c_2, c_3, c_4 is determined by the shape of element and can be derived from the following equation:

$$\begin{pmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{pmatrix} = \frac{1}{a_j b_k - a_k b_j} \begin{bmatrix} b_j - b_k & 0 & b_k & 0 & -b_j & 0 \\ a_k - a_j & 0 & -a_k & 0 & a_j & 0 \\ 0 & b_j - b_k & 0 & b_k & 0 & -b_j \\ 0 & a_k - a_j & 0 & -a_k & 0 & a_j \end{bmatrix} \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \\ u_{2a2-1} \\ u_{2a2} \\ u_{2a3-1} \\ u_{2a3} \end{pmatrix} \dots\dots\dots (2.11)$$

From equation (2.10) and (2.11), we get a general interpolation form for any point located in the boundary element. Because the values of c_1, c_2, c_3, c_4 are determined linearly only by properties of the corresponding boundary element, we could conclude that the u_{2b1-1} and u_{2b1} in (2.10) are linear functions of $u_{2a1}, u_{2a1-1}, u_{2a2},$ and $u_{2a2-1}, u_{2a3}, u_{2a3-1}$ correspondingly.

The transformation matrix S_1 , S_2 in the upper equation is decided by the corresponding boundary element and the x , y coordinate value of node “b1”. After we figure out all of the equations (2.12) for all boundary nodes on component B, we combine them and get the equation (2.9).

$$\begin{aligned} u_{2b1-1} &= S_1 \cdot (u_{2a1-1} \quad u_{2a2-1} \quad u_{2a3-1})^T \\ u_{2b1} &= S_2 \cdot (u_{2a1} \quad u_{2a2} \quad u_{2a3})^T \dots\dots\dots (2.12) \end{aligned}$$

In this paper, we are discussing the straight-line boundary problem. So all boundary node “b1”, “b2”, “bM” is locate on the edge of boundary elements. In this situation, the equation (2.12) could be simplified and decided only by nodal value on boundary edge [Segerlind, L. J., 1984].

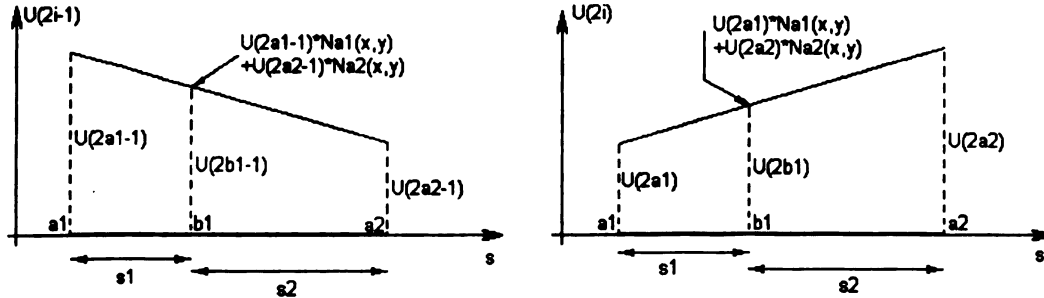


Figure 2.8. Linear transformation on the boundary elements (case 1: interpolation)

Figure 2.7 shows the linear interpolation of node “b1”. Suppose the distance from node “b1” and node “a1” is s_1 , and from node “b1” and node “a2” is s_2 , which could be derived from the their coordinate values, we can get the linear interpolation equation from the shape function [Langhe, K.De., 1995]:

$$\begin{aligned} u_{2b1-1} &= N_{a1}(s) \cdot u_{2a1-1} + N_{a2}(s) \cdot u_{2a2-1} = \frac{s_2}{s_1 + s_2} \cdot u_{2a1-1} + \frac{s_1}{s_1 + s_2} \cdot u_{2a2-1} \\ u_{2b1} &= N_{a1}(s) \cdot u_{2a1} + N_{a2}(s) \cdot u_{2a2} = \frac{s_2}{s_1 + s_2} \cdot u_{2a1} + \frac{s_1}{s_1 + s_2} \cdot u_{2a2} \dots\dots\dots (2.13) \end{aligned}$$

The distance from node “b1” and node “a1”, which is s1, and from node “b1” and node “a2”, which is s2, could be derived from following coordinate computation.

$$\begin{aligned} s_1 &= \sqrt{(x_{b1} - x_{a1})^2 + (y_{b1} - y_{a1})^2} \\ s_2 &= \sqrt{(x_{b1} - x_{a2})^2 + (y_{b1} - y_{a2})^2} \quad \dots\dots\dots (2.14) \end{aligned}$$

It is necessary to mention that the s1 and s2 values are always positive in case 1.

CASE 2. In this case, node “b1” is located on the straight extended line of the edge of corresponded boundary element. This happens because node “b1” is the terminal point on the boundary, and it does not contact with the corresponded component. In the following figure (2.8), we can see that equation (2.13) is still valid. The only thing that needs to be modified is that we should to re-define s1 and s2 according to the situation of “b1” location. If the “b1” location comes with the sequence of “a1, a2, b1”, the s1 and s2 should be:

$$\begin{aligned} s_1 &= \sqrt{(x_{b1} - x_{a1})^2 + (y_{b1} - y_{a1})^2} \\ s_2 &= -\sqrt{(x_{b1} - x_{a2})^2 + (y_{b1} - y_{a2})^2} \quad \dots\dots\dots (2.15) \end{aligned}$$

If the “b1” location comes with the sequence of “b1, a1, a2”, the s1 and s2 should be:

$$\begin{aligned} s_1 &= -\sqrt{(x_{b1} - x_{a1})^2 + (y_{b1} - y_{a1})^2} \\ s_2 &= \sqrt{(x_{b1} - x_{a2})^2 + (y_{b1} - y_{a2})^2} \quad \dots\dots\dots (2.16) \end{aligned}$$

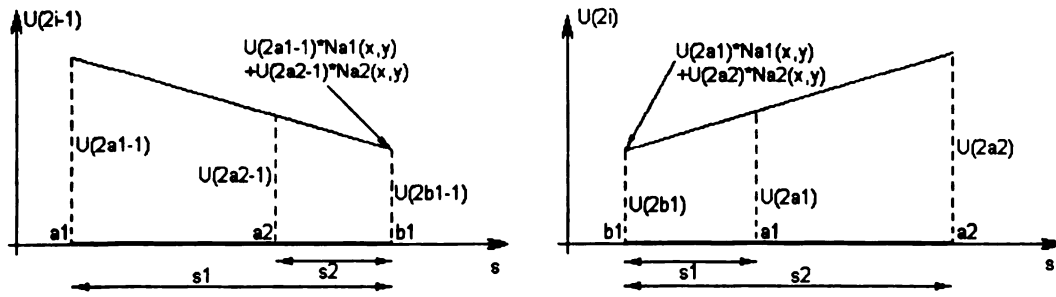


Figure 2.9. Linear transformation on the boundary elements (case 2: extrapolation)

CASE 3. In this case, node “b1” is located on the exact same nodal point of corresponded element. Clearly, this is a “nodal compatibility” situation. It is easy to conclude that if “b1” is located on same nodal point of “a1”

$$\begin{aligned}
 u_{2b1-1} &= u_{2a1-1} = 1 \cdot u_{2a1-1} + 0 \cdot u_{2a2-1} = \frac{s_2}{s_1 + s_2} \cdot u_{2a1-1} + \frac{s_1}{s_1 + s_2} \cdot u_{2a2-1} \\
 u_{2b1} &= u_{2a1} = 1 \cdot u_{2a1} + 0 \cdot u_{2a2} = \frac{s_2}{s_1 + s_2} \cdot u_{2a1} + \frac{s_1}{s_1 + s_2} \cdot u_{2a2} \\
 s_1 &= 0, \dots\dots\dots (2.17)
 \end{aligned}$$

If “b1” is located on same nodal point of “a2”

$$\begin{aligned}
 u_{2b1-1} &= u_{2a2-1} = 0 \cdot u_{2a1-1} + 1 \cdot u_{2a2-1} = \frac{s_2}{s_1 + s_2} \cdot u_{2a1} + \frac{s_1}{s_1 + s_2} \cdot u_{2a2} \\
 u_{2b1} &= u_{2a2} = 0 \cdot u_{2a1} + 1 \cdot u_{2a2} = \frac{s_2}{s_1 + s_2} \cdot u_{2a1} + \frac{s_1}{s_1 + s_2} \cdot u_{2a2} \\
 s_2 &= 0, \dots\dots\dots (2.18)
 \end{aligned}$$

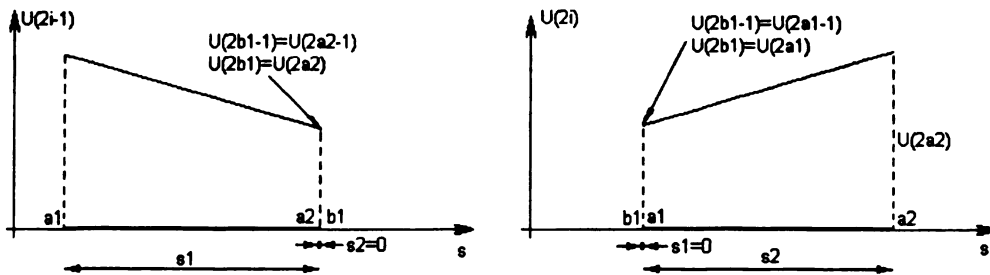


Figure 2.10. Linear transformation on the boundary elements (case 3: equality)

Therefore, equation (2.13) is valid too for this case. When the condition of “s1=0 or s1=0” was represented in (2.13), the output gets the simplified result as equation (2.17) or (2.18). Overall, we could conclude that in the 2-dimensional elastic problem, the displacement constraint equation along the straight-line boundary is equation (2.13). Based on different nodal location, the interpolation parameter s1 and s2 should be computed according to equations (2.14), (2.15), (2.16), (2.17) or (2.18).

From these equations, we could get all interpolation functions for “b1, b2,...bM” in figure 2.4. In equation (2.19), sb11 and sb12 represent the s1 and s2 value for “b1”, and Lb1 represent s1+s2 value for “b1”,

$$\begin{aligned}
 u_{2b1-1} &= \frac{s_2^{b1}}{L_{b1}} \cdot u_{2a1-1} + \frac{s_1^{b1}}{L_{b1}} \cdot u_{2a2-1}, & u_{2b1} &= \frac{s_2^{b1}}{L_{b1}} \cdot u_{2a1} + \frac{s_1^{b1}}{L_{b1}} \cdot u_{2a2} \\
 u_{2b2-1} &= \frac{s_2^{b2}}{L_{b2}} \cdot u_{2a2-1} + \frac{s_1^{b2}}{L_{b2}} \cdot u_{2a3-1}, & u_{2b2} &= \frac{s_2^{b2}}{L_{b2}} \cdot u_{2a2} + \frac{s_1^{b2}}{L_{b2}} \cdot u_{2a3} \\
 &\vdots & &\vdots \\
 u_{2bM-1} &= \frac{s_2^{bM}}{L_{bM}} \cdot u_{2a(N-1)-1} + \frac{s_1^{bM}}{L_{bM}} \cdot u_{2aN-1}, & u_{2bM} &= \frac{s_2^{bM}}{L_{bM}} \cdot u_{2a(N-1)} + \frac{s_1^{bM}}{L_{bM}} \cdot u_{2aN} \quad \dots\dots (2.19)
 \end{aligned}$$

If we define displacements by node as following:

$$\begin{aligned}
 \mathbf{u}_{b1} &= \begin{pmatrix} u_{2b1-1} \\ u_{2b1} \end{pmatrix}, & \mathbf{u}_{b2} &= \begin{pmatrix} u_{2b2-1} \\ u_{2b2} \end{pmatrix}, & \dots\dots \mathbf{u}_{bM} &= \begin{pmatrix} u_{2bM-1} \\ u_{2bM} \end{pmatrix} \\
 \mathbf{u}_{a1} &= \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \end{pmatrix}, & \mathbf{u}_{a2} &= \begin{pmatrix} u_{2a2-1} \\ u_{2a2} \end{pmatrix}, & \dots\dots \mathbf{u}_{aN} &= \begin{pmatrix} u_{2aN-1} \\ u_{2aN} \end{pmatrix}
 \end{aligned}$$

Equation group (2.19) could be gotten if all boundary nodal coordinate values are available. Also, the equations of (2.19) could be re-written as the boundary displacement constraint equation (2.20).

$$\begin{pmatrix} \mathbf{u}_{b1} \\ \mathbf{u}_{b2} \\ \vdots \\ \mathbf{u}_{bM} \end{pmatrix} = \begin{pmatrix} \mathbf{C}_{b1^*a1} & \mathbf{C}_{b1^*a2} & 0 & \vdots & 0 & 0 \\ 0 & \mathbf{C}_{b2^*a2} & \mathbf{C}_{b2^*a3} & \vdots & 0 & 0 \\ \dots & \dots & \dots & \vdots & \dots & \dots \\ 0 & 0 & 0 & \vdots & \mathbf{C}_{bM^*a(N-1)} & \mathbf{C}_{bM^*aN} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{u}_{a1} \\ \mathbf{u}_{a2} \\ \mathbf{u}_{a3} \\ \vdots \\ \mathbf{u}_{a(N-1)} \\ \mathbf{u}_{aN} \end{pmatrix} \quad \dots \quad (2.20)$$

In upper constraint equation, the constraint elements was defined as:

$$\mathbf{C}_{b1^*a1} = \begin{pmatrix} \frac{s_2^{b1}}{L_{b1}} & 0 \\ 0 & \frac{s_2^{b1}}{L_{b1}} \end{pmatrix}; \quad \mathbf{C}_{b1^*a2} = \begin{pmatrix} \frac{s_1^{b1}}{L_{b1}} & 0 \\ 0 & \frac{s_1^{b1}}{L_{b1}} \end{pmatrix}; \quad \dots \dots \mathbf{C}_{bX^*aY} = \begin{pmatrix} \frac{s_i^{bX}}{L_{bX}} & 0 \\ 0 & \frac{s_i^{bX}}{L_{bX}} \end{pmatrix};$$

$i=1 \text{ or } 2$; decided by aY position

Or the equation (2.20) could be write as:

$$\begin{pmatrix} u_{2b1-1} \\ u_{2b1} \\ u_{2b2-1} \\ u_{2b2} \\ \vdots \\ u_{2bM-1} \\ u_{2bM} \end{pmatrix} = \begin{pmatrix} \frac{s_2^{b1}}{L_{b1}} & 0 & \frac{s_1^{b1}}{L_{b1}} & 0 & 0 & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & \frac{s_2^{b1}}{L_{b1}} & 0 & \frac{s_1^{b1}}{L_{b1}} & 0 & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & 0 & \frac{s_2^{b2}}{L_{b2}} & 0 & \frac{s_1^{b2}}{L_{b2}} & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{s_2^{b2}}{L_{b2}} & 0 & \frac{s_1^{b2}}{L_{b2}} & \dots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots & \frac{s_2^{bM}}{L_{bM}} & 0 & \frac{s_1^{bM}}{L_{bM}} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots & 0 & \frac{s_2^{bM}}{L_{bM}} & 0 & \frac{s_1^{bM}}{L_{bM}} \end{pmatrix} \cdot \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \\ u_{2a2-1} \\ u_{2a2} \\ u_{2a3-1} \\ u_{2a3} \\ \vdots \\ u_{2a(N-1)-1} \\ u_{2a(N-1)} \\ u_{2aN-1} \\ u_{2aN} \end{pmatrix}$$

2.3.3. Applying Boundary Forces Constraints

Boundary force constraints analysis is required to solve the system equation. From the description of §2.3.1, we know that after boundary displacement constraint equations convert the system equation to full rank, the boundary force constraint equations and initial conditions are also required to solve the system. In fact, the boundary force constraints equation is very important to enable the system equation to get the final solution. The corresponding mathematical description is, if we need to solve $\mathbf{K}'_{m \times n} \cdot \mathbf{u}_{b \times m} = \mathbf{f}_{b \times m}$; $n > m$. After applying the system boundary displacement constraints equation, the equality of the rank of stiffness matrix $\mathbf{K}'$ and $\mathbf{f}'$ can only guarantee the solution exist, but it can not guarantee that the solution is unique. It is absolutely necessary to have the force constraint equations for the system equation.

The physical sense of this statement is: both the boundary displacement and force constraint equations are required to ensure that the total power (for dynamic model) or work (for static model) flow through the boundary is conserved. Modular modeling ports provide the power constraint. Power is conserved across modular ports because modular ports are power (or work) transfer mechanisms. The power at the connected modular modeling ports always sum to zero. For the elastic stress-strain problem, we need not make dynamic analysis work. Therefore, the work at the connected modular modeling ports always sum to zero.

The work along the boundary defined as : $\delta W = \sum_{i=1}^{aN} W_i = \sum_{i=1}^{aN} (\mathbf{f}_i \cdot \mathbf{u}_i) = \mathbf{f}^T \cdot \mathbf{u}$

Because work across the connection of component A and component B will be

$$\delta W_A + \delta W_B = 0 \quad \Rightarrow \quad \delta W_A = -\delta W_B \dots\dots\dots(2.21)$$

$$\therefore (\mathbf{f}^T \cdot \mathbf{u})_A + (\mathbf{f}^T \cdot \mathbf{u})_B = \mathbf{f}_A^T \cdot \mathbf{u}_A + \mathbf{f}_B^T \cdot \mathbf{u}_B = 0$$

from Equation (2.20), we get

$$(\mathbf{f}_A^T \cdot \mathbf{u}_A) + (\mathbf{f}_B^T \cdot \mathbf{C} \cdot \mathbf{u}_A) = 0 \Rightarrow (\mathbf{f}_A^T + \mathbf{f}_B^T \cdot \mathbf{C}) \cdot \mathbf{u}_A = 0$$

Because the power conservation is valid for any displacement under elastic range, we can say that the equation (2.22) is valid for any of those displacement.

Therefore, we could conclude

$$\mathbf{f}_A^T + \mathbf{f}_B^T \cdot \mathbf{C} = 0 \quad \text{and} \quad \mathbf{f}_A + \mathbf{C} \cdot \mathbf{f}_B = 0$$

Especially, for 2-dimensional case:

$$\delta W_A = \Sigma \mathbf{f}_{AM} \cdot \mathbf{u}_{AN} = \begin{pmatrix} f_{2a1} \\ f_{2a1-1} \\ f_{2a2-1} \\ f_{2a2} \\ \vdots \\ f_{2aN-1} \\ f_{2aN} \end{pmatrix}^T \cdot \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \\ u_{2a2-1} \\ u_{2a2} \\ \vdots \\ u_{2aN-1} \\ u_{2aN} \end{pmatrix}; \quad \delta W_B = \Sigma \mathbf{f}_{BM} \cdot \mathbf{u}_{BM} = \begin{pmatrix} f_{2b1} \\ f_{2b1-1} \\ f_{2b2-1} \\ f_{2b2} \\ \vdots \\ f_{2bM-1} \\ f_{2bM} \end{pmatrix}^T \cdot \begin{pmatrix} u_{2b1-1} \\ u_{2b1} \\ u_{2b2-1} \\ u_{2b2} \\ \vdots \\ u_{2bM-1} \\ u_{2bM} \end{pmatrix}$$

The equation (2.1) should be the beginning step of the constraint equation for nodal forces. From equation (2.1), we get

$$\delta W_A + \delta W_B = \Sigma \mathbf{f}_{AN} \cdot \mathbf{u}_{AN} + \Sigma \mathbf{f}_{BM} \cdot \mathbf{u}_{BM} = \begin{pmatrix} f_{2a1} \\ f_{2a1-1} \\ f_{2a2-1} \\ f_{2a2} \\ \vdots \\ f_{2aN-1} \\ f_{2aN} \end{pmatrix}^T \cdot \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \\ u_{2a2-1} \\ u_{2a2} \\ \vdots \\ u_{2aN-1} \\ u_{2aN} \end{pmatrix} + \begin{pmatrix} f_{2b1} \\ f_{2b1-1} \\ f_{2b2-1} \\ f_{2b2} \\ \vdots \\ f_{2bM-1} \\ f_{2bM} \end{pmatrix}^T \cdot \begin{pmatrix} u_{2b1-1} \\ u_{2b1} \\ u_{2b2-1} \\ u_{2b2} \\ \vdots \\ u_{2bM-1} \\ u_{2bM} \end{pmatrix} = 0$$

When we place equation (2.20) into the above equation,

$$\delta W_A + \delta W_B = \begin{pmatrix} f_{2a1} \\ f_{2a1-1} \\ f_{2a2-1} \\ f_{2a2} \\ \vdots \\ f_{2aN-1} \\ f_{2aN} \end{pmatrix}^T \cdot \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \\ u_{2a2-1} \\ u_{2a2} \\ \vdots \\ u_{2aN-1} \\ u_{2aN} \end{pmatrix} + \begin{pmatrix} f_{2b1} \\ f_{2b1-1} \\ f_{2b2-1} \\ f_{2b2} \\ \vdots \\ f_{2bM-1} \\ f_{2bM} \end{pmatrix}^T \cdot \mathbf{S} \cdot \begin{pmatrix} u_{2a1-1} \\ u_{2a1} \\ u_{2a2-1} \\ u_{2a2} \\ \vdots \\ u_{2aN-1} \\ u_{2aN} \end{pmatrix} = 0 \quad \dots\dots\dots (2.22)$$

in which $\mathbf{S} = \begin{pmatrix} \frac{s_2^{b1}}{L_{b1}} & 0 & \frac{s_1^{b1}}{L_{b1}} & 0 & 0 & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & \frac{s_2^{b1}}{L_{b1}} & 0 & \frac{s_1^{b1}}{L_{b1}} & 0 & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & 0 & \frac{s_2^{b2}}{L_{b2}} & 0 & \frac{s_1^{b2}}{L_{b2}} & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{s_2^{b2}}{L_{b2}} & 0 & \frac{s_1^{b2}}{L_{b2}} & \dots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots & \frac{s_2^{bM}}{L_{bM}} & 0 & \frac{s_1^{bM}}{L_{bM}} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots & 0 & \frac{s_2^{bM}}{L_{bM}} & 0 & \frac{s_1^{bM}}{L_{bM}} \end{pmatrix}$

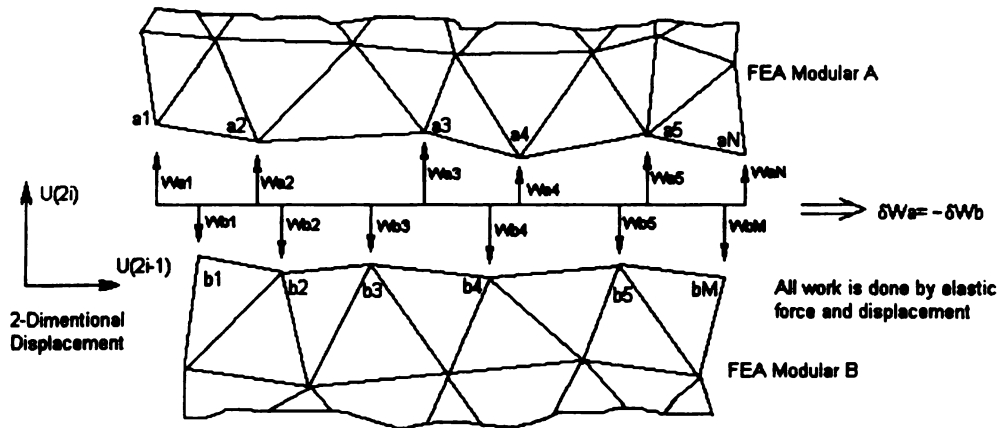


Figure 2.11. The power conservation along the Boundary

Figure 2.10 shows the sum of work transferred along the boundary equal ZERO under the situation of the elastic range.

2.3.4. Combine the Supplemental Boundary Constraint

From the boundary constraint equations (2.20) and (2.23), we get the total $2M$ constraint equations (if $M < N$) or $2N$ constraint equations (if $M > N$). From figure 2.12, we can see that when the system DOF is 0, there must be $M > 3$ and $N > 3$, and the system equation is mathematically non-singular.

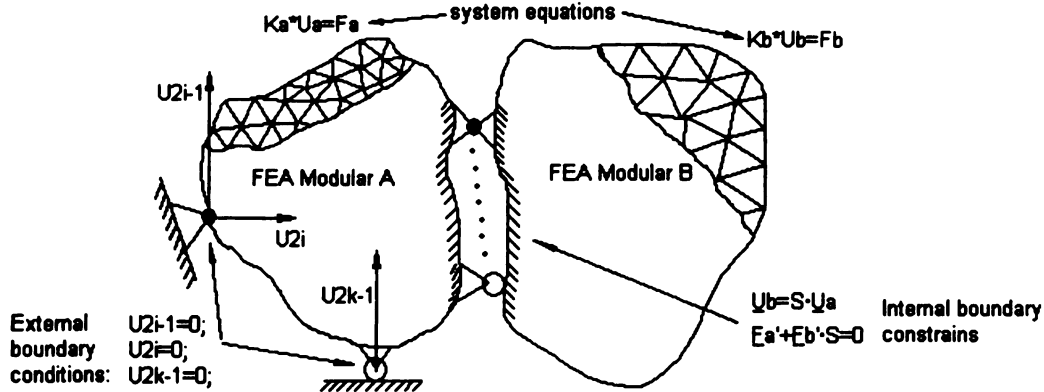


Figure 2.12. The “solvable system” and related equations.

For any solvable system like figure 2.12, the system equation, boundary condition and boundary constraint equation should be combined through the following procedures.

Step 1. Use the FEA models to get the combined system equation (2.5), the boundary displacement constraint equation (2.20), and the boundary force constraint equation (2.23). All of these constraint equations could come out from nodal coordinate values and I/O port shape. If K_A is $p \times p$, K_B is $q \times q$ matrix, then

$$\begin{pmatrix} \mathbf{K}_A & \vdots & \mathbf{0} \\ \cdots & \cdot & \cdots \\ \mathbf{0} & \vdots & \mathbf{K}_B \end{pmatrix} \begin{pmatrix} \mathbf{u}_A \\ \cdots \\ \mathbf{u}_B \end{pmatrix} = \mathbf{K}_{AB} \cdot \begin{pmatrix} \mathbf{u}_A \\ \cdots \\ \mathbf{u}_B \end{pmatrix} = \begin{pmatrix} \mathbf{f}_A \\ \cdots \\ \mathbf{f}_B \end{pmatrix} \dots\dots\dots (2.23)$$

K_{AB} is a $(p+q) \times (p+q)$ matrix.

Step 2. Apply the external boundary conditions and boundary displacement constraint equation (2.20) to transform system equation. This step will replace

$$\begin{aligned} [u_{2i-1} \quad u_{2i} \quad u_{2k-1}] &= [0 \quad 0 \quad 0]; \quad \text{and} \\ [U_{2b1-1} \quad U_{2b1} \quad U_{2b2-1} \quad U_{2b2} \quad \dots \quad U_{2bM-1} \quad U_{2bM}] &\int_B^T \\ &= S \cdot [U_{2a1-1} \quad U_{2a1} \quad U_{2a2-1} \quad U_{2a2} \quad \dots \quad U_{2aN-1} \quad U_{2aN}] \int_A^T \end{aligned}$$

the displacement of some boundary nodes. For example in figure 2.11 the external boundary conditions and boundary displacement constraint equation is

Using upper equations, we can get rid of u_{2i-1} ; u_{2i} ; u_{2k-1} ; and u_{2b1-1} ; u_{2b1} ; u_{2b2-1} ; u_{2b2} ; ... U_{2bM-1} and U_{2bM} from U_A and U_B , and get following equation (assume $N > M$):

$$\begin{pmatrix} u_A \\ \dots \\ u_B \end{pmatrix} = T_U \cdot \begin{pmatrix} u'_A \\ \dots \\ u'_B \end{pmatrix} \Rightarrow \begin{pmatrix} K_A & \vdots & 0 \\ \dots & \cdot & \dots \\ 0 & \vdots & K_B \end{pmatrix} \cdot T_U \cdot \begin{pmatrix} u'_A \\ \dots \\ u'_B \end{pmatrix} = K'_{AB} \cdot \begin{pmatrix} u'_A \\ \dots \\ u'_B \end{pmatrix} = \begin{pmatrix} f_A \\ \dots \\ f_B \end{pmatrix} \quad \dots \quad (2.24)$$

In which T_U is a $(p+q-3-M) \times (p+q)$ transformation matrix, and K'_{AB} is a $(p+q-3-M) \times (p+q)$ matrix that is transformed from K_{AB} .

Step 3. Move the unknown forces (external and internal boundary forces) to the left side, and leave the right side of the equation as a totally known force vector. For example, if f_j, f_k, f_l, f_m in $[F_A \ F_B]^T$ is unknown, we can simply replace it by "0". This will make the system equation become a standard linear algebra equation, in which T_F is a $(2M+3) \times (p+q)$ matrix and K''_{AB} is a $(p+q+M) \times (p+q)$ matrix.

$$K'_{AB} \cdot \begin{pmatrix} u'_A \\ \dots \\ u'_B \end{pmatrix} = \begin{pmatrix} f_A \\ \dots \\ f_B \end{pmatrix} \Rightarrow (K'_{AB} \quad \vdots \quad T_F) \cdot \begin{pmatrix} u'_A \\ \dots \\ u'_B \\ \dots \\ f'_{A+B} \end{pmatrix} = K''_{AB} \cdot \begin{pmatrix} u'_A \\ \dots \\ u'_B \\ \dots \\ f'_{A+B} \end{pmatrix} = \begin{pmatrix} f'_A \\ \dots \\ f'_B \end{pmatrix} \quad \dots \quad (2.25)$$

Step 4. Apply the boundary force constraint equation. Because all boundary forces in equation(2.23) is unknown forces that have been moved to left side, we can directly apply the boundary force constraint equation into the system equation.

$$\begin{pmatrix} \mathbf{f}_{2a1} \\ \mathbf{f}_{2a2} \\ \vdots \\ \mathbf{f}_{2aN} \end{pmatrix} + \begin{pmatrix} \mathbf{f}_{2b1} \\ \mathbf{f}_{2b2} \\ \vdots \\ \mathbf{f}_{2bM} \end{pmatrix} \cdot \mathbf{S}_F = \underline{\mathbf{0}} \Rightarrow \begin{pmatrix} \mathbf{K}'_{AB} & \vdots & \mathbf{T}_F \\ \cdots & \cdot & \cdots \\ \mathbf{0} & \vdots & \mathbf{S}'_F \end{pmatrix} \cdot \begin{pmatrix} \mathbf{u}'_A \\ \cdots \\ \mathbf{u}'_B \\ \cdots \\ \mathbf{f}'_{A+B} \end{pmatrix} = \mathbf{K}''_{AB} \cdot \begin{pmatrix} \mathbf{u}'_A \\ \cdots \\ \mathbf{u}'_B \\ \cdots \\ \mathbf{f}'_{A+B} \end{pmatrix} = \begin{pmatrix} \mathbf{f}'_A \\ \cdots \\ \mathbf{f}'_A \\ \cdots \\ \mathbf{0} \end{pmatrix} \cdots (2.26)$$

Because the boundary force could provide M constraint, the system equation matrix becomes a square matrix with full rank, which $\mathbf{S}'_F$ is a $(2M+3) \times (M)$ matrix, and $\mathbf{K}''_{AB}$ is a $(p+q+M) \times (p+q+M)$ matrix. Therefore, the final integrated system is

$$\begin{pmatrix} \mathbf{K}'_{AB} & \vdots & \mathbf{T}_F \\ \cdots & \cdot & \cdots \\ \mathbf{0} & \vdots & \mathbf{S}'_F \end{pmatrix} \cdot \begin{pmatrix} \mathbf{u}'_A \\ \cdots \\ \mathbf{u}'_B \\ \cdots \\ \mathbf{f}'_{A+B} \end{pmatrix} = \begin{pmatrix} \mathbf{f}'_A \\ \cdots \\ \mathbf{f}'_B \\ \cdots \\ \mathbf{0} \end{pmatrix} \cdots \cdots \cdots (2.27)$$

By solving equation (2.27), the system internal forces $\mathbf{f}'_{A+B}$ and displacement distribution $\mathbf{u}'_A$ and $\mathbf{u}'_B$ are available. In following chapter, I am going to demonstrate this approach with a example. Through this example, the system modeling and integrating details could be shown clearly.

2.4. Result and Discussion

Modular modeling system FEA equations can be developed from the subsystem equations, external boundary conditions, and internal constraint equations. If enough constraint equations are available and the DOF is zero, the modular modeling FEA equations can be solvable. The general form of the combined system equation (2.27) has been developed here for the two-dimensional elastic problem.

Modular modeling FEA equations could be realized automatically with a proper algorithm. The reason is that all of parameters used in combined systems come from individual FEA model and connection description, such as nodal coordinates and boundary conditions. This point is valuable for the modular modeling method because this means that the system could be integrated by the user's selection of parts and connection type.

Solution of the modular modeling formulation can be shown to be equivalent to solving the associated Lagrangian formulation [Appendix I, Radcliffe and Diaz, 2000]. Although the modular modeling formulation is derived based on only work constraint at the boundary, its solution is identical to the solution of the Lagrangian formulation which analysis minimum potential energy for the system. It is gratifying that the solution for the modular modeling method coincides with the traditional energy-based Lagrangian approach.

TWO-DIMENSIONAL EXAMPLE FOR ELASTIC STRAIN-STRESS PROBLEM

3.1. FEA static stress-strain analysis by modular modeling

To demonstrate the last chapter work about integrating the modular modeling FEA systems, a simple example for a 2-D elastic plane force-displacement problem will be solved by the approach and procedures in the chapter 2. After the combined system has been solved, the result was compared with same object but derived from FEM using two different mesh densities.

In figure 3.1, there is a combined plane will be a subject of load. The two triangular components have FEA models: plane A has a FEA stiffness matrix K_A , plane B has a FEA stiffness matrix K_B ,

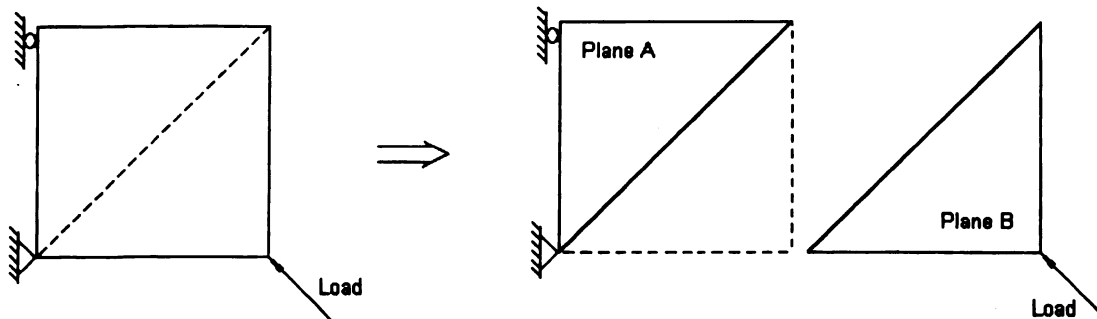


Figure 3.1. Two-Dimensional Example for Elastic Strain-Stress Problem

From figure 3.1, the subject has 3 external boundary conditions so it is 0 DOF, using the approach in last chapter and original equation (3.1) and (3.2), applying the external and internal boundary constraint (port) equations. The system equations are singular.

3.1.1 2-D Elastic Plane A:

We need to get the individual plane meshing and system equation. From the following procedures, it could be solved like any traditional examples [L. Segerlind, 1984]. The related system properties are: Young's modulus: $E = 15 \times 10^6$ (N/cm²); Poisson's ratio: $\mu = 1/4$; Plane thickness: $t = 0.1$ cm.

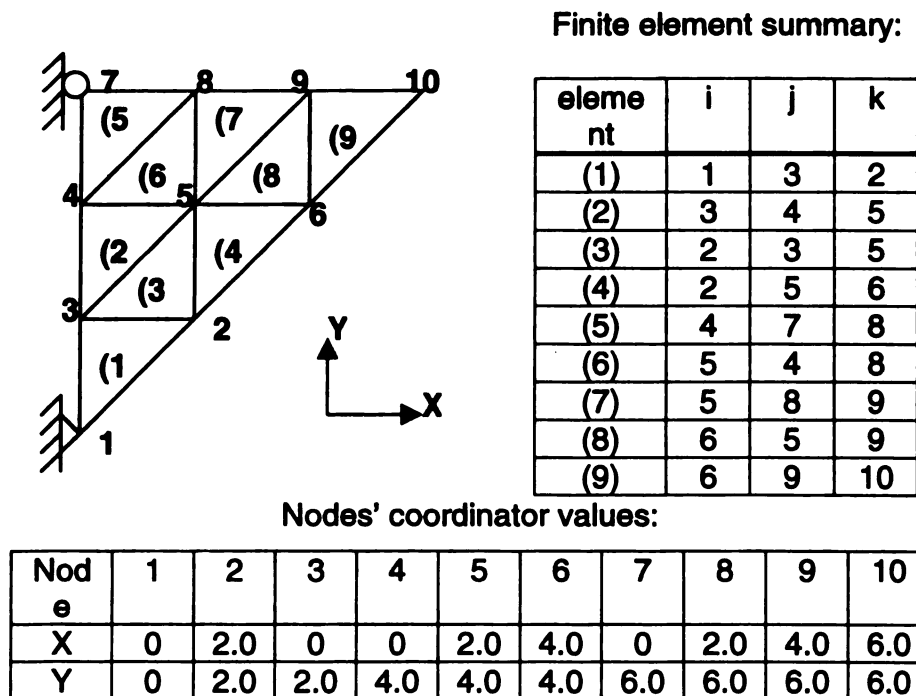


Figure 3.2 Plane A meshing details and nodal distribution

By the traditional method [L. Segerlind, 1984], we get all element stiffness matrixes in plane A: $[K(1)]$, $[K(2)]$, $[K(3)]$, $[K(4)]$, $[K(5)]$, $[K(6)]$ and combine them together as following. See the Appendix II for the detail of computer software.

$$K_A = \begin{bmatrix} 3 & 0 & 0 & 3 & -3 & -3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 8 & 2 & 0 & -2 & -8 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2 & 22 & 5 & -16 & -5 & 0 & 0 & -6 & -5 & 0 & 3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 3 & 0 & 5 & 22 & -5 & -6 & 0 & 0 & -5 & -16 & 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -3 & -2 & -16 & -5 & 22 & 5 & -3 & -3 & 0 & 5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -3 & -8 & -5 & -6 & 5 & 22 & -2 & -8 & 5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -3 & -2 & 22 & 5 & -16 & -5 & 0 & 0 & -3 & -3 & 0 & 5 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -3 & -8 & 5 & 22 & -5 & -6 & 0 & 0 & -2 & -8 & 5 & 0 & 0 & 0 & 0 \\ 0 & 0 & -6 & -5 & 0 & 5 & -16 & -5 & 44 & 10 & -16 & -5 & 0 & 0 & -6 & -5 & 0 & 5 & 0 \\ 0 & 0 & -5 & -16 & 5 & 0 & -5 & -6 & 10 & 44 & -5 & -6 & 0 & 0 & -5 & -16 & 5 & 0 & 0 \\ 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & -16 & -5 & 22 & 5 & 0 & 0 & 0 & 0 & -6 & -5 & 0 \\ 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 & -5 & -6 & 5 & 22 & 0 & 0 & 0 & 0 & -5 & -16 & 2 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & -2 & 0 & 0 & 0 & 0 & 11 & 5 & -8 & -3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & -8 & 0 & 0 & 0 & 0 & 5 & 11 & -2 & -3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 5 & -6 & -5 & 0 & 0 & -8 & -2 & 22 & 5 & -8 & -3 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 5 & 0 & -5 & -16 & 0 & 0 & -3 & -3 & 5 & 22 & -2 & -3 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 5 & -6 & -5 & 0 & 0 & -8 & -2 & 22 & 5 & -8 & -3 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 5 & 0 & -5 & -16 & 0 & 0 & -3 & -3 & 5 & 22 & -2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & -8 & -2 & 8 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 & -3 & -3 & 0 \end{bmatrix} \times 10^5$$

3.1.2. 2-D Elastic Plane B:

The meshing and system equation for plane B could be done by similar way. It has same static properties as plane A: Young's modulus: $E = 15 \times 10^6$ (N/cm²); Poisson's ratio: $\mu = 1/4$; Plane thickness: $t = 0.1$ cm.

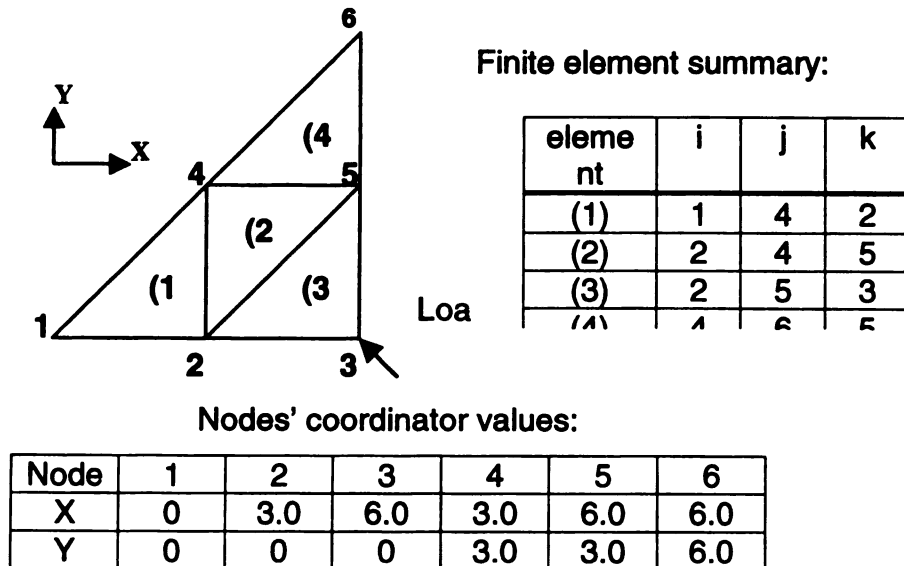


Figure 3.3. Plane B meshing details and nodal distribution

By the traditional method [L. Segerlind, 1984], we could get all element stiffness matrixes in plain A: [K(1)], [K(2)], [K(3)], [K(4)], and combine them together as following. See the Appendix III for the detail of computer software.

$$K_B = \begin{bmatrix} 8 & 0 & -8 & -2 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 \\ 0 & 3 & -3 & -3 & 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 \\ -8 & -3 & 22 & 5 & -8 & -2 & -6 & -5 & 0 & 5 & 0 & 0 \\ -2 & -3 & 5 & 22 & -3 & -3 & -5 & -16 & 5 & 0 & 0 & 0 \\ 0 & 0 & -8 & -3 & 11 & 5 & 0 & 0 & -3 & -2 & 0 & 0 \\ 0 & 0 & -2 & -3 & 5 & 11 & 0 & 0 & -3 & -8 & 0 & 0 \\ 0 & 3 & -6 & -5 & 0 & 0 & 22 & 5 & -16 & -5 & 0 & 2 \\ 2 & 0 & -5 & -16 & 0 & 0 & 5 & 22 & -5 & -6 & 3 & 0 \\ 0 & 0 & 0 & 5 & -3 & -3 & -16 & -5 & 22 & 5 & -3 & -2 \\ 0 & 0 & 5 & 0 & -2 & -8 & -5 & -6 & 5 & 22 & -3 & -8 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 3 & -3 & -3 & 3 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & -2 & -8 & 0 & 8 \end{bmatrix} \times 10^5$$

3.1.3. Using Modular Modeling To Integrate 2-D Elastic Components

We are using this example to demonstrate the details of the modular modeling FEA approach. We need to follow the steps one by one according to the description in previous chapter. The final solvable system equation is too large to be shown here. For detailed computing procedure and result, please review Appendix V.

3.2. Error Estimation and Analysis

To validate the result from §3.1, we need to compare it with a traditional model. Both higher and lower meshing resolution models are needed for the comparison to verify the modular modeling result. For a linear problem, higher resolution meshes yield a more accurate result. The modular model mesh has more resolution than then low resolution reference model but less resolution than the higher resolution model. We expect the modular model to be more accurate than the lower resolution mesh model and less accurate than the higher resolution mesh model. Comparing these models will provide useful information about the detail in the model and component connections at the boundary.

3.2.1. Standard Reference Models for Accuracy Comparison

Two standard reference FEA models (Fig. 3.4) will be used to evaluate the accuracy of the 13 element modular modeling example (Fig. 3.1). The higher resolution standard reference model has 72 elements while the lower resolution standard reference model has only eight elements. We use the same coordinate systems, load and FEA equations to ensure the accuracy of the formulations are directly comparable.

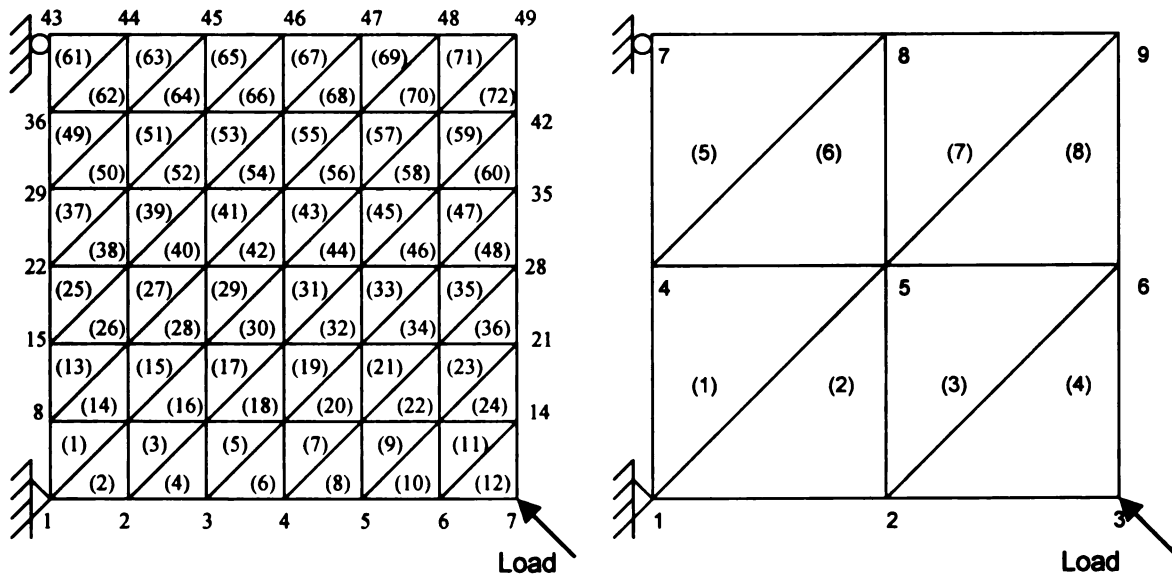


Figure 3.4. Re-meshed reference model with higher or lower meshing resolution

3.2.2. Modular Modeling Method Accuracy Comparison

Displacement contour lines were used for comparing the overall system accuracy. Displacement means the total displacement at every node. If a node has “a” displacement on X direction and “b” displacement on “Y” direction, its total displacement is $\left(\sqrt{a^2 + b^2}\right)$. Because we are focusing on the linear case and the load is same, the linear static elastic problem should have the same and repeatable result. Therefore this comparing is reliable.

Predicted displacement for the modular model were compared to the displacements predicted by the higher resolution standard reference model. EXCEL was used to compute and plot total model displacements (Fig. 3.5). The 3-D displacement difference surface (Fig. 3.6) shows the largest model differences are less than 10%. In particular, note that the modular model results are almost universally smaller displacement than the higher resolution reference model. This indicates that the modular model is stiffer than the higher resolution reference model.

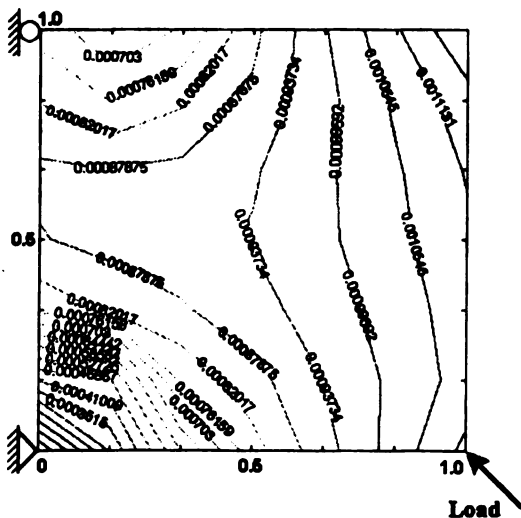


Figure 3.5(a). Displacement contour lines for reference model (7x7).

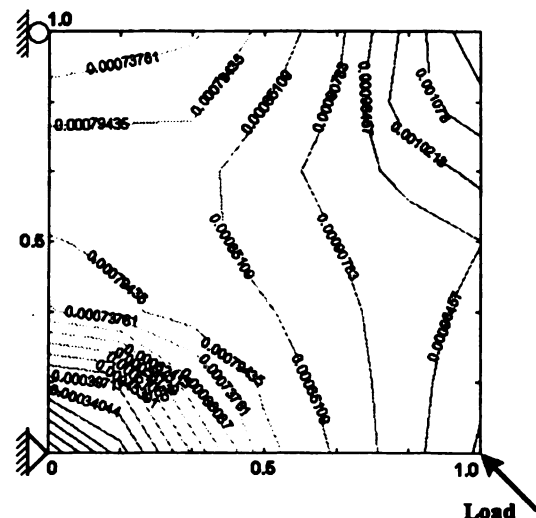


Figure 3.5(b). Displacement contour lines for the Joint Modular Model

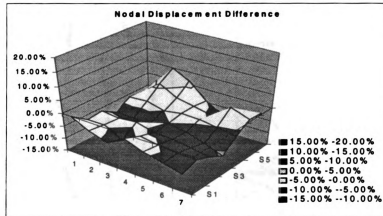


Figure 3.6. Displacement difference between the Joint Modular Model and reference model (7x7)

Predicted displacement for the modular model were next compared to the displacements predicted by the lower resolution standard reference model. EXCEL was again used to compute and plot total model displacements (Fig. 3.7). The 3-D displacement difference surface (Fig. 3.8) shows that in this case, the largest model differences approach 25%. In particular, note that the modular model results are almost universally larger displacement than the lower resolution reference model. In contrast to the comparison with the higher resolution reference model, this results indicates that the modular model is less stiff than the lower resolution reference model.

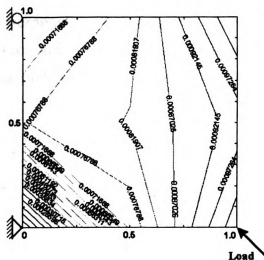


Figure 3.7(a). Displacement contour lines for reference model (3x3).

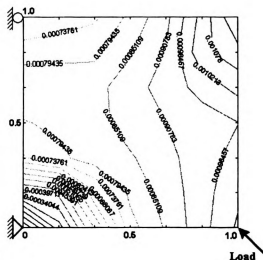


Figure 3.7(b). Displacement contour lines for the Joint Modular Model

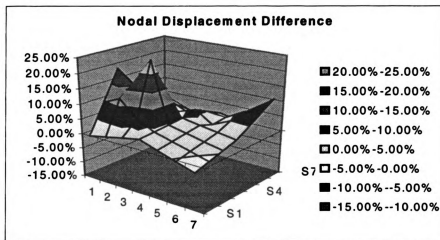


Figure 3.8. Displacement difference between the Joint Modular Model and reference model (3x3)

Comparison of the three models (fig. 3.9) indicates that nodal displacements along the boundary of plane A and B uniformly converge to the higher resolution standard model result as the number of elements in the model increase. Because the number of elements in the modular model lies between the lower and higher resolution reference models, its displacements are bracketed by the

two reference model displacements. As expected, the model stiffness uniformly decreases as the number of elements is increased, indicating the modular model accuracy is comparable to traditional FEM models with the same mesh resolution.

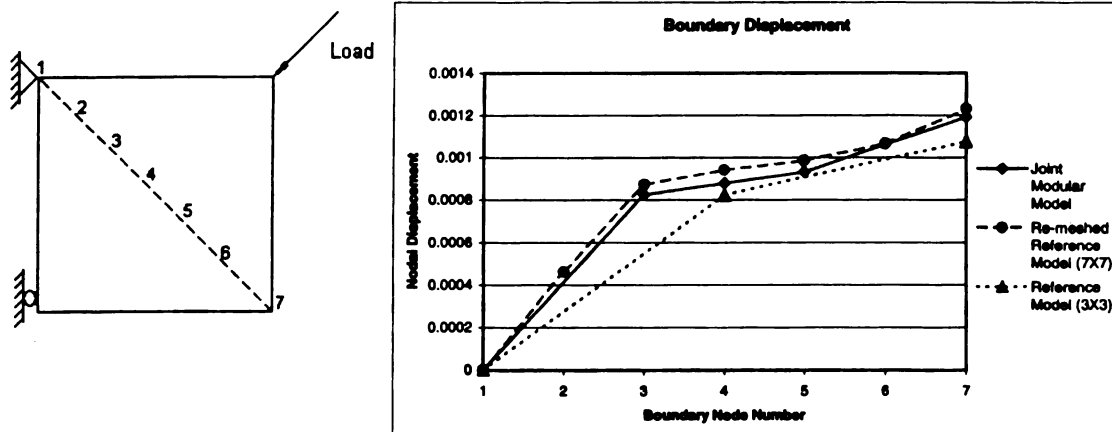


Figure 3.9. The difference of nodal displacements along the boundary

CONCLUSIONS

4.1 Contributions

Application of the modular modeling method in FEA has been discussed; its advantages and efficiency have been shown in examples and numerical results. Most importantly, the 2-dimensional boundary constraint equation that is used to connect FEA modules has been developed and verified in examples. Modular modeling method is shown to use the original FEA data, removing the need for re-meshing and verification, so it saves cost and time.

From a general description and example, we believe that the equation (2.5), (2.20) and (2.23) is the proper way to get the combined system equation and boundary constraint equations. For actual engineering problem, just like the example, we can sequentially use the equations (2.5), (2.24), (2.25), (2.26) and (2.27) to get the solution. For 2-Dimensional elastic stress-strain problem, if the system have 0 DOF, we can conclude that the developed system equation of FEA modular model is non-singular and could be solved with the boundary constraint equations, external boundary conditions and loads.

The modular model has good computational accuracy and quality that is close to the re-meshed higher resolution model and is much better than the relatively low-resolution model. The modular modeling method was shown to provide accuracy comparable to traditional FEA models with comparable mesh resolution. The ability to achieve this accuracy without remeshing is a significant improvement to FEA modeling technology.

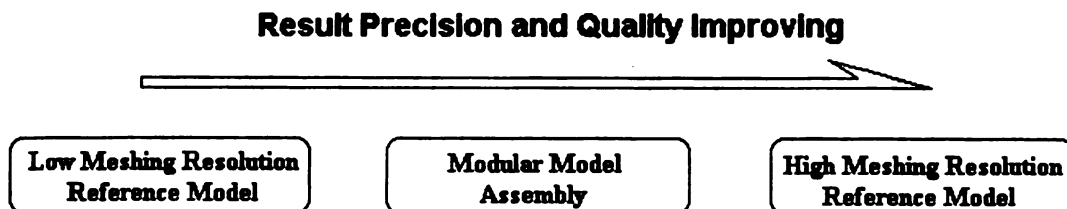


Figure 4.1. Comparing the Modular Model and other models

4.2 Future work

The equation developed in this paper can be only applied to a simple 2-D problem, the accuracy of this approach will be decided by the equality of boundary equation to make this approach practically useful in industrial applications, we must develop the following topics:

- Convergence of the stiffness matrix
- All kinds of meshes
- The 3 dimensional flat boundary
- The 3 dimensional arbitrary boundaries
- The Boundary Constraint Equation development automation,
- Apply modular modeling method into the nonlinear FEA problem.

There are 2 ways to develop the boundary constraint equation. From equations (2.20) and (2.23), we can see that component B was treated as the subject of interpolation from component A. It is easy to understand, the contrary way could also provide similar constraint equations. Those equations should also be enough to make the system solvable. But the problem is, “which one is better?” I believe that this is a more complex and practical problem and needs research through great number of actual case, cause the meshing and constraint have so many kind of types that it is very hard to simply identify them.

Theoretically, the boundary interface input / output equation between components will be only decided by their shape function. Therefore, the automatically formulation should be practically realized by proper programming. This feature should be one of the critical block in the future modular modeling FEA software.

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APPENDIX

- I. Lagrangian FEM Modular Modeling notes
- II. The Mathematica® program for plant A. system stiffness matrix computation.
Filename: *Plane_a.nb*
- III. The Mathematica® program for plant B. system stiffness matrix computation.
Filename: *Plane_b.nb*
- IV. The Mathematica® program for re-meshed plant. Work as the reference for error analysis.
Filename: *Plane_all.nb*
- V. The Mathematica® program for combined plant. System integration.
Filename: *Plane_a+b_connect_6.nb*
- VI. The MS Excel data table for the result of FEA computation and error analysis.

TO: Mr. Xin Li
 From: Clark Radcliffe
 Subject: Lagrangian FEM Modular Modeling
 Date: April 7, 1999

Per discussion with A. Diaz 4/6/00 at your MS defense, I have outlined below the relationship between the Modular Modeling Method and the Lagrangian approach to the joining of two bodies with constraints.

The Modular FEM Modeling Problem (Xin Li MS thesis, 4/6/00) is formulated as

$$\begin{bmatrix} \mathbf{K}_A & \mathbf{0} \\ \mathbf{0} & \mathbf{K}_B \end{bmatrix} \cdot \begin{bmatrix} \mathbf{u}_A \\ \mathbf{u}_B \end{bmatrix} = \begin{bmatrix} \mathbf{f}_A \\ \mathbf{f}_B \end{bmatrix} \quad (1a)$$

subject to the constraints

$$\hat{\mathbf{u}}_B = \hat{\mathbf{S}} \cdot \hat{\mathbf{u}}_A \text{ and} \quad (1b)$$

$$\hat{\mathbf{W}}_A + \hat{\mathbf{W}}_B = \mathbf{0} \quad (1c)$$

where $\mathbf{u}_A = \begin{bmatrix} \bar{\mathbf{u}}_A \\ \hat{\mathbf{u}}_A \end{bmatrix}$, $\mathbf{u}_B = \begin{bmatrix} \bar{\mathbf{u}}_B \\ \hat{\mathbf{u}}_B \end{bmatrix}$, $\mathbf{f}_A = \begin{bmatrix} \bar{\mathbf{f}}_A \\ \hat{\mathbf{f}}_A \end{bmatrix}$, $\mathbf{f}_B = \begin{bmatrix} \bar{\mathbf{f}}_B \\ \hat{\mathbf{f}}_B \end{bmatrix}$

and the work done on the plates A and B by the connection are $\hat{\mathbf{W}}_A = \hat{\mathbf{f}}_A \cdot \hat{\mathbf{u}}_A$ and $\hat{\mathbf{W}}_B = \hat{\mathbf{f}}_B \cdot \hat{\mathbf{u}}_B$ respectively. Substitution of the connection forces and displacements into (1c) yields the relationship between internal joint forces

$$\hat{\mathbf{f}}_A = -\hat{\mathbf{S}}^T \cdot \hat{\mathbf{f}}_B \quad (1c)'$$

All variables denoted "bar" variables ($\bar{x}$) refer to "external" quantities not involved in a joint while all "hat" variables ($\hat{x}$) refer to "internal" joint variables. Partitioning the equations and rewriting them

$$\begin{bmatrix} \bar{\mathbf{K}}_A & \hat{\mathbf{K}}_A & 0 & 0 \\ \hat{\mathbf{K}}_A^T & \hat{\mathbf{K}}_A & 0 & 0 \\ 0 & 0 & \bar{\mathbf{K}}_B & \hat{\mathbf{K}}_B \\ 0 & 0 & \hat{\mathbf{K}}_B^T & \hat{\mathbf{K}}_B \end{bmatrix} \cdot \begin{bmatrix} \bar{\mathbf{u}}_A \\ \hat{\mathbf{u}}_A \\ \bar{\mathbf{u}}_B \\ \hat{\mathbf{u}}_B \end{bmatrix} = \begin{bmatrix} \bar{\mathbf{f}}_A \\ \hat{\mathbf{f}}_A \\ \bar{\mathbf{f}}_B \\ \hat{\mathbf{f}}_B \end{bmatrix} \quad (2a)$$

subject to constraints $\hat{\mathbf{u}}_B = \hat{\mathbf{S}} \hat{\mathbf{u}}_A \quad (2b)$

and $\hat{\mathbf{f}}_A = -\hat{\mathbf{S}}^T \hat{\mathbf{f}}_B \quad (2c)$

Question: Is the solution of this problem equivalent to the solution of the Lagrangian formulation?

The Lagrangian Problem (A. Diaz, 4/6/00)
 Given the potential energy function

$$\Pi = \frac{1}{2} \mathbf{u}_A^T \mathbf{K}_A \mathbf{u}_A - \bar{\mathbf{u}}_A^T \bar{\mathbf{f}}_A + \frac{1}{2} \mathbf{u}_B^T \mathbf{K}_B \mathbf{u}_B - \bar{\mathbf{u}}_B^T \bar{\mathbf{f}}_B \quad (3a)$$

subject to displacement constraints

$$\hat{\mathbf{u}}_B = \hat{\mathbf{S}} \cdot \hat{\mathbf{u}}_A \quad (3b)$$

The Lagrangian is

$$\begin{aligned} \mathcal{L} &= \Pi + (\hat{\mathbf{u}}_B - \hat{\mathbf{S}} \cdot \hat{\mathbf{u}}_A)^T \cdot \lambda \\ &= \frac{1}{2} \mathbf{u}_A^T \mathbf{K}_A \mathbf{u}_A - \bar{\mathbf{u}}_A^T \bar{\mathbf{f}}_A + \frac{1}{2} \mathbf{u}_B^T \mathbf{K}_B \mathbf{u}_B - \bar{\mathbf{u}}_B^T \bar{\mathbf{f}}_B + (\hat{\mathbf{u}}_B - \hat{\mathbf{S}} \cdot \hat{\mathbf{u}}_A)^T \cdot \lambda \end{aligned} \quad (4)$$

The solution of (3a) and (3b) is a stationary point of (4) that minimizes the potential energy (3a).

Taking the derivative of $\mathcal{L}$ with respect to each of the independent variables $\mathbf{u}_A, \mathbf{u}_B, \lambda$.

$$\frac{\partial \mathcal{L}}{\partial \bar{\mathbf{u}}_A} = 0 = \bar{\mathbf{K}}_A \bar{\mathbf{u}}_A + \hat{\mathbf{K}}_A \hat{\mathbf{u}}_A - \bar{\mathbf{f}}_A \quad (4a)$$

$$\frac{\partial \mathcal{L}}{\partial \hat{\mathbf{u}}_A} = 0 = \hat{\mathbf{K}}_A^T \bar{\mathbf{u}}_A + \hat{\mathbf{K}}_A \hat{\mathbf{u}}_A - \hat{\mathbf{S}}^T \lambda \quad (4b)$$

$$\frac{\partial \mathcal{L}}{\partial \bar{\mathbf{u}}_B} = 0 = \bar{\mathbf{K}}_B \bar{\mathbf{u}}_B + \hat{\mathbf{K}}_B \hat{\mathbf{u}}_B - \bar{\mathbf{f}}_B \quad (4c)$$

$$\frac{\partial \mathcal{L}}{\partial \hat{\mathbf{u}}_B} = 0 = \hat{\mathbf{K}}_B^T \bar{\mathbf{u}}_B + \hat{\mathbf{K}}_B \hat{\mathbf{u}}_B - \lambda \quad (4d)$$

$$\frac{\partial \mathcal{L}}{\partial \lambda} = 0 = \hat{\mathbf{u}}_B - \hat{\mathbf{S}} \hat{\mathbf{u}}_A \quad (4e)$$

Putting (4a)-(4e) into matrix form yields,

$$\begin{bmatrix} \bar{\mathbf{K}}_A & \hat{\mathbf{K}}_A & 0 & 0 & 0 \\ \hat{\mathbf{K}}_A^T & \hat{\mathbf{K}}_A & 0 & 0 & -\hat{\mathbf{S}}^T \\ 0 & 0 & \bar{\mathbf{K}}_B & \hat{\mathbf{K}}_B & 0 \\ 0 & 0 & \hat{\mathbf{K}}_B^T & \hat{\mathbf{K}}_B & \mathbf{I} \\ 0 & -\hat{\mathbf{S}} & 0 & \mathbf{I} & 0 \end{bmatrix} \cdot \begin{bmatrix} \bar{\mathbf{u}}_A \\ \hat{\mathbf{u}}_A \\ \bar{\mathbf{u}}_B \\ \hat{\mathbf{u}}_B \\ \lambda \end{bmatrix} = \begin{bmatrix} \bar{\mathbf{f}}_A \\ 0 \\ \bar{\mathbf{f}}_B \\ 0 \\ 0 \end{bmatrix} \quad (5)$$

Note that (5) is a symmetric formulation.

It can be shown that solutions to (4a-e) rewritten as (5) are also solutions to (2a-c). First note that the Lagrange formulation equation (4e) is equivalent to Modular Modeling constraint (2b).

$$0 = \hat{\mathbf{u}}_B - \hat{\mathbf{S}} \hat{\mathbf{u}}_A \Leftrightarrow \hat{\mathbf{u}}_B = \hat{\mathbf{S}} \hat{\mathbf{u}}_A \quad (6)$$

Let the Lagrange multiplier, $\lambda = -\hat{\mathbf{f}}_B$, the internal boundary force on plate B. Then, the Lagrange formulation equation (4d) is equivalent to the fourth equation in the Modular Modeling matrix equation (2a).

$$\mathbf{0} = \hat{\mathbf{K}}_B^T \bar{\mathbf{u}}_B + \hat{\mathbf{K}}_B \hat{\mathbf{u}}_B + \lambda \Leftrightarrow \hat{\mathbf{K}}_B^T \bar{\mathbf{u}}_B + \hat{\mathbf{K}}_B \hat{\mathbf{u}}_B = \hat{\mathbf{f}}_B \text{ for } \lambda = -\hat{\mathbf{f}}_B \quad (7)$$

The Lagrange formulation equation (4b) is equivalent to the second equation in the Modular Modeling matrix equation (2a). Use the force constraint (2c) from the Modular Modeling formulation to redefine the internal force $\hat{\mathbf{f}}_A$ appearing in the second equation in the matrix equation (2a) to write

$$\mathbf{0} = \hat{\mathbf{K}}_A^T \bar{\mathbf{u}}_A + \hat{\mathbf{K}}_A \hat{\mathbf{u}}_A - \hat{\mathbf{S}}^T \lambda \Leftrightarrow \hat{\mathbf{K}}_A^T \bar{\mathbf{u}}_A + \hat{\mathbf{K}}_A \hat{\mathbf{u}}_A = \hat{\mathbf{f}}_A \text{ for } \hat{\mathbf{f}}_B = -\hat{\mathbf{S}}^T \hat{\mathbf{f}}_B = \hat{\mathbf{S}}^T \lambda \quad (8)$$

The Lagrange formulation equations (4a,c) are identical to the first and third equations in the Modular Modeling matrix equation (2a).

$$\mathbf{0} = \bar{\mathbf{K}}_A \bar{\mathbf{u}}_A + \hat{\mathbf{K}}_A \hat{\mathbf{u}}_A - \bar{\mathbf{f}}_A \Leftrightarrow \bar{\mathbf{K}}_A \bar{\mathbf{u}}_A + \hat{\mathbf{K}}_A \hat{\mathbf{u}}_A = \bar{\mathbf{f}}_A \quad (9)$$

$$\mathbf{0} = \bar{\mathbf{K}}_B \bar{\mathbf{u}}_B + \hat{\mathbf{K}}_B \hat{\mathbf{u}}_B - \bar{\mathbf{f}}_B \Leftrightarrow \bar{\mathbf{K}}_B \bar{\mathbf{u}}_B + \hat{\mathbf{K}}_B \hat{\mathbf{u}}_B = \bar{\mathbf{f}}_B \quad (10)$$

The conclusion is that the two formulations have identical solutions with the recognition that the Lagrange multiplier $\lambda = -\hat{\mathbf{f}}_B$, the joint internal force.

General Form of the Displacement Constraint

A more general form of the displacement constraint is the linear relationship

$$\mathbf{S}_A \mathbf{u}_A = \mathbf{S}_B \mathbf{u}_B \quad (11)$$

with this relationship, the Lagrangian formulation becomes

$$\Pi = \frac{1}{2} \mathbf{u}_A^T \mathbf{K}_A \mathbf{u}_A - \mathbf{u}_A^T \mathbf{f}_A + \frac{1}{2} \mathbf{u}_B^T \mathbf{K}_B \mathbf{u}_B - \mathbf{u}_B^T \mathbf{f}_B \quad (12a)$$

subject to displacement constraints

$$\mathbf{S}_A \mathbf{u}_A = \mathbf{S}_B \mathbf{u}_B \quad (12b)$$

Note that in this case, I have not partitioned the displacements $\mathbf{u}_A$ and $\mathbf{u}_B$. The body forces $\mathbf{f}_A$ and $\mathbf{f}_B$ represent all external forces applied to the two bodies and are simply a superset of the previous external forces $\bar{\mathbf{f}}_A$ and $\bar{\mathbf{f}}_B$ that also allow external tractions to be applied at the position of joint connections. The Lagrangian is

$$\begin{aligned} \mathbf{L} &= \Pi + (\mathbf{S}_B \mathbf{u}_B - \mathbf{S}_A \cdot \mathbf{u}_A)^T \cdot \lambda \\ &= \frac{1}{2} \mathbf{u}_A^T \mathbf{K}_A \mathbf{u}_A - \mathbf{u}_A^T \mathbf{f}_A + \frac{1}{2} \mathbf{u}_B^T \mathbf{K}_B \mathbf{u}_B - \mathbf{u}_B^T \mathbf{f}_B + (\mathbf{S}_B \mathbf{u}_B - \mathbf{S}_A \cdot \mathbf{u}_A)^T \cdot \lambda \end{aligned} \quad (13)$$

whose stationary value occurs at the solution to

$$\frac{\partial L}{\partial \mathbf{u}_A} = \mathbf{0} = \mathbf{K}_A \mathbf{u}_A - \bar{\mathbf{f}}_A - \mathbf{S}_A^T \lambda \quad (14a)$$

$$\frac{\partial L}{\partial \mathbf{u}_B} = \mathbf{0} = \mathbf{K}_B \mathbf{u}_B - \bar{\mathbf{f}}_B - \mathbf{S}_B^T \lambda \quad (14b)$$

$$\frac{\partial L}{\partial \lambda} = \mathbf{0} = \mathbf{S}_B \mathbf{u}_B - \mathbf{S}_A \mathbf{u}_A \quad (14c)$$

Putting (14a)-(14c) into matrix form yields,

$$\begin{bmatrix} \mathbf{K}_A & \mathbf{0} & -\mathbf{S}_A^T \\ \mathbf{0} & \mathbf{K}_B & \mathbf{S}_B^T \\ -\mathbf{S}_A & \mathbf{S}_B & \mathbf{0} \end{bmatrix} \begin{bmatrix} \mathbf{u}_A \\ \mathbf{u}_B \\ \lambda \end{bmatrix} = \begin{bmatrix} \mathbf{f}_A \\ \mathbf{f}_B \\ \mathbf{0} \end{bmatrix} \quad (15)$$

Note that (15) is a symmetric formulation with $\mathbf{u}_A$ and $\mathbf{u}_B$ defined as in (1c).

With this definition, the joint constraint matrices in (15) can be used to define the original Modular Modeling formulation by defining,

$$\mathbf{S}_A = \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \hat{\mathbf{S}}_A \end{bmatrix} \quad (16a)$$

and
$$\mathbf{S}_B = \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{I} \end{bmatrix} \quad (16b)$$

APPENDIX II

```

EE = 15000000; u =  $\frac{1}{4}$ ; t = 0.1;
"System properties: Young's modulus(N/cm2); Poisson's ratio; Plane thickness(cm); ";

NI = {1, 3, 2, 2, 4, 5, 5, 6, 6}; NJ = {3, 4, 3, 5, 7, 4, 8, 5, 9};
MK = {2, 5, 5, 6, 8, 8, 9, 9, 10};
"mesh for plane A: nodes distribution, pick up the i,j,k; of #n element by: ";
"in=NI[[n]]; jn=NJ[[n]]; kn=MK[[n]]";

X = {0, 2, 0, 0, 2, 4, 0, 2, 4, 6}; Y = {6, 4, 4, 2, 2, 2, 0, 0, 0, 0};
"X,Y value for plane A nodes, pick up the Xn & Ynj of #n node by: ";
"Xn=X[[n]]; Yn=Y[[n]] ";

Array[Bi, 9]; Array[Bj, 9]; Array[Bk, 9]; Array[Ci, 9]; Array[Cj, 9]; Array[Ck, 9];
Do[ Bi[n] = (Y[[{NJ[[n]]}]] - (Y[[{MK[[n]]}]]), {n, 9}];
Do[ Bj[n] = (Y[[{MK[[n]]}]] - (Y[[{NI[[n]]}]]), {n, 9}];
Do[ Bk[n] = (Y[[{NI[[n]]}]] - (Y[[{NJ[[n]]}]]), {n, 9}];
Do[ Ci[n] = (X[[{MK[[n]]}]] - (X[[{NJ[[n]]}]]), {n, 9}];
Do[ Cj[n] = (X[[{NI[[n]]}]] - (X[[{MK[[n]]}]]), {n, 9}];
Do[ Ck[n] = (X[[{NJ[[n]]}]] - (X[[{NI[[n]]}]]), {n, 9}];
"Define Bi,Bj,Bk,Ci,Cj,Ck for all elements, compute every value by: ";
"Bi[n]=Yj[n]-Yk[n];  Bj[n]=Yk[n]-Yi[n];  Bk[n]=Yi[n]-Yj[n]";
"Ci[n]=Xk[n]-Xj[n];  Cj[n]=Xi[n]-Xk[n];  Ck[n]=Xj[n]-Xi[n]";

"t = Table[{X[[{MK[[n]]}]],Y[[{MK[[n]]}]]},
{n,9}];ListPlot[t,PlotJoined->True, PlotLabel->"e"];

Do[ Print["Bi(", i, ")=", (Bi[i]), "      Bj(", i, ")=",
(Bj[i]), "      Bk(", i, ")=", (Bk[i]), "      Ci(", i, ")=", (Ci[i]),
"      Cj(", i, ")=", (Cj[i]), "      Ck(", i, ")=", (Ck[i]), {i, 9}];

Do[A =  $\frac{1}{2} \times \text{Det} \left[ \begin{pmatrix} 1 & X[[{NI[[n]]}]] & Y[[{NI[[n]]}]] \\ 1 & X[[{NJ[[n]]}]] & Y[[{NJ[[n]]}]] \\ 1 & X[[{MK[[n]]}]] & Y[[{MK[[n]]}]] \end{pmatrix} \right]$ ; "Print[A]", {n, 9}];

Array[BB, 9]; Do[ BB[n] =  $\frac{1}{2A} \begin{pmatrix} Bi[n] & 0 & Bj[n] & 0 & Bk[n] & 0 \\ 0 & Ci[n] & 0 & Cj[n] & 0 & Ck[n] \\ Ci[n] & Bi[n] & Cj[n] & Bj[n] & Ck[n] & Bk[n] \end{pmatrix}$ ;

"Print[BB[n]]", {n, 9}]; DD =  $\frac{EE}{1-u^2} \begin{pmatrix} 1 & u & 0 \\ u & 1 & 0 \\ 0 & 0 & \frac{1-u}{2} \end{pmatrix}$ ;

Array[Kelement, 9];
Do[KK = (Transpose[{BB[n]}]).DD];
KK2 = KK.BB[n]; K = t.*A.*KK2; Kelement[n] = Rationalize[K];
"Print[Kelement[n]]", {n, 9}];

```

$$Kall = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix},$$

```

Array[ijk, 6];
"2i-1=ijk[1],2i=ijk[2],2j-1=ijk[3],2j=ijk[4],2k-1=ijk[5],2k=ijk[6],";
Do[Print["element: ", e]; ijk[1] = (2×NI[[e]]) - 1;
  ijk[2] = (2×NI[[e]]); ijk[3] = (2×NJ[[e]]) - 1; ijk[4] = (2×NJ[[e]]);
  ijk[5] = (2×NK[[e]]) - 1; ijk[6] = (2×NK[[e]]); Ktemp = Kelement[e];
  Do [Do[Kall[[ijk[i], ijk[j]]] = Kall[[ijk[i], ijk[j]]] + Ktemp[[i, j]], {i, 6}],
    {j, 6}], {e, 9}];
"Fill the element K into the integral K(Kall), based on every coordinated nodes";

```

Kall
100000

3	0	0	3	-3	-3	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	8	2	0	-2	-8	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	2	22	5	-16	-5	0	0	-6	-5	0	3	0	0	0	0	0	0	0	0
3	0	5	22	-5	-6	0	0	-5	-16	2	0	0	0	0	0	0	0	0	0
-3	-2	-16	-5	22	5	-3	-3	0	5	0	0	0	0	0	0	0	0	0	0
-3	-8	-5	-6	5	22	-2	-8	5	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	-3	-2	22	5	-16	-5	0	0	-3	-3	0	5	0	0	0	0
0	0	0	0	-3	-8	5	22	-5	-6	0	0	-2	-8	5	0	0	0	0	0
0	0	-6	-5	0	5	-16	-5	44	10	-16	-5	0	0	-6	-5	0	5	0	0
0	0	-5	-16	5	0	-5	-6	10	44	-5	-6	0	0	-5	-16	5	0	0	0
0	0	0	2	0	0	0	0	-16	-5	22	5	0	0	0	0	-6	-5	0	3
0	0	3	0	0	0	0	0	-5	-6	5	22	0	0	0	0	-5	-16	2	0
0	0	0	0	0	0	-3	-2	0	0	0	0	11	5	-8	-3	0	0	0	0
0	0	0	0	0	0	-3	-8	0	0	0	0	5	11	-2	-3	0	0	0	0
0	0	0	0	0	0	0	5	-6	-5	0	0	-8	-2	22	5	-8	-3	0	0
0	0	0	0	0	0	5	0	-5	-16	0	0	-3	-3	5	22	-2	-3	0	0
0	0	0	0	0	0	0	0	0	5	-6	-5	0	0	-8	-2	22	5	-8	-3
0	0	0	0	0	0	0	0	5	0	-5	-16	0	0	-3	-3	5	22	-2	-3
0	0	0	0	0	0	0	0	0	0	0	2	0	0	0	0	-8	-2	8	0
0	0	0	0	0	0	0	0	0	0	3	0	0	0	0	0	-3	-3	0	3

APPENDIX III

```

KE = 15000000; u =  $\frac{1}{4}$ ; t = 0.1;
"System properties: Young's modulus(N/cm2); Poisson's ratio; Plane thickness(cm); ";

NI = {1, 2, 2, 4}; NJ = {4, 4, 5, 6}; NK = {2, 5, 3, 5};
"mesh for plane A: nodes distribution, pick up the i;j;k; of #n element by: ";
"in=NI[[n]]; jn=NJ[[n]]; kn=NK[[n]]";

X = {0, 3, 6, 3, 6, 6}; Y = {0, 0, 0, -3, -3, -6};
"X,Y value for plane A nodes, pick up the Xn & Ynj of #n node by: ";
"Xn=X[[n]]; Yn=Y[[n]] ";

Array[Bi, 4]; Array[Bj, 4]; Array[Bk, 4]; Array[Ci, 4]; Array[Cj, 4]; Array[Ck, 4];
Do[Bi[n] = (Y[[{NJ[[n]]}]] - (Y[[{NK[[n]]}]]), {n, 4}];
Do[Bj[n] = (Y[[{NK[[n]]}]] - (Y[[{NI[[n]]}]]), {n, 4}];
Do[Bk[n] = (Y[[{NI[[n]]}]] - (Y[[{NJ[[n]]}]]), {n, 4}];
Do[Ci[n] = (X[[{NK[[n]]}]] - (X[[{NJ[[n]]}]]), {n, 4}];
Do[Cj[n] = (X[[{NI[[n]]}]] - (X[[{NK[[n]]}]]), {n, 4}];
Do[Ck[n] = (X[[{NJ[[n]]}]] - (X[[{NI[[n]]}]]), {n, 4}];
"Define Bi,Bj,Bk,Ci,Cj,Ck for all elements, compute every value by: ";
"Bi[n]=Yj[n]-Yk[n]; Bj[n]=Yk[n]-Yi[n]; Bk[n]=Yi[n]-Yj[n];
"Ci[n]=Xk[n]-Xj[n]; Cj[n]=Xi[n]-Xk[n]; Ck[n]=Xj[n]-Xi[n];

"t = Table[{X[[{NJ[[n]]}]],Y[[{NJ[[n]]}]]},
{n,4}];ListPlot[t,PlotJoined->True, PlotLabel->"e"];

Do[Print["Bi(", i, ")=", (Bi[i]), " Bj(", i, ")=", (Bj[i]),
" Bk(", i, ")=", (Bk[i]), " Ci(", i, ")=", (Ci[i]),
" Cj(", i, ")=", (Cj[i]), " Ck(", i, ")=", (Ck[i]), {i, 4} ]

Do[A =  $\frac{1}{2} \times \text{Det} \begin{bmatrix} 1 & X[[{NI[[n]]}]] & Y[[{NI[[n]]}]] \\ 1 & X[[{NJ[[n]]}]] & Y[[{NJ[[n]]}]] \\ 1 & X[[{NK[[n]]}]] & Y[[{NK[[n]]}]] \end{bmatrix}$ , {"Print[A]", {n = 4}];

Array[BB, 9]; Do[BB[n] =  $\frac{1}{2A} \begin{pmatrix} Bi[n] & 0 & Bj[n] & 0 & Bk[n] & 0 \\ 0 & Ci[n] & 0 & Cj[n] & 0 & Ck[n] \\ Ci[n] & Bi[n] & Cj[n] & Bj[n] & Ck[n] & Bk[n] \end{pmatrix}$ ;

"Print[BB[n]]", {n, 4} ]; DD =  $\frac{KE}{1-u^2} \begin{pmatrix} 1 & u & 0 \\ u & 1 & 0 \\ 0 & 0 & \frac{1-u}{2} \end{pmatrix}$ 

Array[Kelement, 4];
Do[KK = (Transpose[BB[n]].DD);
KK2 = KK.BB[n]; K = t + A + KK2; Kelement[n] = Rationalize[K];
"Print[Kelement[n]]", {n, 4} ]

```

$$Kall = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix};$$

```

Array[ijk, 6];
"2i-1=ijk[1],2i=ijk[2],2j-1=ijk[3],2j=ijk[4],2k-1=ijk[5],2k=ijk[6],";
Do[Print["element: ", e]; ijk[1] = (2×NI[[e]]) - 1;
  ijk[2] = (2×NI[[e]]); ijk[3] = (2×NJ[[e]]) - 1; ijk[4] = (2×NJ[[e]]);
  ijk[5] = (2×NK[[e]]) - 1; ijk[6] = (2×NK[[e]]); Ktemp = Kelement[e];
  Do [Do[Kall[[ijk[i], ijk[j]]] = Kall[[ijk[i], ijk[j]]] + Ktemp[[i, j]], {i, 6}],
    {j, 6}], {e, 4}];
"Fill the element K into the integral K(Kall), based on every corrdated nodes";

```

Kall/100000

$$\begin{pmatrix} 8 & 0 & -8 & -2 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 \\ 0 & 3 & -3 & -3 & 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 \\ -8 & -3 & 22 & 5 & -8 & -2 & -6 & -5 & 0 & 5 & 0 & 0 \\ -2 & -3 & 5 & 22 & -3 & -3 & -5 & -16 & 5 & 0 & 0 & 0 \\ 0 & 0 & -8 & -3 & 11 & 5 & 0 & 0 & -3 & -2 & 0 & 0 \\ 0 & 0 & -2 & -3 & 5 & 11 & 0 & 0 & -3 & -8 & 0 & 0 \\ 0 & 3 & -6 & -5 & 0 & 0 & 22 & 5 & -16 & -5 & 0 & 2 \\ 2 & 0 & -5 & -16 & 0 & 0 & 5 & 22 & -5 & -6 & 3 & 0 \\ 0 & 0 & 0 & 5 & -3 & -3 & -16 & -5 & 22 & 5 & -3 & -2 \\ 0 & 0 & 5 & 0 & -2 & -8 & -5 & -6 & 5 & 22 & -3 & -8 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 3 & -3 & -3 & 3 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & -2 & -8 & 0 & 8 \end{pmatrix}$$

APPENDIX IV

KE = 15000000; $u = \frac{1}{4}$; $t = 0.1$;

"System properties: Young's modulus(N/cm²); Poisson's ratio; Plane thickness(cm); ";

NI = {1, 1, 2, 2, 3, 3, 4, 4, 5, 5, 6, 6, 8, 8, 9, 9, 10, 10, 11, 11, 12, 12, 13, 13, 15, 15, 16, 16, 17, 17, 18, 18, 19, 19, 20, 20, 22, 22, 23, 23, 24, 24, 25, 25, 26, 26, 27, 27, 29, 29, 30, 30, 31, 31, 32, 32, 33, 33, 34, 34, 36, 36, 37, 37, 38, 38, 39, 39, 40, 40, 41, 41};

NJ = {8, 9, 9, 10, 10, 11, 11, 12, 12, 13, 13, 14, 15, 16, 16, 17, 17, 18, 18, 19, 19, 20, 20, 21, 22, 23, 23, 24, 24, 25, 25, 26, 26, 27, 27, 28, 29, 30, 30, 31, 31, 32, 32, 33, 33, 34, 34, 35, 36, 37, 37, 38, 38, 39, 39, 40, 40, 41, 41, 42, 43, 44, 44, 45, 45, 46, 46, 47, 47, 48, 48, 49};

NK = {9, 2, 10, 3, 11, 4, 12, 5, 13, 6, 14, 7, 16, 9, 17, 10, 18, 11, 19, 12, 20, 13, 21, 14, 23, 16, 24, 17, 25, 18, 26, 19, 27, 20, 28, 21, 30, 23, 31, 24, 32, 25, 33, 26, 34, 27, 35, 28, 37, 30, 38, 31, 39, 32, 40, 33, 41, 34, 42, 35, 44, 37, 45, 38, 46, 39, 47, 40, 48, 41, 49, 42};

"mesh for plane A: nodes distribution, pick up the i;j;k; of #n element by: ";
"in=NI[[n]]; jn=NJ[[n]]; kn=NK[[n]]";

X = {0, 1, 2, 3, 4, 5, 6, 0, 1, 2, 3, 4, 5, 6, 0, 1, 2, 3, 4, 5, 6, 0, 1, 2, 3, 4, 5, 6, 0, 1, 2, 3, 4, 5, 6};

Y = {6, 6, 6, 6, 6, 6, 5, 5, 5, 5, 5, 5, 4, 4, 4, 4, 4, 4, 4, 3, 3, 3, 3, 3, 3, 3, 2, 2, 2, 2, 2, 2, 1, 1, 1, 1, 1, 1, 0, 0, 0, 0, 0, 0, 0};

"X,Y value for plane A nodes, pick up the Xn & Ynj of #n node by: ";
"Xn=X[[n]]; Yn=Y[[n]] ";

Array[Bi, 72]; Array[Bj, 72]; Array[Bk, 72];

Array[Ci, 72]; Array[Cj, 72]; Array[Ck, 72];

Do[Bi[n] = (Y[[{NJ[[n]]}]] - (Y[[{NK[[n]]}]]), {n, 72}];

Do[Bj[n] = (Y[[{NK[[n]]}]] - (Y[[{NI[[n]]}]]), {n, 72}];

Do[Bk[n] = (Y[[{NI[[n]]}]] - (Y[[{NJ[[n]]}]]), {n, 72}];

Do[Ci[n] = (X[[{NK[[n]]}]] - (X[[{NJ[[n]]}]]), {n, 72}];

Do[Cj[n] = (X[[{NI[[n]]}]] - (X[[{NK[[n]]}]]), {n, 72}];

Do[Ck[n] = (X[[{NJ[[n]]}]] - (X[[{NI[[n]]}]]), {n, 72}];

"Define Bi,Bj,Bk,Ci,Cj,Ck for all elements, compute every value by: ";

"Bi[n]=Yj[n]-Yk[n]; Bj[n]=Yk[n]-Yi[n]; Bk[n]=Yi[n]-Yj[n]";

"Ci[n]=Xk[n]-Xj[n]; Cj[n]=Xi[n]-Xk[n]; Ck[n]=Xj[n]-Xi[n]";

Do[Print["Bi(", i, ")=", (Bi[i]), " Bj(", i, ")=", (Bj[i]), " Bk(", i, ")=", (Bk[i]), " Ci(", i, ")=", (Ci[i]), " Cj(", i, ")=", (Cj[i]), " Ck(", i, ")=", (Ck[i]), {i, 72}]

Do[A = $\frac{1}{2} \times \text{Det} \left[\begin{pmatrix} 1 & X[[{NI[[n]]}]] & Y[[{NI[[n]]}]] \\ 1 & X[[{NJ[[n]]}]] & Y[[{NJ[[n]]}]] \\ 1 & X[[{NK[[n]]}]] & Y[[{NK[[n]]}]] \end{pmatrix} \right]$, "Print[A]", {n, 72}]

```
Array[BB, 72]; Do[BB[n] =  $\frac{1}{2\lambda}$   $\begin{pmatrix} Bi[n] & 0 & Bj[n] & 0 & Bk[n] & 0 \\ 0 & Ci[n] & 0 & Cj[n] & 0 & Ck[n] \\ Ci[n] & Bi[n] & Cj[n] & Bj[n] & Ck[n] & Bk[n] \end{pmatrix}$ , {n, 72}];
```

$$DD = \frac{KE}{1-u^2} \begin{pmatrix} 1 & u & 0 \\ u & 1 & 0 \\ 0 & 0 & \frac{1-u}{2} \end{pmatrix}$$

$$\begin{pmatrix} 16000000 & 4000000 & 0 \\ 4000000 & 16000000 & 0 \\ 0 & 0 & 6000000 \end{pmatrix}$$

```
Array[Kelement, 72];
Do[KK = (Transpose[BB[n]]) . DD;
KK2 = KK.BB[n]; K = t + λ + KK2; Kelement[n] = Rationalize[K];
Print[Kelement[n]], {n, 72}]
```

$$Kall = \begin{pmatrix} 0 & : & 0 \\ .. & : & .. \\ 0 & : & 0 \end{pmatrix}_{98 \times 98};$$

```
Array[ijk, 6];
"2i-1=ijk[1],2i=ijk[2],2j-1=ijk[3],2j=ijk[4],2k-1=ijk[5],2k=ijk[6],";
Do[Print["element: ", e]; ijk[1] = (2×NI[[e]]) - 1;
ijk[2] = (2×NI[[e]]); ijk[3] = (2×NJ[[e]]) - 1; ijk[4] = (2×NJ[[e]]);
ijk[5] = (2×NK[[e]]) - 1; ijk[6] = (2×NK[[e]]); Ktemp = Kelement[e];
Do[Do[Kall[[ijk[i], ijk[j]]] = Kall[[ijk[i], ijk[j]]] + Ktemp[[i, j]], {i, 6}],
{j, 6}], {e, 72}];
"Fill the element K into the integral K(Kall), based on every coordinated nodes";
Do[Print[Kall[[i]]], {i, 98}];
"use to list all Kall element";
```

$$U = \begin{pmatrix} U1 \\ U2 \\ U3 \\ : \\ : \\ U12 \\ U13 \\ U14 \\ U15 \\ : \\ : \\ U84 \\ U85 \\ U86 \\ : \\ : \\ U98 \end{pmatrix}; F = \begin{pmatrix} f1 \\ f2 \\ 0 \\ : \\ : \\ 0 \\ -100 \\ -100 \\ 0 \\ : \\ : \\ 0 \\ f85 \\ 0 \\ : \\ : \\ 0 \end{pmatrix};$$


```

Solve[{Kall.U == F, U1 == 0, U2 == 0, U85 == 0},
{f1, f2, f85, U1, U2, U3, U4, U5, U6, U7, U8, U9, U10, U11, U12, U13, U14, U15, U16,
U17, U18, U19, U20, U21, U22, U23, U24, U25, U26, U27, U28, U29, U30, U31, U32, U33,
U34, U35, U36, U37, U38, U39, U40, U41, U42, U43, U44, U45, U46, U47, U48, U49, U50,
U51, U52, U53, U54, U55, U56, U57, U58, U59, U60, U61, U62, U63, U64, U65, U66,
U67, U68, U69, U70, U71, U72, U73, U74, U75, U76, U77, U78, U79, U80, U81, U82,
U83, U84, U85, U86, U87, U88, U89, U90, U91, U92, U93, U94, U95, U96, U97, U98}]

{{f1 → 0, f2 → 100, f85 → 100, U1 → 0, U2 → 0,
U3 →  $-\frac{2407592046919015049412373020803462794172449814748552125670710887246817496493651}{347542805637642019646677141081930240935493453365697259040357719134473019431083040000}$ ,
U4 →  $-\frac{1240118079660558731107203417149144524903115498671995043630743888585903681914279}{6516427605705787868375196395286192017540502250606823607006707233771369114332807000}$ ,
U5 →  $-\frac{2059543348633698800959755489598990805529647954073087896933575375917428236382627}{5792380093960700327444619018032170682258224227616209840059619855745503238513840000}$ ,
U6 →  $-\frac{14818379281897120298598434753367048018546061903207525588802944669254888535911839}{52131420845646302947001571162289536140324018004854588856053657870170952914662456000}$ ,
U7 →  $-\frac{22137170005762172957942048025131657966255020675267314683432452737745739697826513}{347542805637642019646677141081930240935493453365697259040357719134473019431083040000}$ ,
U8 →  $-\frac{1022794547568881169902438630364855134411101948110483020096642720950587875331467}{2896190046980350163722309509016085341129112111380810492002980992787275161925692000}$ ,
U9 →  $-\frac{8159494195900848084706599775574783566907639926889935255106568031324591080886601}{86885701409410504911669285270482560233873363341424314760089429783618254857770760000}$ ,
U10 →  $-\frac{3626936133420891916259467703259369739392842964358328380235199402938555582338309}{8688570140941050491166928527048256023387336334142431476008942978361825485777076000}$ ,
U11 →  $-\frac{3112614005966334782011607673001983864863519559362672015080213011272729412897601}{23169520375842801309778476072128682729032896891046483936023847942298201295405536000}$ ,
U12 →  $-\frac{2123383606361286489971817918389014384523235983528035869708374677817852767676479}{4344285070470525245583464263524128011693668167071215738004471489180912742888538000}$ ,
U13 →  $-\frac{33692102364413670352972673709569061317902354000881057659814443992150994713623803}{173771402818821009823338570540965120467746726682848629520178859567236509715541520000}$ ,
U14 →  $-\frac{1777472296077346438132223937951948762833}{2982651209893062328533406309694265650712000}$ ,
U15 →  $-\frac{10279347345505920387782176010704113180633680566364116940940140541349400676203011}{104262841691292605894003142324579072280648036009709177712107315740341905829324912000}$ ,
U16 →  $-\frac{33743948374049089314060247323220758465272862625047116106898080268663218604372589}{347542805637642019646677141081930240935493453365697259040357719134473019431083040000}$ ,
U17 →  $-\frac{117947659526741963745845406028882056214757888385323187667415408875588265304610003}{1042628416912926058940031423245790722806480360097091777121073157403419058293249120000}$ ,
U18 →  $-\frac{213339190069810914120041910443423915327120387149070243633644408859972757170772107}{1042628416912926058940031423245790722806480360097091777121073157403419058293249120000}$ ,
U19 →  $-\frac{15053847884544485964477221380063016033688175936937182069317454636531416384727727}{130328552114115757367503927905723840350810045012136472140134144675427382286656140000}$ ,
U20 →  $-\frac{302539156432708362842088963187000869264877392172223109936881021805737136971417647}{1042628416912926058940031423245790722806480360097091777121073157403419058293249120000}$ ,
U21 →  $-\frac{130177900967636184385373402405152572901289037448588859903504928231878516958476929}{1042628416912926058940031423245790722806480360097091777121073157403419058293249120000}$ ,
U22 →  $-\frac{41302365069422922908959207825707820490596016876953136039477074853308481029922423}{115847601879214006548892380360643413645164484455232419680119239711491006477027680000}$ ,
U23 →  $-\frac{72852482455170235647385151398271334582444842343627923547161990691485780222679101}{521314208456463029470015711622895361403240180048545888560536578701709529146624560000}$ ,
U24 →  $-\frac{11164325855744279166436188539559201070554814753817030072565585032757399046996393}{26734061972126309203590549313994633918114881028130558387719824548805616879314080000}$ 

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6534906666763785455329635084548986875311314279710360477248414730575955298715143
 U25 → - 41705136676517042357601256929831628912259214403883671084842926296136762331729964800 ,
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 U26 → - 115847601879214006548892380360643413645164484455232419680119239711491006477027680000 ,
 3332472480233073332785420299315720513650659813258550538757304405385767539397179
 U27 → - 20050546479094731902692911985495975438586160771097918790789868411604212659485560000 ,
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 U28 → - 8416103878199922520901048415374801133493646702405462737530720000 ,
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 U29 → - 1629106901426446967093799098821548004385125562651705901751676808442842278583201750 ,
 29331382326278761244717801491485661579522416819081876822557832756103433442920039
 U30 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
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 U31 → - 1042628416912926058940031423245790722806480360097091777121073157403419058293249120000 ,
 119360537160532845405338788678346836247742454898149789880532456464282008592271787
 U32 → - 521314208456463029470015711622895361403240180048545888560536578701709529146624560000 ,
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 U33 → - 40101092958189463805385823970991950877172321542195837581579736823208425318971120000 ,
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 U34 → - 40101092958189463805385823970991950877172321542195837581579736823208425318971120000 ,
 172258763200080732822085473323487565487873681384802749943038955223056060979337759
 U35 → - 1042628416912926058940031423245790722806480360097091777121073157403419058293249120000 ,
 61255812647354562237994479654706518751284055803039417542784485565367218972801889
 U36 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
 338729778637942638895637920852370964903421303067527532032543334409084774597671
 U37 → - 20050546479094731902692911985495975438586160771097918790789868411604212659485560000 ,
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 U38 → - 19307933646535667758148730060107235607527414075872069946686539951915167746171280000 ,
 36104812635842126964911830173314115967865492388956046941073884832955188359327749
 U39 → - 208525683382585211788006284649158144561296072019418355424214631480683811658649824000 ,
 4207371182161304451506833187243985717110173296494463209871320873923016814060071
 U40 → - 9145863306253737359123082660050795814091932983307822606325203135117711037660080000 ,
 89986670214447333434445684987128919247461387734983151063053774042967763159563569
 U41 → - 521314208456463029470015711622895361403240180048545888560536578701709529146624560000 ,
 86303717143006777665096798036634252486294163071228952332937565728607499364890589
 U42 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
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 U43 → - 34754280563764201964667714108193024093549345336569725904035771913447301943108304000 ,
 2592473509336535747200566127539168128060876049280499572061618379516178779498789
 U44 → - 115847601879214006548892380360643413645164484455232419680119239711491006477027680000 ,
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 U45 → - 347542805637642019646677141081930240935493453365697259040357719134473019431083040000 ,
 261849897663306982570276318948405757435128503076143612810929237915994346033466861
 U46 → - 1042628416912926058940031423245790722806480360097091777121073157403419058293249120000 ,
 16356096827774845307641326560853193641363115602755890710806956817582979944427523
 U47 → - 86885701409410504911669285270482560233873363341424314760089429783618254857770760000 ,
 307243747584624899308598237530244139919922723439417144044467342038171786071629021
 U48 → - 1042628416912926058940031423245790722806480360097091777121073157403419058293249120000 ,
 65180635574278047374675685565629304138723620186076554274934705704064262923842723
 U49 → - 347542805637642019646677141081930240935493453365697259040357719134473019431083040000 ,
 120242772876703683830313069620685351949470662277945898242934388777976820993328987
 U50 → - 347542805637642019646677141081930240935493453365697259040357719134473019431083040000 ,
 10864960485416698843397382362362667294003058836891391030235754248629037634775559
 U51 → - 5792380093960700327444619018032170682258224227616209840059619855745503238513840000 ,

U52 → - $\frac{10622986420959809679966565871589323923111425272507632309340362775797318291117619}{26734061972126309203590549313994633918114881028130558387719824548805616879314080000}$,
 U53 → - $\frac{104197892519990118153048183107902329413401529514642686594114166448548835448621}{556068489020227231434683425731088385496789525385115614464572350615156831089732864}$,
 U54 → - $\frac{153620712632911483813983634848578240108898976430806547534238673121718086856181607}{347542805637642019646677141081930240935493453365697259040357719134473019431083040000}$,
 U55 → - $\frac{502375029820938063123880210372885914633774698519709455014933401037934949562787}{2715178169044078278489665164702580007308542604419509836252794680738070464305336250}$,
 U56 → - $\frac{54569418839793723608839296273310983778475856721124419056235282228659359181907949}{115847601879214006548892380360643413645164484455232419680119239711491006477027680000}$,
 U57 → - $\frac{32802769230862019067622768352095441995373570867715003541871424440731405519091}{181011877936271885232644344313505333820569506961300655750186312049204697620355750}$,
 U58 → - $\frac{23334421602208627137667188614407646234756650913811311735704306697045571346251719}{86885701409410504911669285270482560233873363341424314760089429783618254857770760000}$,
 U59 → - $\frac{187166009385152522628716269048623218506174688001165686073862615524702587940621}{990150443412085526058909233851653108078328926967798458804437946252059884419040000}$,
 U60 → - $\frac{69736674015738563518992457174669853821091945124770471978262095166173834807040137}{260657104228231514735007855811447680701620090024272944280268289350854764573312280000}$,
 U61 → - $\frac{33676509916384171380182634558744486189368057605028803901880398283207229145207081}{173771402818821009823338570540965120467746726682848629520178859567236509715541520000}$,
 U62 → - $\frac{370465798531755351641488100359750799510741127350219553695837092107646252325599}{1253159154943420743918306999093498464911635048193619924424366775725263291217847500}$,
 U63 → - $\frac{1758846966300419369049817823495618690389658017909903685245274916412983046348487}{8911353990708769734530183104664877972704960342710186129239941516268538959771360000}$,
 U64 → - $\frac{9842168516776578838301008230960056632444743284025625881593548325317620458776343}{28961900469803501637223095090160853411291121113808104920029809927872751619256920000}$,
 U65 → - $\frac{17365748488109925029153220874059270795600089500752498143336820685147702331812201}{86885701409410504911669285270482560233873363341424314760089429783618254857770760000}$,
 U66 → - $\frac{33662301624615491112327411998360574033372080324044804413403558448832488241449209}{86885701409410504911669285270482560233873363341424314760089429783618254857770760000}$,
 U67 → - $\frac{4646708357288605126794855208184974842610778951262886665042679597926069431628793}{23169520375842801309778476072128682729032896891046483936023847942298201295405536000}$,
 U68 → - $\frac{37258995748189859724797941642052847058073323863856330971957776054848256671469209}{86885701409410504911669285270482560233873363341424314760089429783618254857770760000}$,
 U69 → - $\frac{34623129341501359717008923557308795897371679545611760630370476975129315867640763}{173771402818821009823338570540965120467746726682848629520178859567236509715541520000}$,
 U70 → - $\frac{19813311467682074619613044826639884691334519981927871095177613899728324617962977}{43442850704705252455834642635241280116936681670712157380044714891809127428885380000}$,
 U71 → - $\frac{569491900759919109408737464336831491527178453972908287620625211427807006623957}{3861586729307133551629746012021447121505482815174413989337307990383033549234256000}$,
 U72 → - $\frac{21505592277973830952031830044283220947056365271440607304402493869186865692082941}{69508561127528403929335428216386048187098690673139451808071543826894603886216608000}$,
 U73 → - $\frac{19319182827010288570102868537440440735752278910585267827931016430176483074792267}{115847601879214006548892380360643413645164484455232419680119239711491006477027680000}$,
 U74 → - $\frac{57897718357627718592769238134274778080604825169951266972605382449380409049959107}{208525683382585211788006284649158144561296072019418355424214631480683811658649824000}$,
 U75 → - $\frac{426691171992021156295314415236899703463399846293895939873149196353976450089421}{2286465826563434339780770665012698953522983245826955651581300783779427759415020000}$,
 U76 → - $\frac{61633912587799874770528044217820582467076154761957619890131678483232456394992711}{208525683382585211788006284649158144561296072019418355424214631480683811658649824000}$,
 U77 → - $\frac{70393685154110708973348588959486050341962532365889873475218820631416619256045083}{347542805637642019646677141081930240935493453365697259040357719134473019431083040000}$,
 U78 → - $\frac{23363563748425462001166282985384432574609487107112025803332032179276923505785437}{69508561127528403929335428216386048187098690673139451808071543826894603886216608000}$

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U79 → - 12317807864873283420847671569598044351538415567453216770606309250696341927002029
57923800939607003274446190180321706822582242227616209840059619855745503238513840000 ,
U80 → - 26537967226336380256932992618279352433146668932511164339174346119851319850152669
69508561127528403929335428216386048187098690673139451808071543826894603886216608000 ,
U81 → - 120746789515786911786578163981978144763398211020798512940026909535096184188325
556068489020227231434683425731088385496789525385115614464572350615156831089732864 ,
U82 → - 9771811120624939598269543137505526241339130585276754933023531585728313997936339
23169520375842801309778476072128682729032896891046483936023847942298201295405536000 ,
U83 → - 6301005586964814724576850947136148515072766951307856649171101893054787039542873
28961900469803501637223095090160853411291121113808104920029809927872751619256920000 ,
U84 → - 762904414743221335785898096792034908599766709485547776983386950482238862103661
1695330759208009851935010444302098736270699772515596385562720581143770826493088000 , U85 →
63259036520945883900629510579151873688648222856618449293660762566574808542809717
0, U86 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
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U87 → - 3696747318729169550998017741225752060941858080000 ,
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U88 → - 40101092958189463805385823970991950877172321542195837581579736823208425318971120000 ,
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U89 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
155904276176392353915865267715108669026569681354812131179811791541002846036440981
U90 → - 521314208456463029470015711622895361403240180048545888560536578701709529146624560000 ,
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U91 → - 347542805637642019646677141081930240935493453365697259040357719134473019431083040000 ,
6510960936844358465641208684332022443587042597229899955906195366695392712818783
U92 → - 19307933646535667758148730060107235607527414075872069946686539951915167746171280000 ,
6893658753687449628666166111877476854846277936739632581378288457829111739680707
U93 → - 28961900469803501637223095090160853411291121113808104920029809927872751619256920000 ,
66237696282206757128954283656200263032009795990597732415862727484765276353672537
U94 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
413857284088650504348038792669732326250329859144444776998109692535591123855871
U95 → - 1695330759208009851935010444302098736270699772515596385562720581143770826493088000 ,
73015642602569888258996879153544375079176858001122228649545896040126032829329627
U96 → - 173771402818821009823338570540965120467746726682848629520178859567236509715541520000 ,
4716793426599042231765905883351650821203483553971295207642733604473733627338467
U97 → - 19307933646535667758148730060107235607527414075872069946686539951915167746171280000 ,
26025299126653694461722510173328405199981188065579673797672572561775087656937209
U98 → - 57923800939607003274446190180321706822582242227616209840059619855745503238513840000 }}

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Ulist =

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{U1, U2, U3, U4, U5, U6, U7, U8, U9, U10, U11, U12, U13, U14, U15, U16, U17, U18, U19, U20,
U21, U22, U23, U24, U25, U26, U27, U28, U29, U30, U31, U32, U33, U34, U35, U36, U37, U38,
U39, U40, U41, U42, U43, U44, U45, U46, U47, U48, U49, U50, U51, U52, U53, U54, U55,
U56, U57, U58, U59, U60, U61, U62, U63, U64, U65, U66, U67, U68, U69, U70,
U71, U72, U73, U74, U75, U76, U77, U78, U79, U80, U81, U82, U83, U84, U85,
U86, U87, U88, U89, U90, U91, U92, U93, U94, U95, U96, U97, U98} /. %

```

"Give U(n) value to list of Ulist. Ulist[[1,n]]=Un";

"double check if Ulist is correct. "; Ulist[[1, 3]]

```

TD = Table[Sqrt[(Ulist[[1, (2 × (7 × (j - 1) + 1))]]^2 + (Ulist[[1, (2 × (7 × (j - 1) + 1) - 1)])^2,
{j, 7}], {i, 7}];

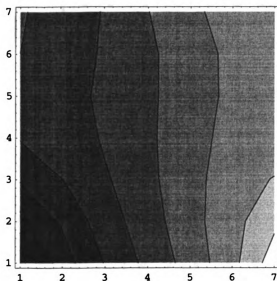
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"TD[[i,j]] is the displacement at position {i,j}";

N[%]

0	0.000397902	0.000636563	0.000793737	0.000913013	0.00101909	0.00112839
0.00055708	0.000663627	0.000770334	0.000857346	0.000936392	0.00101069	0.00107818
0.000776718	0.000833543	0.000872305	0.000914931	0.000965958	0.00102364	0.0010854
0.00087507	0.00089985	0.000915151	0.000942044	0.000984314	0.00103857	0.00110172
0.000890793	0.000878236	0.000893197	0.000931388	0.000985405	0.00104844	0.00111826
0.000816166	0.000759257	0.000811204	0.000892521	0.00097927	0.00106492	0.00114791
0.000740533	0.000549315	0.000709462	0.000864815	0.000992764	0.00110837	0.00123026

ListContourPlot[TD]



- ContourGraphics -

APPENDIX V

$$K_A = \begin{pmatrix} 3 & 0 & 0 & 3 & -3 & -3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 8 & 2 & 0 & -2 & -8 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2 & 22 & 5 & -16 & -5 & 0 & 0 & -6 & -5 & 0 & 3 & 0 & 0 & 0 & 0 & 0 \\ 3 & 0 & 5 & 22 & -5 & -6 & 0 & 0 & -5 & -16 & 2 & 0 & 0 & 0 & 0 & 0 & 0 \\ -3 & -2 & -16 & -5 & 22 & 5 & -3 & -3 & 0 & 5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -3 & -8 & -5 & -6 & 5 & 22 & -2 & -8 & 5 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -3 & -2 & 22 & 5 & -16 & -5 & 0 & 0 & -3 & -3 & 0 & 5 & 0 & 0 \\ 0 & 0 & 0 & 0 & -3 & -8 & 5 & 22 & -5 & -6 & 0 & 0 & -2 & -8 & 5 & 0 & 0 & 0 \\ 0 & 0 & -6 & -5 & 0 & 5 & -16 & -5 & 44 & 10 & -16 & -5 & 0 & 0 & -6 & -5 & 0 & 5 \\ 0 & 0 & -5 & -16 & 5 & 0 & -5 & -6 & 10 & 44 & -5 & -6 & 0 & 0 & -5 & -16 & 5 & 0 \\ 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & -16 & -5 & 22 & 5 & 0 & 0 & 0 & 0 & -6 & -5 \\ 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 & -5 & -6 & 5 & 22 & 0 & 0 & 0 & 0 & -5 & -16 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & -2 & 0 & 0 & 0 & 0 & 11 & 5 & -8 & -3 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & -8 & 0 & 0 & 0 & 0 & 5 & 11 & -2 & -3 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 5 & -6 & -5 & 0 & 0 & -8 & -2 & 22 & 5 & -8 & -3 \\ 0 & 0 & 0 & 0 & 0 & 0 & 5 & 0 & -5 & -16 & 0 & 0 & -3 & -3 & 5 & 22 & -2 & -3 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 5 & -6 & -5 & 0 & 0 & -8 & -2 & 22 & 5 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 5 & 0 & -5 & -16 & 0 & 0 & -3 & -3 & 5 & 22 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & -8 & -2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 & -3 & -3 \end{pmatrix}$$

$$K_B = \begin{pmatrix} 8 & 0 & -8 & -2 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 \\ 0 & 3 & -3 & -3 & 0 & 0 & 3 & 0 & 0 & 0 & 0 & 0 \\ -8 & -3 & 22 & 5 & -8 & -2 & -6 & -5 & 0 & 5 & 0 & 0 \\ -2 & -3 & 5 & 22 & -3 & -3 & -5 & -16 & 5 & 0 & 0 & 0 \\ 0 & 0 & -8 & -3 & 11 & 5 & 0 & 0 & -3 & -2 & 0 & 0 \\ 0 & 0 & -2 & -3 & 5 & 11 & 0 & 0 & -3 & -8 & 0 & 0 \\ 0 & 3 & -6 & -5 & 0 & 0 & 22 & 5 & -16 & -5 & 0 & 2 \\ 2 & 0 & -5 & -16 & 0 & 0 & 5 & 22 & -5 & -6 & 3 & 0 \\ 0 & 0 & 0 & 5 & -3 & -3 & -16 & -5 & 22 & 5 & -3 & -2 \\ 0 & 0 & 5 & 0 & -2 & -8 & -5 & -6 & 5 & 22 & -3 & -8 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 3 & -3 & -3 & 3 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & -2 & -8 & 0 & 8 \end{pmatrix} \times 100000;$$

$$K_{combined} = \begin{pmatrix} K_A & : & 0 \\ .. & : & .. \\ 0 & : & K_B \end{pmatrix}$$

U1	f1	U1	U3
U2	f2	U2	U4
U3	f3	U3	U5
U4	f4	U4	U6
U5	f50	U5	U7
U6	f60	U6	U8
U7	f70	U7	U9
U8	f80	U8	U10
U9	f90	U9	U11
U10	f100	U10	U12
U11	f11	U11	U13
U12	f12	U12	U14
U13	f13	U13	U15
U14	f140	U14	U16
U15	f150	U15	U17
U16	f160	U16	U18
U17	f170	U17	U19
U18	f180	U18	U20
U19	f19	U19	V3
U20	f20	U20	V4
V1	g1	V3	V5
V2	g2	V4	V6
V3	g30	V5	V9
V4	g40	V6	V10
V5	g50	V9	
V6	g60	V10	
V7	g7		
V8	g8		
V9	g90		
V10	g100		
V11	g11		
V12	g12		

UI = ; F = ; UII = ; Uall = ;

U3	0
U4	0
U5	0
U6	0
U7	0
U8	0
U9	0
U10	0
U11	0
U12	0
U14	0
U15	0
U16	0
U17	0
U18	0
U19	0
U20	0
V3	0
V4	0
V5	0
V6	0
V9	0
V10	0
f1	0
f2	-100
f3	-100
f4	0
f11	0
f12	0
f13	0
f19	0
f20	0
g1	0
g2	0
g7	0
g8	0
g11	0
g12	0

UF = ; Zero = ; KIII =

0	3	-3	-3	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
2	0	-2	-8	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
22	5	-16	-5	0	0	-6	-5	0	3	0	0	0	0	0	0	0	0	0
5	22	-5	-6	0	0	-5	-16	2	0	0	0	0	0	0	0	0	0	0
-16	-5	22	5	-3	-3	0	5	0	0	0	0	0	0	0	0	0	0	0
-5	-6	5	22	-2	-8	5	0	0	0	0	0	0	0	0	0	0	0	0
0	0	-3	-2	22	5	-16	-5	0	0	-3	0	5	0	0	0	0	0	0
0	0	-3	-8	5	22	-5	-6	0	0	-8	5	0	0	0	0	0	0	0
-6	-5	0	5	-16	-5	44	10	-16	-5	0	-6	-5	0	5	0	0	0	0
-5	-16	5	0	-5	-6	10	44	-5	-6	0	-5	-16	5	0	0	0	0	0
0	2	0	0	0	0	-16	-5	22	5	0	0	0	-6	-5	0	3	0	0
3	0	0	0	0	0	-5	-6	5	22	0	0	0	-5	-16	2	0	0	0
0	0	0	0	-3	-2	0	0	0	0	5	-8	-3	0	0	0	0	0	0
0	0	0	0	-3	-8	0	0	0	0	11	-2	-3	0	0	0	0	0	0
0	0	0	0	0	5	-6	-5	0	0	-2	22	5	-8	-3	0	0	0	0
0	0	0	0	5	0	-5	-16	0	0	-3	5	22	-2	-3	0	0	0	0
0	0	0	0	0	0	0	5	-6	-5	0	-8	-2	22	5	-8	-3	0	0
0	0	0	0	0	0	5	0	-5	-16	0	-3	-3	5	22	-2	-3	0	0
0	0	0	0	0	0	0	0	0	2	0	0	0	-8	-2	8	0	0	0
0	0	0	0	0	0	0	0	3	0	0	0	0	-3	-3	0	3	0	0
0	1	0	0	0	0	0	0	0	1	0	0	0	0	0	0	0	-8	-2
$\frac{3}{2}$	0	0	0	0	0	0	0	$\frac{3}{2}$	0	0	0	0	0	0	0	0	-3	-3
-3	$-\frac{5}{2}$	0	0	0	0	0	0	-3	$-\frac{5}{2}$	0	0	0	0	0	0	0	22	5
$-\frac{5}{2}$	-8	0	0	0	0	0	0	$-\frac{5}{2}$	-8	0	0	0	0	0	0	0	5	22
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-8	-3
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-2	-3
11	$\frac{5}{2}$	0	0	0	0	0	0	11	$\frac{5}{2}$	0	0	0	0	0	0	0	2	-6
$\frac{5}{2}$	11	0	0	0	0	0	0	$\frac{5}{2}$	11	0	0	0	0	0	3	0	-5	-16
-8	$-\frac{5}{2}$	0	0	0	0	0	0	-8	$-\frac{5}{2}$	0	0	0	0	0	-3	-2	0	5
$-\frac{5}{2}$	-3	0	0	0	0	0	0	$-\frac{5}{2}$	-3	0	0	0	0	0	-3	-8	5	0
0	$\frac{3}{2}$	0	0	0	0	0	0	0	$\frac{3}{2}$	0	0	0	0	0	3	0	0	0
1	0	0	0	0	0	0	0	1	0	0	0	0	0	0	0	8	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0

[illegible]

$$\begin{aligned} & \{ \{ U3 \rightarrow -\frac{-2041150230705}{13317750352961700}, U4 \rightarrow -\frac{25769698430570}{31723750352961700}, U5 \rightarrow -\frac{8662215172565}{56517750352961700}, \\ & U6 \rightarrow -\frac{79359420331675}{110981252941347500}, U7 \rightarrow -\frac{-12067138811082}{66588751764808500}, U8 \rightarrow -\frac{91853071812403}{110981252941347500}, \\ & U9 \rightarrow -\frac{-105415086089231}{543943758824042500}, U10 \rightarrow -\frac{269461324093314}{332943758824042500}, U11 \rightarrow \\ & -\frac{-247924400314217}{867943758824042500}, U12 \rightarrow -\frac{866230479658018}{975943758824042500}, U14 \rightarrow -\frac{666645560737}{55490626470673750}, \\ & U15 \rightarrow -\frac{-19912355731059}{110981252941347500}, U16 \rightarrow -\frac{397587017928351}{563943758824042500}, U17 \rightarrow \\ & -\frac{-26416321906499}{110981252941347500}, U18 \rightarrow -\frac{564259691405924}{656943758824042500}, U19 \rightarrow -\frac{-282201884005113}{987943758824042500}, \\ & U20 \rightarrow -\frac{500819713669427}{432943758824042500}, V3 \rightarrow -\frac{-678632499652849}{10654200282369360000}, V4 \rightarrow \\ & -\frac{1641723224448110}{2130840056473872000}, V5 \rightarrow -\frac{66686394550089}{343943758824042500}, V6 \rightarrow -\frac{2697536739171}{2663550070592340}, \\ & V9 \rightarrow -\frac{-10056869673993300}{54354200282369360000}, V10 \rightarrow -\frac{83087993445792600}{87867200282369360000}, f1 \rightarrow \\ & \frac{84346766723995}{88785002353078}, f2 \rightarrow \frac{8621943225249440}{78646502353078}, f3 \rightarrow -\frac{183863918674800}{44392501176539}, f4 \rightarrow \\ & -\frac{93998899125853}{4345401176539}, f11 \rightarrow \frac{694453503415314}{54654654176539}, f12 \rightarrow -\frac{703562101587}{142634576539}, f13 \rightarrow \\ & -\frac{469089364343894000}{4756754745176539}, f19 \rightarrow \frac{7609309641345}{456456456539}, f20 \rightarrow \frac{1721670875290200}{88957644176539}, \\ & g1 \rightarrow -\frac{31790674735717}{33463465653078}, g2 \rightarrow -\frac{9733417591862140}{88785002353078}, g7 \rightarrow \frac{55046800022816}{6534524876539}, \\ & g8 \rightarrow \frac{87351542020393}{6576579176539}, g11 \rightarrow -\frac{9463422511822380}{567678345353078}, g12 \rightarrow -\frac{19307074104135}{997584296102} \} \} \end{aligned}$$

N[%]

```
{U3→ -0.00015326539217269, U4→ -0.00015326539217269, U5→  
-0.00015326539217269, U6→ -0.00071475324324991, U7→ -0.0001812188769314,  
U8→ -0.00024856457649176903, U9→ -0.0001937977674698, U10→  
-0.000285645697424133, U11→ -0.000285645697424133, U12→  
-0.000285645697424133, U14→ -0.00036403594316897, U15→ -0.00017942089500091,  
U16→ -0.00029906009398439, U17→ -0.000238025082672836,  
U18→ -0.00038117719721275, U19→ -0.000285645697424133,  
U20→ -0.000285645697424133, V3→ -0.00006369624013694, V4→  
-0.00035315173762, V5→ -0.0001938874971248, V6→ -0.0005959370274947,  
V9→ -0.00018502470134319, V10→ -0.00047104487235471, f1→  
0.95001142635067, f2→ 109.62907398655632, f3→ -4.141778764472281, f4→  
-21.631811496107005, f11→ 12.70620981649933, f12→ -4.932619555920048, f13→  
98.61541943476531, f19→ 16.670395461247498, f20→ 19.353827219991384, g1→  
-0.95001142635067, g2→ -109.62907398655632, g7→ 8.423994255565799, g8→  
13.2822155220159, g11→ -16.670395461247498, g12→ -19.353827219991384}}
```

U1 = 0; U2 = 0; U13 = 0; V1 = 0; V2 = 0; V7 = (U3 + U11) / 2;

V8 = (U4 + U12) / 2; V11 = U19; V12 = U20;

"Boundary Conditions and Displacement Constrain Equation";

Ulist = {U1, U2, U3, U4, U5, U6, U7, U8, U9,

U10, U11, U12, U13, U14, U15, U16, U17, U18, U19, U20};

Vlist = {V1, V2, V3, V4, V5, V6, V7, V8, V9, V10, V11, V12};

"Give displacement value to list of Ulist and Vlist";

"Compute the general displacement at nodal position");

UD = Table[$\left[\sqrt{(Ulist[[1, (2 \times n) - 1]])^2 + (Ulist[[1, (2 \times n)])^2}\right], \{n, 1, 10\}];$

N[%]

```
{0.000000000000000000, 0.000826648034066665,  
0.000731311222584203, 0.000847252275786886,  
0.000832209342095123, 0.000932414037579913, 0.000684510360745850,  
0.000727484384202640, 0.000891287339200160, 0.001191523504541550}
```

VD = Table[$\left[\sqrt{(Vlist[[1, (2 \times n) - 1]])^2 + (Vlist[[1, (2 \times n)])^2}\right], \{n, 1, 6\}];$

N[%]

```
{0.000000000000000000, 0.000773086716148328, 0.001031152270031470,  
0.000879531000000000, 0.000963540372081521, 0.001191523504541550}
```

APPENDIX VI

The Result of FEA computation and Displacement Difference Analysis

Reference Model Result: (From re-meshed 7x7 modal, 72 Elements)

0	0.0003	0.00059	0.00079	0.00091	0.00102	0.001128
0.00042	0.00046	0.00077	0.00086	0.00094	0.00101	0.001078
0.00078	0.00083	0.00087	0.00091	0.00097	0.00102	0.001085
0.00088	0.0009	0.00092	0.00094	0.00098	0.00104	0.001102
0.00089	0.00088	0.00089	0.00093	0.00099	0.00105	0.001118
0.00082	0.00076	0.00081	0.00089	0.00098	0.00106	0.001148
0.00074	0.00065	0.00071	0.00086	0.00099	0.00111	0.00123

Modular Model Result (interpolation value in yellow area).

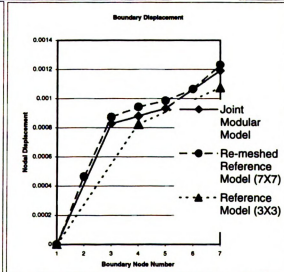
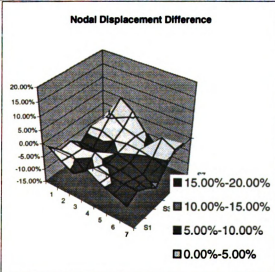
0	0.00026	0.00052	0.00077	0.00086	0.00095	0.001031
0.00037	0.00041	0.00067	0.00081	0.00088	0.00094	0.001009
0.00073	0.00078	0.00083	0.00084	0.00089	0.00094	0.000986
0.00079	0.00081	0.00083	0.00088	0.00091	0.00094	0.000964
0.00085	0.00084	0.00083	0.00088	0.00093	0.00099	0.00104
0.00077	0.00077	0.00078	0.00085	0.00091	0.00106	0.001116
0.00068	0.00071	0.00073	0.00081	0.00089	0.00104	0.001192

Boundary Displacement:

0
4
7
1
3
4
5
7
1
2
3
4
5
6
7

Displacement difference: (result-reference)/MAX(reference) (%)

0.00%	-3.27%	-5.79%	-1.68%	-4.38%	-6.01%	-7.91%
-4.18%	-4.09%	-8.07%	-3.97%	-4.97%	-5.59%	-5.66%
-3.69%	-4.44%	-3.71%	-5.76%	-6.06%	-6.90%	-8.07%
-6.97%	-7.36%	-6.97%	-5.08%	-6.24%	-8.38%	-11.23%
-3.54%	-3.13%	-4.96%	-3.99%	-4.31%	-5.08%	-6.40%
-4.09%	1.11%	-2.55%	-3.79%	-5.48%	-0.24%	-2.63%
-4.55%	4.61%	1.47%	-4.51%	-8.25%	-5.44%	-3.15%



The Result of FEA computation and Displacement Difference Analysis (2)

Reference Model Result: (From 3x3 (8 elements) model).

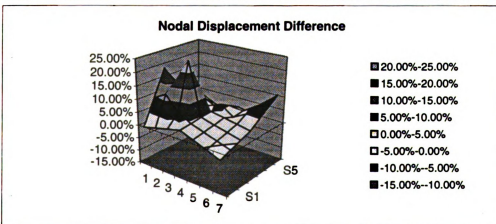
0	0.00024	0.00049	0.00073	0.00084	0.00095	0.00106
0.00026	0.00027	0.00059	0.00076	0.00085	0.00093	0.00102
0.00051	0.0006	0.00055	0.00079	0.00086	0.00092	0.00098
0.00077	0.00078	0.0008	0.00082	0.00086	0.0009	0.00094
0.00073	0.00076	0.00079	0.00081	0.00091	0.00093	0.00099
0.0007	0.00073	0.00077	0.0008	0.00088	0.00099	0.00103
0.00066	0.00071	0.00075	0.00079	0.00089	0.00098	0.00108

Modular Model Result (interpolation value in yellow area).

0	0.00026	0.00052	0.00077	0.00086	0.00095	0.00103
0.00037	0.00041	0.00067	0.00081	0.00088	0.00094	0.00101
0.00073	0.00078	0.00083	0.00084	0.00089	0.00094	0.00099
0.00079	0.00081	0.00083	0.00088	0.00091	0.00094	0.00096
0.00085	0.00084	0.00083	0.00088	0.00093	0.00099	0.00104
0.00077	0.00077	0.00078	0.00085	0.00091	0.00106	0.00112
0.00068	0.00071	0.00073	0.00081	0.00089	0.00104	0.00119

Displacement difference: (result-reference)/MAX(reference) (%)

0.00%	1.16%	2.31%	3.47%	1.52%	-0.43%	-2.38%
9.66%	12.13%	6.73%	3.95%	2.31%	0.68%	-0.95%
19.32%	15.24%	24.25%	4.42%	3.11%	1.79%	0.48%
2.11%	2.17%	2.22%	4.90%	3.90%	2.90%	1.91%
10.22%	7.13%	4.03%	5.97%	2.19%	5.00%	4.66%
6.16%	3.59%	1.03%	3.62%	2.80%	6.18%	7.42%
2.09%	0.06%	-1.97%	1.27%	0.27%	5.22%	10.17%



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