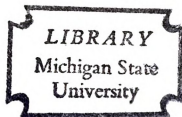


A STOCHASTIC APPROACH TO
POPULATION BALANCE MODELS

Thesis for the Degree of Ph. D.
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ANILKUMAR NARAYAN KHARKAR
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This is to certify that the
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ABSTRACT

A STOCHASTIC APPROACH TO POPULATION BALANCE MODELS

By

Anilkumar Narayan Kharkar

Population balance models are extremely important in the quantitative analysis of chemical process systems characterized by aggregates of matter, as well as ecological systems. A principal aim of this work is to cast these models into a very general stochastic framework.

By and large the development of population balance models in chemical engineering thus far has been based upon purely deterministic considerations and analogies with molecular processes. In this dissertation a probabilistic approach has been presented to develop a common basis for the population balance models used in the residence time distribution analysis of process vessels, various particulate processes and a large class of problems in ecological systems.

The development of the population balance models is based upon the theory of stochastic population processes. The state of each entity in a population is represented by a point in a n -dimensional Euclidean space and the dynamics of the population is characterized by birth, death and movement of the entities in a closed domain in this state space.



In the context of populations of reproducing entities, the process of crystal growth with secondary nucleation and the dynamics of biological populations with reproduction are of particular interest. A common feature of the crystallization process with the state of a crystal described by its characteristic length and many biological populations with individuals characterized by a measure of their maturity is that all the entities in the population are in a particular fixed state at the instant of their first appearance in the population as a result of reproduction, irrespective of the state of the parent. To account for the fact that identical entities do not necessarily grow at the same rate, the growth of each entity is characterized by a nonhomogeneous diffusion process on the interval $[0, \bar{a}]$, where \bar{a} can possibly approach infinity. Some types of boundary behavior of practical interest are discussed. Backward diffusion equations for the probability generating functional and the first two factorial moment distributions, as well as forward diffusion equations for the first two factorial moment densities are derived for a population of reproducing entities with no external input.

The diffusion equations for the moment distributions have been analyzed for some simple cases. The following results have been obtained under mild conditions: i) a general solution for the first moment density in terms of an infinite series of eigenfunctions of the diffusion operator; ii) after a sufficiently large time the first and second factorial moment of the number of entities in any subinterval in $[0, \bar{a}]$ grow exponentially with parameters α_1 and $2\alpha_1$ respectively, where α_1 is the dominant eigenvalue of the

diffusion operator; and iii) the coefficient of variation of the number of entities in any subinterval in $[0, \bar{a}]$ reaches the same constant value after a sufficiently large time. Moreover, this asymptotic value in (iii) has been shown to be inversely proportional to the square root of the initial number of entities in the entire population. An interesting facet of the diffusion process in the context of reproducing populations is that under certain conditions the location of the boundary \bar{a} will determine whether the population will increase or decrease. A graphical solution for the determination of this critical value of \bar{a} has been obtained in terms of dimensionless parameters for the case where \bar{a} is an absorbing barrier. The diffusion equation for the probability generating functional has been solved numerically to obtain the extinction probability of a population. These results had not been obtained thus far for the particular cases considered.

In general, a population will also have an input of entities from a source external to the population. Expressions have been derived for the probability generating functional and the first two factorial moment distributions for a population with a nonhomogeneous Poisson input. It is shown that in the case of a population of non-reproducing entities if the birth, death, and movement of entities in a n -dimensional spatial domain can be characterized by a Markov process, and if the population receives an external input in the form of a nonhomogeneous Poisson process, then the number of entities in any set in the spatial domain at any time is a Poisson-distributed random variable. The use of this result in validating some assumptions behind the models developed in this work has been demonstrated by means of a simple experiment.

The theoretical results obtained in this work are of considerable importance in a broad class of problems related to chemical engineering and ecological systems. Several examples of practical interest are also discussed along with the results.

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By

Anilkumar Narayan Kharkar

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To VAISHALI,
our little angel.



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TABLE OF CONTENTS

	Page
DEDICATION	ii
ACKNOWLEDGEMENTS	iii
LIST OF TABLES	vi
LIST OF FIGURES	vii
CHAPTER	
I INTRODUCTION	1
II MATHEMATICAL BACKGROUND	16
2.1 Point Processes	17
2.2 Counting Measures	19
2.3 Moment Measures and Moment Densities	20
2.4 An Illustrative Example	22
2.5 Generating Functionals	25
2.6 Stochastic Population Processes	28
III THE IDENTIFICATION PROBLEM	29
3.1 Residence Time Distribution of Flow Vessels ...	30
3.2 Life-span Distributions in Biological Populations	32
3.3 A Dispersion Phenomenon in Crystallization ...	35
3.4 Generalizations	37
3.5 Selection of a Proper Model	39
IV DERIVATION OF THE DIFFUSION EQUATIONS	44
4.1 Description of the Process	45
4.2 Boundary Conditions	47
4.3 The Backward Diffusion Equations	50
4.4 The Forward Diffusion Equations	61
4.5 Comparison of the Backward and Forward Equations	72
V SOLUTION OF THE DIFFUSION EQUATIONS	76
5.1 General Considerations	76
5.2 Solution of the Diffusion Equations for the First Moment	79

CHAPTER		Page
	5.3 Solution of the Diffusion Equations for the Second Moment	90
	5.4 The Problem of "Critical Length"	102
	5.5 Solution of the Diffusion Equation for the PGF -- Computation of Extinction Probabilities	111
VI	POPULATIONS WITH AN EXTERNAL INPUT	130
	6.1 Description of the Process	131
	6.2 The Probability Generating Functional of the Process	135
	6.3 Moment Distributions of the Population	141
	6.4 External Input in Populations of Nonreproducing Entities	144
	6.5 Applications	148
	6.6 Justification of a Poisson Input	154
VII	AN EXPERIMENT BASED ON THE THEORY	161
	7.1 Motivation	161
	7.2 Methods and Materials	162
	7.3 Results and Discussion	170
VIII	CONCLUSIONS	181
	8.1 Summary of the Contributions of this Dissertation	181
	8.2 Areas for Future Research	183
	REFERENCES	187
	APPENDIX	
A	SIMPLIFICATION OF THE DIFFUSION MODEL FOR SOME BIOLOGICAL POPULATIONS	192
B	NUMERICAL COMPUTATION OF THE "CRITICAL LENGTH"	199
C	COMPUTATION OF EXTINCTION PROBABILITIES FOR A POPULATION OF REPRODUCING ENTITIES	204
D	ANALYSIS OF EXPERIMENTAL DATA	208

LIST OF TABLES

Table	Page
5.1 Model Parameters for the Dynamics of Populations of Bluegills (<i>Lepomis macrochirus</i>)	110
5.2 Simulation Parameters for the Computation of Extinction Probability of a Population	112
5.3 Results of Simulation No. 1 -- Extinction Probability for a Population Generated by One Ancestor	116
5.4 Results of Simulation No. 2 -- Extinction Probability for a Population Generated by One Ancestor	118
5.5 Results of Simulation No. 3 -- Extinction Probability for a Population Generated by One Ancestor	120
5.6 Results of Simulation No. 4 -- Extinction Probability for a Population Generated by One Ancestor	122
5.7 Results of Simulation No. 5 -- Extinction Probability for a Population Generated by One Ancestor	124
5.8 Results of Simulation No. 6 -- Extinction Probability for a Population Generated by One Ancestor	126
5.9 Results of Simulation No. 7 -- Extinction Probability for a Population Generated by One Ancestor	128
6.1 Statistical Analysis of Cereal Leaf Beetle Oviposition Data with Mated Females	158
6.2 Statistical Analysis of Cereal Leaf Beetle Oviposition Data with Egg-laying Females	159
7.1 Important Steps in the Experiments with CLB Populations	168
7.2 Cereal Leaf Beetle Population Data from Experiment A ..	171
7.3 Cereal Leaf Beetle Population Data from Experiment B ..	173
7.4 Statistical Analysis of Cereal Leaf Beetle Population Data from Experiment A	175
7.5 Statistical Analysis of Cereal Leaf Beetle Population Data from Experiment B	177

LIST OF FIGURES

Figure		Page
2.1a	Counting Measure for a Sample Realization of a Point Process	24
2.1b	Product Counting Measure for the Sample Realization of the Point Process in Figure 2.1a	24
3.1	An Example of Lumping of Variables in the Individual State Space for a Biological Population	40
5.1	Graphical Solution of Equation (5.64) for "Critical Length"	107
6.1	A Typical Sequence of the Daily Egg Input by a Single Cereal Leaf Beetle Female	156
7.1	Comparison of Plant Condition between Experiments A and B	166
B.1	Iteration Scheme for Obtaining a Root of Equation (5.64)	201



CHAPTER I

INTRODUCTION

The use of mathematical models in the analysis of dynamic systems has become a common practice since the advent of high-speed large-capacity computers. The mathematical models used in chemical engineering practice can be classified into three broad categories: (a) Transport phenomena models, (b) Population balance models and (c) Empirical models. The transport phenomena models use physico-chemical properties of matter such as diffusivity, thermal conductivity and viscosity to obtain a dynamic description of the system in the form of equations for conservation of mass, momentum and energy. Systems comprised of ensembles of discrete entities are described by the population balance models. The dynamic description of these systems involves conservation equations for the number of entities in the population. The empirical models are comprised of purely empirical equations or regression analysis of data to describe the system. In particular, the applications of the population balance models include the residence time distribution analysis of imperfectly mixed process vessels (e.g., Levenspiel and Bischoff, 1963; Bischoff, 1966), and various particulate processes such as crystallization (Randolph and Larson, 1971), size reduction (Randolph and Durando, 1971), particle agglomeration (Hulburt and Katz, 1964), fermentation (Tsuchiya et al., 1966) and liquid-liquid extraction (Ramakrishna, 1972).



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The polymerization processes can also be described by population balance models (e.g., Funderburk, 1969).

Typically, the population balance models describe the entities in a population as being distributed in a n -dimensional Euclidean space E according to a density $x(z_1, z_2, \dots, z_n, t)$ at time t . The spatial co-ordinates z_1, \dots, z_n represent the relevant properties of the entities such as geographical location, size, shape, activity in case of catalyst particles, etc. Thus, $x(z_1, \dots, z_n, t) dz_1 \dots dz_n$ represents the number of entities in the infinitesimal volume element $dz_1 \dots dz_n$ located at the point (z_1, \dots, z_n) in E at time t . The "movement" of the entities along the i -th co-ordinate ($i = 1, \dots, n$) of E is characterized by a rate of change of position $v_i(z_1, \dots, z_n, t)$ in that direction. The rate of addition of entities to the infinitesimal volume element $dz_1 \dots dz_n$ from a source internal or external to the population is expressed as the birth rate $B(z_1, \dots, z_n, t)$ of entities per unit time per unit volume of E . Similarly, the rate of change in the number of entities in the infinitesimal volume element representing a permanent loss to the population is characterized as the death rate $D(z_1, \dots, z_n, t)$. If a number balance is written for the entities in the infinitesimal volume element at (z_1, \dots, z_n) at time t by equating the net rate of change in the number of entities due to birth, "movement" along any of the co-ordinates of E , and death to the net rate of ^(6.3) accumulation of entities, the following conservation equation results (Himmelblau and Bischoff, 1968):

$$\frac{\partial x(z_1, \dots, z_n, t)}{\partial t} + \sum_{i=1}^n \frac{\partial [v_i(z_1, \dots, z_n, t)x(z_1, \dots, z_n, t)]}{\partial z_i} = B(z_1, \dots, z_n, t) - D(z_1, \dots, z_n, t) \quad (1.1)$$

The movement of the entities is restricted to a closed domain in E , enclosed by a boundary Γ . The behavior of the entities at the boundary determines the boundary conditions and the density $x_0(z_1, \dots, z_n, 0)$ at time zero describes the initial condition for (1.1).

In general, the quantities x , v_i , B and D will be functions of other variables as well, such as temperature, concentration of a chemical in the medium surrounding the entities, etc. In addition, B will be a function of x when the entities themselves act as sources (reproducing populations). The quantities v_i and D may also depend on x . It should be noted that in the derivation of (1.1) it is assumed that all the entities at a given point in E move in any given direction at essentially the same rate.

In the context of residence time distribution of a process vessel, (1.1) represents a case analogous to a plug flow vessel, where all the fluid elements at a given cross-section of the vessel move toward the outlet at essentially the same velocity. The study of nonideal vessels, i.e., vessels which cannot be characterized by perfect mixing or plug flow, commonly involves a tracer analysis of the flow system. A known input such as a pulse of a tracer is introduced in the inlet stream of the flow system or its scale model and the concentration of the tracer in the outlet stream is measured as a function of time. A mathematical model is then constructed to give the same input-output relation as that obtained

in the tracer experiment. Two distinct methods are available to achieve this: The first method characterizes the flow system as a combination of a number of perfectly mixed vessels connected in series (and parallel, if necessary). This approach is commonly known as the mixing cell approach. The second method assumes that the flow is essentially a modification of a plug flow caused by an effective "diffusion" of the fluid elements in the axial and transverse directions in the flow vessel. The models used in this method are commonly called the dispersed plug flow models. The equations for the concentration of the tracer in the outlet stream are

$$\frac{dc_i}{dt} = k(c_{i-1} - c_i), \quad i = 1, 2, \dots, m, \quad \begin{matrix} \boxed{1} \\ \boxed{2} \\ \boxed{3} \\ \boxed{4} \\ \boxed{5} \end{matrix} \quad (1.2)$$

for the mixing cell model with only a series combination of the perfectly mixed (hypothetical) vessels of equal volume, and

$$\begin{aligned} \frac{\partial c(z_1, z_2, t)}{\partial t} = & \frac{\sigma_L^2}{2} \frac{\partial^2 c(z_1, z_2, t)}{\partial z_1^2} + \frac{\sigma_R^2}{2z_2} \frac{\partial}{\partial z_2} \left[\frac{1}{z_2} \frac{\partial c(z_1, z_2, t)}{\partial z_2} \right] - \\ & - r_L \frac{\partial c(z_1, z_2, t)}{\partial z_1} - r_R \frac{\partial c(z_1, z_2, t)}{\partial z_2} - D \end{aligned} \quad (1.3)$$

for the case of a dispersed plug flow model for a cylindrical vessel. In (1.2) c_i represents the concentration of the tracer in the i -th vessel, t denotes time, k is the mean residence time of the fluid in a single vessel and c_0 is the inlet concentration of the tracer for the first vessel. In the mixing cell models all the vessels in a series combination are often assumed to be of equal volume to keep the model as simple as possible. In (1.3) the quantities z_1 and z_2 represent the axial and radial coordinates respectively, $c(z_1, z_2, t)$ denotes the concentration of

the tracer at a given point in the vessel at time t , $\frac{\sigma_L^2}{2}$ and $\frac{\sigma_R^2}{2}$ are the dispersion coefficients in the axial and radial directions respectively, r_L and r_R are the velocities in the two directions, D denotes the rate of loss of the tracer per unit volume per unit time due to chemical reaction, absorption, etc., and t denotes time. It should be noted that (1.1) thru (1.3) are essentially deterministic in nature.

Since (1.1) represents the conservation equation for the members of the population in a very general sense, it can be seen that the equation can be used to describe any biological population as well. Indeed, the equation has been used to characterize populations of daphnia (Sinko and Streifer, 1969), a species of worms which divide by fission (Sinko and Streifer, 1971), shrimp (Billups, et al., 1971) as well as an insect pest population (Barr, Kharkar and Lee, 1972). Equations similar to (1.2) and (1.3) have also been proposed (Weiss, 1968; Kendall, 1948) and used (Takahashi, 1968; Stuart and Merkle, 1965) to describe cell populations. A detailed discussion of the applicability of the "mixing cell" and "dispersion" analogies in describing a general population is given later in the text. Mathematical modeling of biological populations has become all the more important today in view of the pressing need for mankind to engineer the effects of urbanization and industrialization on the ecological systems. Equations (1.1) thru (1.3) are commonly used in the modeling of chemical engineering systems and the methods of parameter estimation are quite well established. A direct analogy of these models with the models for biological systems makes it possible to think of analogous methods of parameter estimation as well.

Literature in the fields of ecology, applied probability and mathematical biology abounds with works on the dynamics of biological populations. The models for single-species population dynamics used in these fields can be classified into three broad categories:

- (i) models which treat all the members of a population as identical entities. These models usually consist of a single difference or ordinary differential equation;
- (ii) models which divide the population into distinct groups according to age, maturity, etc. (i.e., define a discrete state space to describe an entity) -- these models commonly use a system of difference or ordinary differential equations;
- (iii) models which allow the individuals in the population to occupy any arbitrary point in a n -dimensional Euclidean space E with the relevant descriptors such as age, location, etc. as co-ordinates (i.e., define a continuous state space to describe an entity), as is done in the development of (1.1). These models involve the use of integral or partial differential equations.

All the models can be further classified into stochastic and deterministic types. In the partial differential equation approach for a deterministic analysis, use is made of a form of (1.1) or (1.3). The integral equation approach uses the renewal theory (see e.g., Feller, 1966) quite extensively and is a very convenient alternative to models using (1.1) for the case where $n = 1$ and v_1 is a constant. In more complex situations the integral equations for the system lose their simple character. As far as the

deterministic models are concerned, models of types (i) and (ii) can be shown (Kharkar, 1971) to be approximations of the partial differential equation approach. As regards the stochastic models, much of the work deals with characterizations of type (i), (ii), or simple cases of (iii) in which the individuals are characterized by only one descriptor, namely, the age of an individual. The use of integral equations is very convenient for obtaining a stochastic description of the age distribution of a population (in terms of moments of the age distribution), and is used in almost all cases.

Relatively little work has been done in the stochastic modeling of biological populations characterized as in case (iii) with descriptors other than age. Adke and Moyal (1963) developed a stochastic model for a population diffusing on the real line to characterize the spatial migration of the population. Birth and death rates were assumed to be constant. Adke (1964a) extended the model to the case where the birth and death rates can be functions, and later (Adke, 1964b) to the case where the movement of the diffusing individuals is restricted by absorbing or reflecting barriers. Sevast'yanov (1958, 1961) studied the extinction probability of a population diffusing in a compact region with an absorbing boundary for the case where the individuals in the population reproduce by a branching process. Davis (1965, 1967a, 1967b) studied the asymptotic properties of a population diffusing in an abstract space, and multiplying according to a branching process, i.e., an individual was assumed to produce a random number of offspring only at the end of its life. As in the model of Adke and Moyal (1963), Davis assumed that the offspring are in the same

state as their parent. Radcliffe (1972) extended the results of Davis (1965) by allowing immigration into the population in the form of a nonhomogeneous Poisson process.

Population models are also of relevance in some areas of physics. A great deal of work has been done on stochastic descriptions of processes involving transport of elementary particles (see e.g., Brockwell, 1966 for a class of related problems and bibliography). In a sense, these processes are analogous to the branching processes discussed in the previous paragraph: an elementary particle in motion imparts some or all of its momentum to other particles in its way. The new particles produced are in the same location as the parent particle. When a particle with energy ϵ sets k new particles in motion after a collision, the sum of energies of all the new particles must be equal to ϵ . Thus even in the case where the parent elementary particle does not come to rest, it may be necessary to assume that the parent is replaced by a "new" particle with an appropriate energy level, when the energy distribution of the particles is being studied.

While some problems of chemical engineering interest can possibly be described as branching processes and some phenomena in the biological world besides spatial migration can also be cast in this framework, this analysis is often too cumbersome to use. This is true of a large class of problems related to biological and chemical engineering systems, namely, the processes in which the parent and the offspring are in different states and the parent does not die while giving birth to an offspring. In a biological population of higher organisms this is the case when the individuals

are described by their maturity in terms of their size, weight, etc. An example of a chemical process of this type is the production of nuclei in a crystallizer. Under certain conditions the crystals present in a crystallizer produce new nuclei ("offspring") without any effect on the size of the parent crystal. Theoretical studies on populations often concentrate on the asymptotic properties of the population, i.e., the behavior of the population for large values of time, because an analytical treatment of the problem for short times is not possible in many cases with the currently available methods. Also, the dynamics of the system are often described by assigning stationary transition probabilities to the dynamic changes in the population, thus resulting in a model with time-independent parameters. In contrast, the short-term behavior of a system may be of crucial importance in a problem of engineering interest when the model is used to evaluate control alternatives. In many cases the time-dependence of the phenomena underlying the dynamics of a process cannot be ignored. This is particularly true of ecological systems due to seasonal changes in temperature, solar radiation, etc. Although a general model of a population of engineering interest may be too complex to be solved analytically, a numerical solution of sufficient accuracy can usually be obtained.

The need for population balance models which are cast in a very general stochastic framework and can be directly applied to systems of engineering interest, has motivated this research. Such models are of great value in studying a broad class of problems related to the design and management of ecological systems. In this context the descriptors for the individuals in the population

may be size or weight in the case of fish and plankton in a lake, productive value for trees in a forest or orchard, maturity in case of insect pests and their parasites or predators in a farmland, or the level of acceptance of a new idea or economic status in human populations. To have a greater confidence in the management and control policies derived from the models of these systems one would like to know the expected behavior of the populations as well as a measure of deviations from the expected behavior due to random phenomena. A general stochastic formulation of population balance models is also of value in the analysis of systems encountered in chemical engineering practice. The importance of this approach to particulate processes has been recognized only recently (Curl, 1967; Katz and Shinnar, 1969). Probabilistic models of flow in packed beds, where the fluid elements are assumed to move in discrete jumps, have been used quite successfully for describing the residence time distributions for some time (see e.g., Levenspiel and Bischoff, 1963; Buffham et al., 1970; Schmalzer and Hoelscher, 1971; Srinivasan and Mehata, 1971).

As indicated above, one use of a stochastic model is in predicting the fluctuations in the properties of a population. For small populations the stochastic fluctuations are always important if the rate processes such as birth, growth and death are random in nature. A stochastic model enables one to determine how large a population must be for the random fluctuations to be unimportant. Data on many biological systems shows a considerable scatter and hence the observations are often recorded in the form of means and standard deviations. If the contribution of the errors due to

crudeness of observations, subjective judgement of the observer, etc., to the scatter in the data is known, the variance in the repeated observations due to random phenomena alone can possibly be estimated and used as an additional parameter in testing the mathematical models. This particular facet of stochastic modeling is demonstrated in this work by a controlled experiment, which is discussed later in the text.

Experimental data on particulate processes and residence time distribution analysis of process vessels often exhibits an appreciable amount of scatter. Since the population balance models currently used to characterize these processes are largely based upon purely deterministic considerations and analogies with molecular processes, these models cannot account for the scatter. In molecular processes the number of molecules involved in an experimental setup is almost always very large and consequently, the stochastic fluctuations are insignificant. On the other hand, a very much smaller number of entities (i.e., particles, fluid elements, etc.) are encountered in the experiments with particulate processes and residence time distribution analyses. This fact coupled with a stochastic nature of certain phenomena such as formation and growth of the particles or movement of the eddies of the fluid inside the experimental vessels is responsible for at least a part of the scatter in the data. Stochastic models for such systems provide a useful means for the theoretical analysis of such data.

For the case of reproducing populations with no external input, the coefficient of variation ($\sqrt{\text{variance}}/\text{mean}$) of the

population, which is often taken as a measure of stochastic fluctuations, is always significant if the initial population is small, irrespective of whether the population becomes very large at a later time. Bartlett (1969) indicates this to be the case for age distribution in biological cell populations. A similar result is obtained in this dissertation for a more general description of a population. This fact is quite important in the control of many ecological systems. For example, in the biological control of insect pests, often a small number of parasites or predators is released in the infested area. Control of the pest is usually achieved after a certain time lag required for the parasite (or predator) population to build up to a significant level. A stochastic model for such a system would enable one to estimate a priori the chance of success of such a biological control strategy.

A similar situation may arise in chemical processes of industrial importance. The phenomenon of secondary nucleation, which is a significant factor in industrial crystallizers, serves as a good example to illustrate this. In secondary nucleation new nuclei are generated by the breakup of dendritic growth on the surfaces of a growing crystal or release of micro-clusters of particles during the growth of the parent crystal. Thus, if the initial number of seed crystals is small, the fluctuations in the crystal size distribution will always remain significant. Kane (1971) studied the rates of secondary nucleation of ice crystals in brine. He used a stirred batch crystallizer with liquid isobutylene added directly to the supersaturated brine as a coolant. The experiment

was started with a single ice crystal and the number of crystals in the crystallizer (counted using a photographic technique) and the bulk temperature of the contents of the crystallizer were monitored throughout the experiment. The data on the induction period (i.e., the time required to form a sufficiently large number of crystals from the solution to significantly affect the concentration of the brine so that the freezing point of the solution begins to change) shows a considerable scatter. Expressing the nucleation process as a linear birth process (cf. Parzen, 1962, p. 296), he obtained an estimate of the contribution of the stochastic nature of the process to the scatter in the data. His model was based on the total number of crystals. The rate of formation of new nuclei per parent crystal per unit time was assumed to be a constant. In reality, this rate may be a function of crystal size. The equations developed in this dissertation are applicable to more general situations where the rates of nucleation, growth as well as removal from the crystallizer are functions of size as well as time. By solving these equations one can obtain the mean and variance of crystals in any size range at any point in time. Other situations of industrial importance where similar considerations may apply include failure of a chemical or biochemical reaction due to the catalytic activity of trace quantities of an impurity and growth of mutants of the micro-organisms in a continuous fermenter.

Even in situations where one is not interested in the stochastic fluctuations in the population, a probabilistic approach to the problem offers a certain conceptual convenience in building postulates related to the mechanisms involved in the rate processes,

as well as the relevant properties of the population. In this context, the fraction of the total number of individuals having a certain property may sometimes be interpreted as the probability that an individual has the property. Similarly, the rate of occurrence of a certain event may be looked upon as the rate of the corresponding random phenomenon. In some cases such phenomenological descriptions may not be possible in a deterministic framework. For example, the formulation of (1.1) is true only when the population is large enough to justify the use of the density function $x(z_1, \dots, z_n, t)$. It will be shown later in the text that this restriction can be removed if the quantity $x(z_1, \dots, z_n, t)$ is looked upon as the density of the first moment measure of the population, i.e., when $x(z_1, \dots, z_n, t) dz_1 \dots dz_n$ is taken to be the expected number of individuals at time t in the infinitesimal volume element $dz_1 \dots dz_n$ located at the point (z_1, \dots, z_n) in the space E .

The outline of this thesis is as follows: Chapter II is devoted to the mathematical background necessary for the development of population balance models in a general stochastic framework. In particular, a general probability space is defined for the population. Concepts of counting measures, moment distributions and generating functionals are developed, and the relation between the generating functionals and the moments is discussed.

The analogy between different approaches to the residence time distribution studies and modeling problems related to other populations (e.g., crystals, insects, trees, cells, etc.) is discussed in Chapter III in order to cast all population balance models in a

common framework. Following this discussion, general considerations for the selection of a proper type of model are presented.

A diffusion model for a population in which the individuals are described by one descriptor is derived in Chapter IV. The backward diffusion equation for the probability generating functional, as well as the backward and forward diffusion equations for the first two moments for a reproducing population with no external input are derived and the relevant boundary conditions are discussed. Solutions to these equations for some simple cases are presented in Chapter V.

A general model for a population must also include an input of entities from a source external to the population, i.e., an input other than reproduction. This is dealt with in Chapter VI. Some interesting results for a population with external input and no reproduction are proved and their application to problems of practical interest is discussed.

A controlled experiment performed to demonstrate the use of a result obtained in Chapter VI in testing the validity of some common assumptions behind the mathematical models developed in this thesis is described and discussed in Chapter VII. Chapter VIII is devoted to conclusions and recommendations for future research.



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CHAPTER II

MATHEMATICAL BACKGROUND

To cast the problem of modeling a population into a stochastic framework it is necessary to consider the dynamic changes in the population as a stochastic process. A stochastic process $\{X(t): t \in T\}$ is defined as a family of random variables $X(t)$ indexed by a parameter t varying in an index set T . In the present context, the index t refers to time and T usually refers to the set of nonnegative real numbers or the set of nonnegative integers. When all the entities in the population are identical, as is the case in the population models of type (i) in the previous chapter, $X(t)$ is a nonnegative integer at any given time, representing the total number of entities in the population. Similarly, for models of type (ii) where each entity in the population must belong to one of finitely many categories (such as age groups), $X(t)$ is a finite vector of nonnegative integers defining a measure on the set of age groups. In models of type (iii) each entity can occupy any point in the n -dimensional Euclidean space E , and an analogous description of $X(t)$ will be in terms of measures defined on the subsets of the Euclidean space E . At each point in time the configuration of individuals in the space E is a random distribution. The mathematical tool for the analysis of such processes is the theory of stochastic population processes developed by Moyal (1962).

An excellent review of the work done in this theory is provided by Daley and Vere Jones (1972).

2.1 Point Processes

For the most general description, the state of an entity in the population is described by a point ξ in an abstract space \tilde{E} , called the individual state space. In case of population balance models of type (iii) discussed in the previous chapter the individual state space is an n -dimensional Euclidean space. The population state space \mathcal{G} is the union of disjoint spaces \tilde{E}^k , $k = 0, 1, \dots, \infty$, where \tilde{E}^k is the k -fold cartesian product $\tilde{E} \times \tilde{E} \times \dots \times \tilde{E}$, for $k \geq 1$. By convention \tilde{E}^0 represents a single point, corresponding to zero population. A point $\xi^k \in \tilde{E}^k$ represents a population with k distinguishable individuals. It is a k -tuple (ξ_1, \dots, ξ_k) of points in \tilde{E} . On each \tilde{E}^k is defined the minimal σ -field B^k containing all product sets $e_1 \times e_2 \times \dots \times e_k$, where each set $e_i \in B$, the σ -field of Borel subsets of \tilde{E} . The minimal σ -field defined on the population state space \mathcal{G} containing $\sum_{k=1}^{\infty} B^k$ consists of all subsets e^* of \mathcal{G} whose intersections with each \tilde{E}^k are members of B^k . A measure space $(\mathcal{G}, \mathcal{B})$ can thus be defined from the pair (\tilde{E}, B) and its k -fold cartesian product (\tilde{E}^k, B^k) , $k = 1, \dots, \infty$. A probability measure \mathcal{P} can then be constructed on the measure space $(\mathcal{G}, \mathcal{B})$ either directly, or by constructing measures $P^{(k)}$ on (\tilde{E}^k, B^k) which satisfy $\sum_{k=0}^{\infty} P^{(k)}(\tilde{E}^k) = 1$.

$$p_k = P^{(k)}(\tilde{E}^k) \quad (2.1)$$

is the probability that the population consists of k entities.

The triplet $(\mathcal{S}, \mathcal{B}, \varphi)$ constitutes a model of a stochastic population process, and is called a point process.

Physical interpretations of \mathcal{S} , \mathcal{B} and φ are as follows: the population state space \mathcal{S} is the collection of all possible states of all the entities in the population, the σ -algebra \mathcal{B} refers to the collection of all possible "events" (e.g., all possible combinations of individuals in a biological population being in all possible (denumerable) age groups) and φ denotes the probability measure defined on the collection of all possible events in the population state space \mathcal{S} . Similar interpretations can be given to $\tilde{\mathcal{E}}$, \mathcal{B} , \mathcal{P} and $\tilde{\mathcal{E}}^k$, \mathcal{B}^k , $\mathcal{P}^{(k)}$ in the context of a single entity and a population of k entities respectively.

If the population is comprised of indistinguishable entities, the population state space \mathcal{S}_s consisting of unordered sets of points in $\tilde{\mathcal{E}}$ is a more appropriate basis for the description of the process. The subspace of \mathcal{S}_s of populations with k individuals is denoted by $\tilde{\mathcal{E}}^{(k)}$, and a point in $\tilde{\mathcal{E}}^{(k)}$, which is an unordered set $\xi^{(k)} = \{\xi_1, \xi_2, \dots, \xi_k\}$ represents a population with k indistinguishable entities. The σ -algebras $\mathcal{B}^{(k)}$ and \mathcal{B}_s are comprised of Borel subsets of $\tilde{\mathcal{E}}^{(k)}$ and \mathcal{S}_s respectively. The probability measures $\mathcal{P}_s^{(k)}$ and φ_s on the members of $\mathcal{B}^{(k)}$ and \mathcal{B}_s respectively can be defined in the same way as for the case of distinguishable entities. There is an obvious natural mapping of \mathcal{S} onto \mathcal{S}_s , and to each probability measure φ_s on \mathcal{S}_s there corresponds exactly one symmetric probability measure on \mathcal{S} .



2.2 Counting Measures

An alternative to the point process description of a population is the counting process. For any subset A of the individual state space \tilde{E} , let $N(A)$ be the number of entities in A . Formally, a counting measure N is any non-negative, integer valued measure defined on the class U of all Borel subsets of \tilde{E} . The elements of the space \mathcal{G} define certain counting measures, namely,

$$N(A|\xi^{(k)}) = \sum_{i=1}^k I(A|\xi_i) \quad (2.2)$$

where $I(A|\cdot)$ is the indicator function of the set A :

$$I(A|s) = \begin{cases} 1 & \text{if } x \in A \\ 0 & \text{if } x \notin A. \end{cases}$$

Note that all points in \mathcal{G} which correspond to the same unordered set $\xi^{(k)} = \{\xi_1, \dots, \xi_k\}$ in \mathcal{G}_s lead to the same counting measure (2.2). Moyal (1962) has proved that (2.2) defines a one-to-one correspondence between \mathcal{G}_s , and the space \mathcal{N} of all counting measures on the class U . Every probability space $(\mathcal{G}_s, \mathcal{B}_s, \theta_s)$ has an equivalent counterpart $(\mathcal{N}, \mathcal{B}_N, \theta_N)$.

It is useful to exhibit this last equivalence explicitly. If A_1, \dots, A_r are disjoint subsets of \tilde{E} whose union is \tilde{E} , the probability that there are k_1 entities in A_1 , k_2 entities in A_2, \dots , and k_r entities in A_r , is

$$\begin{aligned} P_N\{N(A_1) = k_1, \dots, N(A_r) = k_r\} &= P_s^{(L)}((A_1^{k_1} \times \dots \times A_r^{k_r})_s) \\ &= \frac{L!}{k_1! \dots k_r!} P_s^{(L)}(A_1^{k_1} \times \dots \times A_r^{k_r}) \quad (2.3) \end{aligned}$$

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where $L = k_1 + \dots + k_r$, $A_i^{k_i}$ is the k_i -fold cartesian product of A_i with itself, and in the first expression on the right hand side, $(C)_s$ is the symmetrization of the set C , i.e., the union of all possible distinct cartesian products $C_1 \times \dots \times C_L$, where k_1 of the C_j 's are equal to A_1 , k_2 are equal to A_2 , etc. There are $L!/k_1!k_2!\dots k_r!$ such possibilities, and thus the second expression on the right hand side is obtained.

2.3 Moment Measures and Moment Densities

The expected number of entities in a measurable set A is

$$M(A) = E(N(A)) = \sum_{L=1}^{\infty} \int_{E^L} \sum_{i=1}^L I(A|\xi_i) P_s^{(L)}(d\xi^{(L)}) = \sum_{L=1}^{\infty} L P_s^{(L)}(A \times \tilde{E}^{L-1}). \quad (2.4)$$

Moyal (1962) has shown that (2.4) can also be derived using (2.3).

Note that $M(A)$ is a non-negative, monotone nondecreasing and countably additive set function, and hence a measure on (\tilde{E}, B) .

$M(A)$ is therefore called the first moment measure or the first moment distribution. When there are only a finite number of entities in every bounded set A with probability one, it is often possible to define a density $m(\cdot)$ of the first moment measure in the case where \tilde{E} is a n -dimensional Euclidean space. Thus, the quantity

$$m(\xi) d\xi = M(d\xi), \quad (2.5)$$

denotes the expected number of entities in the infinitesimal volume element $d\xi = dz_1 dz_2 \dots dz_n$ located at $\xi = (z_1, z_2, \dots, z_n)$ in \tilde{E} .

Higher order moment measures are defined on (\tilde{E}^k, B^k) by taking expected values of the k -th product counting measure

$$N_k(A^{(k)} | \xi^{(L)}) = \sum_{i_1=1}^L \dots \sum_{i_k=1}^L I(A^{(k)} | (\xi_{i_1}, \dots, \xi_{i_k})), \quad k \leq n, \quad (2.6)$$

where $A^{(k)} = A \cap \tilde{E}^{(k)}$ for any $A \subset \tilde{E}$ and $\xi^{(L)} = \{\xi_1, \dots, \xi_L\}$.

The use of these higher-order moment measures is complicated by the fact that they contain "mass" concentrations on subsets of \tilde{E}^k . For this reason it is more convenient to deal with a factorial moment measure $M_{(k)}$ defined as the expected value of

$$\begin{aligned} N_{(k)}(A^{(k)} | \tilde{E}^L) &= \sum_{i_1 \neq i_2 \neq \dots \neq i_k} \sum I(A^{(k)} | \xi_{i_1}, \dots, \xi_{i_k}), \quad L \geq k, \\ &= N(A)(N(A)-1) \dots (N(A)-k). \end{aligned} \quad (2.7)$$

The two types of moment measures may be contrasted as follows: The sum on the right hand side of (2.6) is taken over all samples of size k from the population $\xi^{(L)}$. The sampling is done "with replacement", and (2.6) therefore contains L^k terms. In (2.7) the sum is over all samples "without replacement", and there are $\binom{L}{k} = \frac{L!}{(L-k)!}$ terms. The convenient relationship

$$M_{(k)}(A^{(k)}) = \sum_{j=0}^{\infty} \binom{k+j}{k} P_s^{(k+j)}(A^{(k)} \times \tilde{E}^j) \quad (2.8)$$

holds.

In the case of a n -dimensional Euclidean space, the k -th factorial moment measure often has a density on \tilde{E}^k when the probability of having more than one entity in an arbitrary interval $(\xi, \xi+d\xi)$ is $o(\delta\xi)$, i.e., this probability is of a smaller order of magnitude than $\delta\xi$. When such densities exist, the point process is said to be orderly.

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2.4 An Illustrative Example

At this point it will be worthwhile to consider an example to illustrate some of the concepts developed so far. Consider the case where the entities in the population are described by a single descriptor. To be specific, let the descriptor be the size of the entity, which may increase indefinitely. In this case \tilde{E} is the nonnegative real line and B is the σ -algebra generated by open, closed or half open intervals on the real line (representing ranges of sizes). The product space \tilde{E}^k (or $\tilde{E}^{(k)}$) and the corresponding σ -algebra B^k (or $B^{(k)}$) are necessary to define the joint probability distribution of k distinguishable (or indistinguishable) entities being in k overlapping and/or disjoint size categories. Similarly, the counting measure $N(A)$ denotes the number of entities in the size range A and the first moment measure $M(A)$ represents the expected number of entities in A . Note that for a fixed A , $N(A)$ is a random number. The product counting measure is necessary for defining product moments, as mentioned in the previous section.

To illustrate a product counting measure, consider a sample realization of the process as shown in Figure 2.1a. There are four entities in the population, of which three are in the set A , so that $k = 4$ and $N(A) = 3$. To get the second moment measure, it is necessary to evaluate the expectation $E(N(A)^2) = E(N(A) \times N(A))$. Note that the product counting measure on the set $A \times A$ is equal to $(N(A))^2$. Applying (2.6) to the realization of the point process in Figure 2.1a,

$$N_2(A \times A | \xi^4) = \sum_{i=1}^4 \sum_{j=1}^4 I(A \times A | (\xi_i, \xi_j)) = \sum_{i=1}^4 \sum_{j=1}^4 I(A | \xi_i) I(A | \xi_j) .$$

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The measure $N_2(A \times A | \xi^k)$ thus has a contribution of 1 from every pairwise combination (ξ_i, ξ_j) of the sizes of the entities in a sample realization of the process, when ξ_i and ξ_j both lie in the range A . The combination (ξ_i, ξ_j) refers to the joint occurrence of an entity at ξ_i and another at ξ_j . The sample realization of the process illustrated in Figure 2.1a thus corresponds to $N(A^2) = 9$. Different pairwise combinations (ξ_i, ξ_j) corresponding to this particular realization of the process are illustrated in Figure 2.1b. Note that corresponding to every point in Figure 2.1a there is always a point (ξ_i, ξ_i) in Figure 2.1b. In the particular case where the existence of one entity at ξ_1 is in no way dependent upon the existence of another at ξ_2 , if $f(\xi)d\xi$ denotes the probability that an entity lies in the size range $(\xi, \xi+d\xi)$ for an orderly point process, then the joint probability that an entity lies in $(\xi_1, \xi_1+d\xi_1)$ and another in $(\xi_2, \xi_2+d\xi_2)$ is $f(\xi_1)f(\xi_2)d\xi_1d\xi_2$. Thus, while the probability of having a point on the diagonal of $A \times A$ is $f(\xi)d\xi$, the probability of having a point (ξ_1, ξ_2) in $A \times A$ is only $f(\xi_1)f(\xi_2)d\xi_1d\xi_2 = o(d\xi)$, a quantity of a smaller order of magnitude than the corresponding probability for the diagonal. This results in the "mass" concentration of the second moment measure on a subset of $\tilde{E}^{(2)}$ -- the diagonal of any set $A \times A$. This mass concentration has a density $f(\xi)$ with respect to \tilde{E} , while the points of increase of the second moment measure corresponding to a pairwise occurrence of the entities in the size range A have a density with respect to $\tilde{E}^{(2)}$. The second moment measure can thus be expressed in terms of a sum of two component densities -- a density with respect to \tilde{E}

Figure 2.1

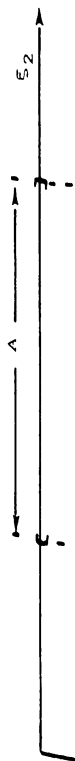


Figure 2.1

$k = \text{total number of entities} = 4$

$N(A) = \text{number of entities in } A = 3$

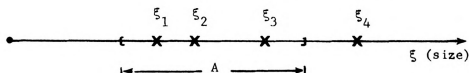


Figure 2.1a Counting measure for a sample realization of a point process

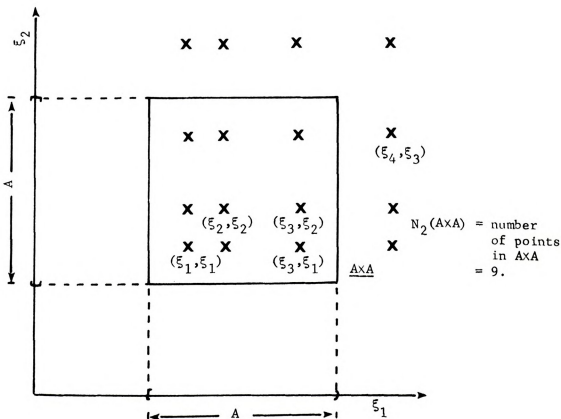


Figure 2.1b Product counting measure for the sample realization of the point process in Figure 2.1a

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representing the mass concentration on the diagonal and a density with respect to $\tilde{E}^{(2)}$. Similar mass concentrations also exist for the higher moments of the counting measure along the diagonals of the corresponding volume elements. For an orderly process the k -th order factorial moment measures ($k = 2, 3, \dots$) do not have such concentrations on subsets of $\tilde{E}^{(k)}$, and thus the use of densities of factorial moment distributions are much more convenient than that of product moment measures. A good discussion of the moment densities is given by Ramakrishnan (1950). When the point process is not orderly, i.e., when the probabilities of having twins, triplets, etc. in $(\xi, \xi + d\xi)$ is of the same order of magnitude as $d\xi$, even the factorial moment distributions have concentrations on the diagonal subspaces of \tilde{E}^k . In this case it is still possible to express the moment measures in terms of densities if each pair, triplet, etc. are treated as distinct populations with each pair, triplet, etc. considered as single entities. Ramakrishnan and Srinivasan (1958) have illustrated this for the case of age distributions in a population. However, even for the case of a population containing only singlets and twins, the mathematics involved was rather elaborate.

2.5 Generating Functionals

A complete description of a point process in the form of the probability distributions $P_s^{(k)}$ or the moment distributions $M_{(k)}$, ($k = 0, 1, \dots, \infty$) can be obtained in terms of a single generating functional -- the probability generating functional. The probability generating functional (PGF) is a natural generalization of

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the probability generating function from an integer-valued random variable to a random variable taking values in a population state space defined in Section 2.1. Analogous extension of the Laplace transform of the probability distribution of an integer-valued random variable (or the moment generating function), namely, the moment generating functional (MGF), and the characteristic functional are also used instead of a PGF.

Let θ be a bounded complex-valued function on \mathcal{S} . Define

$$w(\xi^k) = \theta(\xi_1) \dots \theta(\xi_k)$$

on \tilde{E}^k for every k , and therefore on \mathcal{S} . The PGF

$$G(\theta) = E(w(\xi^k)) = \sum_{k=0}^{\infty} \int_{\tilde{E}^k} \theta(\xi_1) \theta(\xi_2) \dots \theta(\xi_k) P^{(k)}(d\xi^k) \quad (2.9)$$

is defined for all measurable functions θ such that

$\sup_{\xi \in \tilde{E}} |\theta(\xi)| \leq 1$. In (2.9) the product $\theta(\xi_1) \dots \theta(\xi_k)$ is defined

to be equal to one when $k = 0$. The measures $P^{(k)}$ in (2.9) can

be replaced by $P_s^{(k)}$ without changing the result. Moyal (1962)

has proved that there is a one-to-one correspondence between

probability distributions P_s on $(\mathcal{S}_s, \mathcal{B}_s)$ and the PGF of the

form (2.9). If (A_1, \dots, A_r) is a finite measurable partition of

\tilde{E} , then taking

$$\theta(\xi) = \sum_{i=1}^r \zeta_i I(A_i | \xi), \quad (2.10)$$

the PGF reduces to a multivariate probability generating function

of the random variables $N(A_1), \dots, N(A_r)$

$$G(\theta) = E \left[\exp \sum_{i=1}^r \int_{A_i} \log \left\{ \sum_{j=1}^r \zeta_j \delta(A_j | \xi) \right\} N(d\xi) \right] = E \left[\prod_{i=1}^r \zeta_i^{N(A_i)} \right]. \quad (2.11)$$

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The probability distributions $P_s^{(k)}$ and the factorial moment distributions $M_{(k)}$ can be expressed as functional derivatives (variations) of G :

$$P_s^{(k)}(A_1 \times \dots \times A_k) = \frac{1}{k!} \delta_{\theta}^k G(0), \quad (2.12)$$

$$M_{(k)}(A_1 \times \dots \times A_k) = \lim_{\eta \rightarrow 1} \delta_{\theta}^k G(\eta), \quad (2.13)$$

when

$$\theta(\xi) = I(A_1|\xi)I(A_2|\xi)\dots I(A_k|\xi). \quad (2.14)$$

In (2.12) and (2.13) the k -th functional derivative $\delta_{\theta}^k G(\eta)$ is defined as (cf. Hille and Phillips, 1957, p. 111)

$$\delta_{\theta}^k G(\eta) = \frac{d^k [G(\eta + \zeta \theta)]}{d\zeta^k} \Big|_{\zeta=0} \quad (2.15)$$

The MGF (when $\theta(\xi)$ has a nonpositive real part) and the characteristic functional (for all $\theta(\xi)$) are obtained from (2.9) by replacing $\theta(\xi)$ by $\exp(\theta(\xi))$ and $\exp(i\theta(\xi))$ respectively.

Thus, for example, the MGF of a point process is given by

$$\phi(\theta) = E[\exp(\sum_{i=1}^k \theta(\xi_i))] = \sum_{k=0}^{\infty} \int_{\tilde{E}}^k \exp(\sum_{i=1}^k \theta(\xi_i)) P^{(k)}(d\xi^k). \quad (2.16)$$

The summation $\sum_{i=1}^k \theta(\xi_i)$ can also be written as an integral of $\theta(\xi)$ with respect to the counting measure $N(\cdot)$ on \tilde{E} , i.e.,

$$\sum_{i=1}^k \theta(\xi_i) = \int_{\tilde{E}}^k \theta(\xi) N(d\xi). \quad (2.17)$$

The k -th product moment measure M_k is related to the MGF by

$$M_k(A_1 \times \dots \times A_k) = (-1)^k \delta_{\theta}^k \phi(0), \quad (2.18)$$

where

$$\theta(\xi) = I(A_1|\xi)I(A_2|\xi)\dots I(A_k|\xi). \quad (2.19)$$

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6 Stochastic Population Processes

The discussion in this chapter thus far dealt with a probabilistic description of a population only in a static sense. The triplet $(\mathcal{S}, \mathcal{B}, \mathcal{P})$ thus refers to the population state space, the collection of all possible events and the probability measure only at a given point in time. For a complete stochastic description of the dynamic changes in a population one needs to know the joint probability distribution of all finite combinations $\{\xi_{t_1}, \xi_{t_2}, \dots, \xi_{t_i}, \dots, \xi_{t_k}\}$ of sample realizations ξ_{t_i} of the process at times t_i , where $\xi_{t_i} \in \mathcal{S}$. It can be seen that the problem of achieving this is formidable. Moreover, in practice, one is often interested in the first few (usually two) moments of the population at a given subset of the individual state space \tilde{E} , and thus a model describing these moments of a counting measure on \tilde{E} is usually sufficient for most practical purposes. If transition probabilities can be assigned to describe the changes in the state of an entity during a small time interval $(t, t+\delta t)$ (or during the interval $(t, t+1)$ when the time parameter is discrete), it is often possible to write down equations for the moment distributions (or moment densities) of the counting measures on \tilde{E} . The generating functionals prove to be extremely convenient tools in the derivation of such equations. The generating functionals are also useful in evaluating another quantity of interest, namely, the probability of having no entities in a given set $A \in \tilde{E}$ at any time t . For example, this can be obtained from the PGF by setting $\theta(\xi) = I(A|\xi)$.

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CHAPTER III

THE IDENTIFICATION PROBLEM

Before constructing a model for a population it is necessary to identify the relevant descriptors of an entity in the population and thus define the individual state space \tilde{E} . The analysis of change in the descriptors of an entity (e.g., maturity, productivity, etc.) with time, representing the movement of entities in \tilde{E} , is the next step in the modeling process. In a deterministic model, the movement of the entities in \tilde{E} is described by rates of change of the descriptors, whereas in a stochastic description this movement is characterized by transition probabilities when the process is Markovian. Whether the population can be characterized by a continuous movement of the entities in \tilde{E} or whether the population has to be grouped according to location, size, productivity, etc., depends on the particular system itself. The example of residence time distribution analysis of flow vessels may be cited to illustrate this. Equation (1.3) describes the movement of fluid elements inside the flow vessel by a continuous motion, whereas a mixing-cell model described by (1.2) essentially lumps the fluid elements in the vessel into a number of hypothetical "mixing cells". In some situations (1.2) is a more appropriate description of the system than (1.3). Analysis of data collected in tracer experiments as well as the physical structure of the flow vessel are often used to decide the type of model best suited for the purpose. The

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analogies between the models for residence time distributions and the models for other populations, which were mentioned in Chapter I, make it possible to cast the methods of analysis of both the classes of models in a common framework.

1. Residence Time Distribution of Flow Vessels

The residence time distribution (RTD) $C(t)$ of a flow vessel is defined as the age distribution frequency of the fluid elements leaving the vessel under steady flow conditions. Thus, the quantity $C(t)dt$ denotes the fraction of the fluid in the exit stream which has spent between t and $t+dt$ units of time inside the vessel. The RTD can also be interpreted as the probability that the age of a fluid element leaving the vessel is between t and $t+dt$, or the probability that a fluid element entering the vessel would take between t and $t+dt$ units of time to reach the outlet. The RTD $C(t)$ is thus synonymous with the residence time probability density. Sinclair and McNaughton (1965) have used this interpretation in a discussion on the calculation of the RTD of a series-parallel network of flow vessels, when the RTD of each component vessel is known. If two vessels with RTD's $C_1(t)$ and $C_2(t)$ are connected in series and if the residence times are independent of each other, the RTD $C_{12}(t)$ of the combination can be expressed as a convolution

$$C_{12}(t) = \int_0^t C_1(\tau)C_2(t-\tau)d\tau, \quad (3.1)$$

where t is the length of time spent by an element in the composite system and τ and $t-\tau$ are the times spent in each of the vessels. For

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perfectly mixed vessel the exit age of a fluid element has an exponential probability distribution, i.e., the exit age probability density (RTD) is $\frac{\exp(-t/\alpha)}{\alpha}$, where α is the mean residence time.

In the case of a plug flow vessel this probability density is a Dirac function, with a delay α corresponding to the residence time.

Imperfectly mixed vessels have flow patterns intermediate between the two ideal cases described above. The mixing inside such a vessel may be such that a small fraction of the fluid can reach the outlet much faster than the remaining bulk, resulting in bypassing; or a small portion of the fluid may be caught in stagnant zones in the vessel and reach the outlet after a very long time relative to the bulk of the fluid. The magnitude of bypassing and dead space can be quantified by analyzing the data collected in simple tracer experiments performed on the flow system or its scale model (Himmelblau and Bischoff, 1968, p. 71). In flow systems such as packed beds with size of the packing much smaller than the characteristic dimensions of the bed itself, the fluid elements must follow a long tortuous path before reaching the outlet, and the overall effect may be similar to that of a "diffusion" of the fluid elements relative to the mean velocity toward the outlet. Use of the diffusion equation such as (1.3) can be justified when the movement along a direction z is characterized by a Markov process and the incremental distance δz covered along z in a small time interval $(t, t+\delta t)$ can be looked upon as a random variable with

$$\begin{aligned} \text{mean} \quad E(\delta z) &= r(z, t) \delta t + o(\delta t), \\ \text{variance} \quad V(\delta z) &= \sigma^2(z, t) \delta t + o(\delta t), \\ \text{and} \quad E(\delta z^n) &= o(\delta t) \quad \text{for } n \geq 3. \end{aligned} \quad (3.2)$$

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many cases the mixing cell models give a more realistic representation of the RTD. In the mixing cell model, the flow vessel is assumed to be composed of a number of hypothetical perfectly mixed vessels connected in series and/or parallel. In the particular case of k perfectly mixed vessels connected in series the RTD of the combination is a convolution of k exponentials and when all the vessels have the same mean residence time α , the RTD of the combination is the density of the gamma distribution

$$C(t) = \frac{(t/\alpha)^{k-1} \exp(-t/\alpha)}{(k-1)! \alpha} . \quad (3.3)$$

Thus, in this case, the exit age of the fluid elements has a gamma distribution.

Life-Span Distributions in Biological Populations

Probability densities of life-spans of individuals in a wide variety of biological populations show a similarity with the RTD of imperfectly mixed vessels, and analogous models have been proposed and used for some biological populations. To account for the fact that the age at which cells divide is randomly distributed about a mean, Stuart and Merkle (1965) described the dynamics of cell populations by the diffusion equation

$$\frac{\sigma^2}{2} \frac{\partial^2 c(z,t)}{\partial z^2} - r \frac{\partial c(z,t)}{\partial z} = \frac{\partial c(z,t)}{\partial t} , \quad (3.4)$$

where $c(z,t)dz$ is the expected number of cells with physiological age between $(z, z+dz)$, the drift coefficient r denotes the mean rate of maturation and the diffusion coefficient $\frac{\sigma^2}{2}$ accounts for the variability in the rates of maturation among the individual

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cells. The cells were assumed to divide when they reached a certain fixed maturity. Takahashi (1968) divided the complete life cycle of a cell into a number of hypothetical growth stages, and assumed that the time spent by the cell in each of the stages was exponentially distributed with the same mean α . The number of stages and the mean residence time α for an individual stage could then be adjusted to match the output of the model with the corresponding observations in a radioactive tracer experiment. Takahashi's approach can be seen to be an exact analog of a mixing cell model or RTD analysis where vessels with equal mean residence time are connected in series, and thus this model results in a gamma distribution for the life-spans of the cells. Weiss (1968) proposed a general deterministic formulation for cell population dynamics, where a cell was characterized by its chronological as well as physiological age. This model is equivalent to the formulations of Stuart and Merkle as well as Takahashi for the first moment of the population.

As regards other biological populations, a gamma distribution for the maturation periods of insect life-stages was used by Read and Ashford (1968). In a study of productivity of perennial crops such as cocoa and palm, Abkin (1972) divided the life-spans of the trees into a number of stages according to their productive value. All the trees in a given stage were assumed to have the same productivity and the time taken by a tree to mature from one stage to the next was assumed to have a gamma distribution. In both these cases each life stage (of insects and trees) was implicitly taken to be equivalent to a number of identical sub-stages with exponentially

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distributed residence times. Fish populations are known to exhibit an elastic growth. The size distribution of the fish born during a single spawning season widens with time, due to differences in the growth rates among individuals (see e.g., Cooper and Latta, 1954; Cooper, Latta and Schafer, 1956). The magnitude of the spread occurring in a given length of time gets smaller with smaller amounts of food available to the fish (Hall, Cooper and Werner, 1970). No models have been proposed to account for the elastic growth. In general, it can be said that in any vertebrate population the age at which reproduction starts and the age at which an individual dies of old age are randomly distributed about some average values. As in the case of fish, size or biomass of an individual may be used to represent maturity in some cases, or it may be necessary to use an abstract maturity scale $[0,1]$, with the maturity of a newborn individual to be 0 and the end of the life cycle corresponding to the maturity 1 (Stuart and Merkle, 1965). It seems possible to describe these populations by a diffusion equation such as (3.4) with an additional term corresponding to death. In many biological populations the life-spans are randomly distributed due to the genetic variations among individuals. Another factor influencing the variability in growth rates is the availability of food, as mentioned above for the case of fish populations. Leftkovich and Currie (1963) studied the effect of food availability on the larvae of the cigarette beetle, Lasioderma serricornis(F.) and found that the mean as well as the standard deviation of the maturation period of the larval stage of the insect increase with a decreasing supply of food. It is likely that food availability has an effect on the life-span distributions of many other populations as well.

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3.3 A Dispersion Phenomenon in Crystallization

It has been observed that the product crystals from an industrial sugar crystallizer exhibit a much larger spread in the size distribution than that of the seed introduced initially, even when there is no appreciable nucleation during the growth of the seed crystals. Wright and White (1969) and White and Wright (1971) studied the growth rates of sucrose crystals under different conditions to specifically analyze the phenomenon of size dispersion. Starting with a uniform seed size in a batch crystallizer and carefully monitoring the process to see that no nucleation occurred, they measured the crystal size distribution (using characteristic length of a crystal as a basis) at different times. The crystal size distributions obtained under different temperatures, supersaturations and syrup purities were plotted on a normalized scale -- the ratio of the deviation from the mean with the standard deviation of the size distribution. Although the shapes of individual curves were slightly different from one another, the average of twenty six such curves was very close to the density of normal distribution with zero mean and unit variance. The difference in the shapes of the individual size distribution curves was at least partly due to the fact that sieve analysis was used to determine the size distribution for many samples. Since the shapes of individual crystals were non-uniform, conversion of weight-size relation to a number-size relationship resulted in an inaccurate determination of the crystal size distribution. Let the growth of crystals be characterized by (3.4) with $c(z,t)dz$ denoting the expected number of crystals between sizes z and $z+dz$, the drift coefficient r

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representing the (ensemble) mean growth rate and the diffusion coefficient $\frac{\sigma^2}{2}$ being equal to half the value of the variance of individual growth rates. The expected number of crystals crossing a point z per unit time can be shown to be

$$c^*(z,t) = rc(z,t) - \frac{1}{2}\sigma^2 \frac{\partial c(z,t)}{\partial t}, \quad (3.5)$$

(cf. Kimura, 1964, p. 185). Since there is no nucleation, this "flux" across the point $z = 0$ must be zero. Thus, the boundary condition at $z = 0$ is

$$c^*(z,t) = 0. \quad (3.6)$$

Because of the size dispersion phenomenon, the function $c(z,t)$ must be finite for a sufficiently large z at all times. Thus, the second boundary condition is

$$c(\infty,t) = \text{finite}. \quad (3.7)$$

The solution of (3.4), (3.6) and (3.7) with all seed crystals of size z_0 at time zero is approximately proportional to the density of normal distribution for a sufficiently large z_0 (cf. Cox and Miller, 1965, p. 224). Thus, the diffusion model does seem to agree with the experimental observations.

White and Wright postulated that the size dispersion occurs due to an uneven distribution of lattice dislocations on the surfaces of different crystals. These imperfections have a marked effect on growth rates of crystals, thus resulting in a random distribution of individual growth rates.

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3.4 Generalizations

In the context of a general population balance model where each entity is represented as a point in the individual state space \tilde{E} , the random variations among individual rates of movement of the entities along any given direction in \tilde{E} may occur due to three distinct reasons:

- (i) The entities may be present in a uniform homogeneous environment, with the entities themselves having a random distribution in their potential capability to move along any given direction. The random life-span distribution in biological populations occurring due to genetic variability falls in this category.
- (ii) The entities may be present in a uniform homogeneous environment and the process of movement along any direction itself may be random in nature. The phenomenon of size dispersion in sucrose crystals is perhaps of this type, with the number of dislocations on the surfaces of each crystal itself varying randomly as the crystal grows.
- (iii) The entities and the process of movement along any direction (under uniform conditions) may be uniform, but the environment is heterogeneous and the entities spend random lengths of time under different environmental conditions due to the heterogeneity of the environment. This will be the case when there are regions of different supersaturation in a crystallizer and crystals spend random lengths of time in these

regions, giving rise to a size dispersion. A similar situation exists in the RTD analysis of imperfectly mixed vessels -- for example, the RTD of a packed bed differs from that of the ideal cases of perfect mixing and plug flow because the fluid elements spend random amounts of time in the crevices between the packing elements, resulting in little movement for random periods of time.

viously, different combinations of these cases will occur in some situations.

Depending upon the inherent variability among individualities, their rates of movement in \tilde{E} , or the coarseness of the heterogeneity of the environment, situations analogous to dead space bypassing may occur in any general population. In a study of different animal populations in fresh water ponds, Hall, Cooper anderner (1970) introduced approximately a cohort of fish in each of ponds, and the size distributions of the fish were measured at end of the experiment. The size distributions show a long tail with a slight peak in the tail, suggesting the possibility that the fish which get a head-start in their growth continue to grow at a faster rate than the others. This is analogous to the phenomenon of bypassing in a nonideal flow vessel, where a small portion of fluid moves much faster relative to the bulk. Similarly, in a study of economic mobility or acceptance of new ideas in social systems one may encounter a situation where a small portion of the population, possibly representing certain ethnic groups, progress accept new ideas much faster (or much slower) than bulk of the

ulation, representing a situation analogous to bypassing (or
d space) in imperfectly mixed flow vessels.

The nonideal flow vessels and analogous systems involving
er populations may be denoted by a common term -- the nonideal
ulation-flow systems. In the general case, the word "flow"
ers to the movement of entities in the population along any given
ection in \tilde{E} , with the 'directions' specified by the relevant
criptors of an entity, i.e., the co-ordinate axes of \tilde{E} . It
ould be noted that in the general context, a perfectly mixed
essel is analogous to a situation where all the entities contained
a region in \tilde{E} are assumed to be identical in all respects
the time taken by an entity to leave this region has an
ponential distribution. This description thus involves lumping
"averaging" the individuals in the appropriate region of \tilde{E} .
ure 3.1 illustrates this lumping for a biological population
h the relevant domain in \tilde{E} consisting of a two-dimensional
graphical area and a range of maturities corresponding to the
e-span of an individual.

Selection of a Proper Model

The movement of the entities along the i -th co-ordinate axis
of \tilde{E} can be characterized as a diffusion in that direction if
movement of an entity along z_i is Markovian in nature and the
ance δz_i covered in a short time interval δt satisfies (3.2).
implies that the rates of movement of the entitites along z_i
a short period of time should have approximately a normal dis-
ution. It is necessary to study the nature of the individual

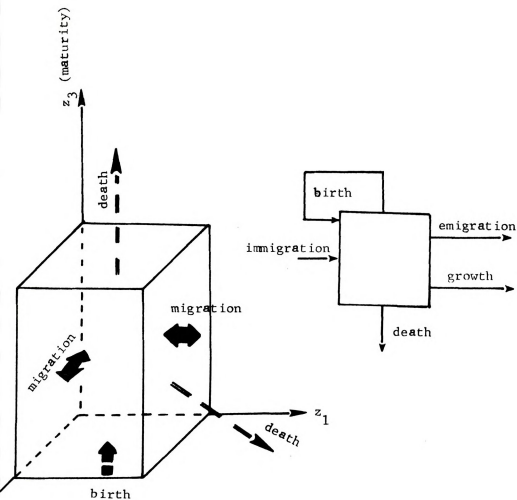


Figure 3.1 An example of lumping of variables in the individual state space for a biological population

variability among the entities, the process of movement along z_i ,
 the heterogeneity of the environment to see whether this con-
 dition is satisfied: When the inherent variability among the
 entities alone is the cause of the randomness in the movement, i.e.,
 when each entity moves at essentially a constant rate at all times
 with a random distribution of such rates among different entities,
 the overall effect of movement of all the entities over a short
 period of time may be analogous to a diffusion if the number of
 entities as well as their individual growth rates are such that
 the assumption of a normal distribution of movement of entities over
 a short time interval can be justified. In the case where the
 process of growth is random in nature, it would be necessary that
 the movement of an individual during a short interval of time be
 characterized by a number of independent random increments, or as
 a continuous stochastic process. In the situation where the
 heterogeneity of the environment is the cause of random movement
 of the entities, it may be necessary to put a constraint on the
 randomness of the heterogeneity for the diffusion model to be
 justified -- for example, while the diffusion equation (1.3) can
 satisfactorily describe the RTD in a packed bed where the char-
 acteristic dimension of the packing is much smaller than that
 of the bed, it is a poor model for short beds with large packing
 particles or for fluidized beds. It should be noted that when the
 diffusion coefficients $\frac{D}{2}$ for the n-dimensional diffusion process
 are zero, the equation for the mean reduces to the case commonly
 used in the population balance models in chemical engineering
 practice, namely, (1.1), where all the entities at a given point

\tilde{E} move along any given direction at essentially the same
e.

When the conditions for a diffusion equation are not
satisfied, it may be necessary to lump the entities in different
regions in \tilde{E} and specify transition probabilities (or rates) for
the movement of entities from one region to another, as mentioned
in the previous section. The mixing cell models used in the RTD
analysis and the model proposed by Takahashi (1968) for cell popula-
tion dynamics are examples of this approach.

Experiments similar to tracer analysis in RTD studies, along
with a careful analysis of the potential causes of random movement
of the entities in the individual state space would usually be
necessary to decide the type of model best suited for the purpose.
For example, the analysis of sucrose crystallization study of
Fitz and Wright (1969, 1971) clearly suggests the applicability of
the diffusion equation, whereas the growth rates of fish may need
a more careful study to decide whether the "bypassing" effect in
the growth of fish populations is of significance.

A stochastic description of a population characterized by
a diffusion process would rely on the theory of stochastic popula-
tion processes. The discrete cases such as the perennial crop
productivity model of Abkin (1972) will involve the use of discrete
continuous parameter Markov chains when the transition proba-
bilities for the movement of the entities from one state (char-
acterized by the productivity, maturity, etc. of the entities in a
state) to another does not depend on the past history of the
entities. For the purpose of the models considered in the next two

chapters it will be assumed that the movement of the entities in
can be described by a diffusion process. A stochastic descrip-
tion of the diffusion of the entities includes the stochastic
counterpart of the process described by (1.1), i.e., the case where
all entities at a point in \tilde{E} move in any given direction at the
same rate, as a special case.

CHAPTER IV

DERIVATION OF THE DIFFUSION EQUATIONS

As a specific example of a general population balance model, the case where the individual state space \tilde{E} consists of the non-negative real line will be considered in this chapter. This case applies to populations in which the entities are characterized by a single descriptor z , denoting a property of an entity such as location, age, maturity or size. The model can be readily extended to the case where \tilde{E} is a n -dimensional Euclidean space for which the properties of an entity representing the co-ordinate axes are independent of each other. The total number of entities in a population can increase due to reproduction by the existing entities in the population, or by virtue of an external input. In the derivation of the equations in this chapter it will be assumed that there is no external input of entities to the population. Moreover, all the "newborn" entities are assumed to appear at the origin of \tilde{E} at the instant of their birth. The descriptions in this chapter will have to be modified slightly in case of populations where the "offspring" is in the same state as the parent, as in the spatial migration of a biological population via a diffusion process (cf. Alderson and Moyal, 1963). The effect of external input will be considered in Chapter VI. In the following discussion the creation of an entity by reproduction will be called birth of the entity and

oval of an entity will be termed as death of the entity.

Description of the Process

At time s the population is started with a single entity at an arbitrary point $y \in \tilde{E}$. The process of movement of an entity from any point $z \in \tilde{E}$ is assumed to be a continuous Markov process, i.e., as the time interval δt becomes small, so also does the distance δz covered by the entity during that interval, and the future locations of the entity depend only on its present location, not the past history. The distance δz covered by an entity in a small time interval $(t, t+\delta t)$ conditional on its survival during this interval is a random variable with a density $g(\delta z; z, t)$ and satisfies the conditions specified by (3.2), i.e.,

$$\begin{aligned} \text{(mean)} \quad E(\delta z) &= r(z, t)\delta t + o(\delta t), \\ \text{(variance)} \quad \text{var}(\delta z) &= \sigma^2(z, t)\delta t + o(\delta t), \\ E(\delta z^k) &= o(\delta t) \quad \text{for } k \geq 3. \end{aligned} \quad (4.1)$$

During the interval $(t, t+\delta t)$ the entity at z has a probability $\lambda_i(z, t)\delta t + o(\delta t)$ of dying and a probability $\lambda_i(z, t)\delta t + o(\delta t)$, $i = 1, \dots, \infty$, of producing i new entities at $z = 0$ with

$$\sum_{i=1}^{\infty} \lambda_i(z, t) \leq \lambda < \infty. \quad (4.2)$$

is,

probability that an entity will not reproduce during $(t, t+\delta t)\} =$

$$1 - \sum_{i=1}^{\infty} \lambda_i(z, t)\delta t + o(\delta t). \quad (4.3)$$

State z of an entity is unaffected by giving birth to an "offspring".

the process of change of position of an entity as well as its reproduction and death proceed independently of other entities in the population. In view of the independence of the entities, the case where the initial population contains k entities ($k > 1$) can be readily modeled by a superposition of k populations, each starting with a single ancestor in the appropriate location corresponding to the initial distribution.

Two obvious cases where this description applies directly are the dynamics of biological populations of individuals characterized by maturity, and crystallization processes involving secondary nucleation. In biological populations the probability of birth refers to reproduction and the probability of death refers to the processes of natural death, predation, parasitism and harvesting. In crystallization processes, reproduction refers to secondary nucleation and death corresponds to precipitation, or loss of the crystals from the crystallizer due to mechanical removal or bulk flow of the magma. Most particulate processes in chemical engineering contain entities which do not reproduce. It can be seen that these processes correspond to the case where the birth rates $\lambda_i(z, t) = 0$, $i = 1, \dots, \infty$. In processes such as degradation of polymer molecules, the reduction of solids, etc., every breakage of an entity normally produces a number of entities of different sizes, thus generating a random size distribution of the pieces. The models derived in this chapter can possibly be generalized to include this case, but the analysis of this general case will be much more complicated.

For the quantitative description of the population, a counting measure $N(A, t|y, s)$, $t \geq s$, is defined as the number of entities

the set $A \subset \tilde{E}$ at time t , given that there was one ancestor at the point y at time s . Obviously,

$$N(A, s|y, s) = I(A|y) \quad (4.4)$$

a probability one. In (4.4), $I(A|\cdot)$ is the indicator function of the set A . Starting with (4.4) the description of the temporal changes in the population can proceed in two directions. One way of describing the population dynamics is to express the counting measure for a fixed arbitrary set A and a fixed time t , as a function of s and y . Since $-\infty \leq s \leq t$, and since the dynamics of the population is expressed relative to the condition specified in (4.4), i.e., with $s = t$ as the reference point in time, it can be seen that this is a description of the population obtained by moving "backward" in time. This description is therefore called a backward description. Another way of describing the changes in the population is to express the counting measure corresponding to an infinitesimal set dz located at $z \in \tilde{E}$ for fixed arbitrary t and s as a function of z and t . Since this description involves moving "forward" in time with respect to the reference point $t = s$, this is called the forward description. Accordingly, (4.4) is a forward description of the population because it describes the population in an infinitesimal interval $(z, z+dz)$ as a function of z and t .

Boundary Conditions

The movement of the entities is restricted to a domain in \tilde{E} called the spatial domain \tilde{D} . The spatial domain may be a subset

\tilde{E} , or the entire space \tilde{E} itself. In the present context, is taken as the closed interval $[0, \tilde{a}]$, where \tilde{a} may possibly approach infinity.

In the derivation of the model, the behavior of entities at boundaries must be considered. The possible modes of boundary behavior of an entity diffusing on the real line were considered Feller (1954). Some possibilities of practical interest are: reaching a boundary, the entity may be instantaneously returned to the interior of \tilde{D} , i.e., the entity is always restricted to stay in the interior of \tilde{D} , resulting in a reflecting barrier. Another possibility is when the entity is removed from the population as soon as it reaches a boundary, representing an absorbing barrier. A third possible type of behavior at the boundary occurs when an entity reaching a boundary stays there for an exponentially distributed random length of time, and then jumps to a point in the interior of \tilde{D} , or to the other boundary. It may be added here that if \tilde{E} is a multidimensional Euclidean space, there is a fourth possibility, namely, the case where an entity reaching a boundary moves along the boundary according to a Markov transition probability.

In the context of biological populations and crystallization processes described above, the boundary at $z = 0$ is a reflecting barrier, because no entity can cross this boundary and moreover, the moment an entity appears at $z = 0$, it starts to "grow", i.e., it moves to the interior of \tilde{D} . When \tilde{D} is a finite interval $[0, \tilde{a}]$, one may encounter a reflecting or an absorbing barrier at \tilde{a} . The third type of boundary behavior, where an entity can take

finite jump into the interior of \tilde{D} may be encountered in rare cases. It is intuitively obvious that as \tilde{a} approaches infinity, the state of the population in any finite interval at any finite time would not be dependent upon the type of boundary behavior at \tilde{a} .

Fish are known to exhibit stunted growth if the population densities are too high (Beckman, 1940). This may be interpreted to mean that there is an upper limit on the size that an individual fish may attain depending upon the availability of food. This limiting size can possibly be viewed as a reflecting barrier restricting further growth of the fish. However, it should be noted that since the stunted growth phenomenon occurs due to a severe interaction among individuals for the available food, the stochastic formulation of the population process may perhaps be only a crude approximation in this case. In an intensively exploited fishery (or a forest resource) few individuals above a certain size (or productive value) may survive, enabling one to imagine an absorbing barrier at a finite size (or productive value). Similar considerations will apply to any intensively exploited biological population. In some biological populations all the individuals above a certain maturity can be considered identical in all respects. In such cases it is possible to consider this maturity to be an absorbing barrier. We assume that an individual reaching this maturity is removed from the original population to become a member of another population consisting of identical individuals. An absorbing barrier can be identified in case of particulate processes in chemical engineering practice when the particles precipitate out of the system, or are removed mechanically by some separation technique when they reach

certain size, weight, etc. The flow of fluids through packed beds offers an interesting possibility of an application of the third type of boundary behavior. The movement of fluid elements near the ends of the bed cannot be characterized by a diffusion equation even in long beds, because the random velocity patterns of the eddies necessary for an effective diffusion are not fully developed in the boundary zones. The distance covered by an eddy at a boundary (i.e., at the inlet or exit end of the bed) in a short time interval will not have a normal distribution as indicated by (4.1) -- instead, the distribution may perhaps look more like an exponential. To account for the end effects it may be worthwhile to consider an idealized flow behavior where a fluid element at a boundary is capable of instantaneously jumping to a point in the interior according to a known probability distribution. When the volume of the headers of a packed vessel is rather large, it may be possible to assume the headers as perfectly mixed (hypothetical) vessels at the boundaries, giving rise to finite exponentially distributed times of stay at the boundaries. A detailed mathematical description and verification of this boundary behavior is beyond the scope of this dissertation.

The Backward Diffusion Equations

A stochastic model for the process described in Section 4.1 may be derived here in the form of a backward diffusion equation for the probability generating functional of the population. Let $n_t(y, s)$, $t \geq s$, denote the number of entities in $[0, z]$ at time t , given that there was one entity at y at time s . Define

$$G(\theta, t|y, s) = E\{\prod_i \theta(z_i), t|y, s\} = E\{\exp \int_{\tilde{D}} [\log \theta(z)] N(dz, t|y, s)\} \quad (4.5)$$

as the conditional probability generating functional for the population at time t , given that there was one entity in the population at y at time s . The product on the right side of (4.5) has one term $\theta(z_i)$ corresponding to each entity at z_i at time t . In view of the independence of the entities in the population a GF conditional on k entities at y_1, \dots, y_k at time s can be shown to be equal to

$$G(\theta, t|y_1, \dots, y_k; s) = \prod_{j=1}^k G(\theta, t|y_j, s). \quad (4.6)$$

The proof of (4.6) follows from the definition of the conditional counting measure with k ancestors

$$N(A, t|y_1, \dots, y_k; s) = \sum_{j=1}^k N(A, t|y_j, s) \quad (4.7)$$

and the second expression on the right side of (4.5) (Moyal, 1964). When there are no entities in the population at time t , $G(\theta, t|y, s)$ equals one (cf. Equation (2.9)). As mentioned before, the spatial domain \tilde{D} is taken as the interval $[0, \tilde{a}]$. The case where \tilde{D} is the semi-infinite interval $[0, \infty)$ is obtained by letting \tilde{a} approach infinity.

Consider the change in $G(\theta, t|y, s-\delta s)$ brought about by a small change in the ancestor during the time interval $(s-\delta s, s)$. In view of the fact that the ancestor can move from y to $y+\delta y$ during the time interval $(s-\delta s, s)$ with probability $g(\delta y; y, s)d(\delta y)$ provided it survived during the interval,

$$\begin{aligned}
 & , t | y, s - \delta s) = \\
 & - \sum_{i=1}^{\infty} \lambda_i(y, s) \delta s (1 - \mu(y, s) \delta s) \int_0^{\bar{a}} g(\delta y; y, s) G(\theta, t | y + \delta y, s) d(\delta y) \\
 & 1 - \mu(y, s) \delta s \sum_{i=1}^{\infty} \lambda_i(y, s) \delta s G(\theta, t | 0, s) \int_0^{\bar{a}} g(\delta y; y, s) G(\theta, t | y + \delta y, s) d(\delta y) \\
 & (y, s) \delta s \cdot 1, \quad (4.8)
 \end{aligned}$$

$$\begin{aligned}
 & \text{at } t | \text{ancestor} \\
 & y \text{ at time } s - \delta s \Big] = \left[\begin{array}{l} P[\text{ancestor neither died nor reproduced} \\ \text{during } (s - \delta s, s)] \{ P[\text{ancestor moved to} \\ y + \delta y \text{ during } (s - \delta s, s)] [PGF \text{ at } t | \text{ancestor} \\ \text{at } y + \delta y \text{ at time } s] \} \end{array} \right] \\
 & \left[\begin{array}{l} P[\text{ancestor did not die during } (s - \delta s, s)] P[\text{ancestor produced} \\ i \text{ offspring during } (s - \delta s, s)] [PGF \text{ at } t \text{ with the } i \text{ offspring} \\ \text{as ancestors at } 0 \text{ at time } s] \{ P[\text{ancestor moved to } y + \delta y \\ \text{during } (s - \delta s, s)] \{ PGF \text{ at } t | \text{ancestor at } y + \delta y \text{ at time } s \} \} \end{array} \right] \\
 & P[\text{the ancestor died during } (s - \delta s, s) \text{ without producing any} \\
 & \text{offspring}] [PGF \text{ for no entities in the population}].
 \end{aligned}$$

probability of the ancestor dying and also producing any off-
 spring in an interval δs is $o(\delta s)^{\#}$, and hence the corresponding
 is omitted from (4.8). Expansion of $G(\theta, t | y + \delta y, s)$ in a
 or series about y gives

$$\begin{aligned}
 t | y + \delta y, s) &= G(\theta, t | y, s) + \frac{\partial G(\theta, t | y, s)}{\partial y} \delta y + \frac{\partial^2 G(\theta, t | y, s)}{\partial y^2} \frac{\delta y^2}{2} \\
 &+ O(\delta y^3)^{\#}. \quad (4.9)
 \end{aligned}$$

the function $f(t) = 0(t)$ if $|f(t)/t| \leq k < \infty$ as $t \rightarrow 0$ (e.g.,
 quantity $E(\delta z)$ in (4.1)). The function $f(t) = o(t)$ if
 $t \rightarrow 0$ as $t \rightarrow 0$, i.e., $f(t)$ is of a smaller order of magnitude
 t , as in case of $E(\delta z^k)$, $k \geq 3$ in (4.1).

stituting (4.9) in (4.8), using (4.1) and omitting terms of order higher than δs , one obtains

$$\begin{aligned} G(\theta, t|y, s-\delta s) &= G(\theta, t|y, s) + \delta s [\mathcal{A}_y G(\theta, t|y, s) + \mu(y, s) \\ &\quad - \sum_{i=1}^{\infty} \lambda_i(y, s) G(\theta, t|y, s) \\ &\quad + \sum_{i=1}^{\infty} \lambda_i(y, s) G(\theta, t|0, s)^i G(\theta, t|y, s)], \end{aligned} \quad (4.10)$$

where the operator \mathcal{A}_y is defined as

$$\mathcal{A}_y = \left[\frac{\sigma^2(y, s)}{2} \frac{\partial^2}{\partial y^2} + r(y, s) \frac{\partial}{\partial y} - \mu(y, s) \right]. \quad (4.11)$$

Subtracting $G(\theta, t|y, s)$ from (4.10), dividing by δs and letting δs approach zero, the backward diffusion equation for the PGF

$$\begin{aligned} \frac{G(\theta, t|y, s)}{\delta s} &= \mathcal{A}_y G(\theta, t|y, s) + \mu(y, s) - \sum_{i=1}^{\infty} \lambda_i(y, s) G(\theta, t|y, s) \\ &\quad + \sum_{i=1}^{\infty} \lambda_i(y, s) G(\theta, t|0, s)^i G(\theta, t|y, s) \end{aligned} \quad (4.12)$$

is obtained.

It may appear that if $E(\delta z^k) = O(\delta t)$ (instead of $o(\delta t)$) (4.1) for say, $k = 3$, then one would be able to get a third order partial differential equation instead of (4.12) by truncating the Taylor series expansion of $G(\theta, t|y+\delta y, s)$ in (4.9) after the third derivative term. However, Pawula (1967) has shown that if $E(\delta z^3) = O(\delta t)$, then $E(\delta z^k) = O(\delta t)$ for all $k > 3$, thus demanding an infinite Taylor series expansion in (4.9) and consequently an infinite-order partial differential equation.

When $y = \tilde{a}$ is an absorbing barrier, in view of the condition that every entity at the absorbing barrier is immediately removed

from the population, i.e., the number of entities at the barrier is always zero with probability one, the corresponding boundary condition is

$$G(\theta, t | \tilde{a}, s) = 1. \quad (4.13)$$

As explained in the previous section, the boundary at $y = 0$ is a reflecting barrier. For a reflecting barrier at $y = 0$ the entity at the boundary is assumed to jump instantaneously to an interior point δy . Thus,

$$G(\theta, t | 0, s) = G(\theta, t | \delta y, s).$$

Hence,

$$G(\theta, t | 0, s) = G(\theta, t | 0, s) + \delta y \frac{\partial G(\theta, t | y, s)}{\partial y} \Big|_{y=0}$$

i.e.,

$$\frac{\partial G(\theta, t | y, s)}{\partial y} \Big|_{y=0} = 0. \quad (4.14)$$

Similarly, if $y = \tilde{a}$ is also a reflecting barrier,

$$\frac{\partial G(\theta, t | y, s)}{\partial y} \Big|_{y=\tilde{a}} = 0. \quad (4.15)$$

Hence there is only one entity at y at time $t = s$,

$\{\Pi\theta(z_i), t | y, t\} = \theta(y)$. Hence, the initial condition is given by

$$G(\theta, t | y, t) = \theta(y). \quad (4.16)$$

If (4.12) can be solved analytically with the initial condition (4.16) and boundary conditions (4.14) and (4.13) or (4.15), the solution $G(\theta, t | y, s)$ would contain a complete stochastic description of the population with one ancestor when there is no

put of entities to the population from an external source. However, because of the nonlinear nature of (4.12) and the occurrence of the arbitrary function $\theta(y)$ in the initial condition, such an analytical treatment is not possible.

Remark 4.1: If there are k entities initially at y_1, y_2, \dots, y_k ; then the complete PGF for the population is obtained by solving the backward diffusion equation for $G(\theta, t | y_j, s)$ for $j = 1, \dots, k$ and substituting these in (4.6).

Remark 4.2: In view of the definition (2.9) for a PGF, the probability of having zero entities in the population (i.e., the extinction probability) is given by the solution of (4.12) with $\theta(y)=0$. Similarly, choosing

$$\theta(y) = \begin{cases} \zeta & \text{if } y \in A \\ 1 & \text{if } y \notin A \end{cases} \quad (4.17)$$

$G(\theta, t | y, s)$ and solving (4.12) one obtains the probability generating function

$$\text{pgf} = E\{\zeta^{N(A, t | y, s)}\} = \sum_{k=0}^{\infty} p_k \zeta^k \quad (4.18)$$

the number of entities in the set $A \subset \tilde{D}$ at time t , when there is one ancestor initially at y . In (4.18) the quantity p_k refers to the probability of having k entities in A at time t , given one ancestor at y at time s .

Remark 4.3: The diffusion equation can also be derived by considering a random walk of the entities in \tilde{E} , where each entity is assumed to jump a random (or fixed) distance forward or backward, or stay in the same position during a small time interval δs with



own probabilities, and taking a limit as the jump size reduces to zero, if the jump probabilities satisfy (4.1) (cf. Cox and Miller, 1955).

It should be noted that a similar set of equations can be derived for the moment generating functional or the characteristic functional as well.

By comparing the population process with a linear birth and death process having birth rate $= \lambda$ and death rate equal to the minimum value taken by $\mu(y,s)$, it can be seen that the first and second factorial moments of populations in any subset of \tilde{E} exist. The backward equations for the first and second factorial moments are readily obtained by using (2.13) and (2.14). Thus, the first moment of the number of entities in the set $A_1 \subset \tilde{D}$ is given by

$$\begin{aligned} \frac{\partial E[\Pi [\eta(z_i) + \zeta \theta(z_i)], t | y, s]}{\partial \zeta} \bigg|_{\zeta=0}^{\eta \rightarrow 1} &= \frac{\partial E[\Pi [\eta(z_i) + \zeta \theta(z_i)], t | y, s]}{\partial \zeta} \bigg|_{\zeta=0}^{\eta \rightarrow 1} \\ &= E \sum_i \theta(z_i) \{ \Pi_{j \neq i} [\eta(z_j) + \zeta \theta(z_j)], t | y, s \} \bigg|_{\zeta=0}^{\eta \rightarrow 1} \\ &= E[N(A_1, t | y, s)] \equiv M_1(A_1, t | y, s) \end{aligned} \quad (4.19)$$

choosing $\theta(z) = I(A_1 | z)$. Similarly, the second factorial moment corresponding to the product counting measure for the set $A_2 \subset \tilde{D} \times \tilde{D}$ (cf. Figure 2.1b), is given by

$$\begin{aligned}
\frac{\partial \{ \prod_{i=1}^k \theta(z_i) \}}{\partial \zeta^2} \Big|_{\substack{\eta \rightarrow 1 \\ \zeta=0}} &= E \left\{ \left[\sum_i \theta(z_i) \right] \frac{\partial \left\{ \prod_{j \neq i} [\eta(z_j) + \zeta \theta(z_j)] \right\}}{\partial \zeta} \Big|_{\substack{\eta \rightarrow 1 \\ \zeta=0}} \right\} \\
&= E \left[\sum_i \theta(z_i) \sum_{j \neq i} \theta(z_j) \right] \\
&= E [N_2((A_1 \times A_2), t | y, s) - N((A_1 \cap A_2), t | y, s)] \\
&= M_{(2)}(A_1 \times A_2, t | y, s), \tag{4.20}
\end{aligned}$$

where $\theta(z) = I(A_1 | z) I(A_2 | z)$, or equivalently, when $\theta(z_i) = I(A_1 | z_i)$ and $\theta(z_j) = I(A_2 | z_j)$. The indices i and j in (4.19) (4.20) refer to all the entities in the population at time t . It can be seen that the quantity $N(A_1 \cap A_2), t | y, s$ is subtracted from $N((A_1 \times A_2), t | y, s)$ to deduct the contribution of the entities corresponding to the terms $\theta(z_i) \theta(z_j)$. When $A_1 = A_2$, (4.20) gives the second factorial moment of the population in the set A_1 , i.e., $E(N(A_1)(N(A_1) - 1))$.

Remark.4.4: It should be noted that in the derivation of (4.19) and (4.20) it was necessary to commute the partial differential operator $\frac{\partial}{\partial \zeta}$ with the expectation operator. This is justified by virtue of the following proof:

$$\begin{aligned}
E \left\{ \prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i)), t | y, s \right\} &= \frac{\partial}{\partial \zeta} \left\{ \sum_{k=0}^{\infty} \int_{\tilde{E}} \prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i)), t | y, s \right\} P^{(k)}(dz^k) \Big\} \\
&= \lim_{\zeta \rightarrow 0} \frac{1}{\zeta} \left\{ \sum_{k=0}^{\infty} \int_{\tilde{E}} \prod_{i=1}^k (\eta(z_i) + (\zeta + \Delta) \theta(z_i)), t | y, s \right\} P^{(k)}(dz^k) \\
&\quad - \sum_{k=0}^{\infty} \int_{\tilde{E}} \prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i)), t | y, s \Big\} P^{(k)}(dz^k) \Big\} \\
&= \lim_{\zeta \rightarrow 0} \frac{1}{\zeta} \left\{ \sum_{k=0}^{\infty} \int_{\tilde{E}} \prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i)), t | y, s \right\} P^{(k)}(dz^k)
\end{aligned}$$

$$\begin{aligned}
& + \sum_{k=0}^{\infty} \int_{\tilde{E}} k [\Delta \frac{\partial}{\partial \zeta} [\prod_{i=1}^k (\eta(z_i) + (\zeta + h\Delta) \theta(z_i), t|y, s) + O(\Delta^2)]] P^{(k)}(dz^k) \\
& - \sum_{k=0}^{\infty} \int_{\tilde{E}} k \prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i), t|y, s) P^{(k)}(dz^k) \}
\end{aligned}$$

where $0 \leq h \leq 1$, by the mean value theorem. Hence, by the dominated convergence theorem (cf. Parzen, 1962, p. 274),

$$\begin{aligned}
& E \left\{ \prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i), t|y, s) \right\} = \sum_{k=0}^{\infty} \int_{\tilde{E}} k \left\{ \lim_{\Delta \rightarrow 0} \left[\frac{\partial}{\partial \zeta} \left(\prod_{i=1}^k (\eta(z_i) \right. \right. \right. \\
& \quad \left. \left. \left. + (\zeta + h\Delta) \theta(z_i), t|y, s) \right) + O(\Delta) \right] P^{(k)}(dz^k) \right\} \\
& = E \left\{ \frac{\partial}{\partial \zeta} \left[\prod_{i=1}^k (\eta(z_i) + \zeta \theta(z_i), t|y, s) \right] \right\}.
\end{aligned}$$

In particular, taking A_1 and A_2 to be the intervals $[0, z_1]$ and $[0, z_2]$ respectively, defining

$$M_1^*(z_1, t|y, s) = M_1([0, z_1], t|y, s), \quad (4.21)$$

$$M_2^*(z_1, z_2, t|y, s) = M_{(2)}([0, z_1] \times [0, z_2], t|y, s) \quad (4.22)$$

and using (4.19) and (4.20), the backward diffusion equations for the first and second factorial moment measures of the population in $[0, z_1]$ and $[0, z_1] \times [0, z_2]$ respectively can be obtained. The backward diffusion equation for the first moment measure is

$$-\frac{\partial M_1^*(z_1, t|y, s)}{\partial s} = \mathcal{Q}_y M_1^*(z_1, t|y, s) + \sum_{i=1}^{\infty} i \lambda_i M_1^*(z_1, t|0, s) \quad (4.23)$$

with the initial condition

$$M_1^*(z_1, t|y, t) = \begin{cases} 0 & \text{for } y > z_1 \\ 1 & \text{for } y \leq z_1 \end{cases} \quad (4.24)$$



and the boundary conditions

$$\frac{\partial M_1^*(z_1, t|y, s)}{\partial y} \Big|_{y=0} = 0, \quad (4.25)$$

$$M_1^*(z_1, t|y, s) \Big|_{y=z_{1+}} = M_1^*(z_1, t|y, s) \Big|_{y=z_{1-}} \quad (4.26)$$

$$\frac{\partial M_1^*(z_1, t|y, s)}{\partial y} \Big|_{y=z_{1+}} = \frac{\partial M_1^*(z_1, t|y, s)}{\partial y} \Big|_{y=z_{1-}} \quad (4.27)$$

d

$$M_1^*(z_1, t|\tilde{a}, s) = 0 \quad (4.28)$$

or an absorbing barrier at $y = \tilde{a}$, or

$$\frac{\partial M_1^*(z_1, t|y, s)}{\partial y} \Big|_{y=\tilde{a}} = 0 \quad (4.29)$$

or a reflecting barrier at $y = \tilde{a}$. The continuity conditions

(4.26) and (4.27) at $y = z_1$ for $M_1^*(z_1, t|y, s)$ and its derivative

with respect to y are necessary in view of the fact that the

initial condition (4.24) essentially divides \tilde{D} into two regions,

$0 \leq y \leq z_1$ and $z_1 < y \leq \tilde{a}$. Conditions (4.26) and (4.27) follow

from the continuity of $G(\theta, t|y, s)$ and its derivative with respect

to y throughout \tilde{D} . Similarly, the diffusion equation for the

second factorial moment measure is

$$\begin{aligned} \frac{\partial M_2^*(z_1, z_2, t|y, s)}{\partial s} &= \mathcal{Q} M_2^*(z_1, z_2, t|y, s) \\ &+ \sum_{i=1}^{\infty} i(i-1) \lambda_i(y, s) M_1^*(z_1, t|0, s) M_1^*(z_2, t|0, s) + \sum_{i=1}^{\infty} i \lambda_i(y, s) M_2^*(z_1, z_2, t|y, s) \\ &+ \sum_{i=1}^{\infty} i \lambda_i(y, s) [M_1^*(z_1, t|y, s) M_1^*(z_2, t|0, s) + M_1^*(z_1, t|0, s) M_1^*(z_2, t|y, s)] \quad (4.30) \end{aligned}$$

with the initial condition

$$M_2^*(z_1, z_2, t|y, t) = 0 \quad (4.31)$$

and boundary conditions

$$\left. \frac{\partial M_2^*(z_1, z_2, t|y, s)}{\partial y} \right|_{y=0} = 0, \quad (4.32)$$

and

$$M_2^*(z_1, z_2, t|\tilde{a}, s) = 0 \quad (4.33)$$

for an absorbing barrier, or

$$\left. \frac{\partial M_2^*(z_1, z_2, t|y, s)}{\partial y} \right|_{y=\tilde{a}} = 0 \quad (4.34)$$

for a reflecting barrier at $y = \tilde{a}$. Since the initial condition does not divide the spatial domain into different regions, the continuity conditions for the points z_1 and z_2 , which follow from (4.19) and (4.20) will always be satisfied, and thus need not be considered separately.

Knowing $M_1^*(z_1, t|y, s)$, the expected number of entities in any arbitrary interval $(\tilde{z}_1, \tilde{z}_2]$ can be calculated from the difference $M_1^*(\tilde{z}_2, t|y, s) - M_1^*(\tilde{z}_1, t|y, s)$. Similarly, the second factorial moment for the number of entities in $(\tilde{z}_1, \tilde{z}_2]$ is given by

$$\begin{aligned} M_{(2)}((\tilde{z}_1, \tilde{z}_2] \times (\tilde{z}_1, \tilde{z}_2], t|y, s) &= M_{(2)}((0, \tilde{z}_2] \times (0, \tilde{z}_2], t|y, s) \\ &+ M_{(2)}((0, \tilde{z}_1] \times (0, \tilde{z}_1], t|y, s) - 2M_{(2)}((0, \tilde{z}_1] \times (0, \tilde{z}_2], t|y, s) \\ &= M_2^*(\tilde{z}_2, \tilde{z}_2, t|y, s) - 2M_2^*(\tilde{z}_1, \tilde{z}_2, t|y, s) \\ &+ M_2^*(\tilde{z}_1, \tilde{z}_1, t|y, s). \end{aligned} \quad (4.35)$$



Knowing the first two factorial moments, the variance of the population in any interval A can be calculated by using the relation

$$\text{var}(N(A)) = M_{(2)}(A \times A) + M_1(A) - (M_1(A))^2.$$

Remark 4.5: When the transition probabilities related to birth and death, as well as the transition probability density for the movement of an entity in \tilde{D} do not depend upon time, i.e., when the stochastic population process is stationary in time, the dependence of $G(\theta, t|y, s)$ on t and s will enter only through the difference $(t-s)$. In view of this,

$$G(\theta, t|y, s) = G(\theta, t-s|y, 0)$$

$$\equiv G(\theta, t-s|y).$$

Similar relationships for the moments follow immediately from the GF. This makes it possible to define a new time scale $\tau = t-s$, so that $\frac{\partial}{\partial s} = -\frac{\partial}{\partial \tau}$. The backward diffusion equation for the PGF for the stationary process and the initial as well as boundary conditions are given by replacing the partial derivative $-\frac{\partial}{\partial s}$ by $\frac{\partial}{\partial \tau}$ and $G(\theta, t|y, s)$ by $G(\theta, \tau|y)$ in (4.12) thru (4.16). Similarly, the corresponding equations for the first two factorial moment measures are given by replacing $-\frac{\partial}{\partial s}$ by $\frac{\partial}{\partial \tau}$, as well as $M_1^*(z, t|y, s)$ and $M_2^*(z_1, z_2, t|y, s)$ by $M_1^*(z, \tau|y)$ and $M_2^*(z_1, z_2, \tau|y)$ respectively, in (4.23) thru (4.34).

4.4 The Forward Diffusion Equations

The derivation of the forward equation for the generating functionals is much more complicated than the backward equations because in the derivation of a forward equation it is necessary to

describe the changes in the generating functional due to the changes in positions of all the entities in the population at time t during a small time interval $(t, t+\delta t)$. On the other hand, only the changes in a single ancestor had to be considered while deriving the backward equation. As indicated in Section 4.1 a forward description is in terms of the number of entities in an infinitesimal interval $(z, z+dz) \subset \tilde{D}$. The forward equation for the moments is thus in terms of densities of the moments. The derivation of the forward diffusion equation therefore depends upon the existence of these densities. In the following derivations the moment generating functional (MGF) will be used for obtaining the forward equations for the moment densities (for the cases where these densities exist). The PGF or the characteristic functional can also be used in place of the MGF.

Kendall (1949) first derived the forward differential equations for the first moment (mean), variance and covariance densities for the age distribution of a population, using a generating functional. The equation for the first moment density was the same as (4.1) with $i = 1$, $v_1 = 1$, and the birth rate appearing only in a boundary condition. The densities were defined as

$$(\text{mean}) \quad E(N(dz, t)) = \tilde{m}_1(z, t) dz + o(dz),$$

$$(\text{variance}) \quad \text{Var}(N(dz, t)) = E(N(dz, t))^2 - (E(N(dz, t)))^2, \\ = \tilde{m}_2(z, t) dz + o(dz)$$

$$\text{and (Covariance)} \quad \text{Cov}(N(dz_1, t)N(dz_2, t)) = E(N(dz_1, t)N(dz_2, t)) \\ - E(N(dz_1, t))E(N(dz_2, t)) = \tilde{m}_{12}(z_1, z_2, t) + o(dz_1, dz_2),$$

where $N(dz, t)$ denotes the number of entities in the interval $(z, z+dz)$ at time t . It can be seen that in the terminology of

the theory of stochastic population processes the individual state space \tilde{E} for Kendall's problem consists of the non-negative real line denoting the age of an individual. The variance density is connected with the density of the second moment measure with respect to the diagonal of $\tilde{E} \times \tilde{E}$ and the covariance density is related to the density of the second moment measure with respect to $\tilde{E} \times \tilde{E}$ itself (cf. Sections 2.3 and 2.4). The method of derivation of the equations in this section will be similar to that of Kendall.

Let $\phi[\theta, t_2 | N_0, t_1]$ denote the MGF for the population at time t_2 given the measure N_0 denoting the distribution of entities in \tilde{D} at time t_1 . In view of (4.4), when there is only one ancestor at y at time s ,

$$\phi[\theta, s | N_0, s] = E[\exp \int_{\tilde{D}} \theta(z) I(dz | y)] = \theta(y). \quad (4.36)$$

It can be seen from the properties of conditional expectations that

$$\begin{aligned} \phi[\theta, t + \delta t | N_0, s] &= E[\exp \int_{\tilde{D}} \theta(z) N(dz, t + \delta t) | N_0, s] \\ &= E[E[\exp \int_{\tilde{D}} \theta(z) N(dz, t + \delta t) | N(z, t)] | N_0, s]. \end{aligned} \quad (4.37)$$

Noting that the integral with respect to the counting measure represents a sum over all the entities in the population,

$$\begin{aligned} E[\exp \int_{\tilde{D}} \theta(z) N(dz, t + \delta t) | N(z, t)] &= E[\prod_{t+\delta t} \exp \theta(z_j) | N(z, t)] \\ &= \prod_{t+\delta t} E[\exp \theta(z_j) | N(z, t)], \end{aligned} \quad (4.38)$$

where the product has one term corresponding to each entity at $z_j \in \tilde{D}$ at time $t + \delta t$, and the second expression on the right follows from the independence of the entities. It is easy to see that

$$\prod_{t+\delta t} E[\exp \theta(z_j) | N(z, t)] = \prod_t \{ (1 - \mu(z_j, t) \delta t) E[\exp(\theta(z_j + \delta z))] + \mu(z_j, t) \delta t \{ 1 - \sum_{i=1}^{\infty} \lambda_i(z_j, t) \delta t + \sum_{i=1}^{\infty} \lambda_i(z_j, t) \delta t \exp(i\theta(0)) \} \}. \quad (4.39)$$

The quantity $E[\exp(\theta(z_j + \delta z))]$ can be written as

$$\begin{aligned} E[\exp(\theta(z_j + \delta z))] &= \int_{\mathbb{D}} g(\delta z; z_j, t) \exp(\theta(z_j + \delta z)) d(\delta z) \\ &= \exp(\theta(z_j)) \{ 1 + \delta t [\theta'(z_j) r(z_j, t) + \frac{\sigma^2(z_j, t)}{2} (\theta''(z_j) + \theta'(z_j)^2)] + o(\delta t) \}. \end{aligned} \quad (4.40)$$

The second expression on the right side of (4.40) is obtained by expanding $\exp(\theta(z_j + \delta z))$ in a Taylor series about z_j and using (4.1). The quantities $\theta'(z_j)$ and $\theta''(z_j)$ denote the first two derivatives of $\theta(z)$ evaluated at $z = z_j$. Substituting (4.39) and (4.40) in (4.38) and simplifying,

$$\begin{aligned} \prod_{t+\delta t} E[\exp(\theta(z_j) | N(z, t))] &\approx \prod_t \exp(\theta^*(z_j)) \\ &= \exp\left[\int_{\mathbb{D}} \theta^*(z) N(dz, t)\right], \end{aligned} \quad (4.41)$$

where

$$\begin{aligned} \theta^*(z) &= \theta(z) + \delta t \{ r(z, t) \theta'(z) + \frac{\sigma^2(z, t)}{2} [\theta''(z) + \theta'(z)^2] \\ &\quad - \mu(z, t) [1 - \exp(-\theta(z))] + \sum_{i=1}^{\infty} \lambda_i(z, t) [\exp(i\theta(0)) - 1] \}. \end{aligned} \quad (4.42)$$

Using (4.41), (4.37) can be written as

$$\phi[\theta, t + \delta t | N_0, s] = \phi[\theta^*, t | N_0, s]. \quad (4.43)$$

Equation (4.43) along with the initial condition (4.36) represents a complete forward description of the population in terms of the MGF. Further simplification of (4.43) is not possible.

The first moment of the population in the set $A \subset \tilde{D}$ is obtained from the MGF by using (2.18) and (2.19):

$$\begin{aligned}\delta_{\theta}\phi(0) &= - \frac{\partial \phi[\zeta\theta, t | N_0, s]}{\partial \zeta} \Big|_{\zeta=0} \\ &= - \frac{\partial}{\partial \zeta} \{E[\exp \int_{\tilde{D}} \zeta\theta(z)N(dz, t) | N_0, s]\} \\ &= - \int_{\tilde{D}} \theta(z)E[N(dz, t) | N_0, s] = \int_A M_1(dz, t | N_0, s), \quad (4.44)\end{aligned}$$

when $\theta(z) = -I(A|z)$. In the derivation of (4.44) it was required to commute the operator $\frac{\partial}{\partial \zeta}$ with the expectation and integral operators. This can be justified by using an argument similar to that in Remark 4.4. In the case where the number of entities in a small arbitrary interval dz is finite with probability one, the first moment density with respect to z exists and $M_1(dz, t | N_0, s)$ can be written as $m_1(z, t | N_0, s)dz$, where $m_1(z, t | N_0, s)$ denotes the density of the first moment measure. When $\theta(z)$ in (4.37) is replaced by $\zeta\theta(z)$, (4.43) changes to

$$\phi[\zeta\theta, t + \delta t | N_0, s] = \phi[\theta^{**}, t | N_0, s], \quad (4.45)$$

where

$$\begin{aligned}\theta^{**}(z) &= \zeta\theta(z) + \delta t \{r(z, t)\zeta\theta'(z) + \frac{\sigma^2(z, t)}{2} [\zeta\theta''(z) + \zeta\theta'^2(z)] \\ &\quad - \mu(z, t)[1 - \exp(-\zeta\theta(z))] + \lambda(z, t)[\exp(\zeta\theta(z)) - 1]\}. \quad (4.46)\end{aligned}$$

Using (4.44) and taking the set A to be an arbitrary interval

$(\tilde{z}_1, \tilde{z}_2]$, (4.45) leads to

$$\int_{\tilde{z}_1}^{\tilde{z}_2} \frac{\partial m_1(z, t | N_0, s) dz}{\partial t} = \int_{\tilde{z}_1}^{\tilde{z}_2} \theta_z^* m_1(z, t | N_0, s) dz, \quad (4.47)$$

where the operator a_z^* is given by

$$a_z^* = \left\{ \frac{1}{2} \frac{\partial^2 [\sigma^2(z, t) \cdot]}{\partial z^2} - \frac{\partial [r(z, t) \cdot]}{\partial z} - [\mu(z, t) \cdot] \right\}. \quad (4.48)$$

It should be noted that although the indicator function $I((\tilde{z}_1, \tilde{z}_2] | z)$ has jump discontinuities at \tilde{z}_1 and \tilde{z}_2 , and thus the derivatives $\theta'(z)$ and $\theta''(z)$ in (4.46) cannot be defined at these points of discontinuity, it is possible to define an infinitely differentiable function $f(z)$ which agrees with the indicator function with an arbitrarily small error $\epsilon > 0$, (defined as the integral of the absolute value of the difference between the two functions) over a closed interval $[\tilde{\alpha}_1, \tilde{\alpha}_2]$ containing $(\tilde{z}_1, \tilde{z}_2]$ and is zero outside $(\beta_1, \beta_2) \subset [\tilde{\alpha}_1, \tilde{\alpha}_2]$ (see e.g., Indritz, 1963; p. 254). The function $f(z)$, instead of the indicator function itself, has to be used in the derivation of (4.47). Since the limits of integration \tilde{z}_1 and \tilde{z}_2 in (4.47) are arbitrary, the integrands must be equal, and hence

$$\frac{\partial m_1(z, t | N_0, s)}{\partial t} = a_z^* m_1(z, t | N_0, s). \quad (4.49)$$

When the interval $(\tilde{z}_1, \tilde{z}_2]$ includes one of the boundaries, i.e., $z = 0$ or $\tilde{\alpha}$, the following boundary conditions result:

$$\begin{aligned} r(0, t) m_1(0, t | N_0, s) - \frac{1}{2} \frac{\partial [\sigma^2(0, t) m_1(0, t | N_0, s)]}{\partial z} \\ = \int_0^{\tilde{\alpha}} \sum_{i=1}^{\infty} i \lambda_i(z, t) m_1(z, t | N_0, s) dz, \end{aligned} \quad (4.50)$$

and

$$r(\tilde{\alpha}, t) m_1(\tilde{\alpha}, t | N_0, s) - \frac{1}{2} \frac{\partial [\sigma^2(\tilde{\alpha}, t) m_1(\tilde{\alpha}, t | N_0, s)]}{\partial z} = 0 \quad (4.51)$$

for a reflecting barrier, or



$$m_1(\tilde{a}, t | N_0, s) = 0 \quad (4.52)$$

for an absorbing barrier. The initial condition is obtained from (4.36) as the Dirac- δ function

$$m_1(z, s | N_0, s) = \delta(z-y) \quad (4.53)$$

when there is one ancestor at $z = y$ at time s . When there are k ancestors located at y_j , $j = 1, 2, \dots, k$, the initial condition will be the sum of k δ -functions

$$m_1(z, s | N_0, s) = \sum_{j=1}^k \delta(z - y_j) \quad (4.54)$$

It can be seen that the forward diffusion equation (4.49)

for the first moment density is the same as the diffusion equation (3.4) with an additional term accounting for the death of entities, and has the same form as (1.3) where $\sigma_R^2 = r_R = 0$ and the death term D replaced by the corresponding term in (4.48). The derivation of (1.3) has been mainly based upon an analogy between mixing of eddies of a fluid in a process vessel and molecular diffusion (cf. Levenspiel and Smith, 1957). Use of the diffusion equation in population balance models in chemical engineering has been restricted to only the residence time distribution analysis of some process vessels. A major reason for not using the diffusion equation to characterize particulate processes such as sucrose crystallization thus far is perhaps a lack of a full appreciation of the stochastic nature of the diffusion process.

The development of (1.1) is based upon purely deterministic considerations, with x denoting the number density of the entities

in the Euclidean space E . To account for the fact that an integer-valued population is characterized by a continuous density function x , it is always assumed that the population must be large in order to justify the "continuum approximation" (Randolph and Larson, 1971; p. 13). Thus, it is assumed that (1.1) is not valid for small populations. Moreover, (1.1) cannot be used to describe any population where the movement of all the entities at a given point in E at time t is not identical in all respects. When the diffusion coefficient $\frac{\sigma^2}{2}$ is zero, (4.49) has the same form as (1.1), and thus (4.49) represents a generalization of (1.1). In view of this, (1.1) can be seen to be rigorously true even for small populations if x is interpreted as the first moment density instead of a number density and other parameters are given the appropriate probabilistic interpretations, and if the initial population is characterized as a sum of Dirac- δ functions instead of a smooth density. Evidently, the assumption of independence of entities underlying the stochastic formulation implies that the resulting partial differential equations and the boundary conditions cannot have any nonlinearities with respect to the first moment density. When the initial population is large, the assumption of a smooth density function may be justified. It is only in this sense that the initial population must be large in order to justify the "continuum approximation" in (1.1).

Derivation of the forward equation for the second moment density is complicated by the fact that the second moment distribution has a "mass concentration" along the diagonal of the domain $\mathcal{D} \times \mathcal{D}$ (see Section 2.4). The second moment of the number of entities in

the set $A_1 \times A_2 \subset \tilde{D} \times \tilde{D}$ can be calculated using (2.18) and (2.19):

$$\begin{aligned} \delta_{\theta}^2(0) &= \frac{\partial^2}{\partial \zeta^2} [E \exp \int_{\tilde{D}} \zeta \theta(z) N(dz, t | N_0, s)] \Big|_{\zeta=0} \\ &= E \int_{\tilde{D}} \theta(z_1) N(dz_1, t | N_0, s) \int_{\tilde{D}} \theta(z_2) N(dz_2, t | N_0, s) \\ &= E \int_{\tilde{D} \times \tilde{D}} \theta(z_1) \theta(z_2) N_2(dz_1 \times dz_2, t | N_0, s) \end{aligned} \quad (4.55)$$

by Fubini's theorem on summations with respect to product measures (see e.g., Feller, 1966; p. 120). As discussed in Sections 2.3 and 2.4, the second product moment can be expressed as a sum of a density $N_2(z, t | N_0, s)$ with respect to the diagonal of $\tilde{D} \times \tilde{D}$ (i.e., the line $z_1 = z_2$) and a density with respect to $\tilde{D} \times \tilde{D}$. Use of these densities in (4.55) yields

$$\begin{aligned} &E \int_{\tilde{D} \times \tilde{D}} \theta(z_1) \theta(z_2) N_2(dz_1 \times dz_2, t | N_0, s) \\ &= \int_{\{z_1=z_2\} \cap \{\tilde{D} \times \tilde{D}\}} \theta(z_1)^2 m_2(z_1, t | N_0, s) dz_1 \\ &+ \int_{\tilde{D} \times \tilde{D}} \theta(z_1) \theta(z_2) m_{12}(z_1, z_2, t | N_0, s) dz_1 dz_2. \end{aligned} \quad (4.56)$$

As in the case of the equation for the first moment density, taking $\theta(z)$ to be the infinitely differentiable function approximating $I(A_1|z)I(A_2|z)$ (or equivalently, taking $\theta(z_1)$ and $\theta(z_2)$ to be the infinitely differentiable functions approximating $-I(A_1|z_1)$ and $-I(A_2|z_2)$ respectively), the second moment of the population on the set $A_1 \times A_2$ can be obtained using (4.45), (4.55) and (4.56). Thus, for the density m_{12} ,

$$\int_{A_1} \int_{A_2} \frac{\partial m_{12}(z_1, z_2, t | N_0, s)}{\partial t} dz_1 dz_2 = \int_{A_1} \int_{A_2} \{a_{z_1}^* + a_{z_2}^*\} m_{12}(z_1, z_2, t | N_0, s) dz_1 dz_2, \quad (4.57)$$

where $a_{z_1}^*$ and $a_{z_2}^*$ denote the operator defined in (4.48). An additional condition that

$$m_1(z_1, t | N_0, s) = m_2(z_1, t | N_0, s) \quad (4.58)$$

needs to be satisfied in the course of the derivation to obtain the equation for the density m_2 :

$$\begin{aligned} & \int_{\{z_1=z_2\} \cap \{A_1 \times A_2\}} \frac{\partial m_2(z_1, t | N_0, s)}{\partial t} \\ &= \int_{\{z_1=z_2\} \cap \{A_1 \times A_2\}} \{a_{z_1}^* m_2(z_1, t | N_0, s) + \mu(z_1, t) [m_1(z_1, t | N_0, s) \\ & \quad - m_2(z_1, t | N_0, s)]\} dz_1. \end{aligned} \quad (4.59)$$

As mentioned in Section 2.4, the second factorial moment of the product counting measure $N_2(A \times A)$ does not have a concentration along the diagonal of $\tilde{D} \times \tilde{D}$ only when the point process describing the population is orderly. Since the second product moment equals the sum of the first moment and the second factorial moment, it can be seen that (4.58) is satisfied only for an orderly process. When this is the case, $m_{12}(z_1, z_2, t | N_0, s)$ completely describes the second factorial moment of the population. In view of this discussion and the fact that A_1 and A_2 are arbitrary, the forward diffusion equation for the second factorial moment density is

$$\frac{\partial m_{12}(z_1, z_2, t | N_0, s)}{\partial t} = \{a_{z_1}^* + a_{z_2}^*\} m_{12}(z_1, z_2, t | N_0, s) \quad (4.60)$$

for the case where the point process describing the population at any given time is orderly. When A_1 or A_2 include the boundaries, the following boundary conditions result:

$$\psi_{z_1}(0, z_2, t) \equiv r(0, t) m_{12}(0, z_2, t | N_0, s) - \frac{1}{2} \frac{\partial [\sigma^2(0, t) m_{12}(0, z_2, t | N_0, s)]}{\partial z_1}$$

$$= \int_0^{\tilde{a}} \sum_{i=1}^{\infty} i \lambda_i(z_1, t) m_{12}(z_1, z_2, t | N_0, s) dz_1 \quad (4.61)$$

$$+ \sum_{i=1}^{\infty} i \lambda_i(z_2, t) m_1(z_2, t | N_0, s),$$

$$\psi_{z_2}(z_1, 0, t) = \int_0^{\tilde{a}} \sum_{i=1}^{\infty} i \lambda_i(z_2, t) m_{12}(z_1, z_2, t | N_0, s) dz_2$$

$$+ \sum_{i=1}^{\infty} i \lambda_i(z_1, t) m_1(z_1, t | N_0, s) \quad (4.62)$$

and

$$\psi_{z_1}(\tilde{a}, z_2, t) = \psi_{z_2}(z_1, \tilde{a}, t) = 0 \quad (4.63)$$

for a reflecting barrier, or

$$m_{12}(\tilde{a}, z_2, t | N_0, s) = m_{12}(z_1, \tilde{a}, t | N_0, s) = 0 \quad (4.64)$$

for an absorbing barrier. In view of (4.58) it is no longer necessary to solve for $m_2(z, t | N_0, s)$ separately. When there is only one ancestor at y at time s , the initial condition will be

$$m_{12}(z_1, z_2, s | N_0, s) = 0. \quad (4.65)$$

If at time s there are k ancestors located at distinct points y_j , $j = 1, \dots, k$, the initial condition will be

$$m_{12}(z_1, z_2, s | N_0, s) = \sum_{\substack{i,j \\ i \neq j}} \delta(z_1^* - y_i) \delta(z_2^* - y_j) + \sum_{\substack{i,j \\ i \neq j}} \delta(z_1^{**} - y_i) \delta(z_2^* - y_j), \quad (4.66)$$

where

$$z_1^*, z_2^* \in [0, z_1] \cap [0, z_2]$$

and

$$z_1^{**} \in \{[0, z_1] \cup [0, z_2] - [0, z_1] \cap [0, z_2]\} \text{ (cf. Figure 2.1b).}$$

Variance of the number of entities in any interval A can be readily calculated by using the relation

$$\begin{aligned} \text{Var}(N(A)) &= \int_A \int_A m_{12}(z_1, z_2, t | N_0, s) dz_1 dz_2 + \int_A m_1(z, t | N_0, s) dz \\ &\quad - \left\{ \int_A m_1(z, t | N_0, s) dz \right\}^2. \end{aligned} \quad (4.67)$$

4.5 Comparison of the Backward and Forward Equations

It is not possible to obtain an explicit forward partial differential equation for a generating functional, whereas the backward equation can be readily derived. Although an analytical solution of this equation cannot be obtained, the extinction probability as well as the probability generating function for the number of entities in any set $A \subset \tilde{D}$ can be calculated numerically. This possibility has broad implications in the management and control of biological populations. For example, in the biological control of insect pests, a small number of parasites or predators are usually released in an infested area. The extinction probability for the controlling population can serve as a measure of the failure probability of the control strategy.

In case of a finite spatial domain \tilde{D} and constant parameters, the forward diffusion equation can be solved by using separation of

variables to yield an infinite series, as illustrated in the next chapter. The coefficients in the infinite series have to be evaluated by using the initial condition. The initial condition is usually in the form of a sum of Dirac- δ functions, and thus cumbersome to handle in general -- when the parameters are constant, as well as in the general case when a numerical solution has to be attempted. In case of a numerical solution the integral in the boundary condition increases the computational effort considerably. As in the case of the corresponding backward equation, the solution of the forward equation with more than one ancestor can also be expressed as a sum of the required number of solutions, each with one ancestor in an appropriate location corresponding to the initial distribution, but this does not simplify the solution to any appreciable extent. When the initial population is large, the assumption of an initial density may be justified and the solution would be much simpler. On the other hand, the backward equation is much simpler to handle when the initial population is small. Although a general analytical solution of the backward equation would be quite difficult, a numerical solution is relatively easier than the forward equation because of the simpler boundary conditions and the fact that the initial condition is a step function rather than a Dirac- δ function. It should be noted that the solution of the backward equation for the first moment gives the value of the expected number of entities in $[0, z]$ for a fixed $z \in \tilde{D}$ at time t , as a function of the initial location of the ancestor and the initial time s . In order to obtain the complete distribution of the expected number of entities in \tilde{D} it is necessary to solve the diffusion

equation for a number of values of z . In addition, when there are more than one ancestor at the initial time s , an appropriate number of solutions each with a single ancestor in a location corresponding to the initial distribution have to be added to obtain the final answer. On the other hand, the forward equation gives the entire distribution of the expected number of entities in the population at time t directly. Thus, when the initial population is large enough to justify a smooth initial density, the forward equation would be computationally more efficient than the backward equation.

The forward equation (4.60) for the second factorial moment density is a linear second order partial differential equation with three independent variables. In spite of its linearity an analytical solution is not possible even in case of constant parameters because of the coupling with the first moment in the boundary conditions. A numerical solution will be cumbersome because of the three independent variables and the integrals in the boundary conditions. Moreover, to get the variance of the population in any interval, the densities have to be integrated over a two dimensional domain (cf. Equation (4.67)). On the other hand, the backward equations for both the moments have only two independent variables and the moments of the population in any interval can be readily computed from two (or three) solutions corresponding to the boundaries of the interval, as discussed earlier in this chapter. Thus, the total computational effort will be much less than that for the forward equations when the initial population is small. Furthermore, the backward equation is valid even if the population is not orderly and does not need the continuity of $r(z,t)$ and $\sigma^2(z,t)$ as well as the continuity of

the first derivative of $\sigma^2(z,t)$. Thus it can be seen that the backward equation is a more convenient tool for obtaining the moments of the population when the initial population is small. When the initial population is large, the forward equation for the first moment is often much easier to handle and since the second moment will always be rather insignificant in this case, its evaluation will no longer be necessary.

CHAPTER V

SOLUTION OF THE DIFFUSION EQUATIONS

Solutions of the diffusion equations for the first two moments and the PGF for a population with a one-dimensional individual state space \tilde{E} are presented in this chapter.

5.1 General Considerations

As in the last chapter, the spatial domain \tilde{D} is taken as the interval $[0, \tilde{a}]$. Analytical solution of the equations for the moments is possible when the parameters are constant. For the analytical treatment, values of the parameters are taken as follows: For the entities at $z \in \tilde{D}$ at time t ,

$$\begin{aligned}\sigma^2(z, t) &= \sigma^2 > 0, \\ r(z, t) &= r, \\ \mu(z, t) &= \mu, \\ \lambda_i(z, t) &= \begin{cases} 0 & \text{for } 0 \leq z < \tilde{b} \\ \lambda_i & \text{for } \tilde{b} \leq z \leq \tilde{a}, i = 1, 2, \dots, \infty. \end{cases}\end{aligned}\tag{5.1}$$

The parameters as given by (5.1) are quite realistic for a broad class of problems. In the crystallization process with z denoting the characteristic length of a crystal, the assumption of a constant growth is very common when the supersaturation of the magma is relatively constant. A constant diffusion coefficient may also be justified under a steady flow of magma into and out

of the crystallizer and constant supersaturation. A constant death rate reflects a perfect mixing in the crystallizer and a constant efflux of the magma, and the assumption of a zero birth rate for $0 \leq z < \tilde{b}$ implies that a nucleus must grow to a certain size before it can breed new nuclei by secondary nucleation.

In biological populations with z representing the weight of an individual, the growth rate as measured in terms of the rate of increase of biomass of an individual is high when the individual is young and reduces progressively as the biomass increases. However, it can be seen that by choosing z to be a proper maturity variable the dependence of the mean growth rate $r(z,t)$ on z can be eliminated. In cold-blooded organisms and plants the metabolic activity is a function of the body temperature of the individual. Due to the diurnal and annual fluctuations in the atmospheric temperature and solar radiation the growth rates of such individuals in a general ecological system are functions of time. In all poikilothermic species (i.e., cold-blooded organisms and plants) the temperature-growth relation shows a characteristic behavior, which can be used to define a physiological time scale from the chronological time-temperature relationship, so as to have a constant mean growth rate with respect to the physiological time. A constant death rate for individuals in a biological population may be a reasonable first approximation in many cases. Similarly, the assumption of zero reproduction rate until a certain value \tilde{b} of z (corresponding to the "age" of puberty) and a constant rate thereafter can also be seen to be a reasonable assumption for many populations. The relationship of physiological and physical activities such as reproduction and locomotion with temperature also shows a pattern similar

to the temperature-growth rate relationship. Hence, it may be assumed that the temperature dependence (and thus, to a large extent, the dependence on chronological time) of these parameters for poikilothermic species can be removed by the use of physiological time.

A detailed discussion of the physiological time is given in Appendix A.

Use of a constant diffusion coefficient in flow through packed beds is very common. Little work has been done so far in the application of the diffusion equations in the modeling of biological populations. A constant diffusion coefficient is usually used in the spatial migration studies. Stuart and Merkle (1965) also used a constant diffusion coefficient in their study of cell dynamics. However, a careful study of the relationship of the diffusion coefficient with temperature and the state z of an individual must be made to check whether the same transformations in z and t to yield new maturity and physiological time variables lead to constant values for all parameters including the diffusion coefficient.

In the case where the diffusion coefficient is zero and the other parameters are given by (5.1), it can be seen that a simple transformation on z by defining $\tilde{z} = z/r$ reduces the rean rate of change r to unity in terms of \tilde{z} . In this case the equations for the moment distributions are the same as those for the age distribution of a biological population, and the results on the age-dependent birth and death processes can be used directly (see e.g., Kendall, 1949; Goodman, 1967; Bartlett, 1969). In the solutions presented in this chapter it is assumed that the diffusion coefficient is always positive.

It should be noted that (5.1) implies stationarity of the diffusion process (i.e., that the parameters are not time-dependent) and the abbreviated notation defined in Remark 4.5 can be used. For abbreviation, the first moment density $m_1(z, t | N_0, s)$ is denoted by $m_1(z, \tau)$, where $\tau = t - s$. The initial distribution N_0 of the entities in \tilde{D} is expressed as a Dirac- δ function for a single ancestor. It may be possible to characterize N_0 by a smooth density function for a large initial population. Also, the quantities $\sum_{i=1}^{\infty} i \lambda_i$ and $\sum_{i=1}^{\infty} i(i-1) \lambda_i$ are abbreviated as λ^* and λ^{**} respectively.

5.2 Solution of the Diffusion Equations for the First Moment

The forward equation is used to obtain a solution for the case where \tilde{D} is finite. For the case of a semi-infinite spatial domain the backward equation is used to obtain an explicit analytical solution when $\tilde{b} = 0$ and the process starts with a single ancestor at $y = 0$.

Case 1. Finite Spatial Domain $[0, \tilde{a}]$ with a Reflecting Barrier at \tilde{a} .

The forward diffusion equation and the corresponding boundary and initial conditions for this case are

$$\frac{\sigma^2}{2} \frac{\partial^2 m_1(z, \tau)}{\partial z^2} - r \frac{\partial m_1(z, \tau)}{\partial z} - \mu m_1(z, \tau) = \frac{\partial m_1(z, \tau)}{\partial \tau}, \quad (5.2)$$

$$r m_1(\tilde{a}, \tau) - \frac{\sigma^2}{2} \frac{\partial m_1(\tilde{a}, \tau)}{\partial z} = 0, \quad (5.2a)$$

$$r m_1(0, \tau) - \frac{\sigma^2}{2} \frac{\partial m_1(0, \tau)}{\partial z} = \lambda^* \int_{\tilde{b}}^{\tilde{a}} m_1(z, \tau) dz \quad (5.2b)$$

and

$$m_1(z, 0) = m_0(z). \quad (5.2c)$$

Let $m_1(z, \tau)$ be expressed as a product

$$m_1(z, \tau) = Z(z)T(\tau), \quad (5.3)$$

where $Z(z)$ and $T(\tau)$ are functions only of z and τ respectively.

Substitution of (5.3) in (5.2), (5.2a) and (5.2b) yields

$$\frac{\alpha}{2} \frac{d^2 Z(z)}{dz^2} - r \frac{dZ(z)}{dz} - \mu Z(z) = \alpha Z(z), \quad (5.4)$$

$$rZ(\tilde{a}) - \frac{\alpha}{2} \frac{dZ(\tilde{a})}{dz} = 0, \quad (5.4a)$$

$$rZ(0) - \frac{\alpha}{2} \frac{dZ(0)}{dz} = \lambda \int_b^{\tilde{a}} Z(z) dz, \quad (5.4b)$$

and

$$\frac{dT(\tau)}{d\tau} = \alpha T, \quad (5.5)$$

where α is a constant. When $\alpha > -(\frac{r^2}{2} + \mu)$, the general solution of (5.4) is given by

$$Z(z) = A_1 \cosh \beta z + A_2 \sinh \beta z, \quad (5.6)$$

where

$$\beta = \frac{\sqrt{r^2 + 2\sigma^2(\mu + \alpha)}}{\sigma}, \quad (5.7)$$

and A_1, A_2 are constants. Similarly, when $\alpha \leq -(\frac{r^2}{2} + \mu)$, the general solution of (5.4) can be written as

$$Z(z) = A_3 \cos \tilde{\beta} z + A_4 \sin \tilde{\beta} z, \quad (5.8)$$

where

$$\beta = \sqrt{\frac{-r^2 - 2\sigma^2(\mu + \alpha)}{\sigma^2}}, \quad (5.8a)$$

and A_3, A_4 are constants. Using boundary condition (5.4a) it is possible to eliminate A_1 and A_3 from (5.6) and (5.8) respectively. Thus, (5.6) yields

$$f_1(z) \equiv Z(z) = A_2 \left[\left(\frac{\sigma^2 + r \tanh(\beta \tilde{a})}{2\sigma \beta \tanh(\beta \tilde{a}) - r} \right) \cosh(\beta z) + \sinh(\beta z) \right] \exp\left(\frac{rz}{\sigma}\right). \quad (5.9)$$

Similarly, (5.8) leads to the expression

$$Z(z) = A_4 \left[\left(\frac{\sigma^2 \tilde{b} - r \tanh(\beta \tilde{a})}{2\sigma \beta \tanh(\beta \tilde{a}) + r} \right) \cos(\beta z) + \sin(\beta z) \right] \exp\left(\frac{rz}{\sigma}\right). \quad (5.10)$$

When the boundary condition (5.4b) is applied to (5.9) and (5.10) to evaluate the admissible values of α , an infinite set $\{\alpha_k\}$ of eigenvalues is obtained. The largest eigenvalue α_1 corresponds to the solution (5.9) and is given by the root of the equation

$$(\alpha + \mu) \sinh(\beta \tilde{a}) = \lambda^* \exp(r\tilde{b}/\sigma^2) \sinh[\beta(\tilde{a} - \tilde{b})] \quad (5.11)$$

with the largest real part. Note that when α_1 is complex, there can possibly be more than one distinct dominant eigenvalues in general. The non-dominant eigenvalues are obtained from (5.10) as the roots of the equation

$$(\alpha + \mu) \sin(\beta \tilde{a}) = \lambda^* \exp(r\tilde{b}/\sigma^2) \sin[\beta(\tilde{a} - \tilde{b})]. \quad (5.12)$$

The solution of (5.5) is

$$T(\tau) = T_0 \exp(\alpha \tau), \quad (5.13)$$

where T_0 is a constant. The eigenfunction corresponding to α_1 is given by substituting the value of α_1 in (5.9). The eigenfunctions

$f_k(z)$ corresponding to the nondominant eigenvalues α_k , $k = 2, 3, \dots, \infty$, are given by substituting the value of α_k in (5.10). The general solution of (5.2) is thus given by the generalized Fourier series (cf. Indritz, 1963)

$$m_1(z, \tau) = \sum_{k=1}^{\infty} a_k f_k(z) \exp(\alpha_k \tau). \quad (5.14)$$

The coefficients a_k have to be evaluated by applying the initial condition (5.2c). Thus,

$$m_0(z) = \sum_{k=1}^{\infty} a_k f_k(z). \quad (5.15)$$

When there is only one ancestor initially at $z = z_0$, $m_0(z) = \delta(z - z_0)$ and (5.15) would represent the generalized Fourier series of the Dirac- δ function. Since the nondominant eigenvalues α_k correspond to the case $\alpha_k \leq -(\frac{r}{2\sigma} + \mu)$, $\alpha_k < 0$, and hence for a sufficiently large value of τ the contributions from only a finite number of terms in (5.14) will be significant, which can be used to determine the criteria for the number of coefficients a_k to be determined in (5.15). For very large values of τ , the contributions from all the terms in (5.14) except those corresponding to the dominant eigenvalues will be negligible, irrespective of the initial distribution. Thus, when the dominant eigenvalue is unique, asymptotically, the first moment density of the number of entities in \tilde{D} will be given by

$$\lim_{\tau \rightarrow \infty} [m_1(z, \tau)] = a_1 f_1(z) \exp(\alpha_1 \tau) \quad (5.16)$$

irrespective of the initial condition. In view of (5.16), it is easy to see that for a sufficiently large τ the first moment density

is always proportional to $f_1(z)$, which is given by (5.9) and (5.11). Similarly, integration of (5.16) over an arbitrary interval $[z_1, z_2] \in \tilde{D}$ leads to the result that when α_1 is real and distinct, asymptotically, the expected number of entities in any subset of \tilde{D} will vary exponentially with the parameter α_1 .

Remark 5.1: It should be noted that the eigenfunctions $f_k(z)$, $k = 1, \dots, \infty$ in (5.14) are linearly independent of each other but not mutually orthogonal. To evaluate a_k , $k = 1, \dots, \infty$, (5.15) is multiplied by $f_j(z)$, $j = 1, \dots, \infty$, and integrated with respect to z over the interval $[0, \tilde{a}]$ to yield an infinite system of linear equations, which has to be solved for a_k . For computational purposes the series has to be truncated after a finite number m of terms and the system of m equations with m unknowns a_k , $k = 1, \dots, m$, has to be solved to give the best approximation to a_k in some sense. If the series (5.15) converges rapidly, only a few terms in the series will be significantly different from zero and the evaluation of a_k is simple and accurate. On the other hand, when the initial condition is in the form of a Dirac- δ function, the series (5.15) will not converge for a finite number of terms, and hence accuracy of such an approximate solution can be expected to be very low.

When $\tilde{b} = 0$, the eigenvalues can be calculated analytically and the solution has a much simpler form. In this case the eigenvalues are

$$\alpha_1 = \lambda^* - \mu$$

$$\alpha_{k+1} = \{-\mu - \frac{r^2}{2\sigma} - \frac{2k^2 \pi^2 \sigma^2}{\tilde{a}^2}\} \quad k = 1, 2, \dots, \infty;$$

and the general solution is

$$m(z, \tau) = \{a_1 \exp[(\lambda^* - \mu)\tau] \left[\left(\frac{\sigma^2 \beta_0 - r \tanh(\beta_0 \tilde{a})}{\sigma^2 \beta_0 \tanh(\beta_0 \tilde{a}) - r} \right) \cosh(\beta_0 z) + s \sinh(\beta_0 z) \right] \right. \\ \left. + \sum_{k=1}^{\infty} a_{n+1} \exp(\alpha_{k+1} \tau) \left[-\frac{2k\pi\sigma^2}{\tilde{a}r} \cos\left(\frac{2k\pi z}{\tilde{a}}\right) + s \sin\left(\frac{2k\pi z}{\tilde{a}}\right) \right] \right\} \exp(rz/\sigma^2), \quad (5.17)$$

where

$$\beta_0 = \sqrt{r^2 + 2\sigma^2 \lambda^*}. \quad (5.18)$$

Since the eigenfunctions are not orthogonal to each other, Remark 5.1 applies in this case as well.

Remark 5.2: Note that in the infinite series solution of the diffusion equation it has been tacitly assumed that the eigenfunctions form a complete set (see e.g., Indritz, 1963) and that all the eigenvalues are distinct. Also, the statement that asymptotically the expected number of entities in any set $A \subset \tilde{D}$ increases exponentially with time is true only when the dominant eigenvalue (i.e., the eigenvalue with the largest real part) is real and distinct. It would be quite difficult to prove the completeness of the set of eigenfunctions and distinctness of the eigenvalues in general. Nevertheless, it is easy to see that for the case of a reflecting barrier at \tilde{a} with $\tilde{b} = 0$ discussed above, the eigenvalues are all real and distinct, and the eigenfunctions form a complete set over a class of functions continuous over a finite interval. Although the hypothesis that the dominant eigenvalue is real and distinct may be very difficult to prove analytically, it is easy to show that there must always be at least one real dominant eigenvalue. The proof of this is as follows: If all the

dominant eigenvalues are distinct and have the form $\alpha_{1R} + i\alpha_{1j}$, $j = 1, \dots, k$; $\alpha_{1j} > 0$ and $i = \sqrt{-1}$, then the asymptotic solution will be of the form $\exp(\alpha_{1R}\tau) \sum_{j=1}^k A_j \exp(i\alpha_{1j}\tau)$, which will be an oscillating function of time assuming negative values during certain time intervals. The same is true even for the case where the dominant eigenvalues are all complex, but not necessarily distinct. Since the first moment distribution must be a nonnegative quantity, at least one of the α_{1j} 's must be zero, and the corresponding coefficient A_j large enough as compared to the other coefficients so as to compensate for the negative contributions from all the terms involving the complex eigenvalues.

Case 2. Finite Spatial Domain $[0, \bar{a}]$ with an Absorbing Barrier at \bar{a} .

The forward diffusion equation, its initial condition and the boundary condition at $z = 0$ will be the same as in Case 1. The boundary condition (5.2a) will be replaced by

$$m_1(\bar{a}, \tau) = 0. \quad (5.2d)$$

Equations (5.3), (5.4), (5.4b) and (5.5) also apply in this case.

Equation (5.4a) will be replaced by

$$Z(\bar{a}) = 0. \quad (5.4c)$$

When $\alpha > -[\frac{r}{2\sigma^2} + \mu]$, the following expression is obtained by applying (5.4c) to the general solution (5.6):

$$A_1 \cosh \beta \bar{a} + A_2 \sinh \beta \bar{a} = 0.$$

Hence,

$$f_1(z) \equiv Z(z) = A_2 [\sinh(\beta z) - \tanh(\beta \bar{a}) \cosh(\beta z)] \exp(rz/\sigma^2). \quad (5.19)$$

Similarly, when $\alpha \leq -(\frac{r^2}{2\sigma^2} + \mu)$,

$$Z(z) = A_4 (\sin(\beta z) - \tan(\beta \tilde{a}) \cos(\beta z)) \exp(rz/\sigma^2). \quad (5.20)$$

When the boundary condition (5.4b) is applied to (5.19), it is found that the dominant eigenvalue is the root of

$$\begin{aligned} (\alpha + \mu) \left\{ \frac{r}{2\sigma} \sinh \beta \tilde{a} + \beta \cosh \beta \tilde{a} \right\} - \lambda^* \exp(r\tilde{b}/\sigma^2) \left\{ \frac{r}{2\sigma} \sinh[\beta(\tilde{a} - \tilde{b})] \right. \\ \left. + \beta \cosh[\beta(\tilde{a} - \tilde{b})] \right\} + \beta \lambda^* \exp(r\tilde{a}/\sigma^2) = 0, \end{aligned} \quad (5.21)$$

with the largest real part. If the dominant eigenvalue is unique and real, it can be seen that as in the case of a reflecting barrier, asymptotically the expected number of entities in any set $A \subset \tilde{D}$ varies exponentially with parameter α_1 . When the boundary condition (5.4b) is applied to (5.20), the following expression for the nondominant eigenvalues results:

$$\begin{aligned} (\alpha + \mu) \left\{ \frac{r}{2\sigma} \sin \beta \tilde{a} + \beta \cos \beta \tilde{a} \right\} - \lambda^* \exp(r\tilde{b}/\sigma^2) \left\{ \frac{r}{2\sigma} \sin[\beta(\tilde{a} - \tilde{b})] \right. \\ \left. + \beta \cos[\beta(\tilde{a} - \tilde{b})] \right\} + \lambda^* \beta \exp(r\tilde{a}/\sigma^2) = 0. \end{aligned} \quad (5.22)$$

It should be noted that Remark 5.2 is applicable in this case also. Similarly, in view of the fact that the eigenfunctions are mutually independent but not orthogonal, Remark 5.1 is also applicable. An explicit analytical solution for the eigenvalues is not possible even when $\tilde{b} = 0$.

Case 3. Semi-Infinite Spatial Domain $[0, \infty)$.

When \tilde{b} is finite and $\tilde{a} \rightarrow \infty$ with an absorbing or reflecting barrier at \tilde{a} , proceeding exactly as in Case 1 and Case 2, it can be shown that regardless of the initial condition, the asymptotic



value of the first moment density is

$$m_1(z, \tau) \sim f_1(z) \exp(\alpha_1 \tau),$$

where α_1 is the root of the equation

$$\alpha + \mu - \lambda^* \exp\left\{\left(\frac{r - \sqrt{r^2 + 2\sigma^2(\mu + \alpha)}}{2\sigma}\right)b\right\} = 0 \quad (5.21a)$$

with the largest real part,

$$f_1(z) = a_1 \exp(\rho z),$$

$$\rho = \frac{r - \sqrt{r^2 + 2\sigma^2(\mu + \alpha)}}{2\sigma},$$

and a_1 is a constant.

In this case it is also possible to obtain an explicit analytical solution for the first moment distribution when $\xi = 0$ and the ancestor is initially at zero (Brockwell, 1972). The backward diffusion equation and the corresponding initial and boundary conditions for the first moment distribution are

$$\frac{\partial M_1^*(z, \tau | y)}{\partial \tau} = \mathcal{A}_y^* M_1^*(z, \tau | y) + \lambda^* M_1^*(z, \tau | 0), \quad (5.23)$$

where the backward diffusion operator \mathcal{A}_y^* is given by (4.11),

$$M_1^*(z, 0 | y) = \begin{cases} 1 & \text{for } y \leq z \\ 0 & \text{for } y > z, \end{cases} \quad (5.23a)$$

$$\frac{\partial M_1^*(z, \tau | 0)}{\partial y} = 0, \quad (5.23b)$$

$$M_1^*(z, \tau | z_+) = M_1^*(z, \tau | z_-), \quad (5.23c)$$

$$\frac{\partial \bar{M}_1^*(z, \tau|y)}{\partial y} \Big|_{y=z_+} = \frac{\partial \bar{M}^*(z, \tau|y)}{\partial y} \Big|_{y=z_-} \quad (5.23d)$$

and

$$\bar{M}_1^*(z, \tau|\infty) = 0. \quad (5.23e)$$

Taking Laplace transforms of (5.23) thru (5.23e) the partial differential equation can be transformed into the following ordinary differential equations and boundary conditions:

$$\begin{aligned} \frac{\sigma}{2} \frac{d^2 \bar{M}_1(z, p|y)}{dy^2} + r \frac{d \bar{M}_1(z, p|y)}{dy} - (\mu + p) \bar{M}_1(z, p|y) \\ + \lambda \bar{M}_1^*(z, p|0) = \begin{cases} 1 & \text{for } y \leq z \\ 0 & \text{for } y > z, \end{cases} \end{aligned} \quad (5.24)$$

$$\frac{d \bar{M}_1(z, p|y)}{dy} \Big|_{y=0} = 0, \quad (5.24a)$$

$$\bar{M}_1(z, p|z_+) = \bar{M}_1(z, p|z_-), \quad (5.24b)$$

$$\frac{d \bar{M}_1(z, p|y)}{dy} \Big|_{y=z_+} = \frac{d \bar{M}_1(z, p|y)}{dy} \Big|_{y=z_-} \quad (5.24c)$$

and

$$\bar{M}_1(z, p|\infty) = 0, \quad (5.24d)$$

where p denotes the Laplace transform variable and $\bar{M}_1(z, p|y)$ is the Laplace transform of $\bar{M}_1^*(z, \tau|y)$. The general solution of (5.24) is

$$\bar{M}_1(z, p|y) = \begin{cases} \frac{\lambda \bar{M}_1^*(z, p|0) + 1}{\mu + p} + a_1 \exp(\beta_1 y) + a_2 \exp(\beta_2 y), & y \leq z; \quad (5.25a) \\ \frac{\lambda \bar{M}_1^*(z, p|0)}{\mu + p} + a_3 \exp(\beta_1 y) + a_4 \exp(\beta_2 y), & y > z; \quad (5.25b) \end{cases}$$

where

$$\left. \begin{aligned} \beta_1 &= \frac{-r + \sqrt{r^2 + 2\sigma^2(\mu+p)}}{\sigma^2} \\ \beta_2 &= \frac{-r - \sqrt{r^2 + 2\sigma^2(\mu+p)}}{\sigma^2} \end{aligned} \right\} \quad (5.26)$$

and a_1, a_2, a_3, a_4 are constants to be determined from the boundary conditions. In addition to these, $\bar{M}_1(z, p|0)$ is to be evaluated by setting $y = 0$ in (5.25a):

$$\bar{M}_1(z, p|0) = a_1 + a_2 + \frac{\lambda^* \bar{M}_1^*(z, p|0) + 1}{\mu + p}. \quad (5.27)$$

In view of (5.24c), $a_3 = 0$. Applying other boundary conditions, one obtains

$$a_4 \exp(\beta_2 y) = \frac{1}{\mu + s} + a_1 \exp(\beta_1 y) + a_2 \exp(\beta_2 y), \quad (5.28)$$

$$\beta_2 a_4 \exp(\beta_2 y) = \beta_1 a_1 \exp(\beta_1 y) + \beta_2 a_2 \exp(\beta_2 y) \quad (5.29)$$

and

$$\beta_1 a_1 + \beta_2 a_2 = 0. \quad (5.30)$$

Solution of (5.27) thru (5.30) for a_1, a_2, a_4 and $\bar{M}_1(z, p|0)$ yields

$$\begin{aligned} \bar{M}_1(z, p|0) &= \frac{1}{p - \lambda^* + \mu} - \frac{\exp(-\beta_1 z)}{p - \lambda^* + \mu} \\ &= \frac{1}{p - \lambda^* + \mu} - \frac{\exp(-rz/\sigma^2) \exp\left(\left[-\sqrt{r^2 + 2\sigma^2(\mu+p)}\right] y/\sigma^2\right)}{p - \lambda^* + \mu}. \end{aligned} \quad (5.31)$$

Taking inverse transform (cf. Cox and Miller, 1965; p. 221, Equations 73 and 75), the analytical solution

$$M_1^*(z, \tau | 0) = \exp[(\lambda^* - \mu^*)\tau] \left\{ 1 - \frac{z}{\sigma} \int_0^\tau \frac{\exp[-\lambda^* u - (z - ru)^2 / 2u\sigma^2]}{\sqrt{2\pi u}} du \right\} \quad (5.32)$$

is obtained.

5.3 Solution of the Diffusion Equations for the Second Moment

As mentioned in the previous chapter, it is not possible to obtain a complete analytical solution of the forward or backward equation for the second moment distribution. However, an asymptotic solution (i.e., a solution for large values of time τ) of the backward equation can be obtained by using Laplace transformations. In this section asymptotic solutions for the second factorial moment distribution will be obtained for the case where there is only one ancestor initially. These can be used to evaluate the second factorial moment distribution for a population with k ancestors initially at y_1, \dots, y_k , from the relation

$$M_2^*(z_1, z_2, \tau | y_1, \dots, y_k) = \sum_{i=1}^k M_2^*(z_1, z_2, \tau | y_i) + \sum_{i=1}^k M_1^*(z_1, \tau | y_i) \sum_{\substack{j=1 \\ j \neq i}}^k M_1^*(z_2, \tau | y_j) \quad (5.33)$$

Equation (5.33) readily follows from (4.6) and (4.20).

Case 1. Finite Spatial Domain $[0, \bar{a}]$ with a Reflecting Barrier at \bar{a} .

The backward diffusion equation and the corresponding boundary conditions are given by

$$\frac{\partial M_2^*(z_1, z_2, \tau | y)}{\partial \tau} = \mathcal{A} M_2^*(z_1, z_2, \tau | y) \quad \text{for} \quad 0 \leq y < \bar{b}, \quad (5.34a)$$

$$\begin{aligned} \frac{\partial M_2^*(z_1, z_2, \tau | y)}{\partial \tau} &= \mathcal{Q}_y M_2^*(z_1, z_2, \tau | y) + \lambda^{**} M_1^*(z_1, \tau | 0) M_1^*(z_2, \tau | 0) \\ &+ \lambda^{**} M_2^*(z_1, z_2, \tau | 0) + \lambda^* [M_1^*(z_1, \tau | 0) M_1^*(z_2, \tau | y) \\ &+ M_1^*(z_2, \tau | 0) M_1^*(z_1, \tau | y)], \text{ for } \tilde{b} \leq y \leq \tilde{a} \end{aligned} \quad (5.34b)$$

where \mathcal{Q}_y is the backward diffusion operator defined in (4.11);

$$\frac{\partial M_2^*(z_1, z_2, y)}{\partial y} \Big|_{y=0} = 0, \quad (5.34c)$$

$$M_2^*(z_1, z_2, \tau | \tilde{b}_-) = M_2^*(z_1, z_2, \tau | \tilde{b}_+), \quad (5.34d)$$

$$\frac{\partial M_2^*(z_1, z_2, \tau | y)}{\partial y} \Big|_{y=\tilde{b}_-} = \frac{\partial M_2^*(z_1, z_2, \tau | y)}{\partial y} \Big|_{y=\tilde{b}_+} \quad (5.34e)$$

and

$$\frac{\partial M_2^*(z_1, z_2, \tau | y)}{\partial y} \Big|_{y=\tilde{a}} = 0. \quad (5.34f)$$

Since λ^* and λ^{**} are step functions defined from $\lambda_i(z, \tau)$ in (5.1) with a jump discontinuity at $y = \tilde{b}$, the backward diffusion equation has to be split into two parts (5.34a) and (5.34b). The conditions (5.34d) and (5.34e) follow from the fact that the second moment distribution and its derivative with respect to y must be continuous on $[0, \tilde{a}]$. Taking Laplace transforms of (5.34a) thru (5.34f) with respect to time, the following set of equations is obtained:

$$\begin{aligned} \frac{\sigma}{2} \frac{d^2 \bar{M}_2(z_1, z_2, p | y)}{dy^2} + r \frac{d \bar{M}_2(z_1, z_2, p | y)}{dy} - (\mu + p) \bar{M}_2(z_1, z_2, p | y) &= 0 \\ \text{for } 0 \leq y < \tilde{b}, \end{aligned} \quad (5.35a)$$

and

$$\begin{aligned}
& \frac{\sigma}{2} \frac{d^2 \bar{M}_2(z_1, z_2, p|y)}{dy^2} + r \frac{d \bar{M}_2(z_1, z_2, p|y)}{dy} - (\mu + p) \bar{M}_2(z_1, z_2, p|y) \\
& + \lambda^{**} \bar{\Phi}_1(z_1, z_2, p|0) + \lambda^* \bar{M}_2(z_1, z_2, p|0) \\
& + \lambda^* [\bar{\Phi}_{12}(z_1, z_2, p|0, y) + \bar{\Phi}_{21}(z_1, z_2, p|y, 0)] \quad \text{for } \bar{b} \leq y \leq \bar{a}, \quad (5.35b)
\end{aligned}$$

where

$$\bar{\Phi}_1(z_1, z_2, p|0) = \mathcal{L}_\tau [M_1^*(z_1, \tau|0) M_1^*(z_2, \tau|0)], \quad (5.36a)$$

$$\bar{\Phi}_{12}(z_1, z_2, p|0, y) = \mathcal{L}_\tau [M_1^*(z_1, \tau|0) M_1^*(z_2, \tau|y)], \quad (5.36b)$$

$$\bar{\Phi}_{21}(z_1, z_2, p|y, 0) = \mathcal{L}_\tau [M_1^*(z_1, \tau|y) M_1^*(z_2, \tau|0)] \quad (5.36c)$$

and $\mathcal{L}_\tau [f(z, \tau|y)] =$ Laplace transform of $[f(z, \tau|y)]$ with respect to τ ,

$$= \int_0^\infty \exp(-pt) f(z, t|y) dt. \quad (5.37)$$

The boundary conditions corresponding to (5.35a) and (5.35b) are

$$\left. \frac{d \bar{M}_2(z_1, z_2, p|y)}{dy} \right|_{y=0} = 0, \quad (5.35c)$$

$$\bar{M}_2(z_1, z_2, p|\bar{b}_-) = \bar{M}_2(z_1, z_2, p|\bar{b}_+), \quad (5.35d)$$

$$\left. \frac{d \bar{M}_2(z_1, z_2, p|y)}{dy} \right|_{y=\bar{b}_-} = \left. \frac{d \bar{M}_2(z_1, z_2, p|y)}{dy} \right|_{y=\bar{b}_+} \quad (5.35e)$$

and

$$\left. \frac{d \bar{M}_2(z_1, z_2, p|y)}{dy} \right|_{y=\bar{a}} = 0. \quad (5.35f)$$

The general solutions of (5.35a) and (5.35b) are

$$\bar{M}_2(z_1, z_2, p|y) = a_1 \exp(\beta_1 y) + a_2 \exp(\beta_2 y) \quad \text{for } 0 \leq y < \bar{b} \quad (5.38a)$$

and

$$\begin{aligned}
\bar{M}_2(z_1, z_2, p|y) &= a_3 \exp(\beta_1 y) + a_4 \exp(\beta_2 y) \\
&+ \frac{\lambda^* \bar{M}_2(z_1, z_2, p|0) + \lambda^{**} \bar{\Phi}_1(z_1, z_2, p|0)}{p + \mu} \\
&+ \frac{2\lambda^* [\Theta_1(z_1, z_2, p|y, 0) \exp(\beta_1 y) - \Theta_2(z_1, z_2, p|y, 0) \exp(\beta_2 y)]}{\sigma^2(\beta_1 - \beta_2)} \\
&\text{for } \tilde{b} \leq y \leq \tilde{a} \quad (5.38b)
\end{aligned}$$

respectively, where

$$\left. \begin{aligned}
\Theta_1(z_1, z_2, p|y, 0) &= \int \exp(-\beta_1 y) [\Phi_{12}(z_1, z_2, p|0, y) \\
&\quad + \Phi_{21}(z_1, z_2, p|y, 0)] dy \\
\Theta_2(z_1, z_2, p|y, 0) &= \int \exp(-\beta_2 y) [\Phi_{12}(z_1, z_2, p|0, y) \\
&\quad + \Phi_{21}(z_1, z_2, p|y, 0)] dy
\end{aligned} \right\} \quad (5.38c)$$

and β_1, β_2 are given by (5.26). The constants a_1, a_2, a_3, a_4 are to be evaluated by using the boundary conditions (5.35c) thru (5.35f). In addition, $\bar{M}_2(z_1, z_2, p|0)$ is evaluated by setting $y = 0$ in (5.38a), i.e.,

$$\bar{M}_2(z_1, z_2, p|0) = a_1 + a_2. \quad (5.35g)$$

As discussed in the previous section, the asymptotic solution is independent of the initial location of the ancestor. It is therefore sufficient to obtain an explicit expression for $\bar{M}_2(z_1, z_2, p|0)$ in order to evaluate the asymptotic solution. Applying the conditions (5.35c) thru (5.35g), the following expression is obtained:

$$\bar{M}_2(z_1, z_2, p|0) = \frac{\left\{ \frac{2\lambda^*}{\sigma^2} (\mu+p) \left\{ \frac{\exp(\beta_1 \bar{a})}{\beta_2} [\Theta_1(z_1, z_2, p|\bar{a}, 0) - \Theta_1(z_1, z_2, p|\bar{b}, 0)] - \frac{\exp(\beta_2 \bar{a})}{\beta_1} [\Theta_2(z_1, z_2, p|\bar{a}, 0) - \Theta_2(z_1, z_2, p|\bar{b}, 0)] \right\} + \lambda^{**} \Phi_1(z_1, z_2, p|0) [\exp[\beta_1(\bar{a}-\bar{b})] - \exp[\beta_2(\bar{a}-\bar{b})]] \right\}}{(\mu+p) [\exp(\beta_1 \bar{a}) - \exp(\beta_2 \bar{a})] - \lambda^* \{ \exp[\beta_1(\bar{a}-\bar{b})] - \exp[\beta_2(\bar{a}-\bar{b})] \}} \quad (5.39)$$

It should be noted that the quantities $\Phi_1(z_1, z_2, p|0)$,

$\Theta_1(z_1, z_2, p|y, 0)$ and $\Theta_2(z_1, z_2, p|y, 0)$ are simply abbreviations of more complicated expressions and would also contribute some terms to the denominator of (5.39).

An important property of the Laplace transforms is that the dominant root of the denominator (i.e., the root with the largest real part, or the dominant pole) of the transform of a function is related to the asymptotic behavior of the function (cf. Carslaw and Jaeger, 1948). For example, if the Laplace transform $\bar{f}(p)$ of a function $f(\tau)$ is given by

$$\bar{f}(p) = \frac{h_1(p)}{h_2(p)} \quad (5.40)$$

where $h_1(p)$ is an analytic function of p and the largest root of $h_2(p)$ is distinct, real and equals q , then for large τ , $f(\tau)$ is given by

$$f(\tau) \sim \frac{(p-q)h_1(q)}{h_2(q)} e^{qt} = \frac{h_1(q)e^{qt}}{\left[\frac{dh_2}{dp} \right]_{p=q}} \quad (5.41)$$

It is easy to see that in the context of the solution of the diffusion equation, the dominant pole of the Laplace transform corresponds to the dominant eigenvalue α_1 of the diffusion operator. In view

of Remark 5.2, at least one dominant eigenvalue must be real. It seems extremely unlikely that the functional expressions (5.11), (5.21) and (5.21a) would have multiple real dominant roots. Uniqueness of the dominant eigenvalue remains to be proved. In the following development it is assumed that α_1 is always real and unique.

As indicated by (5.16) and the discussion following that equation, asymptotically, the first moment distribution $M_1^*(z_1, \tau|y)$ changes exponentially with respect to time, with parameter α_1 . Thus, asymptotically (i.e., as $p \rightarrow \alpha_1$),

$$\bar{M}_1(z, p|y) \sim \frac{F_1(z_1)}{(p - \alpha_1)} \quad \forall y \quad (\text{cf. (5.16)}) \quad (5.42)$$

$$\text{where } F_1(z_1) = \int_0^{z_1} f_1(\xi) d\xi \quad (5.43)$$

and $f_1(z)$ is defined by (5.9). It follows from the asymptotic behavior of $M_1^*(z_1, \tau|y)$ that the asymptotic change in $M_1^*(z_1, \tau|0)M_1^*(z_2, \tau|0)$, $M_1^*(z_1, \tau|0)M_1^*(z_2, \tau|y)$ and $M_1^*(z_1, \tau|y)M_1^*(z_2, \tau|0)$ must be exponential with parameter $2\alpha_1$. The dominant root of denominators of $\bar{\Phi}_{11}$, $\bar{\Phi}_{12}$, and $\bar{\Phi}_{21}$ is therefore equal to $2\alpha_1$. Hence,

$$\begin{aligned} \bar{\Phi}_{11}(z_1, z_2, p|0) &= \mathcal{L}_\tau [M_1^*(z_1, \tau|0)M_1^*(z_2, \tau|0)] \\ &\sim \left\{ \frac{F_1(z_1)F_1(z_2)}{p - 2\alpha_1} \right\} \end{aligned} \quad (5.44)$$

and similarly,

$$\bar{\Phi}_{12}(z_1, z_2, p|0, y) \sim \bar{\Phi}_{21}(z_1, z_2, p|y, 0) \sim \frac{F_1(z_1)F_1(z_2)}{p - 2\alpha_1}. \quad (5.45)$$

Since $M_1^*(z, \tau|y)$ is a bounded function of y for all y and $z \in \bar{D}$, it follows that $\bar{\Phi}_{12}(z_1, z_2, p|0, y)$ and $\bar{\Phi}_{21}(z_1, z_2, p|y, 0)$

are also bounded functions of y , and therefore

$$\Theta_1(z_1, z_2, p|y, 0) = \int \{ \tilde{\Phi}_{12}(z_1, z_2, p|0, y) + \tilde{\Phi}_{21}(z_1, z_2, p|y, 0) \} \exp(-\beta_1 y) dy$$

is bounded for all $y \in \tilde{D}$. Hence, by the dominated convergence theorem (see e.g., Parzen, 1962; p. 274)

$$\begin{aligned} & \lim_{\tau \rightarrow \infty} \{ \mathcal{L}_{\tau}^{-1} [\Theta_1(z_1, z_2, p|y, 0)] \} \\ &= \int \lim_{\tau \rightarrow \infty} \{ \mathcal{L}_{\tau}^{-1} [\tilde{\Phi}_{12}(z_1, z_2, p|0, y) + \tilde{\Phi}_{21}(z_1, z_2, p|y, 0)] \exp(-\beta_1 y) dy \} \quad (5.46) \end{aligned}$$

and therefore,

$$\Theta_1(z_1, z_2, p|y, 0) \sim \frac{-2F_1(z_1)F_1(z_2)\exp(-\beta_1 y)}{(p - 2\alpha_1)\beta_1} \quad (5.47)$$

Similarly,

$$\Theta_2(z_1, z_2, p|y, 0) \sim \frac{-2F_1(z_1)F_1(z_2)\exp(-\beta_2 y)}{(p - 2\alpha_1)\beta_2} \quad (5.48)$$

Substitution of (5.45), (5.47) and (5.48) in (5.39) yields $p = 2\alpha_1$ as the dominant pole of the resulting expression. Hence,

$$\begin{aligned} & M_2^*(z_1, z_2, \tau|y) \\ & \sim \frac{\{ (\lambda^{**} + 2\lambda^*) (\exp[\beta_1^*(\tilde{a} - \tilde{b})] - \exp[\beta_2^*(\tilde{a} - \tilde{b})]) F(z_1) F(z_2) \exp(2\alpha_1 \tau) \}}{(\mu + 2\alpha_1) [\exp(\beta_1^* \tilde{a}) - \exp(\beta_2^* \tilde{a})] - \lambda^* [\exp(\beta_1^* [\tilde{a} - \tilde{b}]) - \exp(\beta_2^* [\tilde{a} - \tilde{b}])]} \quad (5.49) \end{aligned}$$

where β_1^* and β_2^* denote β_1 and β_2 respectively (see Equation (5.26)), evaluated at $p = 2\alpha_1$.

The asymptotic value of the second factorial moment of the number of entities in any given set $\tilde{A} \equiv (\tilde{z}_1, \tilde{z}_2) \subset \tilde{D}$ can be easily obtained from (5.49) by using (4.35). Thus,

$$M_{(2)}(\tilde{A} \times \tilde{A}, \tau | y) \sim \frac{(\lambda^{**} + 2\lambda^*) [F_1(\tilde{z}_2) - F_1(\tilde{z}_1)]^2 [\exp(\beta_1^*(\tilde{a}-\tilde{b})) - \exp(\beta_2^*(\tilde{a}-\tilde{b}))] \exp(2\alpha_1 \tau)}{(2\alpha_1 + \mu) [\exp(\beta_1^*(\tilde{a})) - \exp(\beta_2^*(\tilde{a}))] - \lambda^* [\exp(\beta_1^*(\tilde{a}-\tilde{b})) - \exp(\beta_2^*(\tilde{a}-\tilde{b}))]} \quad (5.50)$$

Similarly, the asymptotic value of the expected number of entities in the set $\tilde{A} \equiv [\tilde{z}_1, \tilde{z}_2]$ is given by

$$M_1(\tilde{A}, \tau | y) \sim [F_1(\tilde{z}_2) - F_1(\tilde{z}_1)] \exp(\alpha_1 \tau) \quad (5.51)$$

The variance of the number of entities in the set \tilde{A} can be readily obtained from the two factorial moments (5.50) and (5.51). The coefficient of variation $CV(\tilde{A})$ of the number of entities in \tilde{A} gives a measure of the stochastic fluctuations in the population $N(\tilde{A})$ of the set \tilde{A} , and is defined as

$$CV(\tilde{A}) = [\text{Var}(N(\tilde{A}))]^{1/2} / E(N(\tilde{A})),$$

where $\text{Var}(N(\tilde{A}))$ is the variance of the number of entities $N(\tilde{A})$.

Thus, using (5.50) and (5.51) the asymptotic value of the coefficient of variation can be seen to be

$$CV(\tilde{A}) \sim \left\{ \frac{(2\alpha_1 + \mu) [\exp(\beta_2^*(\tilde{a})) - \exp(\beta_1^*(\tilde{a}))] + (\lambda^{**} + 3\lambda^*) \{ \exp(\beta_1^*(\tilde{a}-\tilde{b})) - \exp(\beta_2^*(\tilde{a}-\tilde{b})) \}}{(2\alpha_1 + \mu) [\exp(\beta_1^*(\tilde{a})) - \exp(\beta_2^*(\tilde{a}))] - \lambda^* \{ \exp(\beta_1^*(\tilde{a}-\tilde{b})) - \exp(\beta_2^*(\tilde{a}-\tilde{b})) \}} \right\}^{1/2} \quad (5.52)$$

When the initial number of ancestors is k , located at y_1, y_2, \dots, y_k ;

from (5.33) one obtains

$$M_{(2)}(\tilde{A} \times \tilde{A}, \tau | y_1, \dots, y_k) \sim k M_{(2)}(\tilde{A} \times \tilde{A}, \omega | y) + k(k-1) [M_1(\tilde{A}, \omega | y)]^2 \quad (5.53)$$



where $M_{(2)}(\tilde{A} \times \tilde{A}, \omega | y)$ and $M_1(\tilde{A}, \omega | y)$ are given by the expressions on the right side of (5.50) and (5.51) respectively. Since

$$M_1(\tilde{A}, \tau | y_1, y_2, \dots, y_k) \sim k M_1(\tilde{A}, \omega | y), \quad (5.54)$$

it follows from (5.53) and (5.54) that the asymptotic value of the variance of the population in any set $A \subset \tilde{D}$ with k ancestors is k times the corresponding value of the variance with a single ancestor. Hence, when there are k ancestors initially,

$CV(A)$

$$\sim \left\{ \frac{(2\alpha_1 + \mu) [\exp(\beta_2^* \tilde{a}) - \exp(\beta_1^* \tilde{a})] + (\lambda^{**} + 3\lambda^*) \{ \exp(\beta_1^* [\tilde{a} - \tilde{b}]) - \exp(\beta_2^* [\tilde{a} - \tilde{b}]) \}}{k(2\alpha_1 + \mu) [\exp(\beta_1^* \tilde{a}) - \exp(\beta_2^* \tilde{a})] - \lambda^* \{ \exp(\beta_1^* [\tilde{a} - \tilde{b}]) - \exp(\beta_2^* [\tilde{a} - \tilde{b}]) \}} \right\}^{\frac{1}{2}}. \quad (5.55)$$

When $\tilde{b} = 0$, (5.55) assumes a particularly simple form:

$$CV(A) \sim \left\{ \frac{\lambda^{**} + \lambda^* + \mu}{k(\lambda^* - \mu)} \right\}^{\frac{1}{2}}. \quad (5.56)$$

Case 2. Finite Spatial Domain $[0, \tilde{a}]$ with an Absorbing Barrier at \tilde{a} .

The evaluation of the second factorial moment distribution as well as the variance and the coefficient of variation of the number of entities in any given set $A \subset \tilde{D}$ proceeds exactly as in Case 1. The diffusion equation and the boundary conditions are the same except (5.34f) and (5.35f), which must be replaced by

$$M_2^*(z_1, z_2, \tau | a) = 0 \quad (5.34g)$$

and

$$\bar{M}_2(z_1, z_2, p | a) = 0 \quad (5.35h)$$

respectively. The asymptotic solution for $M_{(2)}(\tilde{A} \times \tilde{A}, \tau | y)$, where

$\tilde{A} \equiv (\tilde{z}_1, \tilde{z}_2)$ for this case is given by

$$M_{(2)}(\tilde{A} \times \tilde{A}, \tau | y) \sim \frac{\left\{ [F_1(\tilde{z}_2) - F_1(\tilde{z}_1)]^2 \exp(2\alpha_1 \tau) \{ (\lambda^{**} + 2\lambda^*) [\beta_2^* \exp[\beta_1^* (\tilde{a} - \tilde{b})]] - \beta_1^* \exp[\beta_2^* (\tilde{a} - \tilde{b})] - \beta_2^* + \beta_1^* \} \right\}}{\left\{ (2\alpha_1 + \mu) [\beta_2^* \exp(\beta_1^* \tilde{a}) - \beta_1^* \exp(\beta_2^* \tilde{a})] - \lambda^* [\beta_2^* \exp[\beta_1^* (\tilde{a} - \tilde{b})]] - \beta_1^* \exp[\beta_2^* (\tilde{a} - \tilde{b})] + \lambda^* (\beta_2^* - \beta_1^*) \right\}} \quad (5.57)$$

where α_1 is the dominant eigenvalue given by the root of (5.21).

The asymptotic value of the coefficient of variation of the population in any set $A \subset \tilde{D}$ with k initial ancestors is

$$CV(A) \sim \left\{ \frac{(\lambda^{**} + 3\lambda^*) [\beta_2^* \exp(\beta_1^* \tilde{a} - \tilde{b})] - \beta_1^* \exp(\beta_2^* \tilde{a} - \tilde{b}) - \beta_2^* + \beta_1^*}{k \{ (2\alpha_1 + \mu) [\beta_2^* \exp(\beta_1^* \tilde{a}) - \beta_1^* \exp(\beta_2^* \tilde{a})] - \lambda^* [\beta_2^* \exp[\beta_1^* (\tilde{a} - \tilde{b})]] - \beta_1^* \exp[\beta_2^* (\tilde{a} - \tilde{b})] + \lambda^* (\beta_2^* - \beta_1^*) \}} \right\}^{\frac{1}{2}} \quad (5.58)$$

Case 3. Semi-Infinite Domain $[0, \infty)$.

Proceeding exactly as in Case 1 or Case 2, the following expression is obtained for the Laplace transform of $M_2^*(z_1, z_2, \tau | 0)$:

$$\bar{M}_2(z_1, z_2, p | 0) = \frac{\lambda^{**} \Phi_1(z_1, z_2, p | 0) - \frac{2\lambda^*}{\sigma^2 \beta_2} \Theta_1(z_1, z_2, p | \tilde{b}, 0) \exp(\beta_1 \tilde{b}) (\mu + p)}{(\mu + p) \exp(\beta_1 \tilde{b}) - \lambda^*} \quad (5.59)$$

Noting that α_1 is the dominant eigenvalue given by the root of (5.21a), as before,

$$\Phi_1(z_1, z_2, \tau | 0) \sim \frac{F_1(z_1) F_1(z_2)}{p - 2\alpha_1}.$$

Also, noting that

$$\Theta_1(z_1, z_2, p | \tilde{b}, 0) = \left\{ \int [\tilde{\Phi}_{12}(z_1, z_2, p | 0, y) + \tilde{\Phi}_{21}(z_1, z_2, p | y, 0)] \exp(-\beta_1 y) dy \right\}_{y=\tilde{b}}$$

where $\beta_1 \geq 0$, so that $\exp(-\beta_1 y)$ is bounded on $[0, \infty)$, similar argument as before can be used to show that

$$\Theta_1(z_1, z_2, p | \tilde{b}, 0) \sim \frac{-2F_1(z_1)F_1(z_2)\exp(-\beta_1 \tilde{b})}{(p - 2\alpha_1)\beta_1}.$$

Hence, for the set $\tilde{A} \equiv (\tilde{z}_1, \tilde{z}_2) \subset \tilde{D}$ with one ancestor initially,

$$M_{(2)}(\tilde{A} | \tilde{A}, \tau | y) \sim \frac{[F_1(\tilde{z}_2) - F_1(\tilde{z}_1)]^2 (\lambda^{**} + 2\lambda^*) \exp(2\alpha_1 \tau)}{(\mu + 2\alpha_1) \exp(\beta_1^* \tilde{b}) - \lambda^*} \quad (5.60)$$

and when there are k ancestors initially, the coefficient of variation of the population in any set A in \tilde{D} is given by

$$CV(A, \tau | y) \sim \left\{ \frac{(\lambda^{**} + 3\lambda^*) - (\mu + 2\alpha_1) \exp(\beta_1^* \tilde{b})}{k[(\mu + 2\alpha_1) \exp(\beta_1^* \tilde{b}) - \lambda^*]} \right\}^{\frac{1}{2}}. \quad (5.61)$$

At this point it is worthwhile to summarize the important analytical results obtained thus far regarding the solutions of the diffusion equations for the first and second factorial moment distributions with parameters given by (5.1)

- (i) A complete analytical solution (5.32) has been obtained for the first moment distribution for the case where $\tilde{a} = \infty$, $\tilde{b} = 0$ and at time 0 there is one ancestor at $y = 0$.
- (ii) General solutions for the first moment density in the form of generalized Fourier series of eigenfunctions (assuming that the eigenfunctions form a complete set)

have been obtained for the case of a finite interval $[0, \bar{a}]$ with absorbing and reflecting barrier at \bar{a} .

- (iii) It has been shown that asymptotically the expected number of entities in any set $A \subset \bar{D}$ changes exponentially with a parameter α_1 (cf. (5.16)).

Here the parameter α_1 is given by the dominant root of (5.11), (5.21) and (5.21a) for the cases of a reflecting and an absorbing barrier at a finite \bar{a} and a semi-infinite spatial domain $[0, \infty)$ respectively. Similarly, the asymptotic change in the second factorial moment of the number of entities in any subset of \bar{D} is exponential with respect to time, with parameter $2\alpha_1$ (cf. (5.50), (5.57) and (5.60)).

- (iv) The asymptotic behavior of the coefficient of variation with a given number k of ancestors at time 0, as given by (5.55), (5.58) and (5.61) constitutes a result of particular importance regarding the limiting behavior of the population: The coefficient of variation of the number of entities in any set $A \subset \bar{D}$ attains the same (constant) value depending upon the boundary conditions and the initial number of ancestors. Furthermore, the asymptotic value of the coefficient of variation is inversely proportional to the square root of the initial number of ancestors.

Of course, as indicated before, these results are based upon the hypothesis that α_1 is real and distinct.

In view of the result regarding asymptotic behavior of the coefficient of variation it can be said that if the initial population is small, the stochastic fluctuations in the population of any subset of \tilde{D} will always be significant even if eventually the population becomes very large. On the other hand, when the initial population is large, the stochastic fluctuations will be small and the equations for the first moment alone may be adequate to describe the dynamics of the population. For example, if in the study of the secondary nucleation process in a crystallizer the experiment is started with a small number of crystals, the number of crystals in any size range will always tend to exhibit relatively large fluctuations. This will lead to an appreciable amount of scatter in the data on the induction periods as observed by Kane (1971). Similarly, the results indicate that a biological population started with a small number of ancestors will always tend to exhibit relatively large stochastic fluctuations.

5.4 The Problem of "Critical Length"

Referring to (5.11) and (5.21) it can be seen that when the dominant eigenvalue α_1 is real, for a given set of values of the parameters σ^2 , r , λ^* and \tilde{b} , there exists a critical value \tilde{a}_{cr} of \tilde{a} , such that if $\tilde{a} < \tilde{a}_{cr}$, $\alpha_1 < 0$ and the expected number of entities in the population of any subset of \tilde{D} will always decrease with time. The problem of determining this critical location of the boundary (the "critical length" of \tilde{D}) has important implications in case of biological populations. For example, in an intensively exploited fishery or a forest resource, few individuals

in the population can be expected to survive beyond a certain maturity \tilde{a} . In this context, y and z correspond to the maturity of a tree (measured in terms of its height or productive value) or the length or weight of a fish at times 0 and τ respectively, λ^* denotes the expected number of births per unit time, and μ refers to the death rate due to natural causes as well as harvesting. Of course, it is assumed here that the growth of an individual can be characterized by a diffusion process. In these cases it is reasonable to construct a mathematical model with an absorbing barrier placed at \tilde{a} . As the harvesting pressure increases, the maturity corresponding to the absorbing barrier decreases. If the absorbing barrier corresponds to too short a life-span, the recruitment due to reproduction would be too low to have a self-sustaining population. Evaluation of the maturity corresponding to this critical life-length would be of a great help in determining the necessary management and control policies such as the legal size-limits on fish so as to avoid extinction of the species.

An opposite situation may occur in a continuous crystallizer where the crystals from the product stream are classified and all the crystals below a certain size \tilde{a} are returned to the feed stream. In the context of a crystallizer, y and z denote the size of a crystal (usually expressed in terms of a characteristic length of the crystal), λ^* characterizes the expected number of nuclei generated due to secondary nucleation and μ refers to the rate of removal of crystals from the processing system. As in case of the biological populations mentioned above, the crystals are assumed to grow according to a diffusion process. For a stable operation of the

system the number of crystals in the crystallizer must remain constant. To achieve this it may be necessary to regulate the maximum size \bar{a} of the crystals returned as the feedback. Another way of achieving the control would be to fix \bar{a} and regulate the fraction of the total number of crystals of size less than or equal to \bar{a} to be returned. Thus, the problem reduces to that of determining the critical size for a fixed "death rate" or that of evaluating the "death rate" to make a given size \bar{a} a critical one so as to have a stable operation of the crystallizer. In view of these possible applications, the problem of "critical length" was solved for the case of an absorbing barrier with parameters given by (5.1).

In the case of a reflecting barrier at \bar{a} with $\bar{b} = 0$, the criticality of the change in the population will be dependent upon λ^* and μ alone. However, when $b > 0$, for given values of λ^* and μ there may be a critical location \bar{a}_{cr} of the boundary such that for $\bar{a} < \bar{a}_{cr}$, the population will (asymptotically) decrease monotonically with time. The computation of \bar{a}_{cr} in this case will be similar to that given below for an absorbing barrier.

Asymptotically, the expected number of entities in the population changes exponentially with a parameter α_1 (i.e., the largest eigenvalue of the diffusion operator), which is given by the dominant root of (5.21) for the case of an absorbing barrier at \bar{a} . It can be seen from (5.21) that α_1 is a continuous function of the location \bar{a} of the barrier. The critical length is thus given by setting $\alpha = 0$ in (5.21) and solving the resulting expression for \bar{a} , subject to the condition that the root should be greater than or



equal to \tilde{b} . A solution $\tilde{a}_{cr} < \tilde{b}$ would represent a physically impossible situation in view of the fact that no entity in $[0, b)$ is capable of reproduction (cf. (5.1)).

Since (5.21) is nonlinear in \tilde{a} , the solution has to be obtained numerically. Equation (5.21) can also be written as

$$\begin{aligned}
 (\alpha + \mu)[\beta_1 \exp(\beta_2 \tilde{a}) - \beta_2 \exp(\beta_1 \tilde{a})] - \lambda^* [\beta_1 \exp[\beta_2 (\tilde{a} - \tilde{b})] \\
 - \beta_2 \exp[\beta_1 (\tilde{a} - \tilde{b})]] - \lambda^* (\beta_2 - \beta_1) = 0, \quad (5.62)
 \end{aligned}$$

where

$$\beta_1 = \frac{-r + \sqrt{r^2 + 2\sigma^2(\mu + \alpha)}}{\sigma^2} \quad \text{and} \quad \beta_2 = \frac{-r - \sqrt{r^2 + 2\sigma^2(\mu + \alpha)}}{\sigma^2}.$$

Equation (5.62) has to be solved for \tilde{a} after setting $\alpha_1 = 0$ to give \tilde{a}_{cr} . If the solution to the problem of critical length is to be obtained in the form of a set of families of curves, the parameters in the graphical solution will be σ^2 , r , λ^* , μ and \tilde{b} . The total amount of computation required to evaluate the solution over a range of values of all these parameters can be substantially reduced if the diffusion equation for the first moment density is written in terms of the following dimensionless variables:

$$\begin{aligned}
 \hat{z} &= z/\tilde{a} & \hat{\tau} &= r\tau/\tilde{a} \\
 \hat{p} &= 2r\tilde{a}/\sigma^2 & \hat{b} &= \tilde{b}/\tilde{a} \\
 \hat{y} &= y/\tilde{a} & \hat{\mu} &= \mu\tilde{a}/r \\
 \hat{\lambda} &= \lambda^*\tilde{a}/r
 \end{aligned} \quad (5.63)$$

and an expression analogous to (5.62) is derived in terms of these dimensionless variables. This simplification reduces the problem of evaluation of the critical length to that of solving the equation

$$\begin{aligned}
& (\hat{\mu}/\hat{\lambda})[v_1\hat{P}\exp(v_2\hat{P}) - v_2\hat{P}\exp(v_1\hat{P})] + v_2\hat{P}\exp[v_1(\hat{P}-\hat{P}\hat{b})] \\
& - v_1\hat{P}\exp[v_2(\hat{P}-\hat{P}\hat{b})] + v_1\hat{P} - v_2\hat{P} = 0
\end{aligned} \tag{5.64}$$

for \hat{P} . In (5.64)

$$v_1 = \frac{-1 + \sqrt{1 + 4\hat{\mu}/\hat{P}}}{2}$$

and

$$v_2 = \frac{-1 - \sqrt{1 + 4\hat{\mu}/\hat{P}}}{2}.$$

The computer program used to solve (5.64) is given in Appendix B. It should be noted that the quantities $\hat{\mu}/\hat{\lambda}$, v_1 , v_2 and $\hat{P}\hat{b}$ are all independent of \tilde{a} . Figure 5.1 gives the solution in the form of a family of curves for \hat{P} corresponding to the critical value of \tilde{a} , plotted against $\hat{\mu}/\hat{P}$, with $\hat{P}\hat{b}$ and $\hat{\lambda}/\hat{\mu}$ as parameters. Given the parameters $\frac{\sigma^2}{2}$, r , μ , λ^* , and \tilde{b} , the graph can be used to evaluate \tilde{a}_{cr} as follows: First, the dimensionless parameters $\hat{\lambda}/\hat{\mu}$ and $\hat{P}\hat{b}$ are computed and the curve corresponding to these values of the parameters is located on the graph. The value of $\hat{\mu}/\hat{P}$ is calculated and the corresponding value of \hat{P} is obtained from this curve. Knowing the critical value of \hat{P} as well as $\frac{\sigma^2}{2}$ and r , \tilde{a}_{cr} can be readily computed.

As may be expected, the curves indicate that the value of \tilde{a}_{cr} is less sensitive to the changes in the death rate when the ratio of birth rate to the death rate $\hat{\lambda}/\hat{\mu}$ is high or when the value of $\hat{P}\hat{b}$ is low, which would occur if \tilde{b} is close to zero or if the mean rate of movement r toward the absorbing barrier is small as compared to the diffusion coefficient.



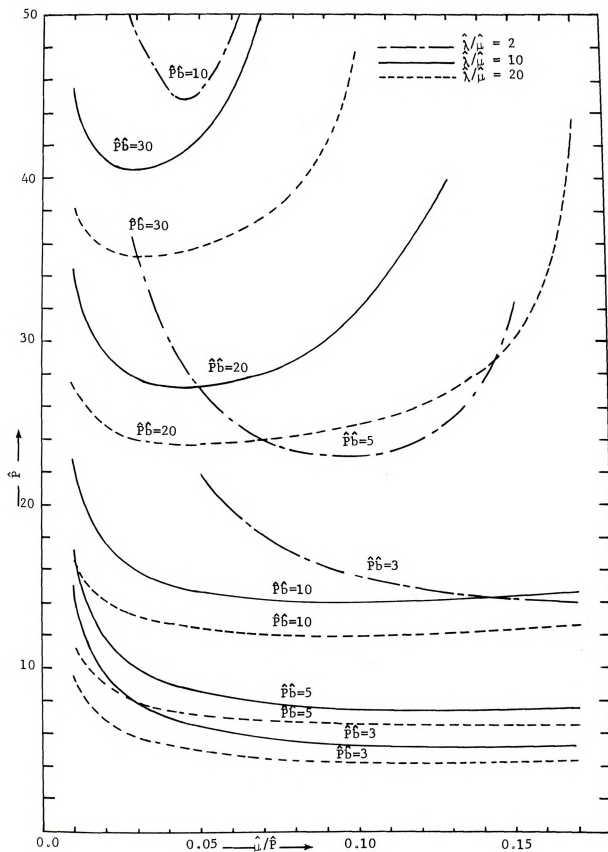


Figure 5.1 Graphical solution of Equation (5.64) for "Critical Length"

Although the main purpose in presenting Figure 5.1 is to demonstrate the method, care was taken to have a realistic range of values of the parameters. The data of Cooper and Latta (1954) and Cooper, Latta and Schafer (1956) on populations of bluegills Lepomis macrochirus was used to obtain crude estimates of the diffusion and drift coefficients and the death rate. Assuming a constant death rate estimates of the yearly total mortalities were obtained by Cooper et. al., which were used to calculate the death parameter μ . The data consisted of the size (length) distribution of the various age classes of the fish for several years. The age classes were characterized according to the age of the fish in years. Since spawning occurs during a relatively short period of time during each year, it was assumed that the size distribution of the fish in each age group had resulted from a large number of identical young ones born at the same instant of time. If the temporal changes in the expected size distribution of the fish can be characterized by a diffusion equation with constant parameters, the size frequency of the fish grown from a cohort should be approximately normal at all times if the effect of the reflecting boundary corresponding to zero size can be neglected. This may be the case when the drift coefficient is much larger than the diffusion coefficient. Under these conditions the changes in the mean and variance of the size frequency distribution during a time interval Δt will be $r\Delta t$ and $\sigma^2\Delta t$ respectively, and thus the drift coefficient r and the diffusion coefficient $\frac{\sigma^2}{2}$ can be readily estimated. The assumption that the size distribution was normal at all times was indeed rather crude, and since the purpose



here is mainly to demonstrate the technique, no further refinements in the estimates of the parameters were attempted.

As discussed by Cooper et. al., horizontal as well as vertical estimates of the parameters were obtained whenever possible. A horizontal estimate is obtained by following the same age cohort year after year, e.g., by considering the two-year old fish in 1954 and the three-year old fish in 1955, assuming that the population estimates for the successive years were obtained with the same accuracy. A vertical estimate uses the different age groups counted during the same season, based on the assumption that the population is in a steady state as regards the yearly recruitment, mortality and growth patterns. The horizontal and vertical estimates were comparable to each other. Different estimates of the parameters are summarized in Table 5.1. It was assumed that the fish were capable of reproduction after they reached a length of 2.5 inches. The values of birth rates were chosen to be simple multiples of the death rates. Based on these estimates the dimensionless parameters for the different curves were chosen to include the ranges of parameters represented in Table 5.1.

Since all the parameters in Figure 5.1 are dimensionless, the same curves can be used to calculate the critical value for any population with any arbitrary units (such as productivity, biomass, etc.) for \bar{a} . However, it must be remembered that in many cases the constant values for σ^2 , r , and λ^* are only an approximation to the real situation and due caution should be exercised in using such charts for practical purpose. For example, it can be seen from Table 5.1 that $\frac{\sigma^2}{2}$ and r are functions of the size of the fish.

TABLE 5.1: Model Parameters for the Dynamics of Populations of
Bluegills (Lepomis macrochirus)

A. Mean Growth Rate r in Inches/Year (Vertical Estimate)

Age Class (years)						Remarks
1	2	3	4	5	6	
2.5	2.2	1.0	0.8	1.0	0.6	Sugarloaf lake, Michigan 1954
2.2	2.0	1.2	1.0	0.9	0.6	Sugarloaf lake, Michigan 1955
2.5	2.3	1.43	1.35	0.78	0.26	Whitmore lake, Michigan 1955
2.4	2.16	1.21	1.05	0.89	0.47	Average of the estimates above

B. Diffusion Coefficient $\sigma^2/2$ in (inches)²/year^{**}

First two years	Third year	Fourth year	Remarks
0.0417	0.0245	0.024	Vertical estimate (1954)
---	0.0201	---	Horizontal estimate (1954-55)

** for Sugarloaf lake, Michigan

C. Average Death Rate μ in (years)⁻¹

Vertical estimate	Horizontal estimate	Remarks
1.11	0.655	Sugarloaf lake, Michigan, 1952, 1954 and 1955.
2.2	1.56	Whitmore lake, Michigan, 1953-1956.
0.892	---	Fine lake, Michigan, 1955
0.415	---	Fife lake, Michigan, 1956

Similarly, many fish spawn only once in a year giving rise to a "pulse" of young ones. In contrast, the model defines $\lambda_1 \delta\tau$ as the probability of producing i offspring in the time interval $(\tau, +\delta\tau)$ irrespective of whether the individual had reproduced at any time in the interval $(-\infty, \tau)$.

5.5 Solution of the Diffusion Equation for the PGF -- Computation of Extinction Probabilities

As mentioned in Remark 4.2, $G(\theta, \tau | y)$ gives the probability that there are no entities in the population, given that there was one ancestor initially at y , if the arbitrary function $\theta(z)$ is taken to be identically zero. To obtain this the backward diffusion equation for the PGF was solved numerically using finite differences. As in the previous section, the main purpose of these simulations was to study the nature of the solution for some simple cases and hence constant values were chosen for $\frac{\sigma^2}{2}$, r and μ from the range of values covered in Table 5.1. Moreover, the probability of having more than one birth in a short time interval $\delta\tau$ was assumed to be negligible, i.e., only λ_1 was assumed to be a significant parameter. Furthermore, $\lambda_1(y, s)$ was assumed to be independent of time. Equation (4.2) was written in the finite difference form as

$$\begin{aligned} \frac{G(\theta, \tau + \delta\tau | y) - G(\theta, \tau | y)}{\delta\tau} = & \frac{\sigma^2}{2} \left[\frac{G(\theta, \tau | y + \delta y) - 2G(\theta, \tau | y) + G(\theta, \tau | y - \delta y)}{\delta y^2} \right] \\ & + r \left[\frac{G(\theta, \tau | y + \delta y) - G(\theta, \tau | y)}{\delta y} \right] + \mu [1 - G(\theta, \tau | y)] \\ & - \lambda_1(y) G(\theta, \tau | y) [1 - G(\theta, \tau | 0)] \end{aligned} \quad (5.65)$$

for $\delta y \leq y \leq \tilde{a} - \delta y$ and $\tau \geq 0$. For $y = 0$ or \tilde{a} , (5.65) has to be modified to include the boundary conditions. For the stability of the numerical computation scheme a ratio of $\frac{\delta y^2}{\delta \tau} = 4$ was chosen. Different time and space increments were tried and values of δt and δy were chosen so as to give results agreeing up to the first three decimal places with those for a much finer discretization. The values of the simulation parameters are summarized in Table 5.2. The results of simulations 1 thru 7 are presented in Tables 5.3 thru 5.9 respectively. A computer program used for the simulations is presented in Appendix C.

It can be seen that in case of a reflecting barrier at \tilde{a} and constant birth and death rates per individual throughout its

TABLE 5.2: Simulation Parameters for the Computation of Extinction Probability for a Population

Simulation No.	Death Rate μ (years) ⁻¹	Birth Rate λ_1 (years) ⁻¹	Nature of Boundary at \tilde{a}	Remarks
1	0.6	0.6	absorbing	For all simulations $\sigma^2/2 = 0.025 \text{ inch}^2/\text{year}$ $r = 1 \text{ inch/year}$ $\delta x = 0.2 \text{ inch}$ $\delta \tau = 0.01 \text{ year}$ $\tilde{a} = 20 \text{ inches}$
2	0.6	0.6	reflecting	
3	0.6	1.2	absorbing	
4	0.6	20sech(y-10)	reflecting	
5	0.6	50sech(y-10)	absorbing	
6	0.6	50sech(y-8)	absorbing	
7	0.6	50sech(y-6)	absorbing	

life-span, the changes in the total population will be the same as that for a linear birth and death process. The extinction probability for the linear birth and death process with one initial ancestor is equal to

$$\frac{\mu - \mu \exp[-(\lambda_1 - \mu)\tau]}{\lambda_1 - \mu \exp[-(\lambda_1 - \mu)\tau]}$$

when $\lambda_1 \neq \mu$ and $\lambda_1 \tau / (1 + \lambda_1 \tau)$ when $\lambda_1 = \mu$; where λ_1 , μ and τ denote the corresponding quantities defined in this section (Cox and Miller, 1965; p. 166). This result for a linear birth and death process checks with that of Simulation 2 for all initial states of the ancestor. Comparing the results of Simulations 1 and 2 it can be seen that when the initial location of the ancestor is sufficiently away from the absorbing boundary (note that this distance from the boundary will be dependent upon $\frac{g^2}{2}$, r and μ) the extinction probability at any time is the same as that for a reflecting boundary. Thus, the effect of the boundary is apparent only near the boundary itself.

Simulations 4 thru 7 were conducted to represent the situation where each individual has a very high reproduction rate for a small portion of its life-span. In Simulations 4 and 5 the peak reproduction was assumed to be at the midpoint of an individual's life-span. The results of both simulations show that although the extinction probability immediately following the introduction of the ancestor in its peak reproductive state is quite low, it increases to a high value as time progresses. The results of Simulations 6 and 7 are more interesting. Simulation 6 shows that if the ancestor is "too young" or "too old" at the time of its introduction, the probability of extinction of the population is very high. In Simulation 7 a steady state was reached after a period of about ten years. Similar steady state solutions can be obtained for other simulations if the period of simulation is long enough. In general, it can be said from Simulations 4 thru 7 that for a population with constant diffusion and drift coefficients, a constant death rate

and the reproduction rate per individual having a functional form similar to that used in Simulations 4 thru 7, the extinction probability at any given time is the lowest if the ancestor is in its peak reproductive state at the time of introduction. As mentioned in Section 4.5, an important application of the extinction probability is in the biological control of pest species. Simulations 4 thru 7 indicate that the maturity of the parasites or predators introduced for achieving the control may be a crucial factor in the survival of the controlling population and hence the chance of success of the biological control strategy.

When the parameters in the diffusion equation are not constant the method described in the previous section cannot be used to find the "critical length". In such cases computation of the extinction probability can be used to find the critical length by using the criterion that when $\tilde{a} < \tilde{a}_{cr}$, the probability of ultimate extinction is one.

Remark 5.3: Application of the results of this chapter to biological populations involve description of the growth of an individual by a diffusion process. The individuals in the population are often characterized by their maturity or a measure of their size such as a characteristic length, which are essentially non-decreasing quantities for each individual in the population. On the other hand, realizations of a diffusion process are not strictly monotone in nature. Thus, for example, when the growth of a fish in terms of its length is represented by a diffusion process, there will be a nonzero probability that the length of the fish as described by the model will decrease sometime during its lifetime. In this

sense the diffusion representation is essentially approximate in nature, and it is instructive to know the probability of a given decrease in the length, maturity, etc. of an individual if a diffusion model is used. Brockwell (1972b) has derived an expression for the probability of such deviations from monotonicity for a diffusion process on the real line $(-\infty, \infty)$. In particular, let the growth of an individual in the maturity interval $[0, \tilde{a}]$ be characterized by a diffusion process on the real line with constant diffusion and drift coefficients $\frac{\sigma^2}{2}$ and r respectively, the death parameter $\mu = 0$ and let z denote the location of the individual in $[0, \tilde{a}]$ at time t . If T_a denotes the time taken by the individual to reach \tilde{a} for the first time (i.e., the first passage time) and $M(t) = \max_{0 \leq \tau \leq t} z(\tau)$, then $\max_{0 \leq t \leq T_a} (M(t) - z(t))$ represents the maximum decrease in maturity experienced^a by an individual during its lifetime. The distribution of this quantity is given by

$$F_c(\omega) = P\left[\max_{0 \leq t \leq T_a} a^{-1}(M(t) - z(t)) \leq \omega\right] = \exp\left[-c^{-2}(\exp(\omega c^{-2}) - 1)^{-1}\right], \quad (5.66)$$

where $c = \frac{\sigma}{(2\tilde{a}r)^{\frac{1}{2}}}$. Thus, for the example of fish population characterized by the parameters \tilde{a} , r , and $\frac{\sigma^2}{2}$ in Table 5.2, (5.66) predicts that while the probability of a maximum decrease in length of greater than one inch during the lifetime of a fish is close to zero, the probability that a fish will decrease in length by more than one fourth inch during its lifetime is 0.036. This probability changes drastically with small changes in c . It is interesting to note that if the diffusion coefficient in the example just considered was 0.05 instead of 0.025, with the other parameters unchanged, so that the parameter c is 0.05 instead of 0.03536, the probabilities of a decrease in length greater than one inch and one fourth inch would be 8.24×10^{-7} and 0.93 respectively.

Table 5.3. (continued)

Time (years) ↓	Y = Initial Length of the Ancestor (inches)										20.00
	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00	
Birth Rate	0.60	0.60	0.60	0.60	0.60	0.60	0.60	0.60	0.60	0.60	0.60
12.25	0.880	0.880	0.880	0.880	0.881	0.881	0.882	0.885	0.893	0.919	1.000
14.25	0.895	0.895	0.895	0.895	0.896	0.896	0.897	0.899	0.907	0.930	1.000
16.25	0.907	0.907	0.907	0.907	0.907	0.907	0.908	0.910	0.917	0.937	1.000
18.25	0.916	0.916	0.916	0.916	0.917	0.917	0.917	0.919	0.925	0.944	1.000
20.25	0.924	0.924	0.924	0.924	0.924	0.924	0.925	0.927	0.932	0.949	1.000
22.25	0.930	0.930	0.930	0.930	0.930	0.931	0.931	0.933	0.938	0.953	1.000
24.25	0.936	0.936	0.936	0.936	0.936	0.936	0.936	0.938	0.943	0.957	1.000

Table 5.5. Results of Simulation No. 3 -- Extinction Probability for a Population Generated by One Ancestor.

Diffusion coefficient $\sigma^2/2 = 0.025$ (inch)²/year, Mean growth rate $r = 1$ inch/year,
 Death rate $\mu = 0.6$ (year)⁻¹, Birth rate $\lambda_1 = 1.2$ (year)⁻¹,
 Absorbing Barrier at $\bar{a} = 20$ inches.

Time		Y = Initial Length of the Ancestor (inches)											
(years)		0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00	20.00	
Birth Rate	→	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20	
0.00	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	
0.25	0.123	0.123	0.123	0.123	0.123	0.123	0.123	0.123	0.123	0.123	0.123	0.123	
0.50	0.207	0.207	0.207	0.207	0.207	0.207	0.207	0.207	0.207	0.207	0.207	0.207	
0.75	0.267	0.267	0.267	0.267	0.267	0.267	0.267	0.267	0.267	0.267	0.269	0.269	
1.00	0.312	0.312	0.312	0.312	0.312	0.312	0.312	0.312	0.312	0.312	0.322	0.322	
1.25	0.346	0.346	0.346	0.346	0.346	0.346	0.346	0.346	0.346	0.346	0.370	0.370	
1.50	0.373	0.373	0.373	0.373	0.373	0.373	0.373	0.373	0.373	0.373	0.412	0.412	
1.75	0.395	0.395	0.395	0.395	0.395	0.395	0.395	0.395	0.395	0.395	0.446	0.446	
2.25	0.426	0.426	0.426	0.426	0.426	0.426	0.426	0.426	0.426	0.427	0.489	0.489	
2.75	0.447	0.447	0.447	0.447	0.447	0.447	0.447	0.447	0.447	0.450	0.513	0.513	
3.25	0.462	0.462	0.462	0.462	0.462	0.462	0.462	0.462	0.462	0.467	0.527	0.527	
3.75	0.472	0.472	0.472	0.472	0.472	0.472	0.472	0.472	0.472	0.479	0.537	0.537	
4.25	0.480	0.480	0.480	0.480	0.480	0.480	0.480	0.480	0.480	0.487	0.543	0.543	
4.75	0.485	0.485	0.485	0.485	0.485	0.485	0.485	0.485	0.486	0.493	0.548	0.548	
5.25	0.489	0.489	0.489	0.489	0.489	0.489	0.489	0.489	0.490	0.497	0.551	0.551	
5.75	0.494	0.494	0.494	0.494	0.494	0.494	0.494	0.494	0.495	0.502	0.556	0.556	
6.25	0.497	0.497	0.497	0.497	0.497	0.497	0.497	0.497	0.498	0.504	0.558	0.558	
7.25	0.498	0.498	0.498	0.498	0.498	0.498	0.498	0.498	0.499	0.506	0.559	0.559	
8.25	0.498	0.498	0.498	0.498	0.498	0.498	0.498	0.498	0.499	0.506	0.560	0.560	
9.25	0.499	0.499	0.499	0.499	0.499	0.499	0.499	0.499	0.500	0.507	0.560	0.560	
10.25	0.499	0.499	0.499	0.499	0.499	0.499	0.499	0.500	0.500	0.507	0.560	0.560	
11.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.560	0.560	

Table 5.5. (continued)

Time		Y = Initial Length of the Ancestor (inches)									
(years)	↓	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00 20.00
Birth Rate	→	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20	1.20
12.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000
14.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000
16.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000
18.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000
20.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000
22.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000
24.25	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.500	0.501	0.507	0.561 1.000

Table 5.6. Results of Simulation No. 4 -- Extinction Probability for a Population Generated by One Ancestor

Diffusion coefficient $\sigma^2/2 = 0.025$ (inch) ² /year, Death rate $\mu = 0.6$ (year) ⁻¹ , Mean growth rate $r = 1$ inch/year, $\lambda_1 = 1$, Birth rate $\lambda_1 = 20\text{sech}(y-10)$ (year) ⁻¹ , Reflecting Barrier at $\bar{a} = 20$ inches.												
Time (years) ↓	Y = Initial Length of the Ancestor (inches)											
	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00	20.00	
Birth Rate →	0.00	0.01	0.10	0.73	5.32	20.00	5.32	0.73	0.10	0.01	0.00	
0.00	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000
0.25	0.140	0.139	0.138	0.128	0.128	0.081	0.033	0.086	0.130	0.138	0.139	0.140
0.50	0.260	0.259	0.254	0.221	0.104	0.104	0.038	0.123	0.231	0.256	0.259	0.260
0.75	0.363	0.362	0.351	0.288	0.119	0.044	0.044	0.150	0.314	0.356	0.362	0.363
1.00	0.452	0.449	0.432	0.337	0.132	0.050	0.050	0.176	0.385	0.442	0.451	0.452
1.25	0.528	0.525	0.500	0.375	0.146	0.058	0.058	0.204	0.449	0.517	0.527	0.528
1.50	0.594	0.589	0.557	0.406	0.162	0.066	0.066	0.233	0.506	0.581	0.593	0.594
1.75	0.650	0.644	0.604	0.432	0.178	0.076	0.076	0.265	0.558	0.637	0.649	0.651
2.25	0.740	0.731	0.675	0.478	0.212	0.099	0.099	0.333	0.648	0.728	0.740	0.742
2.75	0.807	0.794	0.724	0.520	0.250	0.128	0.128	0.408	0.723	0.796	0.807	0.809
3.25	0.856	0.839	0.759	0.559	0.290	0.164	0.164	0.486	0.784	0.848	0.857	0.858
3.75	0.891	0.871	0.786	0.596	0.332	0.207	0.207	0.563	0.833	0.886	0.894	0.895
4.25	0.917	0.893	0.808	0.631	0.375	0.258	0.258	0.631	0.871	0.915	0.922	0.922
4.75	0.936	0.910	0.828	0.663	0.418	0.314	0.314	0.698	0.901	0.937	0.942	0.943
5.25	0.950	0.921	0.845	0.693	0.461	0.374	0.374	0.753	0.924	0.953	0.957	0.957
5.75	0.966	0.938	0.873	0.745	0.541	0.492	0.492	0.833	0.955	0.974	0.976	0.977
6.25	0.975	0.950	0.896	0.787	0.613	0.590	0.590	0.883	0.972	0.985	0.987	0.987
7.25	0.980	0.959	0.914	0.821	0.673	0.664	0.664	0.913	0.981	0.991	0.993	0.993
8.25	0.984	0.967	0.928	0.849	0.722	0.721	0.721	0.932	0.987	0.995	0.996	0.996
10.25	0.987	0.972	0.940	0.872	0.763	0.766	0.766	0.946	0.991	0.997	0.998	0.998
11.25	0.989	0.977	0.949	0.892	0.798	0.802	0.802	0.956	0.993	0.998	0.999	0.999

Table 5.6. (continued)

Time (years)	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00	20.00
Birth Rate →	0.00	0.01	0.10	0.73	5.32	20.00	5.32	0.73	0.10	0.01	0.00
12.25	0.991	0.981	0.957	0.908	0.827	0.833	0.964	0.994	0.999	0.999	0.999
14.25	0.994	0.986	0.969	0.933	0.873	0.880	0.975	0.996	0.999	1.000	1.000
16.25	0.996	0.990	0.977	0.951	0.907	0.912	0.982	0.997	1.000	1.000	1.000
18.25	0.997	0.993	0.983	0.964	0.931	0.936	0.987	0.998	1.000	1.000	1.000
20.25	0.998	0.995	0.988	0.973	0.949	0.953	0.991	0.999	1.000	1.000	1.000
22.25	0.998	0.996	0.991	0.980	0.962	0.965	0.993	0.999	1.000	1.000	1.000
24.25	0.999	0.997	0.993	0.985	0.972	0.974	0.995	0.999	1.000	1.000	1.000

Table 5.7. (Continued)

Time		Y = Initial Length of the Ancestor (inches)									
(years)		0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	20.00
Birth Rate →	0.00	0.03	0.25	1.83	13.29	50.00	13.29	1.83	0.25	0.03	0.00
12.25	0.985	0.964	0.911	0.790	0.581	0.525	0.867	0.980	0.997	1.000	1.000
14.25	0.988	0.970	0.925	0.820	0.634	0.593	0.893	0.984	0.998	1.000	1.000
16.25	0.990	0.975	0.936	0.843	0.677	0.646	0.911	0.987	0.998	1.000	1.000
18.25	0.992	0.979	0.944	0.862	0.713	0.689	0.925	0.989	0.999	1.000	1.000
20.25	0.993	0.981	0.951	0.878	0.744	0.725	0.936	0.991	0.999	1.000	1.000
22.25	0.994	0.984	0.957	0.891	0.770	0.755	0.944	0.992	0.999	1.000	1.000
24.25	0.995	0.986	0.961	0.902	0.792	0.780	0.950	0.993	0.999	1.000	1.000

Table 5.8. Results of Simulation No. 6 -- Extinction Probability for a Population Generated by One Ancestor.

Diffusion coefficient $\sigma^2/2 = 0.025$ (inch)²/year, Mean growth rate $r = 1$ inch/year, -1 ,
 Death rate $\mu = 0.6$ (year)⁻¹, Birth rate $\lambda_1 = 50 \text{ sech}(y-8)$ (year)⁻¹,
 Absorbing Barrier at $\bar{a} = 20$ inches.

Time (years)	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00	20.00
Birth Rate →	0.03	0.25	1.83	13.29	50.00	13.29	1.83	0.25	0.03	0.00	0.00
0.00	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000
0.25	0.139	0.136	0.114	0.044	0.014	0.049	0.117	0.136	0.139	0.140	1.000
0.50	0.258	0.245	0.177	0.051	0.016	0.059	0.195	0.249	0.258	0.260	1.000
0.75	0.359	0.334	0.215	0.058	0.018	0.069	0.256	0.345	0.361	0.367	1.000
1.00	0.445	0.405	0.242	0.065	0.021	0.080	0.309	0.428	0.449	0.474	1.000
1.25	0.518	0.463	0.265	0.074	0.024	0.092	0.358	0.499	0.525	0.586	1.000
1.50	0.580	0.509	0.285	0.082	0.027	0.105	0.404	0.562	0.590	0.696	1.000
1.75	0.633	0.546	0.305	0.091	0.031	0.120	0.448	0.616	0.647	0.792	1.000
2.25	0.715	0.602	0.343	0.111	0.040	0.134	0.530	0.706	0.741	0.919	1.000
2.75	0.772	0.642	0.379	0.132	0.049	0.193	0.604	0.775	0.820	0.974	1.000
3.25	0.812	0.673	0.413	0.152	0.059	0.234	0.667	0.828	0.887	0.992	1.000
3.75	0.839	0.699	0.442	0.170	0.069	0.275	0.719	0.869	0.937	0.998	1.000
4.25	0.859	0.721	0.468	0.187	0.078	0.314	0.761	0.902	0.969	0.999	1.000
4.75	0.873	0.740	0.490	0.202	0.087	0.349	0.794	0.928	0.985	0.999	1.000
5.25	0.884	0.755	0.509	0.215	0.096	0.381	0.820	0.949	0.993	1.000	1.000
6.25	0.901	0.780	0.539	0.238	0.112	0.435	0.858	0.973	0.997	1.000	1.000
7.25	0.913	0.798	0.563	0.258	0.127	0.480	0.884	0.982	0.998	1.000	1.000
8.25	0.921	0.812	0.582	0.275	0.140	0.515	0.900	0.986	0.998	1.000	1.000
9.25	0.928	0.822	0.598	0.290	0.153	0.544	0.910	0.987	0.998	1.000	1.000
10.25	0.932	0.830	0.610	0.301	0.163	0.566	0.917	0.988	0.998	1.000	1.000
11.25	0.936	0.837	0.621	0.312	0.173	0.584	0.922	0.989	0.998	1.000	1.000

Table 5.8. (continued)

Time		Y = Initial Length of the Ancestor (inches)									
(years)		0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	20.00
Birth Rate	→	0.03	0.25	1.83	13.29	50.00	13.29	1.83	0.25	0.03	0.00
	↓										
12.25	0.939	0.842	0.629	0.320	0.181	0.599	0.926	0.989	0.999	1.000	1.000
14.25	0.943	0.850	0.642	0.334	0.195	0.620	0.931	0.990	0.999	1.000	1.000
16.25	0.946	0.856	0.651	0.343	0.205	0.635	0.935	0.991	0.999	1.000	1.000
18.25	0.948	0.859	0.657	0.350	0.213	0.645	0.937	0.991	0.999	1.000	1.000
20.25	0.949	0.862	0.662	0.355	0.218	0.653	0.939	0.991	0.999	1.000	1.000
22.25	0.950	0.864	0.665	0.359	0.222	0.658	0.940	0.992	0.999	1.000	1.000
24.25	0.951	0.865	0.668	0.362	0.225	0.662	0.941	0.992	0.999	1.000	1.000

Table 5.9. Results of Simulation No. 7 -- Extinction Probability for a Population Generated by One Ancestor.

Diffusion coefficient $\sigma^2/2 = 0.025$ (inch) ² /year, Death rate $\mu = 0.6$ (year) ⁻¹		Mean growth rate $r = 1$ inch/year, ⁻¹ Birth rate $\lambda_1 = 50 \text{sech}(y-6)$ (year) ⁻¹ , Absorbing Barrier at $\bar{a} = 20$ inches.		Time Y = Initial Length of the Ancestor (inches)									
(years)	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	18.00	20.00		
Birth Rate →	0.25	1.83	13.29	50.00	13.29	1.83	0.25	0.03	0.00	0.00	0.00		
0.00	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000
0.25	0.136	0.114	0.044	0.014	0.049	0.117	0.136	0.139	0.140	0.140	0.140	1.000	1.000
0.50	0.245	0.177	0.051	0.015	0.059	0.194	0.249	0.238	0.260	0.260	0.260	1.000	1.000
0.75	0.333	0.214	0.057	0.018	0.067	0.254	0.345	0.361	0.363	0.367	0.367	1.000	1.000
1.00	0.403	0.238	0.062	0.020	0.076	0.304	0.427	0.449	0.452	0.474	0.474	1.000	1.000
1.25	0.458	0.257	0.068	0.022	0.084	0.349	0.497	0.524	0.528	0.586	0.586	1.000	1.000
1.50	0.502	0.272	0.073	0.023	0.092	0.389	0.558	0.589	0.594	0.696	0.696	1.000	1.000
1.75	0.535	0.285	0.078	0.025	0.100	0.425	0.611	0.646	0.651	0.793	0.793	1.000	1.000
2.25	0.582	0.307	0.086	0.028	0.114	0.487	0.697	0.736	0.745	0.920	0.920	1.000	1.000
2.75	0.611	0.323	0.091	0.030	0.125	0.536	0.761	0.802	0.824	0.974	0.974	1.000	1.000
3.25	0.629	0.334	0.095	0.032	0.133	0.573	0.809	0.852	0.891	0.993	0.993	1.000	1.000
3.75	0.642	0.343	0.098	0.033	0.140	0.602	0.845	0.890	0.941	0.998	0.998	1.000	1.000
4.25	0.651	0.349	0.101	0.034	0.145	0.624	0.871	0.920	0.971	0.999	0.999	1.000	1.000
4.75	0.658	0.353	0.102	0.034	0.149	0.640	0.891	0.944	0.987	1.000	1.000	1.000	1.000
5.25	0.662	0.356	0.103	0.035	0.152	0.653	0.906	0.962	0.995	1.000	1.000	1.000	1.000
6.25	0.668	0.361	0.105	0.035	0.156	0.669	0.927	0.984	0.998	1.000	1.000	1.000	1.000
7.25	0.672	0.363	0.106	0.036	0.158	0.678	0.939	0.991	0.999	1.000	1.000	1.000	1.000
8.25	0.673	0.364	0.107	0.036	0.160	0.683	0.945	0.992	0.999	1.000	1.000	1.000	1.000
9.25	0.674	0.365	0.107	0.036	0.160	0.686	0.947	0.993	0.999	1.000	1.000	1.000	1.000
10.25	0.675	0.365	0.107	0.036	0.161	0.688	0.947	0.993	0.999	1.000	1.000	1.000	1.000
11.25	0.675	0.366	0.107	0.036	0.161	0.689	0.948	0.993	0.999	1.000	1.000	1.000	1.000

Table 5.9. (continued)

Time		Y = Initial Length of the Ancestor (inches)									
(years)	↓	0.00	2.00	4.00	6.00	8.00	10.00	12.00	14.00	16.00	20.00
Birth Rate →	0.25	1.83	13.29	50.00	13.29	1.83	0.25	0.03	0.00	0.00	0.00
12.25	0.675	0.366	0.107	0.036	0.161	0.689	0.948	0.993	0.999	1.000	1.000
14.25	0.676	0.366	0.107	0.036	0.161	0.688	0.948	0.993	0.999	1.000	1.000
16.25	0.676	0.366	0.107	0.036	0.161	0.688	0.948	0.993	0.999	1.000	1.000
18.25	0.676	0.366	0.107	0.036	0.161	0.688	0.948	0.993	0.999	1.000	1.000
20.25	0.676	0.366	0.107	0.036	0.161	0.688	0.948	0.993	0.999	1.000	1.000
22.25	0.676	0.366	0.107	0.036	0.161	0.688	0.948	0.993	0.999	1.000	1.000
24.25	0.676	0.366	0.107	0.036	0.161	0.688	0.948	0.993	0.999	1.000	1.000

CHAPTER VI

POPULATIONS WITH AN EXTERNAL INPUT

In general, a population will also have an input other than reproduction by its members. In biological populations this external input usually takes place in the form of an input of individuals across the boundaries of a geographical domain due to locomotion in case of mobile species; or due to passive motion along with the carrying medium such as water for aquatic organisms or air in case of flying insects. The examples of external input in case of particulate processes in chemical engineering systems include an input of particles along with the carrying fluid for suspensions, nucleation of bubbles or crystals at the imperfections on the surfaces of a vessel and a homogeneous nucleation of crystals occurring at very high supersaturations of the magma. Formation of free radicals by the initiation reaction in polymerization processes can also be looked upon as an external input of molecules of zero chain length.

Radcliffe (1972) considered the problem of external input (immigration) in the form of a nonhomogeneous Poisson process into a population where the individuals move within an abstract state space according to a Markov transition probability and multiply according to a branching process in which the offspring are produced at the end of life of the parent and are in the same state as the parent. Radcliffe also assumed that all the individuals entering the state

space at any time due to migration are located at the same point in the state space at the instant of their arrival. This description does not apply to many processes which are considered in the present work. In the processes described in the preceding paragraph the entities may be located at any point in the spatial domain at the instant of arrival and in the general process of reproduction described in Chapter IV an offspring is in a state different from its parent.

6.1 Description of the Process

For the analysis of populations with an external input, the individual state space \tilde{E} for the existing (or "live") entities is taken as an n -dimensional Euclidean space. When the entities are removed from \tilde{E} by the process of "death" they are considered as being absorbed into a single point $\xi^* \notin \tilde{E}$. The union $\tilde{E} \equiv \tilde{E} \cup \xi^*$ constitutes the individual state space for the existing and dead entities together. The movement of entities is restricted to the spatial domain $\tilde{D} \equiv \tilde{D} \cup \xi^*$, where \tilde{D} is a closed nonempty set in \tilde{E} with a boundary Γ (cf. Section 4.2). The location of an entity in \tilde{D} is denoted by ξ . Only the entities in \tilde{D} may possibly move to another point in \tilde{D} , or produce new entities by reproduction.

In the context of a chemical process such as crystallization, \tilde{E} may consist of a four-dimensional Euclidean space with three co-ordinates denoting the geographical location of a crystal and the fourth co-ordinate representing the size of the crystal in terms of its characteristic length. The domain \tilde{D} may refer to the processing system itself, such as the interior of a crystallizer (with

the size co-ordinate possibly varying from 0 to ∞). All the crystals that are removed from the system by flow or by mechanical means (i.e., by "death") are considered as being transferred to the absorbing state $\xi^* \notin \tilde{E}$. Similarly, in the case of a biological population \tilde{E} may consist of a four-dimensional Euclidean space of three geographical co-ordinates and a maturity co-ordinate, \tilde{D} may be a given geographical domain of interest and "death" may be interpreted as the combined effect of natural death, emigration and harvesting.

The process is started at time s_0 with k_0 entities in \tilde{D} . At subsequent times s_j , $j = 1, 2, \dots, r$; k_j entities enter \tilde{D} from a source external to the population. Each entity present in \tilde{D} at time s_0 and the entities entering \tilde{D} due to external input at any time $s_j > s_0$ serve as ancestors of the entities generated due to reproduction. The initial locations of these ancestors are denoted by the superscripted variable ξ^0 . The movement of each entity in \tilde{D} is characterized by a Markov transition probability $\chi(A, t | \xi, t_0)$ denoting the probability that an entity at ξ at time t_0 will be in the set $A \subset \tilde{D}$ at time $t \geq t_0$. Moreover, the probabilities of producing any new entities due to reproduction at any time are also assumed to be Markovian in nature, i.e., the probability that an entity will produce any offspring at any time in future depends only on its present state and not the past history. It is also assumed that the input of entities to \tilde{D} , their movement within \tilde{D} , as well as their reproduction proceed independently of other entities in the population. It can be seen that the process described in this section includes the process discussed in Chapter IV as a special case. In the model considered in Chapter IV, \tilde{E} is

the nonnegative real line, the movement of an entity in \tilde{D} is characterized by a continuous Markov process satisfying (4.1) and the Markov probabilities of producing any offspring, or dying (i.e., of moving to ξ^*) are defined in Section 4.1.

For a complete description of the process the rate of input of entities and the points in \tilde{D} where they first appear must also be known. It is assumed that the arrivals of entities into the population due to external input occur according to k simultaneous independent nonhomogeneous Poisson processes, where k is finite. The Poisson inputs are characterized by the following axioms (see e.g., Parzen, 1962; Chapter 4):

- (i) The process has independent increments, i.e., the numbers of entities arriving during disjoint intervals of time are independent of each other.
- (ii) For any time interval (no matter how small) there is a positive probability that some entities will arrive in \tilde{D} , but it is not certain that an arrival will occur.
- (iii) The probability that one entity appears in \tilde{D} during the time interval $(s, s+\delta s)$ from an external source through the i -th Poisson stream, $i = 1, \dots, k$, is $v_i(s)\delta s + o(\delta s)$, and
- (iv) The probability of more than one arrival in \tilde{D} during the time interval $(s, s+\delta s)$ from any of the input streams is $o(\delta s)$.

It can be seen that in view of axiom (iii), the number of entities k_j appearing in \tilde{D} at time s_j must be equal to one for a finite number of Poisson inputs. When axiom (iv) is dropped so that the

probability of more than one arrival during $(s, s+\delta s)$ is $O(\delta s)$, one obtains a generalized Poisson process. The material in this chapter can possibly be generalized to include an input in the form of a generalized Poisson process, but the analysis will be much more complicated.

The location of an entity at the instant of its arrival in \tilde{D} at time s (which is obviously conditional on the fact that an arrival has taken place) from the i -th Poisson input is described by a probability density $w_i(\xi, s)$. The transition probability $\chi(A, t | \xi^0, s)$ can thus be looked upon as the probability that an entity is in the set $A \subset \tilde{D}$ conditional on its first appearance in \tilde{D} at the point ξ^0 at time s .

Remark 6.1: It can be seen that when the number of Poisson streams k is a constant, the net effect of the k Poisson processes with rates $v_i(s)$ and conditional probability densities $w_i(\xi, s)$, $i = 1, \dots, k$, characterizing the location ξ where an entity entering \tilde{D} via the i -th stream appears first, is equivalent to a single Poisson process with rate

$$v(s) = \sum_{i=1}^k v_i(s) \quad (6.1)$$

and a conditional probability density

$$w(\xi, s) = \frac{\sum_{i=1}^k w_i(\xi, s) v_i(s)}{\sum_{i=1}^k v_i(s)} \quad (6.2)$$

for the initial location of an entity at the time of its arrival.

Similarly, when the process has a random number K of independent

Poisson inputs (fixed for each realization of the process) with all rates $v_i(s)$ and the probability densities $w_i(\xi, s)$ of the initial location of an entity being equal to $\bar{v}(s)$ and $\bar{w}(\xi, s)$ respectively, if $K = k$ is given, the input streams can be replaced by a single stream with rate $k\bar{v}(s)$ and a probability density $\bar{w}(\xi, s)$ for the initial location of an entity at the time of its arrival.

To summarize, the process described in this section is a Markov population process with a Poisson input. The main difference between the description here and the process considered by Radcliffe (1972) is that here an entity is assumed to appear at any point in \tilde{D} at the time of its arrival. Moreover, the locations of the entities produced by reproduction have been left unspecified in this description, and therefore can be arbitrary. As mentioned earlier in this chapter, Radcliffe assumed that all the entities are located at a single point in \tilde{D} at the time of their arrival, and that the entities reproduce according to a branching process.

6.2 The Probability Generating Functional of the Process

For a quantitative description of the process, let $N(A, t | \xi_{10}^0, \xi_{20}^0, \dots, \xi_{k0}^0; \xi_{11}^0, \dots, \xi_{k1}^0; \xi_{1r}^0, \dots, \xi_{kr}^0; s_0, \dots, s_r)$ denote the number of entities in the set $A \subset \tilde{D}$ at time t , given that there was one ancestor arriving at each of the points ξ_{ij}^0 at time s_j , $i = 1, \dots, k_j$; $j = 0, 1, \dots, r$. The points ξ_{i0}^0 denote the initial locations of the entities present in the population at the start of the process at time s_0 and ξ_{ij}^0 , $j = 1, \dots, r$, denotes the initial location of the i -th entity arriving into \tilde{D} at time s_j from an external source. Using the conditional counting measure

$N(A, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; \dots; \xi_{1r}^0, \dots, \xi_{k_r r}^0; s_0, \dots, s_r)$, the PGF for the process can be defined as

$$G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; \dots; \xi_{1r}^0, \dots, \xi_{k_r r}^0; s_0, \dots, s_r) = \\ E \left[\exp \int_D \log \theta(\xi) [N(d\xi, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; \dots; \xi_{1r}^0, \dots, \xi_{k_r r}^0; s_0, \dots, s_r)] \right]. \quad (6.3)$$

In view of independence of the entities,

$$N(A, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; \dots; \xi_{1r}^0, \dots, \xi_{k_r r}^0; s_0, \dots, s_r) = \\ \sum_{j=0}^r \sum_{i=1}^{k_j} N(A, t | \xi_{ij}^0; s_j), \quad (6.4)$$

from which it follows that

$$G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; \dots; \xi_{1r}^0, \dots, \xi_{k_r r}^0; s_0, \dots, s_r) = \\ \prod_{j=0}^r \prod_{i=1}^{k_j} G(\theta, t | \xi_{ij}^0; s_j), \quad (6.5)$$

where $G(\theta, t | \xi_{ij}^0, s_j)$ is the PGF of a population starting with one ancestor at ξ_{ij}^0 at time s_j and allowing no external input. Equations (4.7) and (4.6) are special cases of (6.4) and (6.5) respectively when $r = 0$ and $k_0 = k$. The co-ordinates ξ and ξ^0 in (6.3) thru (6.5) correspond to z and y respectively in (4.5) thru (4.7). Note that when the input consists of a Poisson stream, by Axiom (iv) in the previous section $k_i = 1$ for $i = 1, \dots, r$ in (6.3) thru (6.5). For an input stream in the form of a generalized Poisson process, the numbers k_i , $i = 1, \dots, r$ can be greater than one. The conditional PGF (6.5) has to be summed over all possible initial locations ξ^0 of the externally introduced entities at all instants

of arrival s , $s_0 \leq s \leq t$, and all input streams in order to obtain a complete description of the process.

For the case discussed in Remark 6.1 the external input can be characterized by a single Poisson stream of entities. Thus, for the case where k is a constant, the PGF of the process can be written as

$$\begin{aligned} G(\theta, t | \xi_{10}^0, \dots, \xi_{k0}^0; s_0) &= E\{\exp \int_D \log \theta(\xi) [N(d\xi, t | \xi_{10}^0, \dots, \xi_{k0}^0; s_0)] \\ &= E[\int_D \dots \int_D [\exp \int_D \log \theta(\xi) [N(d\xi, t | \xi_{10}^0, \dots, \xi_{k0}^0; \\ &\quad \xi_{11}^0, \xi_{12}^0, \dots, \xi_{1r}^0; s_0, \dots, s_r)] w(\xi_{11}^0, s_1) w(\xi_{12}^0, s_2) \dots \\ &\quad w(\xi_{1r}^0, s_r) d\xi_{11}^0 d\xi_{12}^0 \dots d\xi_{1r}^0] \\ &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}^0; s_0) E\{\prod_{j=1}^r \int_D G(\theta, t | \xi_{ij}^0; s_j) w(\xi_{ij}^0, s_j) d\xi_{ij}^0\}. \end{aligned} \quad (6.6)$$

Equation (6.6) follows from (6.4) and (6.5).

Remark 6.2: Given the number of arrivals r during the interval $[s_0, t]$ due to a single Poisson stream with rate $v(s)$, the instants at which the arrivals occur can be regarded as independently distributed random variables with the probability density $v(s)/V(s_0, t)$ where $s_0 \leq s \leq t$ and

$$V(s_0, t) = \int_{s_0}^t v(s) ds. \quad (6.7)$$

This statement can be proved as follows:

The joint probability that there will be one arrival each during the infinitesimal intervals $(s_1, s_1 + \delta s_1), \dots, (s_r, s_r + \delta s_r) \subset [s_0, t]$ and none during $[s_0, t] - \{(s_1, s_1 + \delta s_1), \dots, (s_r, s_r + \delta s_r)\}$,

where $s_i, i = 1, \dots, r$, constitute an unordered set of distinct points in $[s_0, t]$, is given by

$$\frac{v(s_1)v(s_2)\dots v(s_r)\delta s_1\delta s_2\dots\delta s_r \exp(V(s_0, t))}{r!} \quad (6.8)$$

Hence, the conditional probability that there will be one arrival each during the intervals $(s_1, s_1 + \delta s_1), \dots, (s_r, s_r + \delta s_r)$ given that r arrivals occur during the interval $[s_0, t]$ is

$$\begin{aligned} & \frac{v(s_1)v(s_2)\dots v(s_r)\delta s_1\delta s_2\dots\delta s_r \exp(V(s_0, t))r!}{r! [V(s_0, t)]^r \exp(V(s_0, t))} \\ &= \frac{v(s_1)v(s_2)\dots v(s_r)\delta s_1\delta s_2\dots\delta s_r}{[V(s_0, t)]^r}, \end{aligned}$$

from which it follows that the joint probability density of one arrival each at times $s_1, s_2, \dots, s_r \in [s_0, t]$ given that r arrivals occurred during $[s_0, t]$ is

$$\frac{v(s_1)v(s_2)\dots v(s_r)}{[V(s_0, t)]^r},$$

which is the joint probability density of r independent random variables with densities $v(s_i)/V(s_0, t)$, $i = 1, \dots, r$. (Q.E.D.)

In view of Remark 6.2, averaging (6.6) over s_j conditional on r , the following equation is obtained for the case where the number of Poisson streams k is a constant:

$$\begin{aligned} & G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) \\ &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}^0; s_0) E \left\{ \prod_{j=1}^r \int_{s_0}^t G(\theta, t | \xi_j^0, s_j) w(\xi_j^0, s_j) d\xi_j^0 \frac{v(s_j) ds_j}{V(s_0, t)} \right\} \\ &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}^0; s_0) E \left\{ \left[\frac{X(\theta, s_0, t)}{V(s_0, t)} \right]^r \right\} \quad (6.9) \end{aligned}$$



where

$$X(\theta, s_0, t) = \int_{s_0}^t \int_{\tilde{D}} G(\theta, t | \xi^0, s) w(\xi^0, s) v(s) d\xi^0 ds. \quad (6.10)$$

Noting that r is a Poisson-distributed random variable with a parameter $V(s_0, t)$, averaging $\left\{ \frac{X(\theta, s_0, t)}{V(s_0, t)} \right\}^r$ over r yields

$$\begin{aligned} E \left[\left[\frac{X(\theta, s_0, t)}{V(s_0, t)} \right]^r \right] &= \sum_{r=0}^{\infty} \frac{[V(s_0, t)]^r \exp[-V(s_0, t)] [X(\theta, s_0, t)]^r}{r! [V(s_0, t)]^r} \\ &= \exp[X(\theta, s_0, t) - V(s_0, t)], \end{aligned}$$

so that

$$G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0, s_0) = \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}^0, s_0) \exp[X(\theta, s_0, t) - V(s_0, t)]. \quad (6.11)$$

When $w(\xi^0, s) = \delta(\xi^0 - \xi_0^0)$, i.e., when all the entities arriving in \tilde{D} first appear at ξ_0^0 and when there are no entities in the population at time s_0 (i.e., $k_0 = 0$), the expression for the PGF reduces to

$$G(\theta, t | s_0) = \exp \left\{ \int_{s_0}^t v(s) [G(\theta, t | \xi_0^0; s_0) - 1] ds \right\}, \quad (6.11a)$$

which is essentially the result obtained by Bartlett (1966) for a population of individuals characterized by their age, as well as that of Radcliffe (1972).

When the number of Poisson streams K is a random variable (fixed for each realization of the process) as specified in Remark 6.1, given $K = k$, the analysis is the same as that for a fixed k with $v(s)$ and $w(\xi, s)$ replaced by $k\tilde{v}(s)$ and $\tilde{w}(\xi, s)$ respectively. Thus,



$$G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) = \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}; s_0) E\left\{\left[\frac{\tilde{X}(\theta, s_0, t)}{\tilde{V}(s_0, t)}\right]^r\right\}, \quad (6.12)$$

where

$$\begin{aligned} \tilde{X}(\theta, s_0, t) &= \int_{s_0}^t \int_{\tilde{D}} G(\theta, t | \xi^0; s) \tilde{w}(\xi^0, s) k \tilde{v}(s) d\xi^0 ds \\ &= k \int_{s_0}^t \int_{\tilde{D}} G(\theta, t | \xi^0; s) \tilde{w}(\xi^0, s) \tilde{v}(s) d\xi^0 ds \\ &\equiv k \tilde{\tilde{X}}(\theta, s_0, t) \end{aligned} \quad (6.13)$$

and

$$\begin{aligned} \tilde{V}(s_0, t) &= \int_{s_0}^t k \tilde{v}(s) ds \\ &\equiv k \tilde{\tilde{V}}(s_0, t). \end{aligned} \quad (6.14)$$

Averaging over r yields

$$\begin{aligned} G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}; s_0) E\{\exp[\tilde{X}(\theta, s_0, t) - \tilde{V}(s_0, t)]\} \\ &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}; s_0) E\{\exp[k(\tilde{\tilde{X}}(\theta, s_0, t) - \tilde{\tilde{V}}(s_0, t))]\} \\ &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}; s_0) E\{[\exp(\tilde{\tilde{X}}(\theta, s_0, t) - \tilde{\tilde{V}}(s_0, t))]^k\}. \end{aligned} \quad (6.15)$$

If K is characterized by a probability generating function

$\Psi(u) = E[(u)^k]$, $0 \leq u \leq 1$, then by averaging over k , (6.15) reduces to

$$\begin{aligned} G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) &= \prod_{i=1}^{k_0} G(\theta, t | \xi_{i0}; s_0) \Psi\{\exp[\tilde{\tilde{X}}(\theta, s_0, t) - \tilde{\tilde{V}}(s_0, t)]\}. \end{aligned} \quad (6.16)$$

6.3 Moment Distributions of the Population

In order to obtain the factorial moment distributions of the population it is necessary to apply (2.13) to (6.11). Thus, when the number of Poisson inputs k is fixed, the expected number of entities in the set $A \subset \tilde{D}$ is given by

$$\begin{aligned} M_1(A, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) &= \frac{\partial}{\partial \zeta} G(\eta + \zeta \theta | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) |_{\eta=1, \zeta=0} \\ &= \sum_{i=1}^{k_0} M_1(A, t | \xi_{i0}^0, s_0) \\ &\quad + \left\{ \int_{s_0}^t v(s) \int_{\tilde{D}} M_1(A, t | \xi^0, s) w(\xi^0, s) d\xi^0 ds \right\}, \quad (6.17) \end{aligned}$$

when $\theta = I(A | \xi)$. Thus, the expected number of entities in $A \subset \tilde{D}$ at time t is simply the sum of the expected number due to the entities present in the population at the initial time s_0 and the expected numbers arising from all the entities entering \tilde{D} during the interval (s_0, t) due to all the individual Poisson streams. The second factorial moment of the number of entities in the set

$A_1 \times A_2 \subset \tilde{D} \times \tilde{D}$ is

$$\begin{aligned} M_{(2)}(A_1 \times A_2, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) &= \frac{\partial^2}{\partial \zeta^2} G(\eta + \zeta \theta | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) |_{\eta=1, \zeta=0} \\ &= \sum_{i=1}^{k_0} M_{(2)}(A_1 \times A_2, t | \xi_{i0}^0; s_0) + \int_{s_0}^t v(s) \int_{\tilde{D}} M_{(2)}(A_1 \times A_2, t | \xi^0; s) w(\xi^0, s) d\xi^0 ds \\ &\quad + \sum_{i=1}^{k_0} M_1(A_1, t | \xi_{i0}^0; s_0) \int_{s_0}^t v(s) \int_{\tilde{D}} M_1(A_2, t | \xi^0; s) w(\xi^0, s) d\xi^0 ds \\ &\quad + \sum_{i=1}^{k_0} M_1(A_2, t | \xi_{i0}^0; s_0) \int_{s_0}^t v(s) \int_{\tilde{D}} M_1(A_1, t | \xi^0; s) w(\xi^0, s) d\xi^0 ds \end{aligned}$$



$$\begin{aligned}
& + \sum_{i=1}^{k_0} \sum_{\substack{j=1 \\ j \neq i}}^{k_0} M_1(A_1, t | \xi_{i0}^0; s_0) M_1(A_2, t | \xi_{j0}^0; s_0) \\
& + \prod_{j=1}^2 \left\{ \int_{s_0}^t v(s) \int_{\tilde{D}} M_1(A_j, t | \xi^0; s) w(\xi^0, s) d\xi^0 ds \right\}. \quad (6.18)
\end{aligned}$$

In the context of the problem considered in Chapter IV, where \tilde{E} is the non-negative real line and the movement of entities in \tilde{E} is characterized by a diffusion process; when a fixed number k of Poisson inputs with rates $v_i(t)$ are added, the forward diffusion equation (4.49) for the first moment density of the population has to be modified to

$$\frac{\partial m_1(z, t | N_0, s)}{\partial t} = \mathcal{A}_z^* m_1(z, t | N_0, s) + v(t)w(z, t), \quad (6.19)$$

where \mathcal{A}_z^* is the forward diffusion operator defined in (4.48), and $v(t)$ and $w(z, t)$ are defined by (6.1) and (6.2) respectively. The initial and boundary conditions will remain unchanged. The equation for the second factorial moment density will be more complicated. The backward diffusion equations for the first and second factorial moment distributions of the population as well as the corresponding initial and boundary conditions will remain unchanged. The corresponding solutions $M_1(A, t | \xi^0, s)$ and $M_{(2)}(A_1 \times A_2, t | \xi^0, s)$ will have to be substituted in (6.17) and (6.18) to obtain complete solutions for the moment distributions.

It can be seen that while the diffusion equation for the first moment of the population is quite simple, the equation for the second factorial moment is rather complicated. However, as far as the asymptotic behavior of the coefficient of variation is concerned, it can be expected that normally, the total number of entities



entering \tilde{D} due to an external input in a large finite time interval will be sufficiently large, and since these entities act as ancestors for the entities generated by reproduction, the asymptotic value of the coefficient of variation of the number of entities in any set $A \subset \tilde{D}$ will always tend to be low.

When there are a random number K of Poisson inputs as defined in Remark 6.1, given that $K = k$, the conditional first and second factorial moment distributions are given by (6.17) and (6.18) respectively by changing $v(s)$ and $w(\xi, s)$ to $k\tilde{v}(s)$ and $\tilde{w}(\xi, s)$ respectively. These have to be averaged over k in order to obtain the desired moment distributions. Thus, if $E(K) = \bar{k}$ and $E(K^2) = \bar{k}$, then it can be seen that

$$M_1(A, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) = \sum_{i=1}^{k_0} M_1(A, t | \xi_{i0}^0; s_0) + \bar{k} \int_{s_0}^t \tilde{v}(s) \int_{\tilde{D}} M_1(A, t | \xi^0; s) d\xi^0 ds \quad (6.17a)$$

and

$$M_{(2)}(A_1 \times A_2, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) = \sum_{i=1}^{k_0} M_{(2)}(A_1 \times A_2, t | \xi_{i0}^0; s_0) + \bar{k} \int_{s_0}^t \tilde{v}(s) \int_{\tilde{D}} M_{(2)}(A_1 \times A_2, t | \xi^0; s) \tilde{w}(\xi^0, s) d\xi^0 ds + \sum_{i=1}^{k_0} M_1(A_1, t | \xi_{i0}^0; s_0) \bar{k} \int_{s_0}^t \tilde{v}(s) \int_{\tilde{D}} M_1(A_2, t | \xi^0, s) \tilde{w}(\xi^0, s) d\xi^0 ds + \sum_{i=1}^{k_0} M_1(A_2, t | \xi_{i0}^0; s_0) \bar{k} \int_{s_0}^t \tilde{v}(s) \int_{\tilde{D}} M_1(A_1, t | \xi^0, s) \tilde{w}(\xi^0, s) d\xi^0 ds + \sum_{i=1}^{k_0} \sum_{\substack{j=1 \\ j \neq i}}^{k_0} M_1(A_1, t | \xi_{i0}^0; s_0) M_1(A_2, t | \xi_{j0}^0; s_0)$$



$$+ \bar{k} \prod_{j=1}^2 \int_{s_0}^t v(s) \int_{\tilde{D}} M_1(A_j, t | \xi^0; s) \tilde{w}(\xi^0, s) d\xi^0 ds \} . \quad (6.18a)$$

The quantitative description of the population in terms of the PGF assumes a particularly simple form in the case where the entities in the population do not reproduce and when the number of entities in the population at the initial time s_0 is zero.

6.4 External Input in Populations of Nonreproducing Entities

As mentioned earlier, a large number of processes of chemical engineering interest involve populations of nonreproducing entities. It will be seen later that some problems related to biological populations can also be looked upon as the dynamics of nonreproducing entities.

Given a nonreproducing entity at $\xi^0 \in \tilde{D}$ at time s , there will be only one entity in \tilde{D} at all times $t \geq s$ in the absence of an external input, because if the entity is not in \tilde{D} , it must be in ξ^* , and thus the entity is always restricted to stay within \tilde{D} . In view of this, the PGF $G(\theta, t | \xi^0, s)$ for the population starting with one nonreproducing entity at ξ^0 at time s and allowing no external input can be written as

$$\begin{aligned} G(\theta, t | \xi^0, s) &= E[\theta(\xi), t | \xi^0, s] \\ &= \int_{\tilde{D}} \theta(\xi) \chi(d\xi, t | \xi^0, s), \end{aligned} \quad (6.20)$$

where ξ denotes the location of the entity (the "ancestor") at time t (cf. Equation (2.9)). Substitution of (6.20) in (6.11) yields



$$\begin{aligned}
& G(\theta, t | \xi_{10}^0, \dots, \xi_{k_0 0}^0; s_0) \\
&= \prod_{i=1}^{k_0} \left(\int_{\tilde{D}} \theta(\xi) \chi(d\xi, t | \xi_i^0, s_0) \right) \left(\exp \left\{ \int_{s_0}^t v(s) \int_{\xi \in \tilde{D}} \int_{\xi \in \tilde{D}} \chi(d\xi, t | \xi_{i0}^0, s_0) \right. \right. \\
&\quad \left. \left. w(\xi^0, s) (\theta(\xi) - 1) d\xi^0 ds \right\} \right). \quad (6.21)
\end{aligned}$$

When there are no entities in the population initially, i.e., when $k_0 = 0$, and the number of Poisson inputs k is a constant, in view of Remark 6.1,

$$G(\theta, t | s_0) = \exp \left\{ \int_{s_0}^t v(s) \left\{ \int_{\xi \in \tilde{D}} \int_{\xi \in \tilde{D}} \chi(d\xi, t | \xi^0, s) w(\xi^0, s) [\theta(\xi) - 1] d\xi^0 ds \right\} \right\} \quad (6.22)$$

When

$$\theta(\xi) = \begin{cases} \zeta & \text{if } \zeta \in A \\ 1 & \text{if } \zeta \notin A \end{cases} \quad (6.22a)$$

for an arbitrary set $A \subset \tilde{D}$, (6.22) reduces to

$$G(\theta, t | s_0) = \exp \left\{ \int_{s_0}^t v(s) \left[\int_{\xi \in \tilde{D}} \chi(A, t | \xi^0, s) w(\xi^0, s) d\xi^0 \right] ds (\zeta - 1) \right\}, \quad (6.23)$$

which is the probability generating function for a Poisson process with mean

$$\int_{s_0}^t v(s) \left[\int_{\xi \in \tilde{D}} \chi(A, t | \xi^0, s) w(\xi^0, s) d\xi^0 \right] ds$$

(see e.g. Parzen, 1962; p. 125). This is an important result and may be stated as

Theorem 6.1: Suppose that the population in the spatial domain \tilde{D} defined in Section 6.1 contains no entities at time s_0 and that



- (a) Entities appear in \tilde{D} according to k simultaneous independent nonhomogeneous Poisson processes with rates $v_i(s)$, $i = 1, \dots, k$ from an external source, where the number k is a constant.
- (b) The location ξ_i^0 at which an entity first appears at time s due to the i -th Poisson stream is distributed in \tilde{D} according to a probability density $w_i(\xi_i^0, s)$ with respect to ξ_i^0 .
- (c) Each entity moves in \tilde{D} independently of others according to a Markov transition probability distribution $\chi(A, t | \xi, s)$ denoting the probability that the entity at ξ at time s will be in the set $A \subset \tilde{D}$ at time $t \geq s$.
- (d) The entities do not reproduce.

Then the number of entities in any set $A \subset \tilde{D}$ at time $t \geq s_0$ is a Poisson-distributed random variable with mean

$$M_1(A, t | s_0) = \int_{s_0}^t v(s) \left[\int_{\xi^0 \in \tilde{D}} \chi(A, t | \xi^0, s) w(\xi^0, s) d\xi^0 \right] ds, \quad (6.24)$$

where $v(s)$ and $w(\xi^0, s)$ are related to $v_i(s)$ and $w_i(\xi_i^0, s)$, $i = 1, \dots, k$, by (6.1) and (6.2) respectively.

Bartlett (1966) has indicated that when a population of non-reproducing individuals has an external input in the form of a Poisson process, the total number of individuals in the population at any time is a Poisson-distributed random variable if there are no individuals in the population initially. This observation corresponds to a special case of Theorem 6.1 with $A \equiv \tilde{D}$ in (6.24).

Remark 6.3: To obtain the r -th factorial moment of the number of entities in the set $A \subset \tilde{D}$, it is necessary to evaluate the r -th derivative of the probability generating function (6.23) at $\zeta = 0$, which yields the result

$$M_{(r)}(A, t | s_0) = [M_1(A, t | s_0)]^r. \quad (6.25)$$

Remark 6.4: By choosing the set A to be the single point ξ^* , it can be seen that when the hypotheses (a) thru (d) in Theorem 6.1 are satisfied, the number of entities which have been removed from \tilde{D} during the time interval (s_0, t) due to death is a Poisson-distributed random variable.

Many problems of practical interest can be looked upon as one of the following special cases of Theorem 6.1:

Case 1: When the initial location of each entity at the time of its first appearance due to any input stream is a fixed point $\xi_0^0 \in \tilde{D}$, the number of entities in any set $A \subset \tilde{D}$ at any time $t \geq s_0$ has a Poisson distribution with mean

$$M_1(A, t | s_0) = \int_{s_0}^t v(s) \chi(A, t | \xi_0^0, s) ds, \quad (6.26)$$

which readily follows from the fact that in this case $w(\xi^0, s) = \delta(\xi^0 - \xi_0^0)$ and hence

$$\int_{\tilde{D}} \chi(d\xi, t | \xi^0, s) w(\xi^0, s) d\xi^0 = \chi(d\xi, t | \xi_0^0, s).$$

Case 2: When the initial location of each entity at the time of its first appearance in \tilde{D} due to the external input from the i -th Poisson stream is located on the boundary Γ of \tilde{D} with a probability density $w_i(\hat{\xi}, s)$, $i = 1, \dots, k$; where $\hat{\xi}$ denotes a point on

Γ , the number of entities in any set $A \subset \tilde{D}$ at any time $t \geq s_0$ has a Poisson distribution with mean

$$M_1(A, t | s_0) = \int_{s_0}^t v(s) \int_{\Gamma} \chi(A, t | \xi, s) w(\xi, s) d\xi ds, \quad (6.27)$$

where $v(s)$ and $w(s)$ are related to $v_i(s)$ and $w_i(s)$ by (6.1) and (6.2) respectively.

In the case where the number of Poisson streams is a random variable with the process parameters $\tilde{v}(s)$ and $\tilde{w}(\xi, s)$ defined in Remark 6.1, substitution of (6.20) in (6.16) yields the expression

$$G(\theta, t | s_0) = \Psi \left\{ \exp \left[\int_{s_0}^t \int_{\xi \in \tilde{D}} \int_{\xi^0 \in \tilde{D}} \tilde{v}(s) \chi(d\xi, t | \xi^0, s) \tilde{w}(\xi^0, s) [\theta(\xi) - 1] ds \right] \right\} \quad (6.28)$$

for the PGF of a population of nonreproducing entities with no entities in the population initially. When $\theta(\xi)$ is given by (6.22a) for the set $A \subset \tilde{D}$, the PGF reduces to the probability generating function $\tilde{\Psi}(\zeta, t | s_0)$ of the number of entities in the set A , given by

$$\tilde{\Psi}(\zeta, t | s_0) = \Psi \left\{ \exp \left[\int_{s_0}^t \int_{\xi^0 \in \tilde{D}} \tilde{v}(s) \chi(A, t | \xi^0, s) \tilde{w}(\xi^0, s) d\xi^0 ds (\zeta - 1) \right] \right\}. \quad (6.29)$$

6.5 Applications

Most problems related to the population balance models of chemical engineering interest can be cast in the theoretical framework of the last section. For example, when a homogeneous nucleation of crystals or other suspended particles (as in a flocculation process) occurs in a process vessel, a Poisson rate of production with a certain probability density for the location of the particle in the vessel at the time of formation may be a reasonable assumption.

When uniform conditions prevail throughout the contents of the vessel, a uniform probability distribution for the initial location of a particle at the instant of its formation can be assumed. A similar description can also be used to characterize atmospheric phenomena such as formation of smog. In many chemical processes such as crystallization, nucleate boiling of fluids, distillation and liquid-liquid extraction the particles (i.e., crystal nuclei, bubbles of a vapor or droplets of a liquid) are formed at the surfaces enclosing the fluid. In these cases a Poisson rate of formation with a certain probability density for the initial location of the particles may be a reasonable description. In case of processes carried out in perfectly mixed vessels, the size or weight of a particle is often the only quantity of interest -- in such cases the individual state space for the particles in the population consists of the non-negative real line, and the formation of particles can be characterized as a Poisson input occurring at a single point in the state space, such as nuclei of zero size appearing in the vessel due to a homogeneous nucleation process. In many cases it may be necessary to consider the input of entities as consisting of a number of independent Poisson processes occurring in parallel. For example, in nucleate boiling of a fluid or in the case of a crystallizer where nucleation is taking place at the surfaces of the vessel, the local temperature at any point on the surface as well as the nature of the surface itself determines the rate of formation of the bubbles or crystal nuclei. Similarly, in processes such as distillation in a plate column the rate of bubble formation at the bubble caps or at the holes in a sieve plate will be dependent on the height of the liquid column

above the bubble cap or the respective set of holes in the sieve plate. Since the height of the liquid over a plate changes from the inlet to the outlet weir on the plate, the rates of bubble formation will be different at different points across the plate.

As discussed earlier in this chapter, the stochastic description of a population with an external input is quite involved in case of reproducing populations and therefore application of the theory of Section 6.3 to a crystallization process will be quite difficult when homogeneous nucleation, nucleation at the surfaces of the vessel, as well as secondary nucleation are occurring simultaneously at significant rates. Most particulate processes of chemical engineering interest involve particles of nonreproducing entities, for which the simpler results of the previous section may apply. Thus, in general, it can be said that if

- (i) At the start of a particulate process there are no particles present in the process vessel.
- (ii) The input of the particles into the process vessel can be characterized by a fixed number k of Poisson streams, with a probability density assigned to the particles in each stream for their states (such as the location inside the vessel as well as other characteristics such as size, shape, weight, etc.) at the time of their first appearance in the vessel,
- (iii) The future states of a particle and also the probability of its removal from the population depend only on its present state and not the past history,



- (iv) The changes in the state of a particle as well as its removal from the population proceed independently of other particles,

then the number of particles in any given set of states such as a size class, a volume element inside the vessel, etc. at any time is a Poisson-distributed random number.

In case of biological populations the external input commonly manifests itself in the form of a migration of individuals. As discussed before, the analysis of the stochastic fluctuations in a population with reproductive as well as migratory inputs is quite complicated. However, the study of migration of a single generation of individuals is also of importance in many problems of practical interest and in these cases the simpler analysis of the previous section can be applied. The example of the cereal leaf beetle (CLB), Oulema melanopus (L.) may be cited in this context. The CLB is an economic pest of the small grain crops in the North American continent. It has only one generation every year. In early spring the hibernating adults emerge from their overwintering sites in wooded areas surrounding grain fields and move to the succulent grain crops. They feed on the plant leaves for some time before starting their mating and egg-laying activity. Most of the damage to the plants is done by the larvae emerging from the eggs. Among other things, a detailed study of the movement patterns of the adults is necessary to arrive at efficient strategies for the control of the pest. Ruesink and Haynes (1972) studied the distribution patterns of the CLB adult populations in grain fields under different population densities. Their observations indicate that the adults are scattered



in a field according to a Poisson distribution only at medium densities. The mean and variance of the number of adults caught in a number of samples in the same field are significantly different from each other at very high and very low densities indicating that the spatial distribution is not Poisson at these densities.

If the rate of emergence of the CLB adults is assumed to occur according to a Poisson process with the points of emergence (i.e., the overwintering sites) distributed along the boundary of the field with a certain probability density and if the movement of a beetle within a field as well as its death are assumed to occur independently of others according to a Markov transition probability, then the number of beetles in any section of the field at any time should have a Poisson distribution. This description may be a good starting point in the quantitative study of the migration of the CLB adults. It is likely that the assumption of a Poisson emergence rate is not justifiable at low densities (cf. Section 6.6) and there may be a significant interaction among the individuals at high densities.

The analysis of external input to populations of nonreproducing entities can also be applied to problems other than spatial migration of biological populations. For example, in the study of the CLB population dynamics, the egg-laying activity of the overwintering adults can be considered as an "external" input to the new generation of the insects. In this case the individual state space \tilde{E} consists of the positive octant of a three dimensional Euclidean space with co-ordinates z_1, z_2 and z_3 ; with z_1 representing the maturity of an individual and z_2, z_3 denoting its geographical location in the field. The egg-laying activity of an adult located at



the point (z_2, z_3) in the geographical domain constitutes an "external" input of individuals of new generation at the boundary point $\xi = (0, z_2, z_3)$. Since the eggs and pupae are not translocated and since the migration of larvae is small relative to the dimensions of the field, the (Markovian) transition probability $\chi(A, t | \xi, s)$ characterizing the movement of an individual in \tilde{E} can be looked upon as the transition probability of movement along the maturity co-ordinate alone. Thus, $\chi(A, t | \xi, s)$ is independent of the geographical co-ordinates of an individual. In view of this, the quantitative description of the population in a fixed geographical area can be written in terms of z_1 alone.

It should be noted that when the diffusion equation can be used to characterize the growth of an individual, the transition probability $\chi(A, t | z_1, s)$ is given by the solution of the backward diffusion equation

$$-\frac{\partial \chi(A, t | z_1, s)}{\partial s} = \mathcal{A}_{z_1} \chi(A, t | z_1, s) \quad (6.30)$$

with the initial condition

$$\chi(A, t | z_1, t) = I(A | z_1)$$

and the boundary conditions

$$\frac{\partial \chi(A, t | z_1, s)}{\partial z_1} \Big|_{z_1=0} = 0$$

and

$$\frac{\partial \chi(A, t | z_1, s)}{\partial z_1} \Big|_{z_1=\tilde{a}} = 0$$

for a reflecting boundary at $z_1 = \tilde{a}$, or

$$\chi(A, t | \tilde{a}, s) = 0$$

for an absorbing boundary at $z_1 = \tilde{a}$. The operator \mathcal{Q}_{z_1} is the backward diffusion operator defined in (4.11). Alternatively, the transition probability can also be expressed in terms of the solution of the corresponding forward diffusion equation for the transition probability density.

The problem is complicated by the fact that the number of egg-laying adult females in any fixed area in a grain field is a random number. If the overall effect of the egg-laying females can be characterized as a random number k of Poisson inputs, each with the same rate $\tilde{v}(s)$, such that the number of inputs k varies from one part of the field to the other, but does not change with time in any fixed area within the field, and if the random number k has the probability generating function Ψ , then the number of individuals of the new generation in any given maturity range is given by (6.29). In the case where it is necessary to assume that the number of Poisson inputs k for any fixed area also changes with time, an additional transition probability describing the change in k must also be included in the description, thus further complicating the problem.

6.6 Justification of a Poisson Input

Throughout this chapter the external input to the population has been assumed to be in the form of a Poisson process. Besides the obvious advantages of its nice analytical properties, the assumption of a Poisson process is often justified on the basis of the fact that a stochastic process formed by a superposition of a number of independent sequences of events converges in distribution (see e.g.,

Feller, 1966; p. 243 for a definition of convergence in distribution) to a Poisson process if the individual sequences satisfy certain simple conditions (Jagers, 1971).

Connin (1971) studied the rate of egg-laying by CLB females under constant environmental conditions. The females were caught in the field and kept under controlled conditions at 80°F, 40% relative humidity and a 16 hour day-8 hour night cycle. The same experiment was also performed using adults reared and mated under laboratory conditions. A typical sequence of the daily egg-input by a CLB female is shown in Figure 6.1. A mated female starts to lay eggs at a relatively small number per day, quickly reaching a plateau, after which it continues to lay eggs at a more or less steady rate for two or three weeks (or even longer) before reaching a low rate again. The rate of egg output of a female is mainly a function of temperature (Yun, 1967) but the average rate during the egg-laying period varies from individual to individual, with an average of about 12 eggs per day at 80°F. It can be seen that under constant temperature conditions, if the female lays ten or twelve eggs during one day, it is very likely to lay about the same number of eggs during the next day. Since the number of eggs laid in disjoint intervals of time are not necessarily independent of each other and since it is possible to predict with certainty that an egg will be laid during certain time intervals, axioms (i) and (ii) of a Poisson process (see Section 6.1) will not be satisfied if the egg input is provided by a single female.

Wellso (1972) studied the egg-laying and feeding behavior of CLB adults. He enclosed the females in individual cages set up on oat plants in a field. The number of eggs laid by each female were

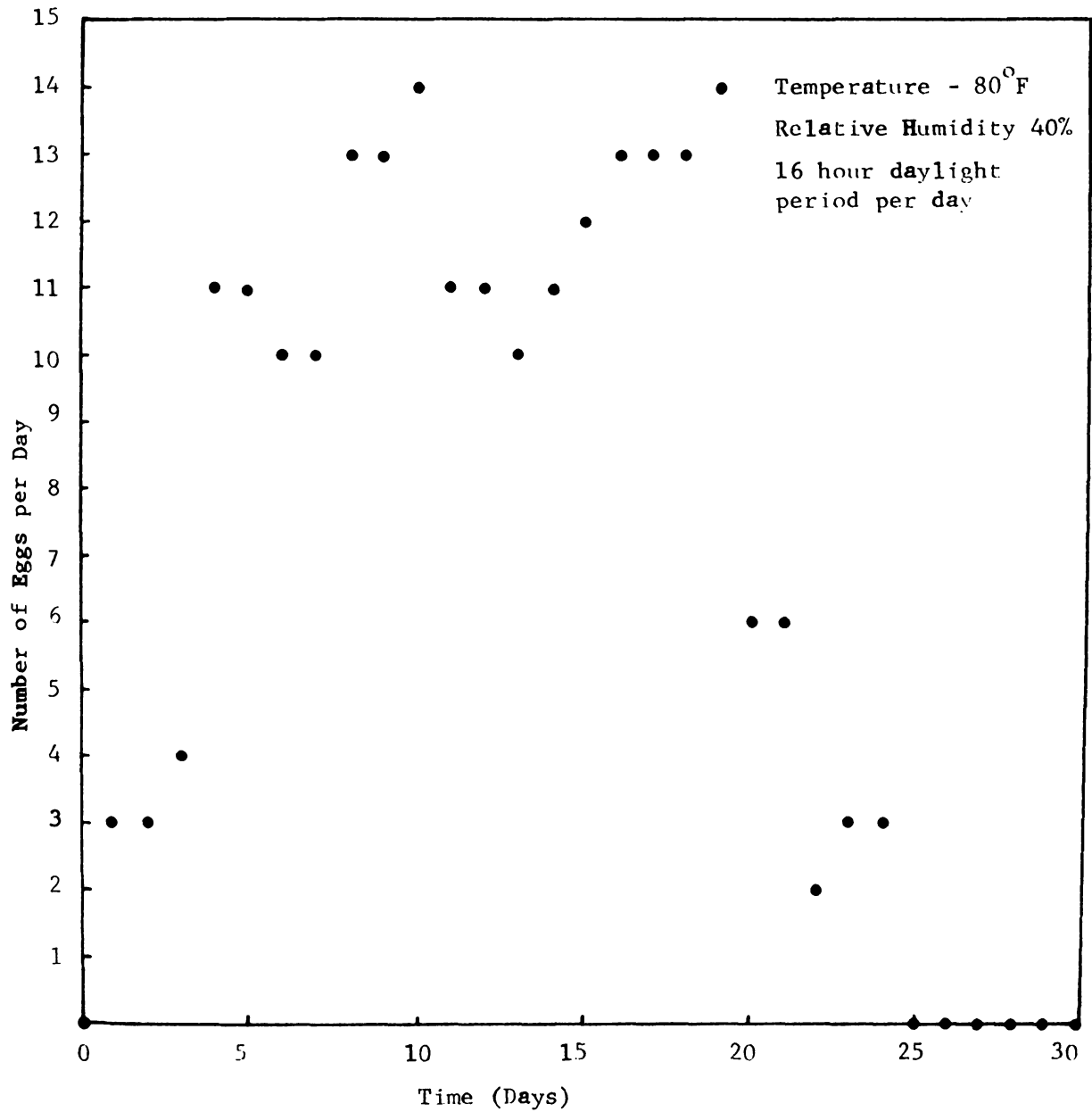


Figure 6.1 A typical sequence of the daily egg-input by a single cereal leaf beetle female. (Source: Connin, 1971)

counted every day. The χ^2 test (see e.g., Himmelblau, 1970; p. 74) was performed on this data to see how well the combined net daily input of eggs by all the females in the experiment can be represented by a Poisson process. The computed values of χ^2 were compared with those for a χ^2 distribution with a 10% level of significance for the appropriate number of degrees of freedom. The results of the statistical analysis of Wellso's data are summarized in Tables 6.1 and 6.2. In the results reported in Table 6.1, all the females taken for the experiment were considered whereas in the results listed in Table 6.2, only the females which laid eggs during each day were analyzed. When all the females in the experiment were considered in the analysis, the computed values of χ^2 were greater than those for the χ^2 distribution in four out of ten cases. When only the egg-laying females were considered, the computed values of χ^2 were greater than those for a χ^2 distribution in only two cases. This indicates that the total egg-input from all the females in the experiment may be only a poor approximation to a Poisson process, but when only the females which lay eggs are considered, the overall process may be a good approximation to the Poisson process. It should be noted that all the females used in the experiment were already mated and ready for oviposition. In the field populations there will also be a population of adults who have just emerged from their hibernation sites and have just begun their feeding and mating activity. However, if the adults are described in terms of their "age" since the time of emergence, it should be possible to differentiate the egg-laying adults from others, and when the number of egg-laying adults in a given area is sufficiently large, a Poisson rate of egg-input can be justified.

Table 6.1. Statistical Analysis of Cereal Leaf Beetle Oviposition Data with Mated Females*

Day	No. of Females	Mean # eggs/day	Variance (No.) ² /day	Degrees of Freedom	χ^2 (computed)	χ^2 (from Tables) 10% significance level
1	24	4.125	8.984	3	2.716	6.25
2	24	2.92	3.56	2	1.631	4.61
3	24	5.042	9.00	3	6.436	6.25
4	24	5.583	9.384	2	4.369	4.61
5	23	2.65	3.51	1	7.68	2.71
6	23	3.52	3.08	2	5.288	4.61
7	23	1.74	2.75	1	5.21	2.71
8	23	3.74	8.02	2	3.175	4.61
9	23	8.217	13.36	2	1.046	4.61
10	22	9.041	16.253	2	2.56	4.61

* The data was taken from Wellso (1972).

Table 6.2. Statistical Analysis of Cereal Leaf Beetle Oviposition Data with Egg-laying Females*

Day	No. of Females	Mean # eggs/day	Variance (No.) ² /day	Degrees of Freedom	χ^2 (computed)	χ^2 (from Tables) 10% significance level
1	22	4.5	8.71	2	4.584	4.61
2	22	3.182	3.013	2	2.942	4.61
3	22	5.5	7.214	2	3.015	4.61
4	22	6.09	7.039	2	3.421	4.61
5	21	2.904	3.091	1	13.388	2.71
6	22	3.682	2.608	2	2.405	4.61
7	15	2.667	1.667	1	3.759	2.71
8	19	4.526	6.041	2	3.197	4.61
9	22	8.591	10.634	2	0.648	4.61
10	22	9.041	16.253	2	2.56	4.61

* This table lists the same data used for Table 6.1, except that only the females which laid eggs during any given day were considered in the analysis.

Remark 6.5: It should be noted that the χ^2 test for goodness of fit of a collection of random numbers to a given distribution is only an approximate test. Thus the results of the test for CLB oviposition data can only be interpreted to mean that there is no reason to reject the hypothesis that the number of eggs laid by a collection of ovipositing females each day does not differ significantly from the Poisson distribution.

In the context of the particulate processes in chemical engineering systems, the number of "active sites" responsible for the generation of the particles (such as the number of surface imperfections on a heating surface in the case of nucleate boiling) is usually much larger than the number of "sources" (i.e., the CLB females) considered in the analysis above. Hence, a similar reasoning of a superposition of a large number of independent sequences of events converging to a Poisson process can be used to justify a Poisson input at least in an approximate sense in many particulate processes as well.

CHAPTER VII

AN EXPERIMENT BASED ON THE THEORY

7.1 Motivation

It was shown in Chapter IV that a diffusion equation can be used to characterize the moment distributions of a population if the movement of each entity in its spatial domain \tilde{D} can be represented by a Markov transition probability satisfying (4.1), the birth and death processes are also Markovian -- defined as in Section 4.1 and the movement of an entity in \tilde{D} as well as its reproduction and death take place independently of others. It is therefore necessary to test the validity of these conditions before the use of the diffusion equation can be justified in the modeling of a population. The interdependence between the members of the population may not necessarily be in the form of a physical interaction. For example, the insects feeding on the leaves of a grain plant may be located on different leaves of the plant with no physical interference between the feeding activity of each other. However, the feeding by each insect affects the condition of the plant in terms of the growth rate and succulence of the plant tissues, which in turn affects the new leaf surface area as well as the quality of food available to all the insects feeding on the same plant. The result obtained in Theorem 6.1 in the last chapter can be used to test whether the conditions of Markov transition probabilities and independence of entities in the population are

satisfied. To demonstrate this with the help of a simple problem of practical interest, a controlled experiment was performed on the cereal leaf beetle (CLB) populations.

The life-cycle of the CLB can be divided into seven distinct stages -- eggs, four stages of development (commonly called "instars" by the entomologists) of larvae, pupae and adults. Most of the damage to the crop is done by the larvae, mainly because they feed on the plant leaves in the early stages of development of the plant, thus severely affecting the plant growth at high larval densities. Thus, for the within-generation dynamics of the CLB population and its impact on the plants the range of maturities of the CLB from the beginning of the egg stage to the end of the fourth larval instar is of interest. The survivals in the pupae and the adults are related to the supply of eggs for the succeeding season.

When the egg-input occurs according to a Poisson process, the maturation and death of any individual are characterized by Markov transition probabilities and the change from one life-stage into another occurs at fixed levels of maturity, then by Theorem 6.1 the number of individuals in any life-stage at any time must be a Poisson-distributed random number. Thus, if a Poisson input of eggs is generated and if the population of any particular life-stage at any time in a large number of replicate experiments do not have a Poisson distribution, then it can be said that the conditions mentioned above are not satisfied for the population of that life-stage.

7.2 Methods and Materials

The CLB eggs are about one millimeter in length. The larvae grow from approximately the size of an egg to a length of about six



millimeters. In an experiment aimed at testing the independence of individuals and Markov transition probabilities it is necessary to have a large number of replicate experiments and to count the populations of different life-stages of the CLB at different points in time without disturbing the plants or the insects. Since the experiment is based on the measurement of the sample realizations of a stochastic process, the accuracy in counting the number of individuals is also very important. For the ease and accuracy of counting and to ensure a better control over the dynamic processes in the population, the experiment was conducted in a controlled atmosphere room in a laboratory. Two sets of experiments were performed: Experiment A and Experiment B. In Experiment A an abundant supply of food was provided to the insects so as to have little interaction between individuals. In Experiment B the food supply was severely restricted so that growth and death of an individual may be significantly influenced by other individuals in the population. Thirty replicates of each experiment were run to have a sufficient number of sample realizations of the process.

Although a large number of beetles laying eggs simultaneously lead to approximately a Poisson input of eggs (cf. Section 6.6), the experiment was performed with an artificially generated Poisson input by manually transferring the newly laid eggs onto the experimental plants to ensure a better control of the input. A sequence of 30 random numbers having a Poisson distribution with mean = 10 was computed using a table of random numbers. Each number in the sequence represented the total number of eggs to be used in a replicate experiment. Adult beetles were allowed to lay eggs on plants enclosed in



a cage in the laboratory for a period of eight hours, thus generating approximately a cohort of eggs. The required number of eggs was then carefully transferred on the experimental plants. The eggs have a layer of an adhesive on the outside, which made the transfer of eggs relatively easy. The same sequence of Poisson-distributed random numbers was used to generate the egg input in Experiments A and B.

Pots with barley plants having a uniform height of about 3.5 inches were chosen for the experiment. The number of stems per pot was chosen to be 40 for Experiment A and 25 for Experiment B. Plaster of Paris suspension was poured in each pot to cover the soil. The larvae have to go underground to form pupae. At the appropriate time a layer of vermiculite (a mineral normally used as a soil conditioner) was spread over the plaster of Paris surface. The plaster of Paris surface acted as a barrier restricting the insects to form pupae in the vermiculite layer. The required number of eggs were transferred to the plants and the pots were kept well separated from each other in large metal trays. Water was added to the trays every day to irrigate the plants through the holes at the bottoms of the pots. The atmospheric conditions in the room were set at 16 hours of daylight provided by fluorescent light and 8 hours of night, with the temperature being 80°F and 76°F during the day and night respectively. The relative humidity in the room was kept at 40%. Under these conditions the eggs began to hatch on the seventh day. Several hours before the actual emergence of the larvae from the eggs, the eggs had turned dark brown and the heads of the larvae could be clearly seen through the shells. After the eggs were ready to hatch the number of individuals in each life-stage were counted every day.

Each larval instar of the CLB is characterized by its head capsule diameter, which changes only during transition from one instar to another, when the insect sheds its head capsule and a part of its skin, and forms a larger head capsule. The head capsule diameter changes from one instar to the next by only a few microns. It is therefore difficult to distinguish between an insect which would molt shortly and one which has just molted without a microscope. Moreover, the molting process lasts several hours. Since the counting was to be done without disturbing the insects or plants, the larvae had to be classified according to the instars only by a visual observation. However, the presence of shedded head capsules and skins on the leaves often provided a clue regarding the number of individuals undergoing transition.

The larvae can crawl from one plant to another. When the larvae grew large enough to crawl from one pot to another, the plants in each pot were enclosed in lantern globes covered with pieces of cloth. Plants in Experiment A were sprayed periodically with water by means of a spray bottle to simulate rain. However, this did not seem to affect the larval mortality. In spite of the limited food availability in Experiment B, the larval mortalities were not noticeably different from that in Experiment A. In view of this the number of stems per pot in Experiment B were reduced progressively to four per pot. At this point in time however, all the available green foliage in almost all the replicate experiments was readily consumed by the larvae. After this stage only a limited amount of food (one stem per day or less) was supplied to the larvae in Experiment B. Figure 7.1 illustrates the difference in the condition of the plants in Experiments

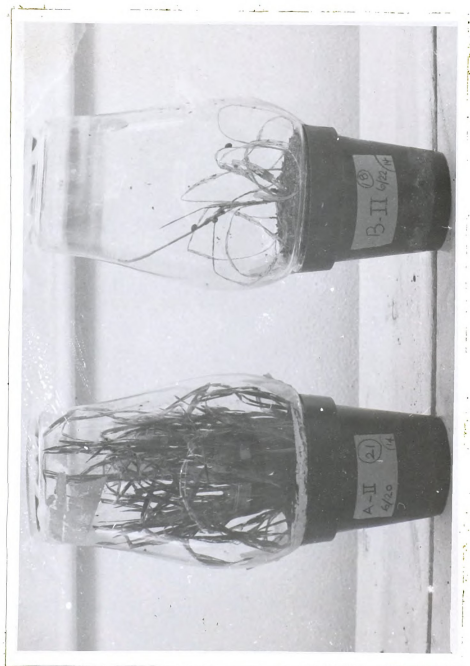


Figure 7.1. Comparison of the Plant Condition between Experiments A and B

A and B on the 14th day for pots, each of which had 14 eggs initially. The laboratory was also being used for other experiments, for which it was necessary to change the environmental conditions in the room. All the pots were therefore moved to growth chambers (i.e., small controlled-atmosphere cabinets) where the same temperature and day-light cycles were used. At the time of this transfer most of the insects had already formed pupae. All the pupae were left undisturbed until the adult emergence was complete, after which the pupal cases were removed from the vermiculite in the pots and opened to make sure that any insects left inside the cocoons were dead. The important steps in the experimental procedure are summarized in Table 7.1.

The daily emergence of the adults did not show the expected behavior -- in some replicates the emergence took place much earlier than the others, while in some replicates there was no emergence at all. The humidity in the growth chambers was not controlled. Moreover, the presence of a stagnant mass of air enclosed by the lantern globe and the covering cloth over each pot created conditions significantly different from those outside the lantern globes. The moisture evaporating from the soil was condensing on the inner surfaces of many lantern globes, thus creating almost a complete retention of moisture inside the globes. The presence of the condensate indicated that the air inside the globes was saturated with moisture. The temperature of vermiculite in the pots was also found to be substantially different from that of the ambient air inside the growth chambers, and varied from pot to pot. The vermiculite temperature was found to be as high as 98°F in some pots. In general, the pots with drier soil tended to have higher temperatures.

Table 7.1. Important Steps in the Experiments with CLB Populations

	Day	Procedural Step
Experiment A	1	Experiment started with 40 barley plants per pot
	7	Eggs began to hatch -- daily counting of populations started
	10	Put lantern globes on all pots
	11	Sprayed water on the plants to simulate rain
	12	Sprayed water on the plants to simulate rain
	13	Sprayed water to simulate rain, covered lantern globes with cloth
	14	Sprayed water to simulate rain
	15	Added vermiculite to all pots for the insects to pupate
	18	Moved all pots to growth chambers
	33	Experiment terminated
Experiment B	1	Experiment started with 25 barley plants per pot
	7	Eggs began to hatch -- daily counting of populations started
	9	Put lantern globes on all pots
	10	Reduced number of plants to 16 per pot
	11	Covered lantern globes with cloth
	12	Reduced number of plants to 9 per pot
	13	Added vermiculite to all pots for the insects to pupate, reduced number of plants to 4 per pot
	14	Added one stem of barley as food
	15	Added one stem of barley as food
	16	Moved all pots to growth chambers
	31	Experiment terminated

The effect of moisture on growth of the CLB is not well understood. The vermiculite in many pots exhibited fungus growth. The fungus growth was more significant in pots with a higher moisture content. The fungus was probably responsible for a part of the pupal mortalities, and thus the nonuniform humidity in different replicate experiments may have resulted in nonuniform mortalities in the pupal populations in the different replicates. However, since Experiments A and B were conducted under almost identical environmental conditions, it is reasonable to assume that if the difference in the food supply to the larvae did not have any effect on growth and mortalities of the insects in the two experiments, the effect of nonuniformities in humidity in the different replicates of Experiments A and B should also be identical.

The growth rate of an individual is directly related to the temperature (cf. Appendix A), which is believed to be a principal reason for the deviations from the expected emergence behavior of the adults. However, the result of Theorem 6.1 is still applicable if the total number of adults that ever emerged from the pupae are chosen as the quantity of interest by considering the range of maturities corresponding to the entire population coupled with an additional absorbing state ξ^{**} to characterize all the individuals that died during the adult stage (cf. Section 6.4) and by taking the time of observation to be such that no live individuals are present in any stage except the adults. In view of this only the total number of adults which emerged from each pot were analyzed instead of the daily number of emerging adults.

7.3 Results and Discussion

The data on the daily population counts of each life-stage in Experiments A and B is listed in Tables 7.2 and 7.3 respectively. This data was analyzed using an IBM 1800 computer. The χ^2 test was used to see whether the observed numbers in each life-stage at any given time had a Poisson distribution. The results of Experiments A and B are summarized in Tables 7.4 and 7.5 respectively. The computer program is listed in Appendix D. It can be seen from Tables 7.4 and 7.5 that the populations of all life-stages from the eggs to the pupae show a good fit to the Poisson distribution with few exceptions. The poor fit in these exceptional cases could be attributed to errors in distinguishing between successive instars. Better results were obtained in some instances when the dark-colored eggs which were ready to hatch were counted as larvae.

Although the degree of starvation of the insects in quite a few replicates of Experiment B was extremely severe, the mortality as well as the χ^2 statistics of all stages up to the fourth instar did not exhibit any significant difference from that in Experiment A. All the fourth instar larvae which went underground were counted as pupae. Since the number of pupal cases recovered at the end of Experiment B was less than the number fourth instar larvae in many replicates, it is likely that some of the starved larvae could not form complete cocoons, or died in the early part of the pupal stage and were completely decomposed by microbial activity. The total pupal mortality in Experiment B was much higher than Experiment A. It can be seen from the results that while the final population of the emerging adults in Experiment A was not a good approximation to a

Table 7.2. Cereal Leaf Beetle Population Data from Experiment A.

Day	Life Stage	-----Populations in the Thirty Replicate Experiments-----																													
1	Eggs	6	9	8	16	11	7	12	9	8	5	4	7	8	5	6	10	10	14	8	9	14	9	9	10	8	9	13	7	5	12
7	Unhatched eggs	5	5	1	12	7	4	9	7	4	4	4	5	8	5	3	10	8	5	3	6	7	2	2	6	2	6	5	2	1	3
7	Hatching eggs	1	4	6	4	4	3	2	2	4	1	0	2	0	0	3	0	1	9	4	3	7	7	7	3	6	1	4	4	4	8
7	Instar I	0	0	0	0	0	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	2	0	0	1
8	Unhatched eggs	1	4	0	1	0	0	2	0	1	2	0	0	2	0	1	3	2	2	1	1	2	1	1	0	2	3	2	1	0	2
8	Hatching eggs	2	0	0	2	0	1	1	0	0	0	0	0	1	1	0	1	0	0	0	0	0	0	0	1	0	0	1	0	0	1
8	Instar I	3	5	7	13	11	6	9	8	7	3	3	7	5	3	4	6	7	12	4	8	11	8	7	9	6	3	9	5	4	10
9	Unhatched eggs	1	4	0	1	0	0	2	0	1	2	0	0	2	0	1	3	2	2	3	1	2	1	1	0	0	1	2	1	0	0
9	Hatching eggs	2	0	0	1	0	1	1	0	0	0	0	0	1	1	0	1	0	0	0	0	0	0	0	0	0	1	0	0	1	0
9	Instar I	3	5	7	14	11	6	8	7	7	3	3	7	5	3	4	6	7	11	4	8	10	8	7	9	6	3	9	5	3	8
10	Unhatched eggs	1	4	0	1	0	0	1	0	1	2	0	0	2	0	1	3	2	2	3	1	2	1	1	0	0	1	2	1	0	0
10	Hatching eggs	2	0	0	1	0	1	2	0	0	0	0	0	1	0	0	1	0	0	1	0	0	0	0	0	0	1	0	0	0	0
10	Instar I	3	5	7	14	11	6	8	7	7	3	3	7	5	3	4	6	7	11	4	8	10	8	7	9	6	3	9	5	3	8
10	Unhatched eggs	1	4	0	1	0	0	1	0	1	2	0	0	2	2	1	1	2	2	0	1	0	1	0	0	0	0	0	0	0	0
10	Hatching eggs	1	0	0	1	0	1	2	0	0	0	0	0	1	0	0	2	0	0	1	0	2	0	1	0	0	1	0	1	0	0
10	Instar I	1	2	1	3	2	0	1	0	2	0	1	2	2	0	0	0	0	1	0	0	0	0	0	0	0	0	1	0	0	0
10	Instar II	2	3	6	11	7	6	7	6	5	3	2	5	3	3	4	6	7	9	3	8	10	8	7	9	6	4	8	4	4	10
11	Instar I	0	1	0	2	0	0	0	0	1	0	0	2	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
11	Instar II	3	4	7	12	9	6	7	6	6	3	3	5	5	3	4	6	7	10	4	8	11	8	6	9	4	3	7	4	2	7
11	Instar III	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
12	Instar I	0	0	0	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
12	Instar II	0	2	4	5	3	0	1	1	1	1	1	2	1	2	1	4	1	2	1	1	4	4	0	0	0	0	2	2	0	0
12	Instar III	3	3	4	6	7	5	6	5	5	2	1	4	4	1	3	2	6	8	3	7	6	3	6	9	6	3	7	2	4	9
13	Instar II	0	0	0	1	1	0	1	0	2	0	1	1	0	0	0	0	0	2	1	0	0	0	0	0	0	0	0	0	2	0
13	Instar III	3	4	7	11	9	5	6	6	4	3	1	5	4	2	4	6	7	8	3	8	10	7	6	9	6	3	9	2	4	9

Table 7.2. (continued)

Life Stage		-----Populations in the Thirty Replicate Experiments-----																											
Day	Life Stage	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
14	Instar II	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
14	Instar III	1	3	4	9	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
14	Instar IV	2	1	3	3	5	4	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
15	Instar III	0	0	0	2	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
15	Instar IV	3	4	7	10	8	5	4	3	1	5	3	2	4	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
16	Instar III	0	0	0	2	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
16	Instar IV	2	4	6	10	7	4	6	5	3	1	5	4	2	4	0	0	0	0	0	0	0	0	0	0	0	0	0	0
16	Pupae	1	0	1	0	1	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
17	Instar III	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
17	Instar IV	0	1	0	5	1	0	1	0	1	0	4	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
17	Pupae	3	3	7	7	7	5	4	2	1	2	3	1	3	3	0	0	0	0	0	0	0	0	0	0	0	0	0	0
18	Instar III	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
18	Instar IV	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
18	Pupae	3	4	7	12	8	5	6	4	3	2	5	4	2	4	4	4	4	4	4	4	4	4	4	4	4	4	4	4
19	Instar III	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
19	Instar IV	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
19	Pupae	3	4	7	12	8	5	6	4	3	2	6	4	2	4	5	7	8	4	4	4	4	4	4	4	4	4	4	4
20	Instar IV	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20	Pupae	3	4	7	12	8	5	6	4	3	2	6	4	2	4	5	7	9	4	4	4	4	4	4	4	4	4	4	4
33	Adults	1	4	3	6	0	3	1	3	2	3	1	1	4	1	3	1	5	6	3	3	4	2	6	0	7	1	2	0



Table 7.3. (continued)

Day	Life Stage	-----Populations in the Thirty Replicate Experiments-----																															
15	Instar III	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
15	Instar IV	4	6	4	12	3	4	5	10	7	6	4	6	8	2	6	5	4	10	5	7	10	6	4	7	7	3	7	11	3	7	0	
15	Pupae	1	1	2	2	1	1	0	0	0	0	0	0	0	1	0	2	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	
16	Instar III	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	1	0	0	0	0	0	0	0	
16	Instar IV	1	0	1	8	1	3	1	5	1	3	3	3	4	2	3	4	3	3	1	7	10	6	2	7	4	2	4	9	1	4	0	
16	Pupae	4	7	5	6	3	2	4	5	6	3	1	3	4	1	3	3	3	7	4	0	1	0	2	0	3	1	3	2	2	3	0	
17	Instar IV	0	0	0	1	1	0	0	4	0	2	0	0	2	0	0	1	2	2	0	2	2	0	0	3	2	0	0	3	0	0	0	
17	Pupae	5	7	6	13	3	5	5	6	7	4	4	6	6	3	6	6	4	8	5	5	9	6	4	5	5	3	7	8	3	7	0	
18	Instar IV	0	0	0	0	0	0	0	2	0	0	0	0	1	0	0	0	1	2	0	1	1	0	0	1	1	0	0	1	0	0	0	
18	Pupae	5	7	6	14	4	5	5	8	7	6	4	6	7	3	6	7	5	8	5	6	10	6	4	7	6	3	7	10	3	7	7	
19	Instar IV	0	0	0	0	0	0	0	2	0	0	0	0	1	0	0	0	0	1	0	1	1	0	0	0	1	0	0	0	0	0	0	
19	Pupae	5	7	6	14	4	4	5	8	7	6	4	6	7	3	6	7	6	9	5	6	10	6	4	8	6	3	7	11	3	7	7	
20	Instar IV	0	0	0	0	0	0	0	2	0	0	0	0	1	0	0	0	0	0	0	0	1	0	0	0	1	0	0	0	0	0	0	
20	Pupae	5	7	6	14	4	4	5	8	7	6	4	6	7	3	6	7	6	10	5	7	10	6	4	8	6	3	7	11	3	7	7	
31	Adults	0	5	5	0	0	1	1	0	3	0	3	2	2	0	5	3	5	3	2	0	1	1	1	1	0	0	1	0	0	0	0	0

Table 7.4. Statistical Analysis of Cereal Leaf Beetle Population Data from Experiment A

Day	Life Stage of the Beetle	Mean	Variance	Degrees of Freedom	χ^2 (computed)	χ^2 (from Tables) 10% significance level
7	Unhatched eggs	5.033	7.137	3	3.295	6.25
7	Hatching eggs	3.467	6.326	3	17.754	6.25
7	All eggs	8.500	7.845	3	5.314	6.25
7	Hatching eggs + instar I	3.600	6.800	3	25.903	6.25
8	Unhatched eggs	1.233	1.151	2	1.227	4.61
8	Hatching eggs	0.400	0.386	1	0.096	2.71
8	Instar I	6.767	8.185	3	6.154	6.25
8	All eggs	1.633	1.551	2	0.211	4.61
8	Hatching eggs + instar I	7.167	8.144	3	2.585	6.25
9	Unhatched eggs	1.100	1.197	1	0.607	2.71
9	Instar I	6.567	7.771	3	9.760	6.25
9	All eggs	1.433	1.909	2	2.498	2.71
9	Hatching eggs + instar I	6.900	7.748	3	5.234	6.25
10	Unhatched eggs	0.733	0.961	1	1.740	2.71
10	Instar II	5.867	6.395	3	3.097	6.25
10	All eggs	1.200	1.269	2	0.566	4.61
10	Hatching eggs + instar I	1.167	1.247	2	2.799	4.61
10	All larvae	6.567	7.357	3	4.543	6.25
11	Instar II	5.967	6.585	3	1.989	6.25
11	All larvae	6.433	7.289	3	4.386	6.25
12	Instar II	1.533	2.120	2	0.706	4.61
12	Instar III	4.667	4.920	3	2.828	6.25
12	All larvae	6.233	6.599	3	3.952	6.25
13	Instar II	0.400	0.455	1	1.762	2.71
13	Instar III	5.700	6.907	3	2.987	6.25
13	All larvae	6.100	6.990	3	5.539	6.25
14	Instar III	3.233	6.461	3	7.397	6.25
14	Instar IV	2.633	4.378	3	1.773	6.25
14	All larvae	5.933	7.651	3	5.057	6.25

Table 7.4. (continued)

Day	Life Stage of the Beetle	Mean	Variance	Degrees of Freedom	χ^2 (computed)	χ^2 (from Tables) 10% significance level
15	Instar IV	5.367	6.309	3	4.483	6.25
16	Instar IV	4.833	6.006	3	1.582	6.25
16	Pupae	0.633	1.137	1	1.538	2.71
16	All larvae + pupae	5.733	6.547	3	1.694	6.25
17	Instar IV	1.433	3.771	2	11.994	4.61
17	Pupae	4.200	5.131	3	4.131	6.25
17	All larvae + pupae	5.733	6.547	3	1.694	6.25
18	Pupae	5.467	6.464	3	1.605	6.25
18	All larvae + pupae	5.733	6.547	3	1.694	6.25
19	Pupae	5.633	6.447	3	1.871	6.25
20	Pupae	5.667	6.644	3	1.874	6.25
33	Adults	2.800	4.372	2	11.513	4.61

Table 7.5. Statistical Analysis of Cereal Leaf Beetle Population Data from Experiment B

Day	Life Stage of the Beetle	Mean	Variance	Degrees of Freedom	χ^2 (computed)	χ^2 (from Tables) 10% significance level
7	Unhatched eggs	1.767	3.978	2	1.131	4.61
7	Hatching eggs	3.200	5.821	4	21.819	7.78
7	Instar I	3.667	6.437	3	7.757	6.25
7	All eggs	4.967	7.206	4	12.486	7.78
7	Hatching eggs + instar I	6.867	8.947	3	2.221	6.25
8	Unhatched eggs	0.733	1.375	1	3.143	2.71
8	Instar I	7.267	7.926	3	0.895	6.25
8	All eggs	0.900	1.610	1	4.776	2.71
8	Hatching eggs + instar I	7.433	8.116	4	3.554	7.78
9	Instar I	2.967	13.758	2	63.476	4.61
9	Instar II	4.200	9.338	3	72.042	6.25
9	Hatching eggs + instar I	3.000	13.586	2	182.436	4.61
9	All larvae	7.167	8.075	4	1.112	7.78
10	Instar II	6.967	7.620	4	1.641	7.78
10	All larvae	7.033	7.826	3	1.371	6.25
11	Instar II	4.200	7.200	3	4.430	6.25
11	Instar III	2.767	4.254	3	5.018	6.25
11	All larvae	6.967	7.275	3	3.266	6.25
12	Instar III	6.400	6.110	3	5.050	6.25
12	All larvae	6.767	6.875	3	3.122	6.25
13	Instar III	4.733	7.926	3	7.702	6.25
13	Instar IV	1.900	2.645	2	4.225	4.61
13	All larvae	6.633	7.206	3	2.716	6.25
14	Instar III	0.800	2.441	1	6.293	2.71
14	Instar IV	5.733	8.133	3	5.189	6.25
14	All larvae	6.533	6.878	3	4.173	6.25
15	Instar IV	6.100	6.507	3	2.239	6.25
15	Pupae	0.400	0.455	1	1.762	2.71
15	Pupae + larvae	6.600	6.662	3	2.124	6.25



Table 7.5. (continued)

Day	Life Stage of the Beetle	Mean	Variance	Degrees of Freedom	χ^2 (computed)	χ^2 (from Tables) 10% significance level
16	Instar IV	3.533	6.602	3	6.016	6.25
16	Pupae	3.033	3.758	3	6.316	6.25
16	All larvae	3.567	6.875	3	6.388	6.25
16	Larvae + pupae	6.600	6.662	3	2.124	6.25
17	Instar IV	0.900	1.403	1	6.357	2.71
17	Pupae	5.700	4.355	3	3.091	6.25
17	Larvae + pupae	6.600	6.662	3	2.124	6.25
18	Pupae	6.233	5.289	3	2.790	6.25
18	Larvae + pupae	6.600	6.662	3	2.124	6.25
19	Pupae	6.333	5.885	3	4.665	6.25
19	Larvae + pupae	6.567	6.806	3	2.670	6.25
20	Pupae	6.400	6.110	3	3.444	6.25
20	Larvae + pupae	6.567	6.806	3	2.670	6.25
31	Adults	1.467	3.085	2	41.810	4.61



Poisson distribution (possibly due to the nonuniformity of temperature, humidity, etc. resulting in different death rates in the replicates), the adult population in Experiment B was even worse. Since the higher pupal mortality in Experiment B could be directly related to the amount of feeding done by the insects during the larval stages, i.e., the past feeding history, it can be said that the probability of death in the pupal stage was not Markovian, and therefore the Poisson character of the adult population was further distorted in Experiment B.

Remark 7.1: In view of Remark 6.5, the goodness of fit of the population data to the Poisson distribution can be inferred only in an approximate sense. The results of the χ^2 test can be taken to mean that the hypothesis regarding the independence of individuals and Markov nature of the dynamic processes in laboratory populations of the CLB may be true for any life stage up to and including the fourth instar even under severely stressed conditions. If growth of an individual can be characterized by a continuous Markov process satisfying (4.1) and the probability of death can be defined as in Section 4.1, then the diffusion equation can be used to describe the changes in the population in the range of maturities from the beginning of the egg stage to the end of the fourth instar. This is the case when the egg input is known and only the impact of the beetle on the crop for a given season has to be studied.

Barr, Kharkar and Lee (1972) have used the diffusion equation to characterize the first moment density of field populations of the CLB. Remark 7.1 will hold for their model if the behavior of the insects in field populations is similar to that observed in the experiments. In practice, the CLB population densities in the grain



fields will never be allowed to reach a stage of a complete defoliation of the plants. Chemical controls will normally be applied at a very much smaller level of damage to the plants than that observed in Experiment B. It is therefore possible that in field populations the conditions of independence of individuals and Markov transition probabilities will continue to hold even for the pupae and the adults at all the densities normally encountered, and the use of the diffusion equation may be justified for describing the entire life-cycle of the insect. However, it is also likely that under field conditions the solar radiation, wind and humidity will play a significant effect on the mortalities of larvae under stressed conditions. When there is a significant defoliation of the plants due to insect feeding, the larvae have to move larger distances on a plant or in a group of neighboring plants to find suitable food, and in this process they are more likely to be blown or shaken off the plant due to the action of wind. This increases the chances of a higher metabolic stress as well as mechanical damage to the tissues, resulting in a greater chance of dying at the time of molting due to inability of a larva to shed its head capsule and skin. Moreover, the larvae on defoliated plants are directly exposed to the sun for longer periods of time and the effect of dehydration due to solar radiation is also likely to be more pronounced, particularly when the humidity of the ambient air is low. A careful investigation of the effect of these factors is needed before a definite conclusion about the justification of the diffusion equation to describe the CLB field populations can be drawn.



CHAPTER VIII

CONCLUSIONS

8.1 Summary of the Contributions of this Dissertation

This work has cast the population balance models in a very general stochastic framework. Applications of the models considered in this dissertation include the residence time distribution analysis (RTD) of fluids in process vessels under steady flow conditions, dynamics of particulate and polymerization processes in chemical engineering systems as well as dynamics of populations in biological and social systems. A probabilistic treatment of the population balance models in chemical engineering has been restricted only to some simple cases thus far (e.g., Kane, 1971; Shinnar and Katz, 1969; Schmalzer and Hoelscher, 1971). The "continuous stirred tank" and "dispersion" approaches in the RTD analysis have been shown to be applicable for a large class of general populations. Although analogous models had been proposed and used for some biological populations, the generality of these approaches had not been fully appreciated so far. Since the methods of parameter estimation are well developed for the RTD analysis, the possibility of use of analogous models to other populations makes it possible to use analogous methods for the evaluation of parameters as well.

Diffusion equations have been derived for the probability generating functional as well as the first two factorial moment



distributions and densities for a population of reproducing entities. These equations represent a more general situation than similar equations previously reported in literature. The backward equations are used only very rarely in engineering applications -- the value of these equations in some cases of practical interest has been demonstrated. It is a common engineering practice to use Monte Carlo methods to characterize the stochastic fluctuations in the model of a process. These methods use models where the parameters are assumed to be random numbers with known distributions (the normal distribution is often used for this purpose). The values of the parameters are generated with the help of a random number generator and a number of solutions are obtained by using different values of the parameters thus generated. The mean and variance of the output of the model are then computed from the replicate solutions. Coulman, Riece and Tummala (1971) have used this method in a model for a species of freshwater shrimp to characterize the mean and variance of the population as functions of time. The equations for the moment distributions derived in the present work represent an alternative to the Monte-Carlo techniques. The simple cases of practical interest for which the analytical and numerical solutions have been obtained in the present work represent new results and demonstrate the value of the stochastic approach to the population balance models.

A very general analysis for populations of reproducing entities has been presented to include the effect of external input on the probability generating functional and the first two factorial moments of the population. In this analysis the external input is assumed to be in the form of a nonhomogeneous Poisson process, the

entities in the population are considered to be independent of one another and the movement of each entity in its state space is characterized by a Markov transition probability. An explicit partial differential equation has been obtained for the first moment density of a population of reproducing entities with an external input, with the movement of each entity in its one-dimensional state space being characterized by a diffusion process. An important result (Theorem 6.1) has been obtained for the particular case of populations of nonreproducing entities, and its application to several cases of practical interest has been discussed.

The models derived in this dissertation are based upon two key assumptions: mutual independence of entities in the population and Markov transition probabilities characterizing the death, reproduction and movement of each entity in its state space. Possibility of using the theory developed in this thesis in validating these assumptions for the particular case of nonreproducing populations has been demonstrated by means of a simple experiment.

8.2 Areas for Future Research

The analytical results in Chapter V were based upon the hypothesis that the dominant eigenvalue of the diffusion operator (with parameters defined by (5.1)) is unique and real. Although the existence of a real eigenvalue has been proved, it remains to be shown that this eigenvalue is the only dominant eigenvalue.

The processes for which the present work does not apply include the size reduction processes in chemical engineering systems and degradation of polymer molecules, where every breakage of an



entity may result in a point process, i.e., a random distribution of sizes of the "pieces" may be generated at every breakage. Possibilities of extending the present analysis to describe these processes need to be explored.

Another situation not covered in this dissertation is the case of a population with an external input in the form of a generalized Poisson process. Incorporation of such an input in the analysis will greatly add to the generality of the present work.

In industrial processes involving particulate matter the number of particles in the system is usually very large and consequently the stochastic fluctuations in the dynamic characteristics of the process system are almost always of a minor importance. However, the data needed for the design of these systems are usually collected in a small experimental setup where the number of particles is rather small, and the scatter in the data is often rather large. Use of stochastic models in the analysis of such systems instead of considering the scatter to be "noise" in the data may lead to a better understanding of these systems. Similarly, in systems such as fluidized beds often it is not possible to use a purely deterministic approach to characterize certain phenomena like the formation and breakup of large bubbles and the relationship of these phenomena to the characteristics of the fluidized particles. The stochastic framework provided in this thesis, or other formulations based upon the theory of stochastic population processes offer alternative ways to describe a broad range of such problems of chemical engineering interest. In order to fully appreciate the value of a stochastic approach such applications need to be explored in detail.



As mentioned before, the stochastic framework for the population balance models makes it possible to formulate the models based entirely upon phenomenological descriptions of the processes involved. In the analysis of crystallization processes the growth and nucleation rates are commonly expressed by empirical kinetic expressions. Hypotheses treating these phenomena as stochastic processes would be a natural next step in gaining a better quantitative understanding of the crystallization processes. It would be beneficial to carry out similar investigations regarding the dynamics of other particulate processes also.

The possibility of an application of the models developed in this thesis to social systems needs to be explored. For example, orthodox or enterprising nature of certain ethnic groups can possibly be quantified by using the concepts developed in Chapter III: Thus, if the post-world war II boom in the U.S. economy is considered as a known change (such as a step or a ramp change) in the available opportunities, a comparison of the economic (or educational) status of different ethnic groups can be made to quantify how well each group has used the opportunities. A group significantly surpassing (or lagging behind) the rest of the population in reaping the benefits from the opportunities can be seen to make the system a "nonideal population-flow system" (see Section 3.4), since all the individuals in the population do not progress at the same rate. The opportunistic groups in the population can be characterized by the "bypassing" phenomenon whereas the groups lagging behind the rest of the population can be described by the phenomenon of "dead space" in the population. The quantitative models for the social systems based

upon such considerations can then be used to derive alternatives to the current means of social control aimed at achieving a more equitable progress in the society.

Of the possible boundary conditions for diffusion equations, the absorbing and reflecting boundaries have been used extensively in various applications. The residence time distribution analysis of flow in short packed beds offers a unique possibility of application of a third type of boundary condition describing the possibility that a fluid element at a boundary can jump to an interior point, as pointed out in Section 4.2. The application of such a boundary condition needs to be studied in greater detail.

The formulations presented in this work are based upon the assumption that the entities in the population are independent of one another. However, many phenomena in the dynamics of biological populations result from interactions among individuals in the form of competition for a certain vital resource. Similarly, chemical engineering systems such as liquid-liquid extraction and gas absorption involve liquid droplets and gas bubbles which coalesce and break-up again, thus violating the assumption of independence of the entities. Application of the theory of stochastic population processes to populations of interacting entities represents an uncharted area in the study of stochastic processes. Development of even approximate techniques for the calculation of the first two moment distributions will represent an important milestone in the mathematical modeling of populations.

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APPENDICES

APPENDIX A

SIMPLIFICATION OF THE DIFFUSION MODEL FOR SOME BIOLOGICAL POPULATIONS

The physiological processes such as growth and reproduction in poikilothermic ("cold-blooded") organisms exhibit a common characteristic dependence on temperature. The rate of these processes is negligible below a certain threshold temperature T_l , which depends upon the species and the process, and increases more or less linearly with temperature up to an upper threshold T_u . The rate decreases with temperature above the upper threshold. However, for a species adopted to a particular climate, the threshold temperatures are often such that the total length of time spent by the organism at temperatures above the upper threshold during a year is rather small. For example, the upper and lower threshold temperatures for the rates of growth and egg-laying for the cereal leaf beetle, a major pest of the small grain crops in midwestern United States, are about 48°F and 90°F respectively (Yun, 1967). The number of days in a year with temperatures in the nineties is indeed small in this part of the United States.

The commonality in the temperature dependence of poikilothermic organisms is used to define a physiological time in terms of "degree days". The use of degree days often leads to clearer biological insights into a system (see e.g., Wang, 1960), and is widely used by biologists to study population data. The profound effect of



temperature on the physiological processes in poikilothermic organisms implies that except for constant temperature studies, the parameters in the mathematical models of such populations must be functions of temperature. Barr, Kharkar and Lee (1972) have shown that the temperature dependence of the model parameters can often be eliminated by using degree days instead of chronological time in the quantitative descriptions. They have discussed this simplification for the forward diffusion equation for the first moment density of a population where each individual is characterized by only one descriptor, namely, maturity. The same argument can be used to show that the diffusion equation for the probability generating functional (PGF) can also be greatly simplified by the use of degree days, and hence the simplification applies to all the moment distributions of the population.

Consider the backward diffusion equation for the PGF with y denoting the initial location of an ancestor at time s :

$$\begin{aligned}
 - \frac{\partial G(\theta, t | y, s)}{\partial s} &= \frac{\sigma^2(y, s)}{2} \frac{\partial^2 G(\theta, t | y, s)}{\partial y^2} + r(y, s) \frac{\partial G(\theta, t | y, s)}{\partial y} \\
 &\quad - \mu(y, s) G(\theta, t | y, s) + \mu(y, s) - \sum_{i=1}^{\infty} \lambda_i(y, s) G(\theta, t | y, s) \\
 &\quad + \sum_{i=1}^{\infty} \lambda_i(y, s) G(\theta, t | 0, s)^i G(\theta, t | y, s)
 \end{aligned} \tag{4.12}$$

with the initial condition

$$G(\theta, t | y, t) = \theta(y) \tag{4.16}$$

and boundary conditions

$$\frac{\partial G(\theta, t | y, s)}{\partial y} \Big|_{y=0} = 0, \tag{4.14}$$



and

$$\left. \frac{\partial G(\theta, t | y, s)}{\partial y} \right|_{y=\tilde{a}} = 0 \quad (4.15)$$

for a reflecting barrier at $y = \tilde{a}$, or

$$G(\theta, t | \tilde{a}, s) = 1 \quad (4.13)$$

for an absorbing barrier at $y = \tilde{a}$. Moreover, let the parameters $\frac{\sigma^2}{2}$, r , μ and λ_i , $i = 1, 2, \dots$ be piecewise linear functions of temperature with the same lower threshold temperature, i.e., for temperatures $T(s)$ below the upper threshold let

$$\frac{\sigma^2(y, s)}{2} = f[T(s)] \frac{\tilde{\sigma}^2(y)}{2}, \quad (A.1)$$

$$r(y, s) = f[T(s)] \tilde{r}(y), \quad (A.2)$$

$$\mu(y, s) = f[T(s)] \tilde{\mu}(y), \quad (A.3)$$

$$\text{and} \quad \lambda_i(y, s) = f[T(s)] \tilde{\lambda}_i(y), \quad i = 1, 2, \dots \quad (A.4)$$

where

$$f[T(s)] = \begin{cases} T(s) - T_\ell & \text{for } T(s) \geq T_\ell \\ 0 & \text{for } T(s) \leq T_\ell \end{cases} \quad (A.5)$$

and define the physiological time variable $\tilde{\tau}$ as

$$\tilde{\tau} = \int_t^s f[T(w)] dw, \quad (A.6)$$

for a fixed t .

Since $d\tilde{\tau} = f[T(s)] ds$, it follows that (4.12) thru (4.16) can be written as



$$\begin{aligned}
-\frac{\partial G(\theta, t|y, \tilde{\tau})}{\partial \tilde{\tau}} &= \frac{\tilde{\sigma}^2(y)}{2} \frac{\partial^2 G(\theta, t|y, \tilde{\tau})}{\partial y^2} + \tilde{r}(y) \frac{\partial G(\theta, t|y, \tilde{\tau})}{\partial y} \\
&- \tilde{\mu}(y) G(\theta, t|y, \tilde{\tau}) + \tilde{\mu}(y) - \sum_{i=1}^{\infty} \tilde{\lambda}_i(y) G(\theta, t|y, \tilde{\tau}) \\
&+ \sum_{i=1}^{\infty} \tilde{\lambda}_i(y) G(\theta, t|0, \tilde{\tau}) {}^i G(\theta, t|y, \tilde{\tau}), \tag{A.7}
\end{aligned}$$

$$G(\theta, t|a, 0) = \theta(y), \tag{A.8}$$

$$\left. \frac{\partial G(\theta, t|a, \tilde{\tau})}{\partial y} \right|_{y=0} = 0, \tag{A.9}$$

$$\left. \frac{\partial G(\theta, t|a, \tilde{\tau})}{\partial y} \right|_{y=\tilde{a}} = 0 \tag{A.10}$$

and

$$G(\theta, t|\tilde{a}, \tilde{\tau}) = 1 \tag{A.11}$$

respectively. Note that the temperature dependence of the parameters has been eliminated in (A.7) thru (A.11) because $f[T(s)]$ has been cancelled from each term in the equations.

In view of the fact that $\lambda_i(y, s)$, $i = 1, 2, \dots$ essentially describe a physiological process, (A.4) is a realistic approximation for reproduction in poikilothermic organisms. The actual temperature dependence of growth or maturation would very likely be such that the random increment δy in the maturity of an arbitrary individual during a time interval δs is a piecewise linear function of temperature. The incremental growth δy is related to the diffusion and drift coefficients by (4.1), from which it follows that r would be characterized by a piecewise linear function as in (A.2) and $\frac{\sigma^2}{2}$ would be a quadratic function of temperature. Analysis of the data of Yun (1967) and Helgesen (1969) for maturation of the cereal leaf beetle lends some support to this conjecture.



In spite of the quadratic dependence of the diffusion coefficient on temperature, it may be possible to obtain a fair piecewise linear approximation for $\frac{\sigma^2}{2}$ as shown in (A.1).

The term $\mu(y,s)ds$ is the probability of an individual of maturity y being removed from the population by death or emigration during the time interval $(s, s+ds)$. In fact, it can be viewed as a sum of three different terms representing loss of individuals to the population from (i) emigration, (ii) death due to old age, disease, etc. and natural enemies such as parasites and predators, and (iii) artificial removal by harvesting, pesticide application, etc. The following discussion considers (i), (ii) and (iii) in sequence as if each occurred alone:

When emigration of individuals occurs solely due to locomotion, (A.3) can be justified in view of the fact that locomotive activity of an individual is essentially a physiological process and hence dependent upon temperature. For natural death and death due to natural enemies, (3.4) does not at first seem to be an accurate representation. However, in some species these mortalities are often considered to occur at those particular discrete maturities which define the transitions between the various life stages of an organism. For such a situation (4.12) may be used to describe the behavior between each of these transitions with $\mu(y,s)$ being 0. Thus (A.3) is trivially satisfied with $\tilde{\mu}(y) = 0$. When this form of analysis is employed, these mortalities are used to determine boundary conditions at each of the transitions. The assumption of mortalities



occurring only at maturity transitions seems to be particularly good for natural death and parasitism, as compared to predators. Nevertheless, since a prey may be less mobile at transitions between life stages, it could be more vulnerable to predators. Artificial removal by harvesting and pesticides is essentially independent of temperature and thus for this case (A.3) is invalid, and the function $\mu(y,s)$ in (A.7) must be replaced by

$$\frac{\mu(y,s)}{f[T(s)]} . \quad (\text{A.12})$$

If $\mu(y,s) = 0$ whenever $f[T(s)] = 0$, then the analytical form of (A.3) and the resulting (A.7) are well posed. This supposition is true when harvesting, pesticide application, etc. do not occur at temperatures below the lower threshold T_l . When $f[T(s)] > 0$, (A.5) describes a one-to-one relationship between s and $\tilde{\tau}$, and thus in (A.12) s may be replaced by $\tilde{\tau}$. The resulting form of (A.7) is simpler than (4.12) in the sense that only one parameter is dependent upon temperature. When (i), (ii) and (iii) occur simultaneously, all the remarks above must hold for the simplification to be possible.

When (A.1) thru (A.4) are valid even after a suitable choice of a maturity variable y to remove the dependence of parameters on y , (A.7) reduces to an equation with constant coefficients, and the analysis of Sections 5.2 and 5.3 can be used to solve the moment equations derived from the diffusion equation for the PGF. In practice, the parameters $\frac{\sigma^2}{2}$, r , μ , and λ_i have to be taken as being independent of maturity over the respective maturity intervals. For example, the data of Yun (1967) and Helgesen (1969) on the growth

characteristics and egg laying behavior were obtained in the form of the mean and standard deviation of the developmental times for each of the life stages, and mean number of eggs per adult per day respectively. In such cases elimination of temperature dependence of the parameters again leads to a diffusion equation with constant or piecewise constant parameters.



APPENDIX B

NUMERICAL COMPUTATION OF THE "CRITICAL LENGTH"

As discussed in Section 5.4, the critical location \tilde{a} of an absorbing barrier for a population of reproducing entities characterized by a diffusion process is given by solving

$$(\hat{\mu}/\hat{\lambda})[\nu_1 \hat{P} \exp(\nu_2 \hat{P}) - \nu_2 \hat{P} \exp(\nu_1 \hat{P})] + \nu_2 \hat{P} \exp[\nu_1(\hat{P} - \hat{P}b)] - \nu_1 \hat{P} \exp[\nu_2(\hat{P} - \hat{P}b)] + \nu_1 \hat{P} - \nu_2 \hat{P} = 0, \quad (5.64)$$

with

$$\nu_1 = \frac{-1 + \sqrt{1 + 4\hat{\mu}/\hat{P}}}{2}$$

and

$$\nu_2 = \frac{-1 - \sqrt{1 + 4\hat{\mu}/\hat{P}}}{2},$$

for $\hat{P} = \frac{2\tilde{a}r}{\sigma^2}$ and substituting for the diffusion coefficient $\frac{\sigma^2}{2}$ and the drift coefficient r . The computer program CRITL was used to obtain the family of curves in Figure 5.1 by solving (5.64) for \hat{P} with a number of values of other parameters. In the program the variables XLM, PB, XMP and X1 correspond to the quantities $\hat{\lambda}/\hat{\mu}$, $\hat{P}b$, $\hat{\mu}/\hat{P}$ and the assumed value of \hat{P} (for an iterative solution) respectively from (5.64).

Since \hat{P} must be larger than PB for a physically meaningful solution (cf. Section 5.4), the iterative procedure in CRITIC is started with $X1 = PB$. The function on the left hand side of (5.64)



is generated by the function subprogram FNC. The root P of (5.64) is computed by using the subroutine CVG1 available on magnetic tape on the CDC 6500 computer at Michigan State University. The subroutine works as follows: X_1 and X_2 are two trial values of a root of (5.64). If both X_1 and X_2 lie on the same side of the root, the values Y_1 and Y_2 of the function FNC evaluated at X_1 and X_2 will have the same sign. When this is the case, the subroutine sets IND to -1 and returns to the main program, to take a new value of X_2 . This procedure is repeated until X_1 and X_2 lie on opposite sides of a root. If this procedure leads to too large a value of X_2 to be physically meaningful, the iteration is stopped and after printing a message to that effect, the program proceeds to the next set of values of parameters. Once a value of X_2 is found such that the root lies between X_1 and X_2 , the subroutine computes the point where the straight line connecting the points (X_1, Y_1) and (X_2, Y_2) crosses the X-axis. The value of X corresponding to this point is taken as the new value of X_2 and the iteration is continued until Y_1 or Y_2 is sufficiently close to zero. This iterative procedure is illustrated in Figure B.1.

Since the subroutine CVG1 gives only one root of the function depending on the initial guesses of X_1 and X_2 and the iteration scheme for changing the values of X_2 , the reliability of this computer program hinges on the uniqueness of the solution to the problem of critical length. The shape of the function was studied for a large number of values of the parameters using an IBM 1800 computer coupled with a plotter and it was found that only one root \hat{P} of (5.64) occurred for $\hat{P} > \hat{P}_b$. Although the uniqueness of the solution to the problem of critical length was not proved theoretically, these results were taken as an indication of the uniqueness.



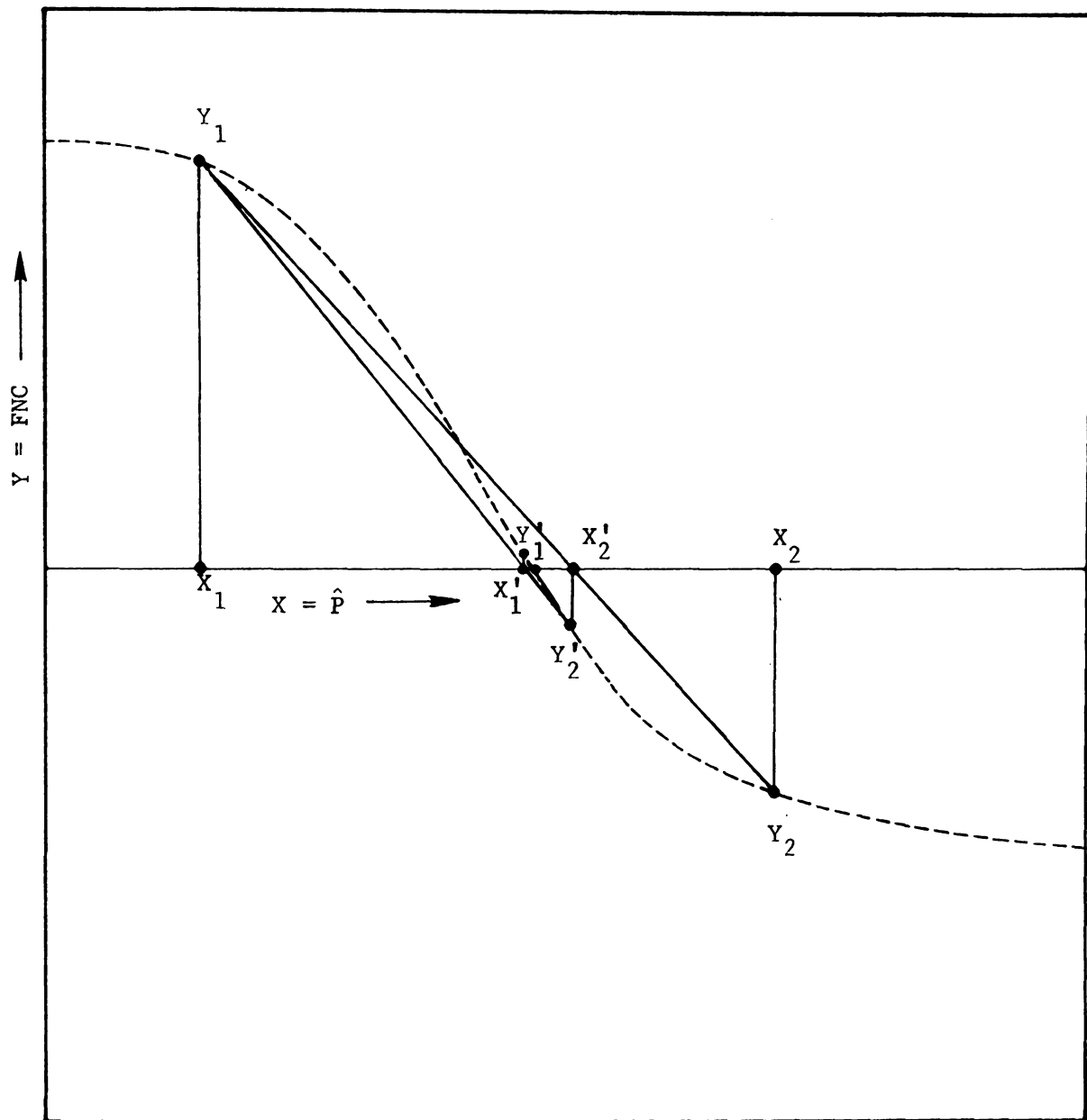


Figure B.1 Iteration scheme for obtaining a root of Equation (5.64)



```

PROGRAM CRITIC(INPUT,OUTPUT)
DIMENSION XLM(3),XMP(17),PB(5)
DATA XLM/2.,10.,20./
DATA PB/3.,5.,10.,20.,30./
DATA XMP/.01,.02,.03,.04,.05,.06,.07,.08,.09,.1,.11,.12,.13,.14,.1
15,.16,.17/
PRINT 100
100 FORMAT(*1 BIRTH/DEATH P BETA DEATH/P P*,7X,*CRI
1T LENGTH*)
I=2
J=4
DO 51 K=1,17,8
Y1=XLM(I)
Y2=PB(J)
Y3=XMP(K)
X1=Y2
X2=X1+1.
IND=1
20 ROOT=FNC(Y1,Y2,Y3,X1)
23 CALL CVG1(X1,ROOT,X2,IND)
24 CRL=2.5*X1/Y2
IF(IND)21,22,20
21 X2=X2+1.0
IND=1
IF(CRL-20.)25,26,26
26 PRINT 102
GO TO 102
25 GO TO 23
22 CONTINUE
PRINT 101,Y1,Y2,Y3,X1,CRL
101 FORMAT(1X,6(G12.5,2X))
102 FORMAT(* CRITICAL LENGTH TOO LARGE*)
51 CONTINUE
52 CONTINUE
STOP
END

```



```

FUNCTION FNC(X,Y,Z,P)
PH=(P-Y)*0.5
SQR=SQR(1.+4.*Z)    $PH1=-1.+SQR    $PH2=-1.-SQR
FNT=PH2*EXP(PH1*P*.5)-PH1*EXP(PH2*P*.5)
FNX=X*(PH2*EXP(PH1*PH)-PH1*EXP(PH2*PH))-PH2+PH1)
FNC=FNT-FNX
PRINT 100,PH1,PH2,FNT,FNX,FNC
100 FORMAT(1X,10(E12.5,1X))
RETURN
END

```



APPENDIX C

COMPUTATION OF EXTINCTION PROBABILITY FOR A POPULATION OF REPRODUCING ENTITIES

As discussed in Section 5.5, the extinction probability for a population of reproducing entities characterized by a diffusion process with parameters as in Table 5.2 (given one ancestor initially at y) was computed by solving the backward diffusion equation for the PGF with the initial condition $G(\theta, 0|y) = 0$. In view of the nonlinear nature of the diffusion equation, the solution was obtained numerically using finite differences. The difference equations for the simulations with conditions described in Table 5.2 are

$$\begin{aligned} \frac{G(\theta, \tau + \delta\tau|i\delta y) - G(\theta, \tau|i\delta y)}{\delta\tau} &= \frac{\sigma^2}{2} \left\{ \frac{G(\theta, \tau|(i+1)\delta y) - 2G(\theta, \tau|i\delta y) + G(\theta, \tau|(i-1)\delta y)}{\delta y^2} \right\} \\ &+ r \left\{ \frac{G(\theta, \tau|(i+1)\delta y) - G(\theta, \tau|i\delta y)}{\delta y} \right\} + \mu[1 - G(\theta, \tau|i\delta y)] \\ &- \lambda_1(i\delta y)G(\theta, \tau|i\delta y)[1 - G(\theta, \tau|0)] \quad \text{for } \delta y \leq i\delta y \leq \tilde{a} - \delta y, \end{aligned} \quad (C.1)$$

$$\begin{aligned} \frac{G(\theta, \tau + \delta\tau|0) - G(\theta, \tau|0)}{\delta\tau} &= \sigma^2 \left\{ \frac{G(\theta, \tau|\delta y) - G(\theta, \tau|0)}{\delta y^2} \right\} + r \left\{ \frac{G(\theta, \tau|\delta y) - G(\theta, \tau|0)}{\delta y} \right\} \\ &+ \mu[1 - G(\theta, \tau|0)] - \lambda_1(0)G(\theta, \tau|0)[1 - G(\theta, \tau|0)] \end{aligned} \quad (C.2)$$

and

$$\begin{aligned} \frac{G(\theta, \tau + \delta\tau|\tilde{a}) - G(\theta, \tau|\tilde{a})}{\delta\tau} &= \sigma^2 \left\{ \frac{G(\theta, \tau|\tilde{a} - \delta y) - G(\theta, \tau|\tilde{a})}{\delta y^2} \right\} \\ &+ r \left\{ \frac{G(\theta, \tau|\tilde{a} - \delta y) - G(\theta, \tau|\tilde{a})}{\delta y} \right\} + \mu[1 - G(\theta, \tau|\tilde{a})] - \lambda_1(\tilde{a})G(\theta, \tau|\tilde{a})[1 - G(\theta, \tau|0)] \end{aligned} \quad (C.3)$$



when there is a reflecting barrier at $y = \tilde{a}$, or

$$G(\theta, \tau | \tilde{a}) = 1 \quad (C.4)$$

for all $\tau > 0$, when there is an absorbing barrier at \tilde{a} . As mentioned above, the initial condition for the difference equations is

$$G(\theta, 0 | y) = 0. \quad (C.5)$$

Equations (C.1) thru (C.5) were solved using program PROB. The notation for PROB is as follows: The arrays $G(I)$ and $XG(I)$ contain values of $G(\theta, \tau | i\delta y)$ and $G(\theta, \tau + \delta\tau | i\delta y)$ respectively. Similarly, the quantities $BH(I)$, S , R , DH , DT and DX denote $\lambda_1(i\delta y)$, $\frac{\sigma^2}{2}$, r , μ , $\delta\tau$, and δy respectively in (C.1) thru (C.5). The number M represents the total number of mesh points in the discretizing scheme with $M-1$ equal divisions of the interval $[0, \tilde{a}]$. The program PROB given on the following pages characterizes only one set of parameters from Table 5.2.



```

PROGRAM PROB(INPUT,OUTPUT)
DIMENSION G(201),XG(201),BH(201)
DIMENSION XT(11)
SECH(Z)=2./(EXP(Z)+EXP(-Z))
C*****DATA
DATA S,R,DH,DX,DT/.025,1.0,.6,.2,.01/
TM=0.
M=101
M1=M-1
NK=25
N=30
DTX=DT/(DX*DX)
DTY=DT/DX
XY=20.*N/FLOAT(M1)
C*****SET INITIAL CONDITION.
DO 15 I=1,M
15 G(I)=0.
C*****DEFINE ARRAY BH(I)
DO 10 I=1,M
Z=(I-1)*DX-6.
10 BH(I)=SECH(Z)*50.
C*****COMPUTATION OF THE PROFILE AT TIME T+DT. THE VALUES ARE PRINTED AFTER
C*****EVERY NK ITERATIONS
XT(1)=0.
DO 5 I=2,11
5 XT(I)=XT(I-1)+10.*DX
PRINT 100,DH
PRINT 99
PRINT 98
PRINT 97,(XT(I),I=1,11)
PRINT 99
PRINT 95
PRINT 96,(BH(I),I=1,101,10)
PRINT 94

```



```

DO 71 KQK=1,4
DO 70 JK=1,7
PRINT 101,7M,(G(I),I=1,M,10)
DO 65 J=1,NK
  G11=1.-G(1)
C*****XG(1) INCORPORATES THE REFLECTING BOUNDARY CONDITION AT Y = 0.
  XG(1)=2.*(G(2)-G(1))*DTX*S+R*DTY*(G(2)-G(1))+G11*DT*(DH-BH(1))*G(1)
    1)+G(1)
DO 25 I=2,M1
  XG(I)=DTX*S*(G(I+1)-2.*G(I)+G(I-1))+R*DTY*(G(I+1)-G(I))+DH*DT*(1.-
    1G(I))-BH(I)*G(I)*G11*DT+G(I)
25 CONTINUE
  XG(M)=1.
C*****UPDATE THE ARRAY G(I)
DO 26 I=1,M
  G(I)=XG(I)
  TM=TM+DT
65 CONTINUE
70 CONTINUE
  NK=NK+NK
71 CONTINUE
94 FORMAT(* -----*)
1-----
95 FORMAT(17X,*BIRTH*)
96 FORMAT(17X,*RATE= *,11(F5.2,1X))
97 FORMAT(22X,11(1X,F5.2))
98 FORMAT(/,17X,* T .....X = INITIAL MATURITY OF THE ANCESTOR
  1R.....*)
99 FORMAT(/,* -----*)
1-----
100 FORMAT(#1*,///,12X,*EXTINCTION PROBABILITY FOR THE POPULATION WHE
  1N THE DEATH RATE IS *,F5.2,*YEARS .*)
101 FORMAT(17X,F5.2,11(1X,F5.3))
      STOP
      END

```



APPENDIX D

ANALYSIS OF EXPERIMENTAL DATA

The experimental data on the CLB populations was tested for goodness of fit with a Poisson distribution using the χ^2 test. The program CHSQR was written to do this analysis on an IBM 1800 computer equipped with a keyboard, typewriter, card reader and printer. The logical unit numbers 1, 2, 3 and 6 in the READ and WRITE statements refer to the typewriter, card reader, line printer and keyboard respectively. The function IDTSW takes signals from the data switches on the main console to go from one part of the program to another.

The program CHSQR calculates the χ^2 statistics for N replicate populations in Experiments A and B to test the goodness of fit of the data with the Poisson distribution having the same mean as the data. The program was thus used to compute

$$\chi^2 = \sum \frac{(n_i - n_i^*)^2}{n_i^*} \quad (D.1)$$

where n_i = the observed frequency of the number i (or a combination of numbers i_1, i_2, \dots , etc.) and

n_i^* = the frequency of occurrence of the number i (or the same combination of numbers as in n_i) computed using the distribution against which the sequence of numbers is to be tested.



This value of χ^2 was then compared with that from a standard table of χ^2 distribution for a 10% level of significance and $k-2$ degrees of freedom, where k is the number of terms in the summation in (D.1). If the computed value of χ^2 was greater than that from the tables, the hypothesis that the sequence of numbers (representing the populations in different replicates) has a Poisson distribution was rejected. For a greater reliability, the numbers in the sequence to be tested are usually combined so that each n_i represents one or a combination of more than one numbers in the sequence such that each n_i is always greater than or equal to five. This convention was followed for the computations with few exceptions, in which one of the n_i 's was equal to four.

In program CHSQR $IX(I,J)$ contains a two-dimensional array of populations of different life stages on a given day, with I denoting the serial number of the replicate experiment and J representing the life stages of the CLB. The array $IX(I,J)$ is to be read in as data from cards. The working array $X(I)$ consists of populations of a particular life stage or a combination of the numbers of individuals in different life stages in the replicate experiments. The program uses subrouting MOMNT to compute the first four central moments (i.e., mean, variance, skewness and kurtosis, cf. Himmelblau, 1970) of the numbers in $X(I)$. Next, the frequencies of occurrence of various numbers in the sequence $X(I)$ are computed -- $IKQ(J)$ represents the number of times $KQ(J) = 0, 1, 2, \dots$, etc. occurs in the sequence $X(I)$. This as well as the calculation of the (Poisson) probabilities of occurrence $POIS(I)$ of $KQ(I)$ is done using subrouting POISN. To evaluate χ^2 as in



(D.1), different terms from the sequence $KQ(I)$ are to be combined to make each n_i greater than or equal to five. This is done by selecting the appropriate numbers from the sequence $KQ(I)$ and feeding these as data for evaluation of each n_i using the keyboard. The individual terms in the summation in (D.1) are computed and printed as $CHI(I)$. The variable $CHISQ$ denotes the final value of χ^2 .



```

*****
**STATISTICAL ANALYSIS OF CEREAL LEAF BEETLE POPULATION DATA WITH POISSON INPUT
**
DIMENSION IX(30,8),X(30),KQ(20),IKQ(30),POIS(20),KQX(10),CHI(10)
DIMENSION AXI(80)
N=30
MN=17

*****CONTAINS THE INSTRUCTIONS REGARDING
*****MANIPULATION OF DATA SWITCHES, OR ANY OTHER RELEVANT COMMENT. AL=A FOR
*****EXPERIMENT A, AL=8 FOR EXPERIMENT B. IDAY = DAY, AS COUNTED FROM THE START
*****OF THE EXPERIMENT, JB SPELLS OUT THE LIFE STAGE AND JA = THE LOCATION OF
*****THE LIFE STAGE IN THE ARRAY IX(I,J).IDA=IDAY, JB=JA

      1 READ(2,50)(AXI(I),I=1,80)
      WRITE(1,50)(AXI(I),I=1,80)
      2 READ(2,52)A,IDAY,B,JA
      READ(2,53)AL,IDA,JB,(IX(I,JA),I=1,N)
*****CHECK WHETHER THE DATA CARDS ARE IN PROPER ORDER.
      IF(IDA-IDAY)99,9,99
      9 IF(AL-A)99,10,99
      10 IF(JB-JA)99,11,99
      11 IF(IDAY)12,32,12
      12 WRITE(3,54)A,B,IDAY
      DO 13 I=1,N
      13 X(I)=IX(I,JB)
      WRITE(1,57)B,IDAY
      212 WRITE(3,55)(X(I),I=1,N)
*****COMPUTE FIRST FOUR CENTRAL MOMENTS
      CALL MOMNT(X,N,AM)
*****CALCULATE FREQUENCIES AND CORRESPONDING POISSON PROBABILITIES
      CALL POISN(X,M,MN,POIS,KQ,IKQ,N,AM)
      KKK=0

```



```

C*****CALCULATE CHI SQUARE
C*****I IS THE NO. OF TERMS IN ABOVE FREQUENCIES TO ADJUST THE FREQUENCY TO (II)
C*****FOR A GOOD APPROXIMATION TO CHI SQUARE DISTRIBUTION (II) SHOULD NOT BE LES
C*****S THAN FIVE.
213 READ(6,56)I,II,III
    IF(III)16,214,16
214 IF(I)14,16,14
    14 KKK=KKK+1
    READ(6,56)(KQX(J),J=1,I)
    WRITE(3,58)(KQX(J),J=1,I)
    WRITE(3,59)II,KKK
    SUM=0.
    DO 15 J=1,I
    IJ=KQX(J)
15 SUM=SUM+POIS(IJ)
    FJ =N*SUM
    CHI(KKK)=((FLOAT(II)-FJ)**2.)/FJ
    GO TO 213
16 CHISQ=0.
    DO 17 I=1,KKK
17 CHISQ=CHISQ+CHI(I)
    WRITE(3,102)
    WRITE(3,55)(CHI(I),I=1,KKK)
    WRITE(3,103)CHISQ
C*****RECALCULATE CHI SQUARE WITH DIFFERENT COMBINATIONS OF THE DATA
C
C
C*****COMBINE UNHATCHED AND HATCHING EGGS
18 IF(IDTSW(0))19,21,19
19 DO 20 I=1,N
20 X(I)=IX(I,1)+IX(I,2)
    WRITE(3,104)
    WRITE(1,104)
    GO TO 212

```



```

C*****COMBINE HATCHING EGGS WITH FIRST INSTAR POPULATION
21 IF(IDTSW(1))22,24,22
22 DO 23 I=1,N
23 X(I)=IX(I,2)+IX(I,3)
  WRITE(3,105)
  WRITE(1,105)
  GO TO 212
C*****COMBINE ALL LARVAL INSTARS TOGETHER
24 IF(IDTSW(2))25,27,25
25 DO 26 I=1,N
  X(I)=0.
  DO 26 J=3,6
26 X(I)=IX(I,J)+X(I)
  WRITE(3,106)
  WRITE(1,106)
  GO TO 212
C*****COMBINE ALL LARVAL INSTARS AND HATCHING EGGS TOGETHER
27 IF(IDTSW(3))28,30,28
28 DO 29 I=1,N
  X(I)=0.
  DO 29 J=2,6
29 X(I)=X(I)+IX(I,J)
  WRITE(3,107)
  WRITE(1,107)
  GO TO 212
C*****COMBINE ALL LARVAE AND PUPAE TOGETHER
30 IF(IDTSW(5))34,36,34
34 CONTINUE
  DO 35 I=1,N
  X(I)=0.
  DO 35 J=2,7
35 X(I)=X(I)+IX(I,J)
  WRITE(3,109)
  WRITE(1,109)
  GO TO 212

```



```

36 IF(IDTSW(4))31,1,31
31 GO TO 32
99 WRITE(3,108)
32 CALL EXIT
50 FORMAT(80A1)
52 FORMAT(A1,I2,A6,I1)
53 FORMAT(A1,I2,I1,I1,30I2)
54 FORMAT('1 EXPERIMENT ',A1,' POPULATION OF ',A6,' ON ',I2,'TH DAY',
1)
55 FORMAT(1X,10F10.5)
56 FORMAT(25(1X,I2))
57 FORMAT(///,1X,A6,1X,I2,'TH DAY')
58 FORMAT(' TERMS COMBINED ARE NO. ',6(1X,I2))
59 FORMAT('+,45X,' TO ADJUST FREQUENCY TO ',I2,'FOR ',I2,'TH TERM IN
1CHI SQUARE TEST')
102 FORMAT('0 INDIVIDUAL TERMS IN SUMMATION FOR CHI SQUARE TEST')
103 FORMAT('0 CHI SQUARE =',F12.5)
104 FORMAT('1 REPEAT CHI SQUARE TEST WITH EGGS (HATCHED AND UNHATCHED)
1LUMPED TOGETHER')
105 FORMAT('1 REPEAT CHI SQUARE TEST WITH 1ST INSTAR LARVAE AND HATCHI
1NG EGGS LUMPED TOGETHER')
106 FORMAT('1 REPEAT CHI SQUARE TEST WITH ALL LARVAL POPULATION LUMPED
1TOGETHER')
107 FORMAT('1 REPEAT CHI SQUARE TEST WITH LARVAE AND HATCHING EGGS LUM
1PED TOGETHER')
108 FORMAT('0DATA CARDS PROBABLY OUT OF SEQUENCE--CHECK')
109 FORMAT(' REPEAT CHI SQUARE TEST WITH LARVAE AND PUPAE LUMPED TOGET
1HER')
END

```



```

SUBROUTINE MOMNT(X,N,AM)
DIMENSION X(1),Y(30),Z(30)
C*****COMPUTE MEAN
SUM=0.
DO 6 I=1,N
6 SUM=SUM+X(I)
AM=SUM/N
DO 7 I=1,N
7 Z(I)=X(I)-AM
7 Y(I)=Z(I)*Z(I)
C*****COMPUTE VARIANCE
SUM=0.
DO 8 I=1,N
8 SUM=SUM+Y(I)
VAR=SUM/(N-1)
C***** COMPUTE SKEWNESS
DO 9 I=1,N
9 SUM=SUM+Y(I)*Z(I)
SKEW=SUM/(N-2)
C*****COMPUTE FOURTH CENTRAL MOMENT (KURTOSIS).
SUM=0.
DO 10 I=1,N
10 SUM=SUM+Y(I)*Y(I)
A4TH=SUM/(N-3)
WRITE(3,102)N,AM,VAR,SKEW,A4TH
102 FORMAT('O SAMPLE SIZE= ',I3,'MEAN= ',F8.5,'VARIANCE= ',F8.5,' THIR
10 MOMENT= ',E12.5,'FOURTH MOMENT= ',E12.5)
RETURN
END

```



```

SUBROUTINE POISN(X,M,MN,MN,POIS,KQ,IKQ,N,AM)
DIMENSION X(1),POIS(1),KQ(1),IKQ(1)
C*****EVALUATE ARRAYS KQ(I) AND IKQ(I).
IJ=0
DO 18 I=1,30
18 IKQ(I)=0
DO 31 I=1,MN
IND=0
DO 30 J=1,N
IJK=X(J)
LX=IJK-I+1
IF(IND)19,19,23
19 IF(LX)23,20,23
20 IND=1
IJ=IJ+1
KQ(IJ)=IJK
23 IF(LX)30,24,30
24 IKQ(IJ)=IKQ(IJ)+1
30 CONTINUE
31 CONTINUE
C*****COMPUTE PROBABILITIES POIS(I) OF KQ(I) FOR A POISSON DISTRIBUTION WITH
C*****MEAN = AM.
AMX=AM
POISX=EXP(-AMX*.5)
IJ=1
IF(KQ(1))32,29,32
29 POIS(IJ)=EXP(-AMX)
IJ=IJ+1
32 DO 38 I=1,100
IF(I-KQ(IJ))38,33,38
33 POIS(IJ)=POISX
II=KQ(IJ)
DO 34 J=1,II
34 POIS(IJ)=POIS(IJ)*AMX/J
POIS(IJ)=POIS(IJ)*POISX
IJ=IJ+1

```



```

38 CONTINUE
KK=0
DO 40 I=1,N
  IF(IKQ(I))39,40,39
39 KK=KK+1
40 CONTINUE
  WRITE(3,110)
110 FORMAT('O POPULATIONS OF SAMPLES, THEIR FREQUENCIES IN 30 REPLICATE
15 AND CORRESPONDING POISSON PROBABILITIES')
  WRITE(3,56)(KQ(I),I=1,KK)
  WRITE(3,56)(IKQ(I),I=1,KK)
  WRITE(1,56)(KQ(I),I=1,KK)
  WRITE(1,56)(IKQ(I),I=1,KK)
56 FORMAT(1X,25(1X,I2))
  WRITE(3,101)(POIS(I),I=1,KK)
101 FORMAT(1X,10F10.5)
  RETURN
END

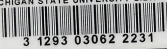
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