



THESIS



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thesis entitled

INDUCED ELECTROMAGNETIC FIELDS IN CONDUCTING
BODIES IRRADIATED BY MAGNETIC FIELDS
AND MICROWAVES

presented by

Jen-Hwang Lee

has been accepted towards fulfillment
of the requirements for

Ph. D. degree in Elect. Engr. & Sys. Sci.

Major professor

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INDUCED ELECTROMAGNETIC FIELDS IN CONDUCTING
BODIES IRRADIATED BY MAGNETIC FIELDS
AND MICROWAVES

By

Jen-Hwang Lee

A DISSERTATION

Submitted to
Michigan State University
in partial fulfillment of the requirements
for the degree of

DOCTOR OF PHILOSOPHY

Department of Electrical Engineering
and Systems Science

1979

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ABSTRACT

INDUCED ELECTROMAGNETIC FIELDS IN CONDUCTING BODIES IRRADIATED BY MAGNETIC FIELDS AND MICROWAVES

By

Jen-Hwang Lee

A new theoretical method for determining the eddy current induced by a uniform RF magnetic field or a beam of RF magnetic field in a biological body is developed in this research. This study was motivated by the fact that more biological research and medical applications utilize the irradiation of RF magnetic fields, and by the need for an efficient method of determining the magnetic mode of the induced electric field in a thick biological body irradiated by HF-VHF EM waves. The body of rotational symmetry is subdivided into a number of circular rings with various radii and cross-sectional areas. The eddy current induced by the impressed magnetic field can be considered, physically, as induced by a circulatory impressed electric field associated with the uniform impressed magnetic field. Induced eddy currents in all the rings maintain a scattered electric field which can be added to the impressed electric field to yield the total induced electric field in the body.

69

At the same time, the complex conductance of these relations between simultaneous equations can be obtained, and with the present magnetic fields, biological bodies have been conventionally used, but fails to yield a 20 MHz, while the accuracy and efficiency of biological bodies

Also included in quantifying the irradiated biological integral equation, the quantification of the irradiated body has been equation methods. The human body is in the range, the body has an EM field concentration on the surface. A volume is inadequate or inefficient is more efficient

At the same time, the eddy current is the product of the complex conductivity and the induced electric field. Using these relations and the point-matching method, a set of simultaneous equations for the induced electric fields can be obtained, and they are subsequently solved numerically. With the present method, the eddy currents induced by magnetic fields of 1 to 200 MHz in various biological bodies have been obtained. It was found that the conventionally used, quasi-static solution for the eddy current fails to yield accurate results for frequencies higher than 20 MHz, while the present method can be used to determine accurately and efficiently the induced eddy currents in biological bodies for frequencies higher than 20 MHz.

Also included in this thesis is a numerical method for quantifying the induced EM field on the surface of an irradiated biological body based on two coupled, surface integral equations. In the field of theoretical dosimetry, the quantification of the induced EM field inside an irradiated body has been performed mainly by volume integral equation methods. When a large biological body such as the human body is irradiated by an EM wave of microwave range, the body becomes electrically large and the induced EM field concentrates mainly in a thin layer of the body surface. A volume integral equation method then becomes inadequate or inefficient to handle this problem, and it is more efficient to quantify the induced EM field on the

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body surface based on a surface integral equation method for this case. After the tangential components of the induced electric and magnetic fields on the body surface are determined, the internal EM field can be calculated. To check the accuracy of this method, electrically small bodies are considered first, and results obtained from the surface integral equation method are compared with that obtained from the tensor integral equation method, a volume integral equation method. It was learned that, in some cases, the surface integral equation method showed advantages in accuracy and computational cost over the volume integral equation method.

To

To my mother, Kim-Chu Hong Lee,
my wife, Ai-Chih,
and my son, Yueh.

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I wish to express my sincere appreciation to my research advisor, Dr. K. M. Chen, for his constant support and guidance throughout the course of this research. I also wish to thank Dr. D. P. Nyquist for his helpful comments and encouragement about this study. Sincere appreciation is also extended to Dr. M. Siegel and Dr. D. Riska for their valuable help.

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Finally, I thank my wife and my son for their patience and love, which made completion of this work possible.

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CHAPTER I

INTRODUCTION

Since electromagnetic radiation and propagation were discovered more than a century ago, electromagnetic waves of various frequencies have been utilized in many ways to benefit the human society. Especially, due to the ingenious contributions of engineers and scientists, such incredible dreams as satellite communications, long-range radar detection, high power-rating microwave ovens, and many other products related to EM technology have been realized. Moreover, high-level beamed microwave power is proposed as a new scheme for transporting energy in the future, and the use of electromagnetically induced hyperthermia as an adjunct in cancer treatment is receiving increased attention from many medical researchers. Indeed, these highly advanced techniques of utilizing energy in EM form do benefit us in many respects; unfortunately, they also bring us certain unwanted potential hazards.

Over a decade ago, medical personnel and public health officials began to suspect that low-level, long-term ionizing radiation could be one of the major causes for several different cancers. Even more astonishingly, it is found

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impossible to completely avoid this low-level radiation since it exists everywhere in space, originating from either man-made or natural sources. Besides the above possible hazards due to low-level radiation, EM radiation in high intensities, under an uncontrolled condition, can be harmful to living systems. The thermal effects associated with such high-intensity radiation can produce burns, cataracts, and chemical changes, etc. All these effects have become the major concerns of the public, especially for those who are subject to possible long-term or high-intensity radiation exposure. As a result, there is strong demand for a more thorough understanding of the biological effects of EM radiation. There is therefore a demonstrated need to determine, for example, the safe power density for long-term human exposure, the correct applications methodology and techniques in utilizing high-intensity EM radiation for medical treatment, etc. However, based on the present state of knowledge in this area, many questions still remain unanswered. As a matter of fact, there exist strong arguments between involved groups about where realistic safety levels for microwave exposure should lie. It is surprising to observe that the maximum allowable safe power density for long-term human exposure varies from 10 mW/cm^2 in this country to as low as 0.01 mW/cm^2 in the USSR.

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To clarify these uncertainties, electromagnetic engineers should assume the responsibility for providing more complete, accurate, quantitative data on biological effects of EM radiation; they belong to the professional groups having adequate knowledge and ability to analyze this problem both theoretically and experimentally.

Under these circumstances, we complete the present research for the purpose of seeking for a better understanding of EM interactions with biological bodies. We study the eddy currents induced by RF magnetic fields inside biological bodies in Chapter II. The RF magnetic field has received increased attention due to its potential usefulness for local heating in cancer treatment. Based on our research, we obtain a better understanding about the mechanism exciting induced currents inside the bodies; furthermore, we are able to predict accurately the current distributions inside the bodies, which has not been accomplished by previous methods. In Chapter III, we study the microwave interactions with biological bodies through the solutions of two coupled, surface integral equations. This numerical technique can be applied to quantify the EM fields induced on the surfaces of arbitrarily shaped biological bodies. Through this study, we expect to obtain a more complete, accurate, quantitative evaluation of microwave interactions. Chapter IV describes

in detail the

problems. Th

in detail the computer programs used for the above two problems. Then, Chapter V summarizes this research.

EDDY

In this case the method for determining the current induced by the RF magnetic field in a body of rotational symmetry is as follows.

The body is divided into a number of circular rings of circular ring areas. The induced current in each ring is numerically determined by the induced potential and the magnetic field and the results are summed to give the total current. The results deviate significantly from the analytical solutions. An analysis of the electric fields in the conducting phase found to be in good agreement with the accuracy of the

CHAPTER II

EDDY CURRENTS INDUCED BY RF MAGNETIC FIELDS INSIDE A FINITE CONDUCTING BODY WITH ROTATIONAL SYMMETRY

In this chapter, we will present a new theoretical method for determining the electric field or the eddy current induced by a uniform RF magnetic field or a beam of RF magnetic field in a finite conducting body with rotational symmetry.

The body being considered is divided into a number of circular rings with various radii and cross-sectional areas. The induced electric field in each ring is then numerically determined based on the theory of vector potential and the moment method. Numerical examples are given and the results based on the present theory are found to deviate significantly from the often-used quasi-static solutions. An experiment was conducted to measure the electric fields induced by a UHF magnetic field in finite conducting phantom models. Experimental results were found to be in a good agreement with the theory. The accuracy of the present theory was also verified by the

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2.1. Introduction

Due to its increasing applications in the biomedical area, the radiation of RF magnetic field is becoming an important new technique which has drawn a large amount of attention from many investigators. Conventionally, the electric field or the eddy current induced by a magnetic field inside a finite conducting body is estimated by the quasi-static solution. Unfortunately, the quasi-static approximation can give accurate results only when the body is relatively small or when the frequency of the applied magnetic field is lower than about 20 MHz. If an electrically large body such as human body is exposed to a RF magnetic field with a frequency higher than 20 MHz, the quasi-static approximation becomes inadequate. It is desirable to solve this practical problem accurately and effectively.

Based on Mie theory, Lin et al. [1] have obtained a quasi-static solution for the magnetic mode of the electric field induced by a HF EM wave inside a spherical model of man. Their theory is valid up to 20 MHz and restricted to a spherical body. Spiegel [2] has used the same quasi-static solution to estimate the electric field induced by the magnetic field of an EHV power line inside a simplified

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model of man. His solution is valid only for the very low frequency range. A theory on the eddy current induced by an ac uniform magnetic field in a conducting sphere was presented by Van Bladel [3]. This theory imposes less approximations but only applies to a spherical geometry (see Section 2.5). There is a need for a new theoretical method which can be used to quantify the electric fields or the eddy currents induced by RF magnetic fields of up to VHF range in finite conducting bodies with geometries more complex than a sphere.

Another motivation of the present study is the need for improving the efficiency of a recent numerical method [4] developed by our group for quantifying the internal electric field induced by an EM wave in a human body. When a thick biological body is exposed to an EM wave of HF-VHF range, the induced electric field or current can be divided into the electric and magnetic modes. The electric mode being linear and relatively independent of the body shape is easily determined by the numerical method. However, the magnetic mode is circulatory and strongly dependent on the body shape. The determination of this magnetic mode with a numerical method often encounters the difficulty of numerical convergence. Thus, a new theoretical method which can efficiently determine the magnetic mode of the induced current in an irradiated biological body is desired.

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2.2. Theory

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The theoretical method presented in this chapter is capable of predicting the electric field or the eddy current induced by a uniform RF magnetic field or a beam of RF magnetic field in a finite conducting body with rotational symmetry. The development of the theory is presented in Section 2.2. Some numerical examples are given in Section 2.3. The convergence of numerical results is discussed in Section 2.4. The comparison of the present solution with a closed form solution for a sphere is made in Section 2.5. The experimental setup is described in Section 2.6, and experimental results are compared with theoretical results in Section 2.7.

2.2. Theory

When a finite conducting body with rotational symmetry is exposed to a uniformly impressed, RF magnetic field $\vec{H}^i$, the induced electric field $\vec{E}$ inside the body is conventionally estimated by Maxwell equation of

$$\oint \vec{E} \cdot d\vec{\ell} = -j\omega\mu_0 \int \vec{H}^i \cdot d\vec{s}. \quad (2.1)$$

The induced electric field $\vec{E}$ determined from the above equation is circulatory and its amplitude increases linearly with the radial distance from the central axis of the body and its phase is 90° out of that of impressed magnetic field $\vec{H}^i$. In this quasi-static solution, the interaction between the induced currents at different

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locations is completely ignored. However, to obtain accurate results, the scattered magnetic field produced by the induced current must be taken into consideration. The quasi-static approximation can predict good results only for frequencies lower than 20 MHz. A theoretical method which can be used to quantify the electric field in a finite conducting body with a rotational symmetry induced by a RF magnetic field with a frequency higher than 20 MHz is developed below.

We consider two different cases separately: (1) a finite conducting body immersed in an impressed, uniform RF magnetic field, and (2) the body irradiated by a cylindrical beam of RF magnetic field.

Physically, we may consider that the electric field in the body is induced by an impressed electric field which is associated with the impressed RF magnetic field.

For the case of an impressed, uniform RF magnetic field $\vec{H}^i$ as shown in Figure 2.1, the associated, impressed electric field $\vec{E}^i$ can be obtained from

$$\oint \vec{E}^i \cdot d\vec{l} = -j\omega\mu_0 \int \vec{H}^i \cdot d\vec{s}$$

as

$$\vec{E}^i = -\frac{j}{2}\omega\mu_0 H^i r \hat{\phi} \quad (2.2)$$

where r is the radial distance from the central axis of the body.



Figure 2.1

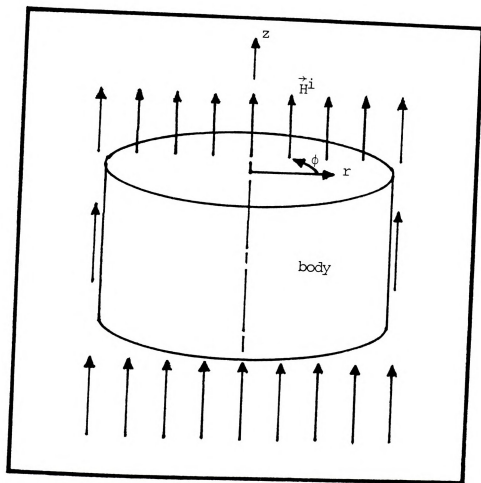


Figure 2.1. A biological body is exposed to a uniformly impressed RF magnetic field.

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For the case of a cylindrical beam of uniform magnetic field $\vec{H}^i$ with a beam radius of "b" as shown in Figure 2.2, the associated, impressed electric field can be found to be

$$\vec{E}^i = \begin{cases} -\frac{j}{2}\omega\mu_0 H^i r \hat{\phi} & \text{for } 0 \leq r < b \\ -\frac{j}{2}\omega\mu_0 H^i \frac{b^2}{r} \hat{\phi} & \text{for } b \leq r. \end{cases} \quad (2.3)$$

These impressed electric fields $\vec{E}^i$ associated with the impressed magnetic fields $\vec{H}^i$ are in the azimuthal direction and they induce circulatory electric fields inside the finite conducting body.

If the rotational symmetry is assumed for the body geometry and the electrical properties of the body, the body can be subdivided into a number (N) of circular rings of various radii and cross-sectional areas as shown in Figure 2.3. We further assume that the induced electric field within each ring is uniform, but it can vary from ring to ring. Let's consider two sample rings shown in Figure 2.3; namely, the nth ring with radius r_n , cross-sectional area S_n and a reference point $(r_n, 0, z_n)$, and the mth ring with radius r_m , cross-sectional area S_m and a reference point $(r_m, 0, z_m)$. The geometries of these two rings are depicted in Figure 2.4, where the reference points of the rings, $\vec{r}_n$ and $\vec{r}_m$, the source point in the nth ring, $\vec{r}_n$, and the distance between the source point



Figure 2.2

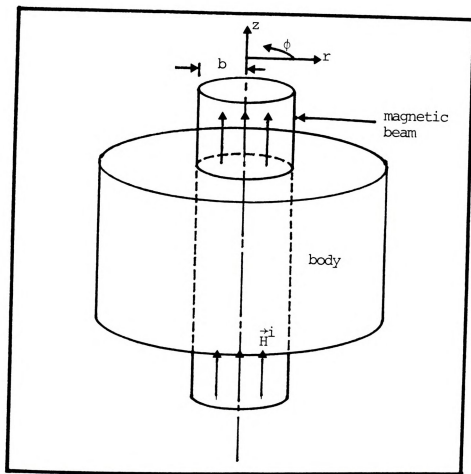


Figure 2.2. A biological body is irradiated by a beam of RF magnetic field.



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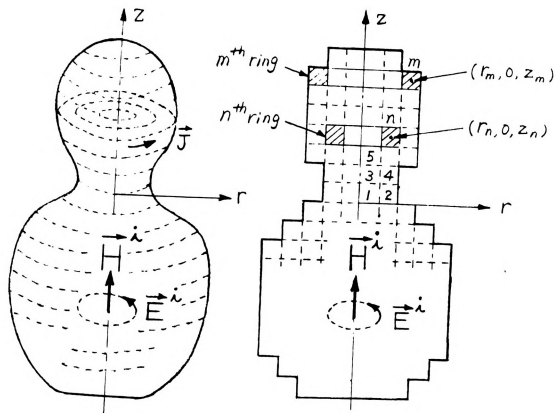
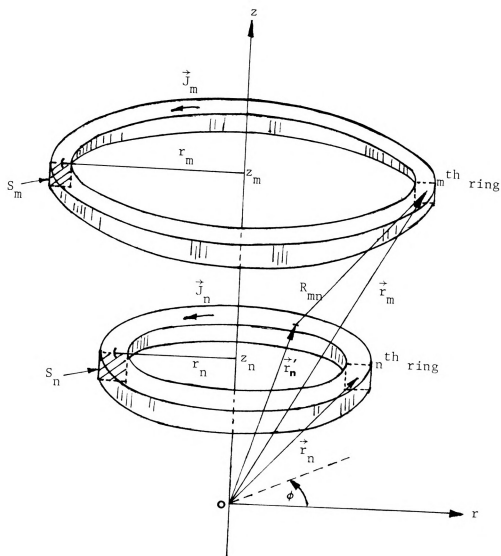


Figure 2.3. A biological body of rotational symmetry irradiated by a uniform RF magnetic field is subdivided into a number of rings of various radii.



$$\begin{aligned} \vec{r}_n &= (r_n, 0, \\ \vec{r}_m &= (r_m, 0, \\ \vec{r}_n' &= (r_n', \phi, \\ R_{nm} &= |\vec{r}_m - \vec{r}_n'| \\ &= \text{distance} \\ &= \text{reference} \end{aligned}$$

Figure 2.4.



- $\vec{r}_n = (r_n, 0, z_n)$ = reference point of the n^{th} ring
 $\vec{r}_m = (r_m, 0, z_m)$ = reference point of the m^{th} ring
 $\vec{r}_n = (r_n, \phi, z_n)$ = source point in the n^{th} ring
 $R_{mn} = |\vec{r}_m - \vec{r}_n| = [r_m^2 + r_n^2 - 2r_m r_n \cos \phi + (z_m - z_n)^2]^{1/2}$
 = distance between the source point in the n^{th} ring and the reference point of the m^{th} ring.

Figure 2.4. Geometries of two rings in the subdivided biological body.

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If the volume density of the induced current in the n th ring is denoted as $\vec{J}_n$, the vector potential at $\vec{r}_m$ of the m th ring maintained by $\vec{J}_n$ of the n th ring, $\vec{A}_{mn}$, can be determined as

$$\vec{A}_{mn} = \frac{\mu_0}{4\pi} \int_{\text{nth ring}} \vec{J}_n \frac{e^{-jk_0 R_{mn}}}{R_{mn}} dv \quad (2.4)$$

$$\text{where } R_{mn} = [r_m^2 + r_n^2 - 2 r_m r_n \cos\phi + (z_m - z_n)^2]^{\frac{1}{2}}$$

$$k_0 = \omega \sqrt{\mu_0 \epsilon_0}.$$

Using $dv = S_n r_n d\phi$ and because of the rotational symmetry, eq. (2.4) can be evaluated as

$$\vec{A}_{mn} = \hat{\phi} \frac{\mu_0}{2\pi} J_n S_n r_n K_{mn} \quad (2.5)$$

$$\text{where } K_{mn} = \int_0^\pi d\phi \cos\phi \frac{e^{-jk_0 R_{mn}}}{R_{mn}}. \quad (2.6)$$

To determine the vector potential $\vec{A}_{mm}$ at $\vec{r}_m$ of the m th ring maintained by $\vec{J}_m$, the induced current in the m th ring, we consider the geometry of the m th ring shown in Figure 2.5. Note that we approximate the square ring cross-section by a circle with the same area. This leads to the following relationship:

top view
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ring

side view
of the mth
ring

Figure 2.

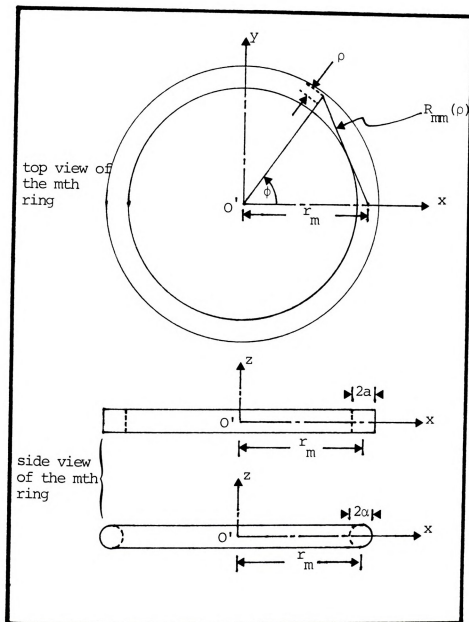


Figure 2.5. Geometries of a single ring inside the subdivided biological body.

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$$\vec{A}_m = \sum_n$$

$$\alpha = \frac{2a}{\sqrt{\pi}}$$

where $2a$ is the dimension of the square cross-section, while α is the radius of the circular cross-section.

The distance between the source point inside the m th ring and the reference point $\vec{r}_m$, R_{mm} , can be expressed as

$$R_{mm}(\rho) \doteq r_m [2(1 - \cos\phi) + \frac{\rho^2}{r_m^2}]^{\frac{1}{2}}.$$

The vector potential $\vec{A}_{mm}$ can be written in terms of R_{mm} as

$$\vec{A}_{mm} = \frac{\mu_0}{4\pi} \int_{\text{ring}} \vec{J}_m \frac{e^{-jk_0 R_{mm}(\rho)}}{R_{mm}(\rho)} dv.$$

After substituting the expression of R_{mm} and some straightforward manipulations, we have the following form for the above equation:

$$\vec{A}_{mm} = \hat{\phi} \frac{\mu_0}{2\pi} J_m S_m r_m K_{mm} \quad (2.7)$$

$$\text{where } K_{mm} = \frac{2\pi}{S_m} [R_{mm}(\alpha) - R_{mm}(0)] \int_0^\pi \cos\phi e^{-jk_0 R_{mm}(0)} d\phi. \quad (2.8)$$

The total vector potential at $\vec{r}_m$ maintained by the induced currents in all the rings, $\vec{A}_m$, can then be obtained as

$$\vec{A}_m = \sum_n \vec{A}_{mn} = \hat{\phi} \sum_n \frac{\mu_0}{2\pi} S_n r_n K_{mn} J_n. \quad (2.9)$$

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At the same time, the induced current $\vec{J}_n$ in the nth ring can be related to the total induced electric field in the nth ring, $\vec{E}_n$, as

$$\vec{J}_n = [\sigma_n + j \omega (\epsilon_n - \epsilon_0)] \vec{E}_n = \tau_n \vec{E}_n. \quad (2.10)$$

Equation (2.10) expresses $\vec{J}_n$ as the sum of the conduction current and the polarization current, and τ_n can be considered as the complex conductivity of the nth ring.

With eq. (2.10), $\vec{A}_m$ can be rewritten as

$$\vec{A}_m = \hat{\phi} \sum_n \frac{\mu_0}{2\pi} S_n r_n K_{mn} \tau_n E_n. \quad (2.11)$$

This $\vec{A}_m$ is related to the scattered electric field $\vec{E}_m^s$ in the mth ring which is maintained by the induced currents in all the rings:

$$\vec{E}_m^s = -j\omega \vec{A}_m = -\hat{\phi} \sum_n j f \mu_0 S_n r_n K_{mn} \tau_n E_n \quad (2.12)$$

where f is the frequency. It should be emphasized that since the induced current in each ring is rotationally symmetrical, no electric charge is induced and, consequently, no scalar potential is maintained.

The total induced electric field $\vec{E}_m$ in the mth ring is the sum of the impressed electric field $\vec{E}_m^i$ at the mth ring and the scattered electric field $\vec{E}_m^s$ obtained in eq. (2.12):

$$\vec{E}_m = \vec{E}_m^i + \vec{E}_m^s \quad (2.13)$$

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With eq. (2.12), eq. (2.13) can be rearranged as

$$E_m + \sum_n jf \mu_0 S_n r_n K_{mn} \tau_n E_n = E_m^i. \quad (2.14)$$

Equation (2.14) can be point-matched at N reference points of the N rings to yield N simultaneous equations in a matrix form as

$$\begin{bmatrix} M_{11} & M_{12} & \dots & M_{1N} \\ M_{21} & M_{22} & \dots & M_{2N} \\ \dots & \dots & \dots & \dots \\ M_{N1} & M_{N2} & \dots & M_{NN} \end{bmatrix} \cdot \begin{bmatrix} E_1 \\ E_2 \\ \dots \\ E_N \end{bmatrix} = \begin{bmatrix} E_1^i \\ E_2^i \\ \dots \\ E_N^i \end{bmatrix} \quad (2.15)$$

where

$$M_{nn} = 1 + jf \mu_0 S_n r_n K_{nn} \tau_n \quad (2.16)$$

$$M_{mn} = jf \mu_0 S_n r_n K_{mn} \tau_n. \quad (2.17)$$

The total induced electric fields, E_1 to E_n , in all the rings can be determined from eq. (2.15) by the matrix inversion technique. Note that if the induced eddy current is needed, it can be obtained readily from eq. (2.10).

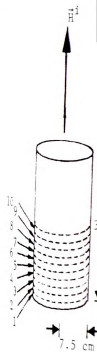
A computer program has been developed based on this method to calculate induced electric fields in various finite conducting bodies. Some numerical results are discussed in the next section.

2.3. Numerical

The first biological body 30 cm being MHz with an have a constant (ϵ_r) of body is subdivided with a square. The amplitude of the field $\vec{E}$ in each of the sections of the radial distance from the different sections is observed in the induced electric field at the radial distance from the solution as a function of the angle of the induced field $\vec{H}^i$ while the angle is -90° at any body, the often estimate the angle of its phase and the critical body is

2.3. Numerical Results

The first numerical example deals with a cylindrical biological body with a diameter of 15 cm and a height of 30 cm being exposed to a uniform magnetic field of 100 MHz with an intensity of 1 Amp/m. The body is assumed to have a conductivity (σ) of 0.889 S/m and a dielectric constant (ϵ_r) of 71.7. In the numerical calculation, the body is subdivided into 100 rings of various radii and with a square cross-sectional area of 1.5 cm x 1.5 cm. The amplitude and phase angle of the induced electric field $\vec{E}$ in each ring were calculated, and the distributions of these quantities are plotted as functions of the radial distance from the cylindrical axis r for 10 different sections of the body as shown in Figure 2.6. It is observed in Figure 2.6 that the amplitude of the induced electric field $|\vec{E}|$ increases linearly with the radial distance; this result is close to the quasi-static solution as can be obtained from eq. (2.1). The phase angle of the induced electric field varies between around -110° and -150° (with respect to the impressed magnetic field $\vec{H}^i$) while the quasi-static solution predicts it to be -90° at any point in the body. For this cylindrical body, the often-used quasi-static solution can be used to estimate the amplitude of the induced electric field but not its phase angle. However, as the size of the cylindrical body is increased, the quasi-static solution is



$\vec{E} = 1 \text{ A/m}$ Ph
 $f = 100 \text{ MHz}$ of
 $\epsilon = 0.889 \text{ S/m}$ $\vec{E}$
 $\theta = 71.7^\circ$ (deg)

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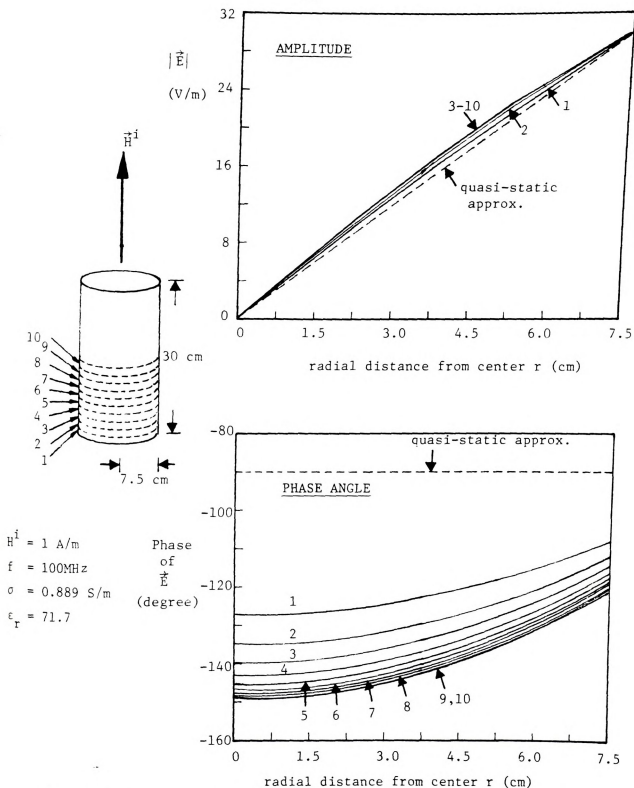


Figure 2.6. Distributions of amplitudes and phase angles of electric fields induced by a 100 MHz magnetic field of 1 A/m in a cylindrical biological body with a diameter of 15 cm and a height of 30 cm.

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found to be completely inadequate, as can be observed in the next example.

In the second example, we consider a cylindrical biological body with the same electric properties ($\sigma = 0.889$ S/m and $\epsilon_r = 71.7$) as the first example, but with a diameter of 36 cm and a height of 90 cm as shown in Figure 2.7. The same magnetic field of 100 MHz with an intensity of 1 Amp/m is impressed on the body. The distributions of the amplitudes and phase angles of the induced electric fields in Figure 2.7 show that the present results deviate greatly from the quasi-static solutions: the amplitude of the induced electric field is two or three times lower than the quasi-static solution, while the phase angles vary between -120° and -290° against the quasi-static solution of -90° . It is clear, from this example, that when a human body is exposed to a magnetic field of VHF range, the often-used, quasi-static solution becomes completely invalid.

The next example as shown in Figure 2.8 deals with the electric fields induced by magnetic fields of unit intensity (1 Amp/m) and of various frequencies (10, 40, 100 and 200 MHz) in a muscle disk with a diameter of 30 cm and a thickness of 1 cm. The electrical properties of the disk are assumed to have $\sigma = 0.625$ to 1.28 S/m and $\epsilon_r = 160$ to 56.5 for the frequency range of 10 to 200 MHz. The amplitudes of the induced electric fields are found to



$\vec{H}^i = 1 \text{ A/m}$ PH
 $f = 100 \text{ MHz}$
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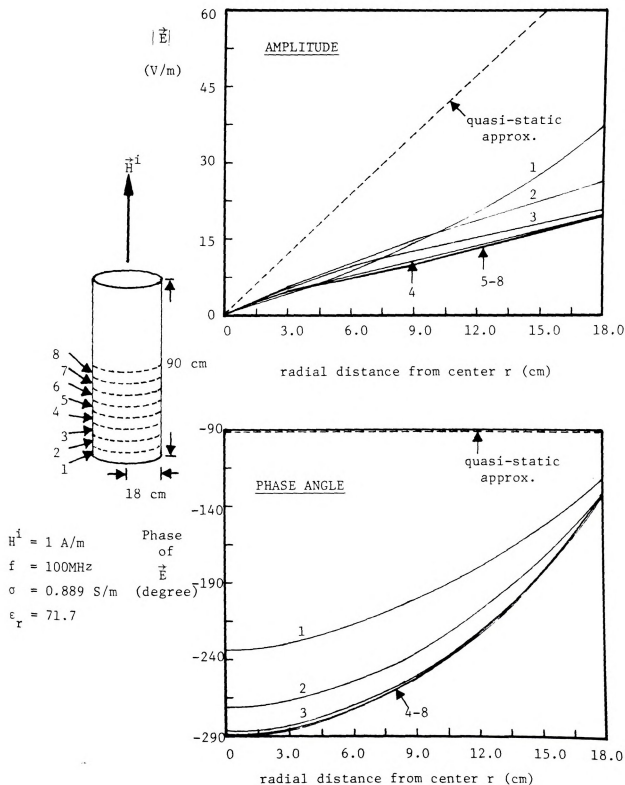


Figure 2.7. Distributions of amplitudes and phase angles of electric fields induced by a 100 MHz magnetic field of 1 A/m in a cylindrical biological body with a diameter of 36 cm and a height of 90 cm.



$B^1 = 1 \text{ A/m}$
 $\sigma = 0.625 \text{ to } 1.28 \text{ S/m}$
 for 10-200 MHz
 $\epsilon_r = 160 \text{ to } 560$
 for 10-200 MHz

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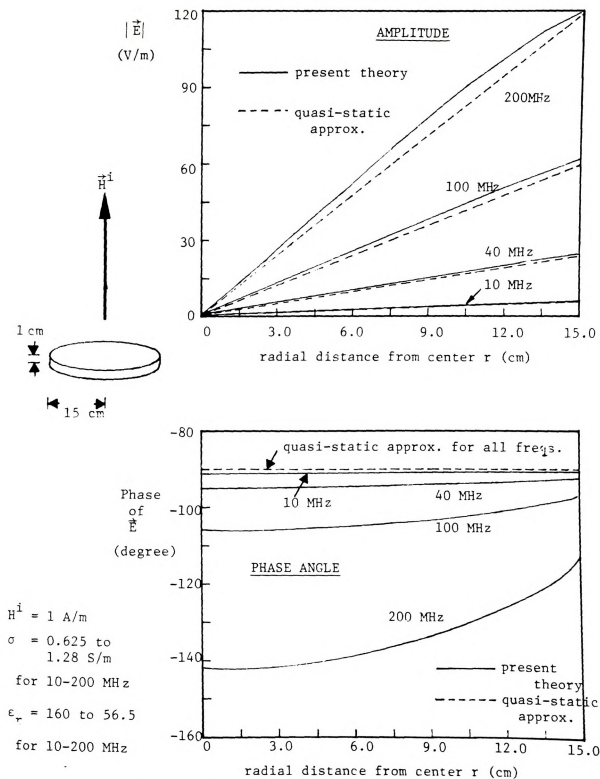


Figure 2.8. Distributions of amplitudes and phase angles of electric fields induced by magnetic fields of various frequencies in a muscle disk with a diameter of 30 cm and a thickness of 1 cm.

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increase linearly with the frequency and with the radial distance from the cylindrical axis; these results deviate slightly from the quasi-static solutions. The present results for the phase angles of the induced electric fields, however, deviate significantly from the quasi-static solutions, especially, for the high frequency cases.

In the example shown in Figure 2.9, we study the effect of the cylindrical height on the induced electric field at the central section of the cylinder. The cylindrical body has a diameter of 36 cm, but its height varies from 6 cm to 90 cm. The impressed magnetic field at 100 MHz has an intensity of 1 Amp/m. The electrical properties of the body at this frequency are assumed to have $\sigma = 0.889$ S/m and $\epsilon_r = 71.7$. The induced electric fields at the central section of the cylinder are determined and the distributions of their amplitudes and phase angles are plotted for the cases of cylindrical heights; $H = 6, 12, 18, 30$ and 90 cm. It is observed in Figure 2.9 that as the cylindrical height is increased, the amplitude of the induced electric field at the central section of the cylinder decreases greatly, and its phase angle varies drastically from the quasi-static solution of -90° . This example indicates that the eddy current induced by a RF magnetic field in a cylindrical biological body is strongly dependent on the cylindrical height.



Phase
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$$H^i = 1 \text{ A/m}$$

$$f = 100 \text{ MHz}$$

$$\sigma = 0.889 \text{ S/}$$

$$\epsilon_r = 71.7$$

Figure 2.9. D
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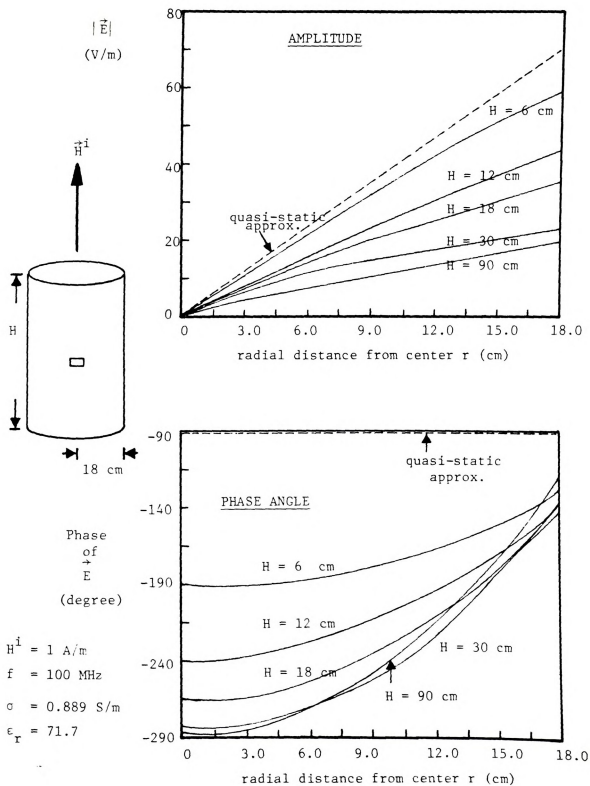


Figure 2.9. Distributions of amplitudes and phase angles of electric fields induced by a 100 MHz magnetic field in the central section of a cylindrical biological body of diameter 36 cm with various heights.

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The case of the eddy current induced by a beam of magnetic field in a biological body is studied next. Figure 2.10 shows the distributions of the amplitudes and phase angles of electric fields induced in a cylindrical body with a diameter of 32 cm and a height of 20 cm, by a beam of 100 MHz magnetic field with a beam diameter of 8 cm and an intensity of 1 Amp/m. The electrical properties of $\sigma = 0.889$ S/m and $\epsilon_r = 71.7$ are assumed for the body. The impressed electric field for this case is calculated from eq. (2.3). The amplitudes and phase angles of the induced electric fields in five different sections of the cylinder are plotted in Figure 2.10. It is observed that the induced electric field or the induced eddy current is zero at the center of the magnetic beam and it increases linearly to a maximum value at the edge of the beam, and then decays down toward the edge of the body. The deviations between the present results and the quasi-static solutions are indicated in Figure 2.10. The application of a beam of RF magnetic field for the purpose of locally heating a biological body in a hyperthermia cancer therapy is feasible if the low heating at the beam center can be compensated by other means of EM heating.

The last numerical example is the quantification of the electric fields induced by a 100 MHz magnetic field with an intensity of 1 Amp/m in a cylindrical model of man with a height of 168 cm and a maximum diameter of



Phase
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$$\begin{aligned} H_z^i &= 1 \text{ A/m} \\ f &= 100 \text{ MHz} \\ \sigma &= 0.889 \text{ S/m} \\ \epsilon_r &= 71.7 \end{aligned}$$

Figure 2.10.

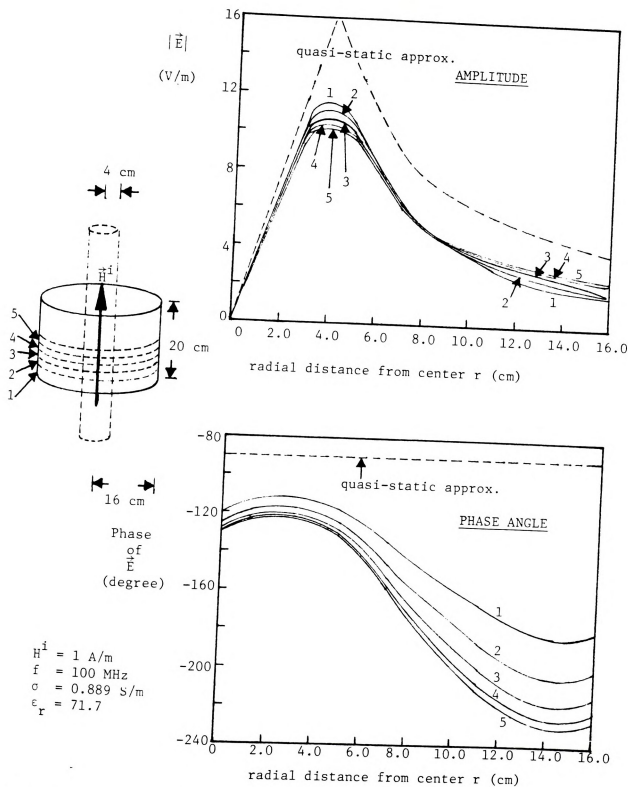


Figure 2.10. Distributions of amplitudes and phase angles of electric fields induced in a cylindrical biological body with a diameter of 32 cm and a height of 20 cm by a beam of 100 MHz magnetic field with a diameter of 8 cm and an intensity of 1 A/m.

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36 cm as shown in Figure 2.11. Again, $\sigma = 0.889 \text{ S/m}$ and $\epsilon_r = 71.7$ are assumed for the body. In the numerical calculation, the body was subdivided into 66 rings of various radii and with a square cross-sectional area of 6 cm x 6 cm. The amplitudes and phase angles of the induced electric fields inside the body are indicated in Figure 2.11. It is noted that these results deviate greatly from the quasi-static solutions. For example, the induced electric field in the head can be greater than that in the middle section of the body even though the radius of the latter is greater than that of the former. The phase angles of the induced electric fields are very different from -90° and they vary widely throughout the body.

2.4. Convergence of Numerical Results

The accuracy and validity of numerical results presented in this chapter depend greatly on the convergence of numerical results as the number of subdivided rings is increased. Fortunately, the present theory gives an excellent convergence of numerical results. This fact is demonstrated in an example shown in Figure 2.12, in which a cylindrical body with a diameter of 36 cm and a height of 18 cm is exposed to a uniform magnetic field of 100 MHz with an intensity of 1 Amp/m. The electrical properties of $\sigma = 0.889 \text{ S/m}$ and $\epsilon_r = 71.7$ are assumed for the body.



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Figure 2.11

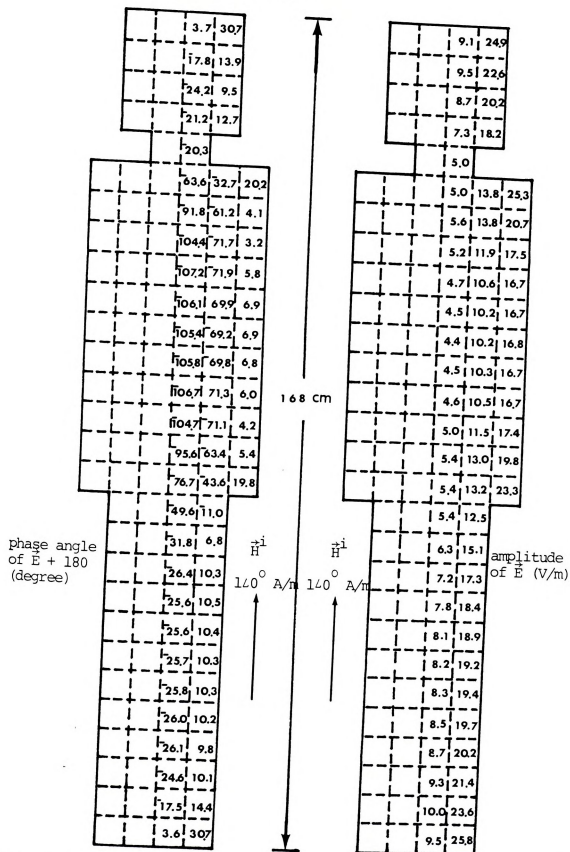


Figure 2.11. Distributions of amplitudes and phase angles of electric fields induced by a uniform magnetic field of 100 MHz in a model of man ($\sigma = 0.889 \text{ S/m}$, $\epsilon_r = 71.7$).

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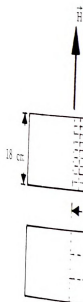
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In the first numerical calculation, the body was subdivided into 9 rings of various radii and with a square cross-sectional area of 6 cm x 6 cm. In the second numerical calculation, the same body was subdivided into 81 rings of various radii and with a square cross-sectional area of 2 cm x 2 cm. The numerical results on the induced electric fields based on these two ring subdivisions are compared in Figure 2.12. It is observed that the agreement between these two sets of numerical results is excellent even though the number of rings in the second subdivision increases by a factor of 9 from that of the first subdivision. This excellent convergence of numerical results implies that it is possible to obtain accurate results on the electric fields or the eddy currents induced by a RF magnetic field in most biological bodies without subdividing the body into too many rings. This advantage provided by the present method can lead to a great saving in computer time and cost.

2.5. Closed Form Solution for Sphere and Comparison With Present Numerical Results

To compare our numerical results with the existing analytical solution [3], we consider the case of a conducting sphere with radius R which is immersed in a uniform time-harmonic magnetic field,

$$\vec{H}^i = \hat{z} H^i = \hat{r} H^i \cos\theta - \hat{\theta} H^i \sin\theta$$



$$\begin{aligned}
 H^i &= 1 \text{ A/m} \\
 f &= 100 \text{ MHz} \\
 \sigma &= 0.889 \text{ S/m} \\
 \epsilon_r &= 71.7
 \end{aligned}$$

Figure 2.12.

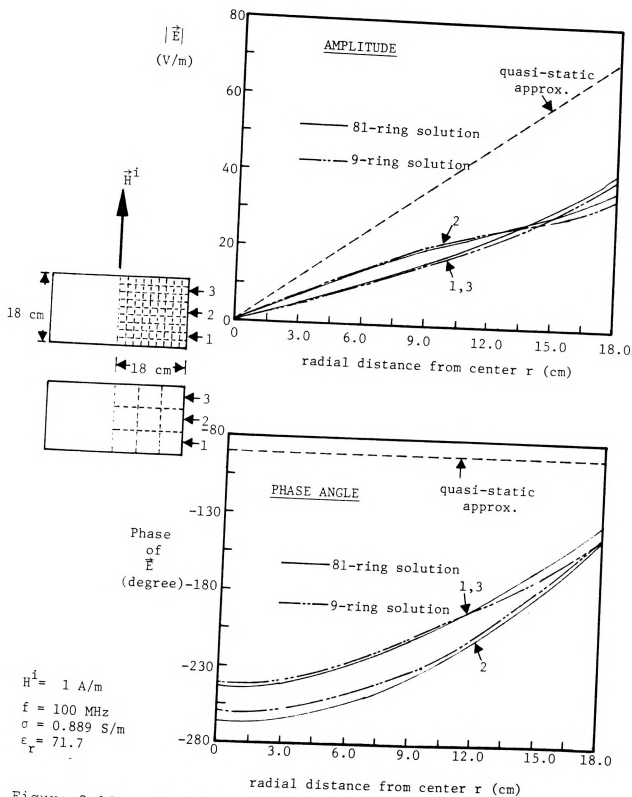


Figure 2.12. Comparison of 81-ring solutions and 9-ring solutions for the electric fields induced in a cylindrical biological body with a diameter of 36 cm and a height of 18 cm by a 100 MHz magnetic field with an intensity of 1 A/m.

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where $\hat{r}$ and $\hat{\theta}$ are unit vectors in spherical coordinates, as shown in Figure 2.13. The vector potentials maintained by the induced currents and the sources for the regions inside and outside the sphere can be written as

$$\vec{A}_{in} = (1/2) \mu_0 H^i \sin\theta \frac{C_1}{\sqrt{r}} I_{3/2} [(j\omega \mu_0 \sigma)^{1/2} r] \hat{\phi} \quad \text{when } r < R \quad (2.18)$$

$$\vec{A}_{out} = (1/2) \mu_0 H^i \sin\theta \left(r + \frac{C_2}{r^2}\right) \hat{\phi} \quad \text{when } r > R \quad (2.19)$$

where $I_{3/2}$ is the spherical Bessel function, and σ is the conductivity of the sphere. The constants C_1 and C_2 can be determined by use of the following boundary conditions at $r = R$,

$$1) \quad A_{in} = A_{out}$$

$$2) \quad \frac{\partial}{\partial r} (r A_{out}) = \frac{\partial}{\partial r} (r A_{in}).$$

After some manipulations and utilizing some identities of spherical Bessel functions we have

$$C_1 = \frac{3 R^{3/2}}{w I_{1/2}(w)} \quad (2.20)$$

$$C_2 = \frac{3 w I_{-1/2}(w) - (3+w^2) I_{1/2}(w)}{w^2 I_{1/2}(w)} R^3 \quad (2.21)$$

where $w = \sqrt{2j} \frac{R}{\delta}$, and $\delta = \left(\frac{2}{\omega \mu_0 \sigma}\right)^{1/2}$ is the skin depth of the sphere. Combining eqs. (2.18) and (2.20), we obtain



Figure

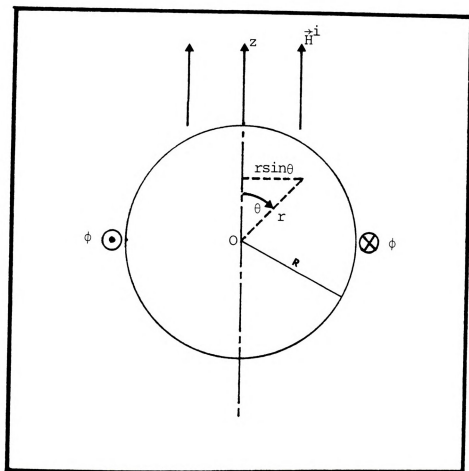


Figure 2.13. A finite conducting sphere with radius R is immersed in a uniform magnetic field polarized in the $+z$ direction.

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the complete expression for $\vec{A}_{in}$,

$$\vec{A}_{in} = (3/2) \mu_0 H^i \sin\theta \frac{\sqrt{Rr}}{\frac{wr}{R}} \frac{I_{3/2}(wr/R)}{I_{1/2}(w)} \hat{\phi}.$$

Then the electric field induced inside the sphere is given by

$$\vec{E} = -j\omega\vec{A}_{in} = - (3/2) j\omega\mu_0 H^i \sin\theta \sqrt{Rr} \frac{R}{wr} \frac{I_{3/2}(wr/R)}{I_{1/2}(w)} \hat{\phi}. \quad (2.22)$$

To compare eq. (2.22) with our numerical results, we assume that the sphere has a radius of 2 cm, a conductivity of 8.0 S/m, and a dielectric constant of 50, which is immersed in a uniform oscillating magnetic field with a frequency of 300 MHz. Under these assumptions the skin depth and the radius of the sphere are in the same order of magnitude. Now for the case of numerical evaluation, the sphere is simulated by a "ring sphere" which is constructed with 30 different rings, as shown in Figure 2.14. The numerically evaluated results of the induced electric fields in these circular rings are then compared with those values obtained from eq. (2.22) at corresponding points inside the sphere. The comparison is made in Table 2.1. Also listed in Table 2.1 are the results obtained by the quasi-static approximation.

It is apparent that the agreement between our numerical solution and the existing closed form solution is



Figure

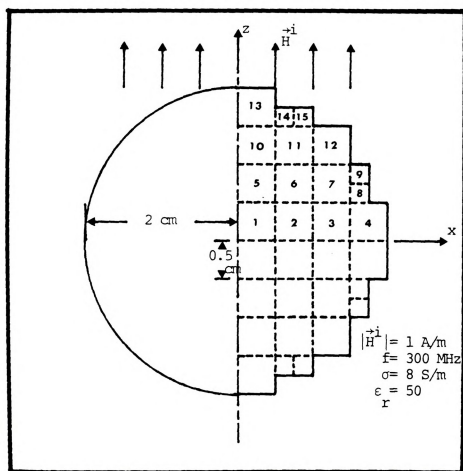


Figure 2.14. A finite conducting sphere with a radius of 2 cm is simulated by a "ring sphere" constructed with 30 circular rings of two different cross-sectional areas. The sphere is assumed to have a conductivity of 8.0 S/m and a dielectric constant of 50, and is immersed in a 300 MHz oscillating magnetic field with an intensity of 1 A/m.

Table 2.1.

Comparison of the present solutions and a closed form solution for the conductivity of a sphere in a conducting sphere with a conductivity of 8.0 S/m and a dielectric permittivity of 1.0 by a 300 MHz magnetic field with an intensity of 1 A/m . Also listed are the solutions obtained by the quasi-static approximation.

Quasi-static

Table 2.1. Comparison of the present solutions and a closed form solution for the electric fields induced in a conducting sphere with a conductivity of 8.0 S/m and a dielectric constant of 50 by a 300 MHz magnetic field with an intensity of 1 A/m . Also listed are the solutions obtained by the quasi-static approximation.

Ring. No.	Quasi-Static Approximation		Closed Form Solution		Present Solution	
	Magnitude (10 V/m)	Phase (Degree)	Magnitude (10 V/m)	Phase (Degree)	Magnitude (10 V/m)	Phase (Degree)
1	0.296	-90	0.230	-25.6	0.259	-26.4
2	0.889	-90	0.689	-31.0	0.770	-31.7
3	1.480	-90	1.162	-41.8	1.269	-42.4
4	2.070	-90	1.681	-57.6	1.787	-58.3
5	0.296	-90	0.230	-30.9	0.256	-32.1
6	0.889	-90	0.692	-36.4	0.763	-37.3
7	1.480	-90	1.173	-47.2	1.264	-48.4
8	1.930	-90	1.557	-56.6	1.647	-58.2
9	1.930	-90	1.582	-60.7	1.658	-62.5
10	0.296	-90	0.232	-41.9	0.253	-43.4
11	0.889	-90	0.704	-47.1	0.759	-48.5
12	1.480	-90	1.206	-57.8	1.271	-59.0
13	0.296	-90	0.241	-57.6	0.255	-60.1
14	0.741	-90	0.600	-56.8	0.635	-59.0
15	1.040	-90	0.849	-60.6	0.894	-62.7

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excellent for both the magnitude and the phase distributions throughout the interior of the sphere. This excellent agreement confirms the accuracy of our numerical technique. Furthermore, the deviation of the quasi-static solution from the other two sets of solutions indicates that the conventional quasi-static solution is usually inadequate.

2.6. Experimental Setup

A series of experiments has been conducted to measure the electric fields induced by UHF magnetic fields inside some experimental models constructed with plexiglass and filled with phantom biological materials for the purpose of substantiating the accuracy of theoretical results presented in this chapter.

The experimental setup is depicted schematically in Figure 2.15. The experiment was conducted inside a large microwave anechoic chamber in which a standing EM wave was created by radiating an EM wave upon a metallic reflector. The electric field of the wave was polarized horizontally and the magnetic field was vertically polarized. The experiment was conducted in the frequency range of 500 to 750 MHz. The phantom body was placed at the location of a maximum magnetic field or a minimum electric field. If the body dimension in the direction of wave propagation is small compared with the wavelength, the



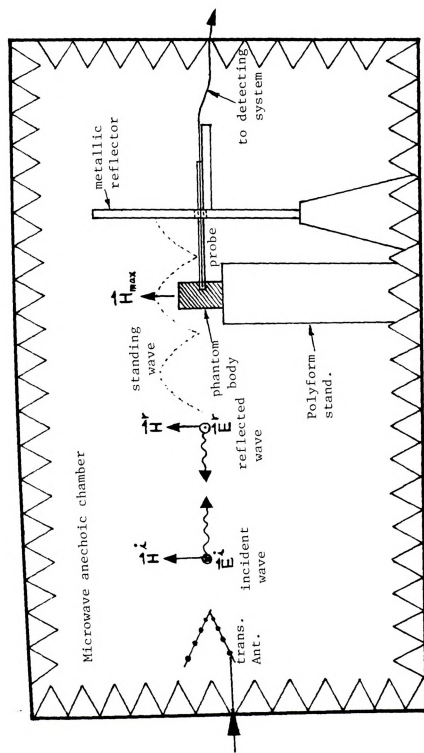


Figure 2.15. Experimental setup for measuring the electric fields inside the phantom biological bodies induced by RF magnetic fields.

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impressed magnetic field at any point inside the phantom body can be assumed to be uniform. The induced electric field inside the phantom body was measured by an implantable electric field probe with an interference-free lead-wire system. This probe was described elsewhere [5]. The output of the probe was connected to a detecting system outside the anechoic chamber.

2.7. Experimental Results and Comparison With Theoretical Results

The first experiment was conducted on a cylindrical phantom body with a diameter of 4 cm and a variable height. The electric properties of the phantom body were estimated to be approximately $\sigma = 5.0$ S/m and $\epsilon_r = 50$. The experiment was conducted at 750 MHz and the cylinder was placed at the location of a maximum magnetic field, one wavelength (40 cm) in front of the metallic reflector. The direction of the magnetic field was in parallel with the cylindrical axis. The induced electric field in the central section of the cylinder was probed for three cases of cylindrical heights; $H = 2, 4$ and 8 cm. The measured amplitudes of the induced electric fields are plotted in Figure 2.16 with dashed lines, in comparison with corresponding theoretical results which are indicated with solid lines. It is noted that with an implantable electric field probe loaded with a microwave diode, only the amplitude of the induced electric field could be measured. In



Figure 2.16.

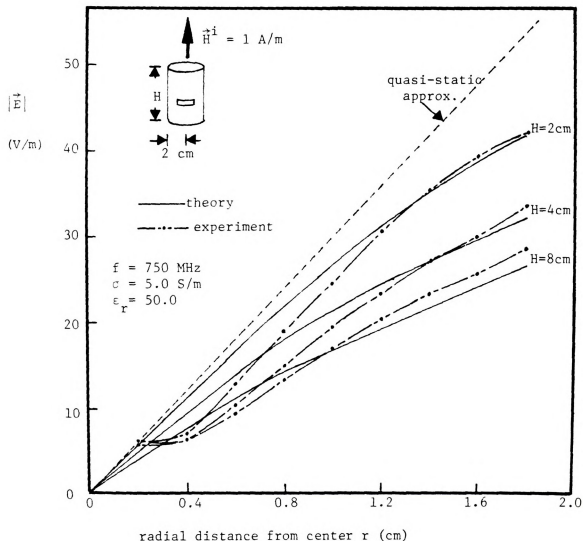


Figure 2.16. Comparison of theoretical and experimental results for the amplitudes of electric fields induced in a phantom biological cylinder of 4 cm diameter with various heights by a 750 MHz magnetic field.

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Figure 2.16, the agreement between experimental and theoretical results is very good except in the region near the cylindrical axis where the perturbation of the probe (about 1 cm in size) on the induced current becomes significant. Another cause of discrepancy between the theory and the experiment can be attributed to the image effect of the metallic reflector; this effect is not severe if the distance between the body and the reflector is about one wavelength or greater. The important finding in Figure 2.16 is that it was observed experimentally and predicted theoretically that as the cylindrical height was increased, the induced electric field in the central section of the cylinder decreased. These results were also found to be significantly different from the quasi-static solutions.

The second experiment was conducted on a phantom cylinder with a diameter of 3.8 cm and a height of 3.8 cm to study the effect of body conductivity on the induced electric field. Three kinds of phantom materials with conductivities of $\sigma = 2.2, 4.5$ and 6.0 S/m were used. The experiment was conducted at 750 MHz, and the cylinder was placed at the location of a maximum magnetic field as in the previous experiment. The induced electric fields at the central section of the cylinder were probed, and the distributions of their amplitudes are plotted in Figure 2.17 in comparison with corresponding theoretical

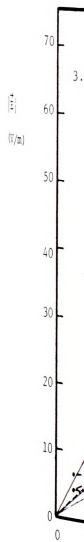


Figure 2.17.

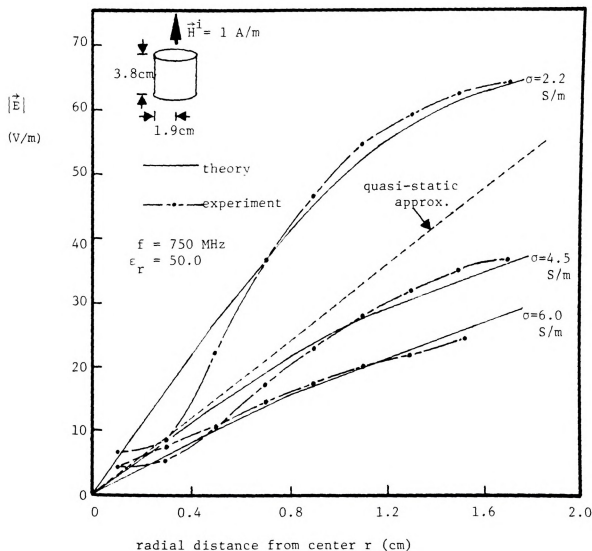


Figure 2.17. Comparison of theoretical and experimental results for the amplitudes of electric fields induced in a phantom biological cylinder with a diameter of 3.8 cm, a height of 3.8 cm and various conductivities by a 750 MHz magnetic field.

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results. It was observed that as the body conductivity is increased, the induced electric field decreases. This effect was accurately predicted by theory and confirmed by experiment. It is also noted that both theoretical and experimental results deviate significantly from the quasi-static solutions. The agreement between theory and experiment as shown in Figure 2.17 is very good.

The third experiment was conducted to measure the induced electric fields inside a phantom model of man. The phantom model of man was constructed with plexiglass and filled with phantom material, and has a height of 22.2 cm and a maximum diameter of 3.8 cm as shown in Figure 2.18. The experiment was conducted at 750 MHz, and electrical properties of the model at this frequency are approximately $\sigma = 5.0$ S/m and $\epsilon_r = 50$. The model was placed upright at the location of a maximum magnetic field with the impressed magnetic field in parallel with the long dimension of the body. The induced electric fields at 32 locations of the body were probed. The measured amplitudes of the induced electric fields are indicated in Figure 2.18 in comparison with the corresponding theoretical results. The agreement between theory and experiment is considered to be very good in this rather complicated body. It is to be noted that experimental and theoretical results shown in Figure 2.18 are very different from the quasi-static solutions.

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Figure 2.18.

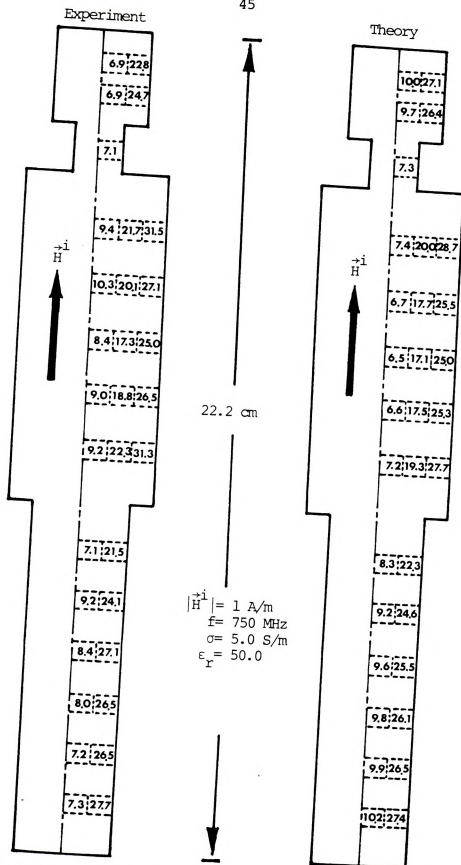


Figure 2.18. Comparison of theoretical and experimental results for the amplitudes of electric fields induced in a phantom model of man (height = 22.2 cm, max. diameter = 3.8 cm) by a 750 MHz magnetic field.

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A very good agreement between theory and experiment as demonstrated in Figures 2.16, 2.17 and 2.18 confirms the validity and accuracy of the theoretical method presented in this chapter.

2.8. Discussion

In the experiment, the body was placed at the location of a maximum magnetic field in a standing EM wave. Since the body dimension in the direction of wave propagation was small compared with the wavelength, the impressed magnetic field inside the body was assumed to be uniform. Under this approximation, the impressed electric field associated with the impressed magnetic field was

$$\vec{E}^i = - \frac{j}{2} \omega \mu_0 H^i r \hat{\phi} \quad (2.2)$$

as has been given before. This $\vec{E}^i$ was used in the numerical calculation of the induced electric field inside the body.

From a different point of view, the body in the experiment was located at a minimum electric field of a standing EM wave, and the actual impressed electric field was

$$\vec{E}^i = \hat{x} j \sqrt{\frac{\mu_0}{\epsilon_0}} H^i \sin ky = \hat{x} j \omega \mu_0 H^i y \quad (2.23)$$

where $y = 0$ corresponds to the center of the body.

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$\vec{E}^i$ given in eq. (2.2) is circulatory and that given in eq. (2.23) is linear and antisymmetrical. Also the amplitude of the latter is twice that of the former along the y axis.

It is important to ask whether these two different $\vec{E}^i$'s can yield the same theoretical values for the induced electric fields in the body? We have numerically proved that, indeed, these two $\vec{E}^i$'s give very similar numerical results on the induced electric fields in the same body. This phenomenon resulted in a very good agreement between experimental and theoretical results reported in this chapter.

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CHAPTER III

SURFACE INTEGRAL EQUATION METHOD FOR
INTERACTION OF MICROWAVE WITH
BIOLOGICAL BODY

Only in recent years has beamed microwave power been proposed as a new technique in transporting high-level energy between two remote transmitting-receiving stations, and has microwave heating been considered as an effective method in hyperthermia cancer therapy. Although EM transmission and scattering phenomena have been major and classical topics in electrodynamics since the end of the last century, uncertainties about the potential hazards of these high-power microwave radiations are still major concerns relating to the acceptability of these newly developed technologies. In order to fully understand the biological effects of microwave radiation, it is necessary to quantify the electric fields induced in the biological bodies when they are irradiated by the incident EM waves.

Existing methods of quantifying the electric fields induced by EM waves are mostly based on the solutions of volume integral equations. Unfortunately, when the body is large compared with the wavelength of the incident

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wave, the volume integral equation method becomes inefficient. Since for this case, the induced electric field is mainly concentrated in the region near the body surface, it is believed that the surface integral equation technique is potentially more efficient.

In this chapter, we develop a new technique based on two coupled surface integral equations. A good agreement between our solutions and the existing results obtained by volume integral equation method has been verified.

3.1. Introduction

High-level, beamed-power microwave transmission between satellites and earth has been proposed as an alternate scheme for providing energy in the future. However, the potential hazards which exist in the interface between the beamed-power microwave and personnel should be carefully considered by the microwave systems design engineers. Also, the thermal-heating effects of intense microwaves are becoming well known, and have been demonstrated effective in hyperthermia cancer therapy. But the allowable field intensity and the permissible level of absorbed power for the human body still remain unknown. It is desirable to obtain some quantitative information about microwave interactions with human or other biological bodies.

Johnson and Durney [6] have investigated the plane wave irradiation of a prolate-spheroid model of man based

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on the perturbation theory described by Van Bladel [7]. Their theory is valid only when the wavelength is long compared to the dimensions of the spheroid. The same approach has been adopted by Massoudi [8] to solve for the internal electric field induced by a plane wave inside an ellipsoidal model of man. Livesay and Chen [9] have developed a theoretical method based on a volume integral equation, the so called Tensor Integral Equation, which can be used to quantify the electric field induced inside an arbitrarily shaped biological body by an incident EM wave. This volume integral equation method removes the restrictions that the irradiated bodies have simple shapes and small electrical dimensions. However, when the body is electrically large, the EM waves cannot penetrate deeply inside the body, and will concentrate near the body surface. Consequently, the volume integral equation method becomes inefficient. For this case, the surface integral equation method seems to be a reasonable alternative approach. A surface integral equation method has been employed by Wu et al. [10] to solve for the EM fields induced inside arbitrary cylinders of biological tissue; only infinitely long cylinders were considered. Another surface integral formulation has been developed by Poggio and Miller [11]; however, in their formulation the derivatives of the surface unknowns are involved in the integrands of the surface integrals. This makes the numerical

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solutions difficult, if not impossible, to implement. Wang [12] has determined the scattering characteristics of conducting bodies by use of the moment-method, which divides the surface of the body into trilateral surface cells, but only results for sphere have been reported. It is desirable to obtain more complete information about the fields induced by EM waves, based on the surface integral equations, for finite conducting bodies with geometries more complicated than that of the sphere.

In this chapter, we will develop a new surface integral equation technique to quantify the induced EM field on the surface of an arbitrarily shaped, finite conducting body. It should be emphasized that the internal field can be readily determined in terms of the surface field, once it is found.

Before the surface integral equations are derived in Section 3.4, two preliminary theorems and a general solution of Maxwell's equations are described in Sections 3.2 and 3.3, respectively. The application of the coupled, surface integral equations to the special case of an infinite interface is considered in Section 3.5. The moment-method of solution of these integral equations is briefly reviewed in Section 3.6. The development of the numerical technique for the solutions of the surface integral equations is presented in Section 3.7. Finally, some numerical examples are given in Section 3.8.

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3.2. The Preliminary Theorems

Before the coupled, surface integral equations are derived, we must introduce two preliminary theorems which consist of several vector integral identities. These theorems are necessary for the derivation, and will be described without proof. The related proof can be found elsewhere [13].

Theorem 1.

Let $\vec{v}$, $\nabla_{\mathbf{O}} \cdot \vec{v}$ be continuous on the closed surface F , and U be a differentiable function. Then

$$\int_F U \nabla_{\mathbf{O}} \cdot \vec{v} dF' = - \int_F \vec{v} \cdot \nabla_{\mathbf{O}} U dF'$$

with the surface divergence $\nabla_{\mathbf{O}} \cdot \vec{v}$ defined as

$$\nabla_{\mathbf{O}} \cdot \vec{v} \equiv \lim_{|S| \rightarrow 0} \frac{1}{|S|} \oint_C (\hat{n}_{\mathbf{O}} \cdot \vec{v}) dl$$

where $|S|$ is the area of a small surface element S enclosed by contour C , and $\hat{n}_{\mathbf{O}}$ is a unit normal vector on C directed exterior to S . It can be shown that $\nabla_{\mathbf{O}} = \nabla - \hat{n} \frac{\partial}{\partial n}$, where ∇ represents the conventional three-dimensional differential del-operator.

Theorem 2.

If the surface field $\vec{j}(\vec{r}')$ is continuous on F , then

$$\hat{n}(\vec{\xi}) \times \int_{F_1} [\vec{j}(\vec{r}') \times \nabla' \Phi(\vec{r}, \vec{r}')] dF' = -2\pi \vec{j}(\vec{\xi}) +$$

$$\hat{n}(\vec{\xi}) \times \int_F [\vec{j}(\vec{r}') \times \nabla' \Phi(\vec{\xi}, \vec{r}')] dF'.$$

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$$\hat{A}(\vec{\xi}) \times \int_{F_e} [\vec{j}(\vec{r}') \times \nabla' \phi(\vec{r}, \vec{r}')] dF' = 2\pi \vec{j}(\vec{\xi}) +$$

$$\hat{A}(\vec{\xi}) \times \int_F [\vec{j}(\vec{r}') \times \nabla' \phi(\vec{\xi}, \vec{r}')] dF'$$

The sign $\int_{F_e}$ implies that $\vec{r}$ is the field point in the integration, and lying in the exterior side of the surface F .

The eventual field point $\vec{\xi}$ and the source point $\vec{r}'$ are both on the surface F . We denote with $\int_{F_i}$ and $\int_{F_e}$ the limits which we obtain if $\vec{r}$ approaches the point $\vec{\xi}$ of the surface F from the interior or exterior, respectively.

3.3. The General Solution of Maxwell's Equations in an Unbounded Homogeneous Space

Let us consider the problem where a finite source region G_s , consisting of electric and magnetic sources with volume current densities of $\vec{j}_e$ and $\vec{j}_m$, respectively, radiates into a lossless, unbounded, homogeneous, linear, and isotropic medium with permeability μ and permittivity ϵ , as shown in Figure 3.1 where we assume that the source region G_s is enclosed by surface F_s .

Based on the continuity equation, the volume charge densities are given as

$$Q_e = \frac{j}{\omega} \nabla \cdot \vec{j}_e \quad Q_m = \frac{j}{\omega} \nabla \cdot \vec{j}_m$$

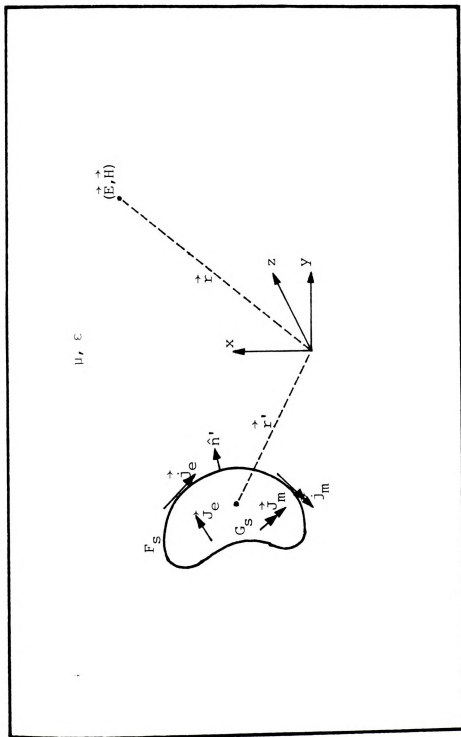


Figure 3.1. A source region G_S enclosed by surface F_S is characterized by volume current densities ($\vec{J}_e, \vec{J}_m$) and equivalent surface current densities ($\vec{j}_e, \vec{j}_m$). These current sources radiate into a homogeneous space with permeability μ and permittivity ϵ , and maintain an EM field (E, H) at point I .

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$$\eta_0 = \sqrt{\mu_0}$$

where a time variation of $e^{j\omega t}$ has been assumed, and ω is the angular frequency of the sources.

From the general solution of Maxwell's equations [14], the electromagnetic field at any point $\vec{r}$ maintained by the above sources can be expressed in terms of $\vec{J}_e$, $\vec{J}_m$, Q_e , and Q_m as

$$\vec{E}(\vec{r}) = \frac{1}{4\pi} \int_{G_S} [-\vec{J}_m \times \nabla' \phi + \frac{Q_e}{\epsilon} \nabla' \phi - j\omega\mu \vec{J}_e \phi] dG' \quad (3.1)$$

$$\vec{H}(\vec{r}) = \frac{1}{4\pi} \int_{G_S} [\vec{J}_e \times \nabla' \phi + \frac{Q_m}{\mu} \nabla' \phi - j\omega\epsilon \vec{J}_m \phi] dG' \quad (3.2)$$

with $\phi = \frac{e^{-jk|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|}$, where $\vec{r}'$ is the position vector locating the source point inside the source region G_S , and $k = \omega\sqrt{\mu\epsilon}$ is the wavenumber in unbounded space.

Now let's define the following source densities of the equivalent surface currents and charges on the closed surface F_S

$$\vec{J}_e \equiv \hat{n}' \times \vec{H}$$

$$\vec{J}_m \equiv -\hat{n}' \times \vec{E}/\eta_0$$

$$q_e = \frac{1}{\omega} \nabla_O \cdot \vec{J}_e$$

$$q_m = \frac{1}{\omega} \nabla_O \cdot \vec{J}_m$$

where $\hat{n}'$ is the outward unit normal vector on F_S , and

$\eta_0 = \sqrt{\mu_0/\epsilon_0}$ is the intrinsic wave impedance in free space.

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The above electromagnetic field ($\vec{E}$, $\vec{H}$) at position $\vec{r}$ can also be expressed in terms of $\vec{j}_e$, $\vec{j}_m$, q_e , and q_m as [14]

$$\vec{E}(\vec{r}) = \frac{1}{4\pi} \int_{F_S} [-\eta_0 \vec{j}_m \times \nabla' \phi + \frac{q_e}{\epsilon} \nabla' \phi - j\omega\mu \vec{j}_e \phi] dF' \quad (3.3)$$

$$\vec{H}(\vec{r}) = \frac{1}{4\pi} \int_{F_S} [\vec{j}_e \times \nabla' \phi + \frac{q_m}{\mu} \nabla' \phi - j\omega\epsilon \eta_0 \vec{j}_m \phi] dF' \quad (3.4)$$

where ϕ has the same form as that of eqs. (3.1) and (3.2). The last two equations are very important, and will serve as the basic tool in the following derivation.

3.4. Derivation of the Coupled Surface Integral Equations

To derive the coupled, surface integral equations, we first consider the geometry shown in Figure 3.2. We assume that a finite conducting body G enclosed by surface F is irradiated by the electromagnetic wave from the source region G_s . The medium in the region external to F is assumed to be free space with permeability μ_0 and permittivity ϵ_0 . The finite conducting body is constructed with a homogeneous material characterized by permeability μ_0 and complex permittivity ϵ_d . The complex permittivity ϵ_d is defined by

$$\epsilon_d = \epsilon_0 (\epsilon_r - j \frac{\sigma}{\omega \epsilon_0})$$

where ϵ_r is the dielectric constant, and σ is the conductivity of the body. The source region G_s is

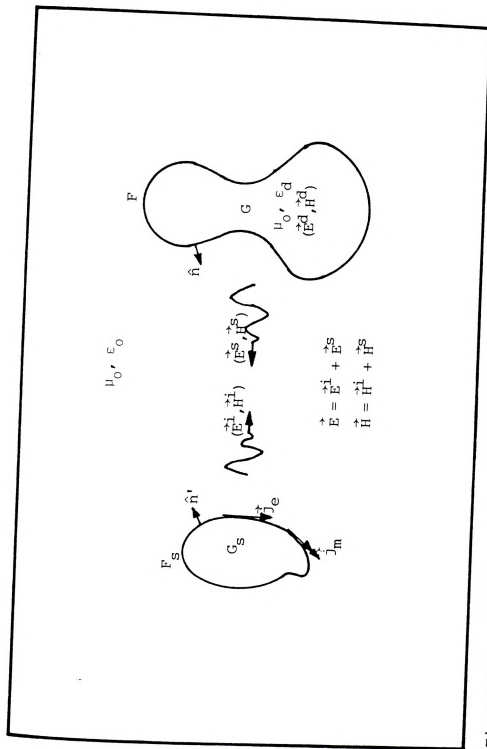


Figure 3.2. A finite conducting body G , immersed in free space and enclosed by surface F , with a permeability μ_0 and a complex permittivity ϵ_d is irradiated by the incident field $(\vec{E}^i, \vec{H}^i)$ radiated from the source region G_s where the equivalent surface current densities $(\vec{j}_e, \vec{j}_m)$ are defined on the equivalent surface F_s . The EM fields scattered from and induced in the body are denoted as $(\vec{E}^s, \vec{H}^s)$ and $(\vec{E}_s, \vec{H}_s)$, respectively.

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Let us denote the incident field which is radiated by the sources in the absence of the finite conducting body by $(\vec{E}^i, \vec{H}^i)$. With the body present, we denote the fields scattered from and induced in it as $(\vec{E}^s, \vec{H}^s)$ and $(\vec{E}^d, \vec{H}^d)$, respectively. Now we are in a position to consider these three sets of EM fields separately.

I) The Scattered Field $(\vec{E}^s, \vec{H}^s)$:

To begin with, let's consider Maxwell's equations in the region internal to F . They are

$$\nabla \times \vec{H}^d = j\omega\epsilon_0 \vec{E}^d + j\omega(\epsilon_d - \epsilon_0) \vec{E}^d \quad (3.5)$$

$$\nabla \times \vec{E}^d = -j\omega\mu_0 \vec{H}^d. \quad (3.6)$$

Defining the last term of eq. (3.5) as an equivalent volume current density, then eq. (3.5) becomes

$$\nabla \times \vec{H}^d = j\omega\epsilon_0 \vec{E}^d + \vec{j}_e^{eq} \quad (3.5')$$

where

$$\vec{j}_e^{eq} \equiv j\omega(\epsilon_d - \epsilon_0) \vec{E}^d. \quad (3.7)$$

After substituting the expression for ϵ_d into eq. (3.7), we may see that $\vec{j}_e^{eq}$ consists of two components: the conduction current and the polarization current. We may think of $\vec{j}_e^{eq}$ as a current existing in free space as shown in

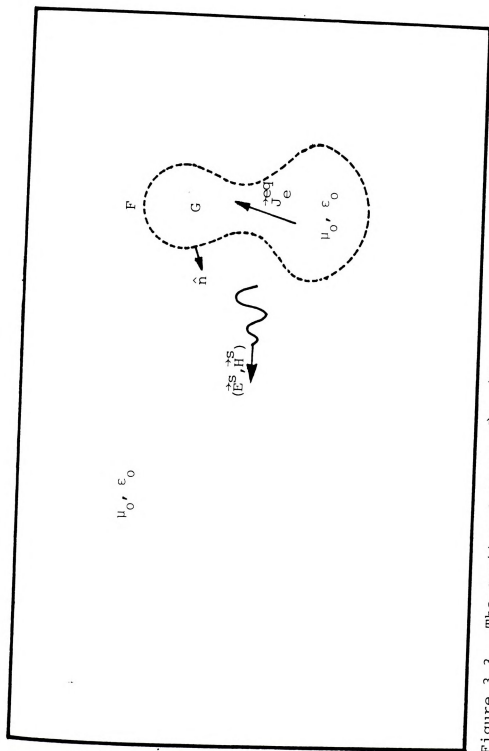


Figure 3.3. The scattered field $(\vec{E}^s, \vec{H}^s)$ can be considered as originated from an equivalent current source $\vec{j}_{eq}$ existing inside the body region which has been replaced by free space. This equivalent current source consists of two components, the conduction current and the polarization current.

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Figure 3.3, since only μ_0 and ϵ_0 appear in eqs. (3.5') and (3.6).

It is to be noted that the source region G_s is not included in Figure 3.3, because $\vec{J}_e^{\text{eq}}$ is the only source of the scattered field ($\vec{E}^s$, $\vec{H}^s$). Furthermore, the equivalent current $\vec{J}_e^{\text{eq}}$ radiates into an unbounded, homogeneous space. Hence, based on eqs. (3.1) and (3.2), $\vec{E}^s$ and $\vec{H}^s$ can be written as

$$\vec{E}^s = \frac{1}{4\pi} \int_G \left[\frac{\vec{Q}_e^{\text{eq}}}{\epsilon_0} \nabla' \phi_f - j\omega\mu_0 \vec{J}_e^{\text{eq}} \phi_f \right] dG'$$

$$\vec{H}^s = \frac{1}{4\pi} \int_G [\vec{J}_e^{\text{eq}} \times \nabla' \phi_f] dG'$$

where $\vec{Q}_e^{\text{eq}} = \frac{1}{\omega} \nabla \cdot \vec{J}_e^{\text{eq}}$ is the equivalent volume charge density,

$\phi_f = \frac{e^{-jk_0 |\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|}$ is the Green's function for free space,

and $k_0 = \omega \sqrt{\mu_0 \epsilon_0}$ is the wavenumber in free space.

Now let's define the equivalent surface current densities in terms of the scattered field ($\vec{E}^s$, $\vec{H}^s$) on the surface F as

$$\vec{J}_e^s \equiv \hat{n} \times \vec{H}^s$$

$$\vec{J}_m^s \equiv -\hat{n} \times \vec{E}^s / \eta_0$$

$$\vec{Q}_e^s = \frac{1}{\omega} \nabla_0 \cdot \vec{J}_e^s$$

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$$\mathbf{q}_m^s = \frac{1}{\omega} \nabla_0 \cdot \mathbf{j}_m^s$$

where $\hat{n}$ is the outward unit normal vector on F .

From the equivalence theorem [15], [16], we know that the equivalent surface currents $(\mathbf{j}_e^s, \mathbf{j}_m^s)$ on F can support $(\mathbf{E}^s, \mathbf{H}^s)$ external to F and zero field internal to F as shown in Figure 3.4. Since the surface currents $(\mathbf{j}_e^s, \mathbf{j}_m^s)$ radiate in an unbounded homogeneous space (free space), we can evaluate the EM fields supported by them with eqs. (3.3) and (3.4). From the uniqueness theorem, we know that the EM fields so calculated will be those postulated by the equivalence theorem. Hence we have

$$\mathbf{E}^s = \frac{1}{4\pi} \int_F [-\eta_0 \mathbf{j}_m^s \times \nabla' \phi_f + \frac{\mathbf{q}_e^s}{\epsilon_0} \nabla' \phi_f - j\omega\mu_0 \mathbf{j}_e^s \phi_f] dF' \quad (3.8)$$

$$\mathbf{H}^s = \frac{1}{4\pi} \int_F [\mathbf{j}_e^s \times \nabla' \phi_f + \frac{\mathbf{q}_m^s}{\mu_0} \nabla' \phi_f - j\omega\epsilon_0 \eta_0 \mathbf{j}_m^s \phi_f] dF' \quad (3.9)$$

external
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as the scattered field at any point in the region external to F . In the region internal to F , the above two integrals vanish, i.e.,

$$0 = \frac{1}{4\pi} \int_F [-\eta_0 \mathbf{j}_m^s \times \nabla' \phi_f + \frac{\mathbf{q}_e^s}{\epsilon_0} \nabla' \phi_f - j\omega\mu_0 \mathbf{j}_e^s \phi_f] dF' \quad (3.10)$$

$$0 = \frac{1}{4\pi} \int_F [\mathbf{j}_e^s \times \nabla' \phi_f + \frac{\mathbf{q}_m^s}{\mu_0} \nabla' \phi_f - j\omega\epsilon_0 \eta_0 \mathbf{j}_m^s \phi_f] dF' \quad (3.11)$$

internal
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because there is a null field internal to F .



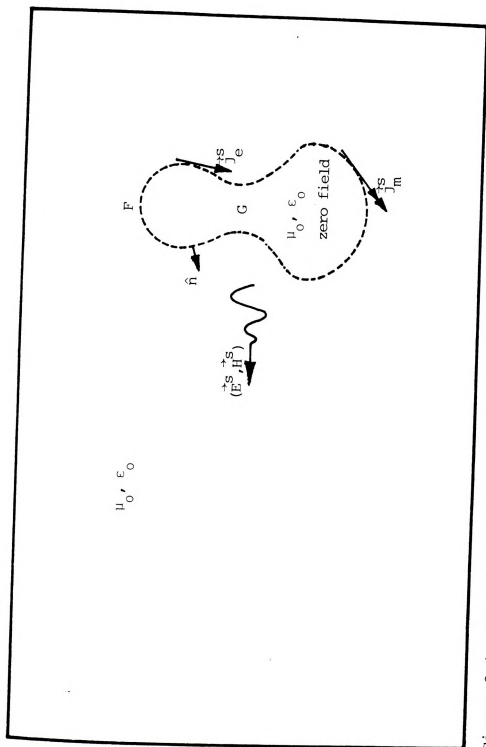


Figure 3.4. The scattered field ($\vec{E}^s, \vec{H}^s$) in the exterior of the body and a zero field in the interior of the body are maintained by the equivalent surface currents ($\vec{j}_e, \vec{j}_m$), which are defined in terms of the scattered field on surface F. The body region G has been replaced by free space.

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II) The Induced Field ($\vec{E}^d, \vec{H}^d$):

To investigate the induced field ($\vec{E}^d, \vec{H}^d$), let us consider Figure 3.2 again. This time we treat the interior region of G as source free with the induced field ($\vec{E}^d, \vec{H}^d$) maintained by the sources existing in the exterior region of G. As for the case of scattered field, we may define a set of equivalent surface current densities ($\vec{j}_e^d, \vec{j}_m^d$) in terms of the surface value of ($\vec{E}^d, \vec{H}^d$):

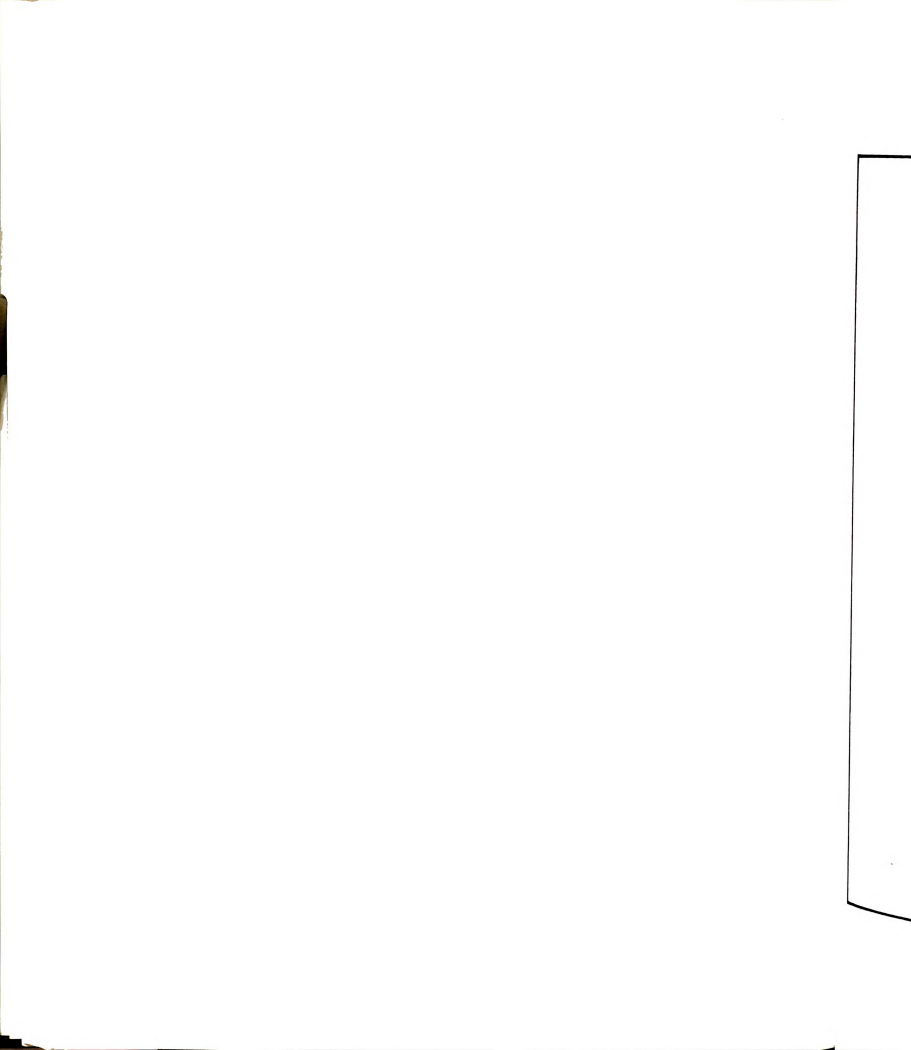
$$\vec{j}_e^d \equiv \hat{n} \times \vec{H}^d$$

$$\vec{j}_m^d \equiv -\hat{n} \times \vec{E}^d / \eta_0$$

$$q_e^d = \frac{1}{\omega} \nabla_{\odot} \cdot \vec{j}_e^d$$

$$q_m^d = \frac{1}{\omega} \nabla_{\odot} \cdot \vec{j}_m^d$$

Again, from the equivalence theorem [15], [16], we know that, with the exterior region of G being replaced by a homogeneous source free region with permeability μ_0 and complex permittivity ϵ_d , the negative of the equivalent surface currents ($\vec{j}_e^d, \vec{j}_m^d$) will support the total induced field ($\vec{E}^d, \vec{H}^d$) in the region internal to F and a null field in the region external to F, as shown in Figure 3.5. Equations (3.3) and (3.4) can be applied to determine these fields; a sign change is necessary since we define ($\vec{j}_e^d, \vec{j}_m^d$)



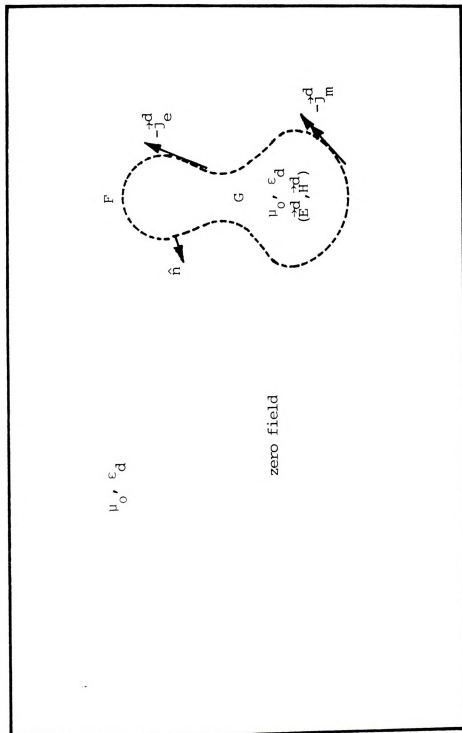


Figure 3.5. The induced field $(\vec{E}^d, \vec{H}^d)$ in the interior of the body and a zero field in the exterior of the body are maintained by the negative of the equivalent surface currents $(\vec{j}_e^d, \vec{j}_m^d)$, which are defined in terms of the induced field on surface F . The region external to F is replaced by a medium with permeability μ_0 and complex permittivity ϵ_d .

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in terms of the unit normal vector $\hat{n}$ which is pointing into the source region. Hence we have

$$\vec{E}^d = \frac{1}{4\pi} \int_F [\eta_o \vec{j}_m^d \times \nabla' \phi_d - \frac{q_e^d}{\epsilon_d} \nabla' \phi_d + j\omega \mu_o \vec{j}_e^d \phi_d] dF' \quad (3.12)$$

$$\vec{H}^d = \frac{1}{4\pi} \int_F [-\vec{j}_e^d \times \nabla' \phi_d - \frac{q_m^d}{\mu_o} \nabla' \phi_d + j\omega \epsilon_d \eta_o \vec{j}_m^d \phi_d] dF' \quad (3.13)$$

internal
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as the induced field at any point in the region internal to the surface F, where $\phi_d = \frac{e^{-jk_d |\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|}$ is the Green's function for the conducting body region, and $k_d = \omega \sqrt{\mu_o \epsilon_d}$ is the complex wavenumber in the body. In the region external to F, the above two integrals vanish, i.e.,

$$0 = \frac{1}{4\pi} \int_F [\eta_o \vec{j}_m^d \times \nabla' \phi_d - \frac{q_e^d}{\epsilon_d} \nabla' \phi_d + j\omega \mu_o \vec{j}_e^d \phi_d] dF' \quad (3.14)$$

$$0 = \frac{1}{4\pi} \int_F [-\vec{j}_e^d \times \nabla' \phi_d - \frac{q_m^d}{\mu_o} \nabla' \phi_d + j\omega \epsilon_d \eta_o \vec{j}_m^d \phi_d] dF' \quad (3.15)$$

external
to F

because there is a null field external to F.

III) The Incident Field ($\vec{E}^i, \vec{H}^i$):

To investigate the incident field ($\vec{E}^i, \vec{H}^i$), let's consider the geometry shown in Figure 3.2 again. Based on ($\vec{E}^i, \vec{H}^i$), we define the following equivalent surface currents on the body surface F:

$$\vec{j}_e^i \equiv \hat{n} \times \vec{H}^i$$

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$$\vec{j}_m^i = -\hat{n} \times \vec{E}_e^i / \eta_0$$

$$q_e^i = \frac{1}{\omega} \nabla_0 \cdot \vec{j}_e^i$$

$$q_m^i = \frac{1}{\omega} \nabla_0 \cdot \vec{j}_m^i$$

From the induction theorem [15], [16], we know that the negative of these equivalent surface currents, radiating in the presence of the conducting body, will support the total induced field ($\vec{E}^d$, $\vec{H}^d$) in the region internal to surface F, and the scattered field ($\vec{E}^s$, $\vec{H}^s$) external to F. This is illustrated in Figure 3.6.

Suppose that we are considering an extreme case with the complex permittivity of the body ϵ_d being replaced by the free space permittivity ϵ_0 , then the negative of the surface currents ($\vec{j}_e^i$, $\vec{j}_m^i$) will support the incident field ($\vec{E}^i$, $\vec{H}^i$) in the region internal to F, and a null field in the region external to F. Because the entire space becomes homogeneous (free space), there exists no real scattering obstacle—the conducting body, so there is no scattered field. This situation is shown in Figure 3.7.

Again, eqs. (3.3) and (3.4) can be used to express the EM fields shown in Figure 3.7. We have

$$\vec{E}^i = \frac{1}{4\pi} \int_F [\eta_0 \vec{j}_m^i \times \nabla' \phi_f - \frac{q_e^i}{\epsilon_0} \nabla' \phi_f + j\omega\mu_0 \vec{j}_e^i \phi_f] dF' \quad \left. \begin{array}{l} \text{internal} \\ \text{to F} \end{array} \right\} \quad (3.16)$$



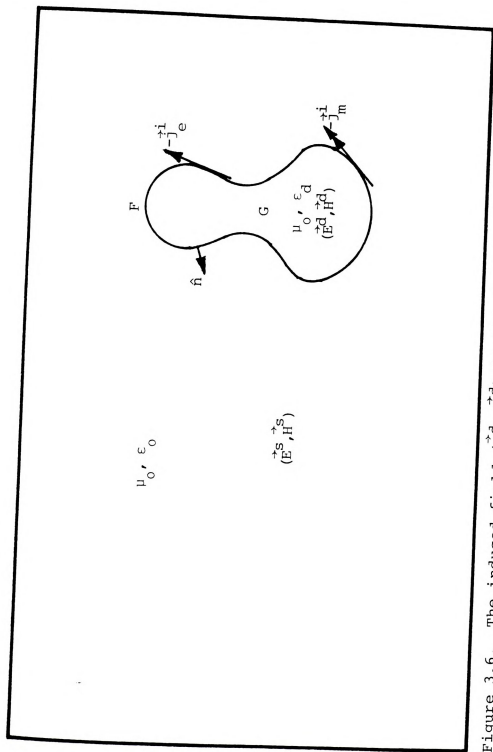
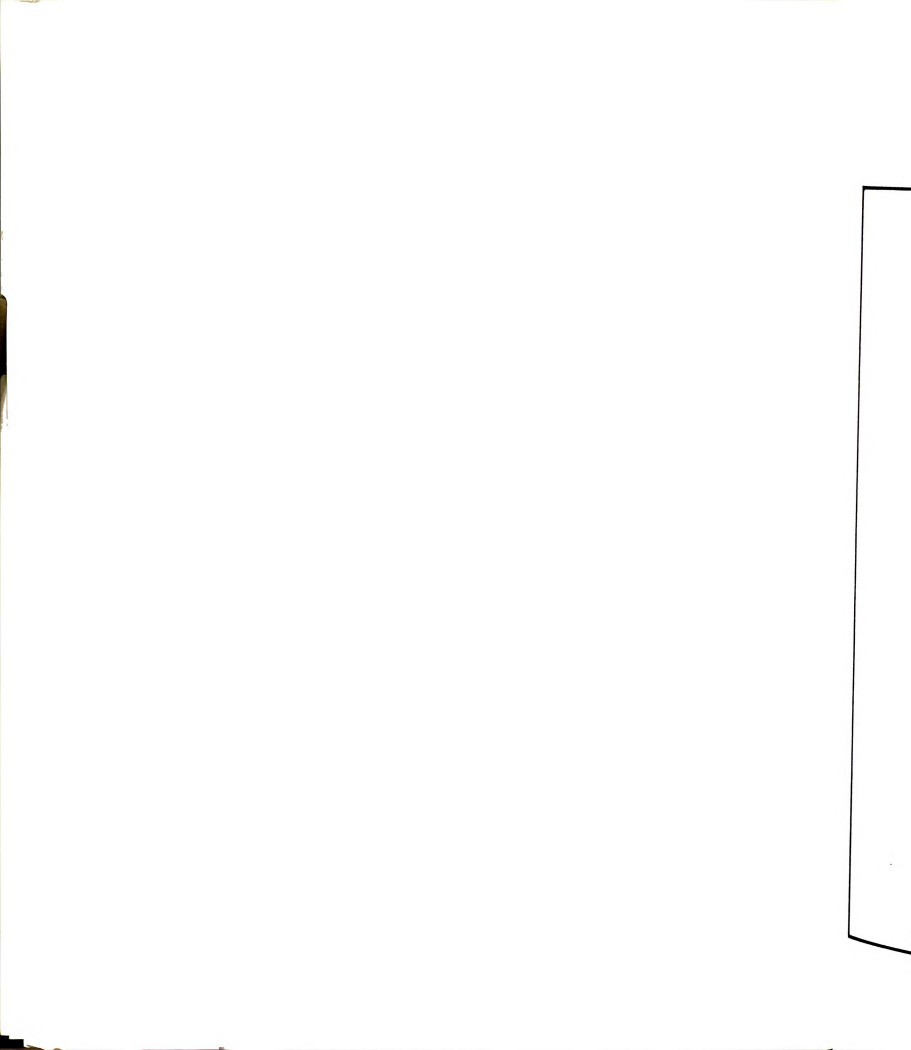


Figure 3.6. The induced field $(\vec{E}^d, \vec{H}^d)$ in the interior of the body and the scattered field $(\vec{E}^s, \vec{H}^s)$ in the exterior of the body are mainly defined by the negative of the equivalent surface currents $(\vec{j}_e, \vec{j}_m)$, which are defined in terms of the incident field on surface F . The body G with permeability μ_0 and complex permittivity ϵ_d is immersed in free space.



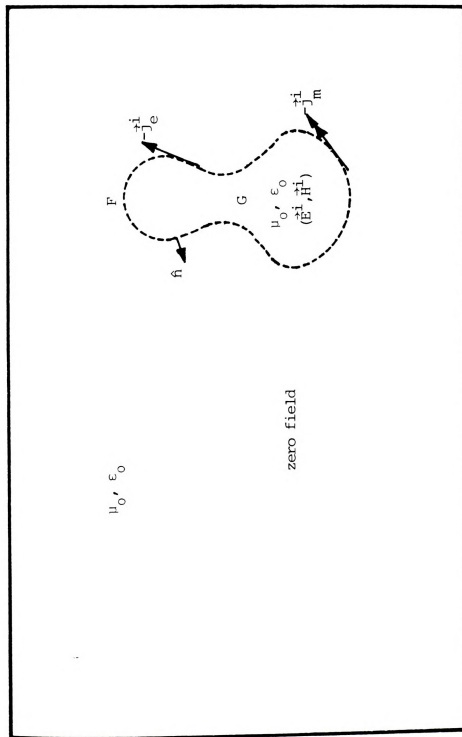


Figure 3.7. With the body region G replaced by free space, the incident field $(\vec{E}^i, \vec{H}^i)$ in the interior of the body and a zero field in the exterior of the body are maintained by the negative of the equivalent surface currents $(\vec{j}_e^i, \vec{j}_m^i)$, which are defined in terms of the incident field on surface F.

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$$\vec{H}^i = \frac{1}{4\pi} \int_F [-\vec{j}_e^i \times \nabla' \phi_f - \frac{q_e^i}{\mu_0} \nabla' \phi_f + j\omega \epsilon_0 \vec{j}_m^i \phi_f] dF' \left\{ \begin{array}{l} \text{internal} \\ \text{to } F \end{array} \right. \quad (3.17)$$

as the incident field at any point in the region internal to F. However, the above two integrals vanish in the exterior region of the body, i.e.,

$$0 = \frac{1}{4\pi} \int_F [\eta_0 \vec{j}_m^i \times \nabla' \phi_f - \frac{q_e^i}{\epsilon_0} \nabla' \phi_f + j\omega \mu_0 \vec{j}_e^i \phi_f] dF' \left\{ \begin{array}{l} \text{external} \\ \text{to } F \end{array} \right. \quad (3.18)$$

$$0 = \frac{1}{4\pi} \int_F [-\vec{j}_e^i \times \nabla' \phi_f - \frac{q_m^i}{\mu_0} \nabla' \phi_f + j\omega \epsilon_0 \eta_0 \vec{j}_m^i \phi_f] dF' \left\{ \begin{array}{l} \text{external} \\ \text{to } F \end{array} \right. \quad (3.19)$$

Up to this point, we have considered the incident field ($\vec{E}^i, \vec{H}^i$), the induced field ($\vec{E}^d, \vec{H}^d$), and the scattered field ($\vec{E}^s, \vec{H}^s$) individually. Now we are ready to derive the coupled surface integral equations based on the results we just obtained.

Since eq. (3.16) is true at any point $\vec{r}$ in the interior region of G, if we let the field point $\vec{r}$ approach a point $\vec{\zeta}$ on the surface F along the interior normal, then the limiting value of eq. (3.16) will be the incident electric field at the point $\vec{\zeta}$ on the surface F. Hence we may write

$$\vec{E}^i(\vec{\zeta}) = \frac{1}{4\pi} \int_{F_i} [\eta_0 \vec{j}_m^i \times \nabla' \phi_f - \frac{q_e^i}{\epsilon_0} \nabla' \phi_f + j\omega \mu_0 \vec{j}_e^i \phi_f] dF'$$

where $\int_{F_i}$ has the same meaning as that described in theorem 2. After taking the cross-product of the above

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$$\begin{aligned} \hat{n} \times \frac{1}{4\pi} \int_{F_i} [\eta_o \vec{j}_m^i \times \nabla' \phi_f - \frac{q_e^i}{\epsilon} \nabla' \phi_f + j\omega \mu_o \vec{j}_e^i \phi_f] dF' \\ = -\eta_o \vec{j}_m^i. \end{aligned} \quad (3.20)$$

Applying the same arguments to eq. (3.17), we have

$$\begin{aligned} -\hat{n} \times \frac{1}{4\pi} \int_{F_i} [-\vec{j}_e^i \times \nabla' \phi_f - \frac{q_m^i}{\mu_o} \nabla' \phi_f + j\omega \epsilon_o \eta_o \vec{j}_m^i \phi_f] dF' \\ = -\vec{j}_e^i. \end{aligned} \quad (3.21)$$

Also, from eqs. (3.10) and (3.11), it is straightforward to obtain

$$\hat{n} \times \frac{1}{4\pi} \int_{F_i} [-\eta_o \vec{j}_m^s \times \nabla' \phi_f + \frac{q_e^s}{\epsilon} \nabla' \phi_f - j\omega \mu_o \vec{j}_e^s \phi_f] dF' = 0 \quad (3.22)$$

and

$$-\hat{n} \times \frac{1}{4\pi} \int_{F_i} [\vec{j}_e^s \times \nabla' \phi_f + \frac{q_m^s}{\mu_o} \nabla' \phi_f - j\omega \epsilon_o \eta_o \vec{j}_m^s \phi_f] dF' = 0. \quad (3.23)$$

Before going any further, we should introduce the boundary conditions which must be satisfied at the surface F . The boundary conditions require that the tangential components of the electric and magnetic field must be continuous across F , or equivalently we may write

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Now, take the difference between eqs. (3.20) and (3.22), and apply the above boundary conditions. We have

$$\begin{aligned} \hat{n} \times \frac{1}{4\pi} \int_{F_i} [\eta_o \vec{j}_m^d \times \nabla' \phi_f - \frac{q_e^d}{\epsilon_o} \nabla' \phi_f + j\omega \mu_o \vec{j}_e^d \phi_f] dF' \\ = -\eta_o \vec{j}_m^i. \end{aligned} \quad (3.24)$$

Similarly, subtracting eq. (3.23) from eq. (3.21), we obtain

$$\begin{aligned} -\hat{n} \times \frac{1}{4\pi} \int_{F_i} [-\vec{j}_e^d \times \nabla' \phi_f - \frac{q_m^d}{\mu_o} \nabla' \phi_f + j\omega \epsilon_o \eta_o \vec{j}_m^d \phi_f] dF' \\ = -\vec{j}_e^i. \end{aligned} \quad (3.25)$$

On the other hand, we may derive the following two equations from eqs. (3.14) and (3.15):

$$\hat{n} \times \frac{1}{4\pi} \int_{F_e} [\eta_o \vec{j}_m^d \times \nabla' \phi_d - \frac{q_e^d}{\epsilon_d} \nabla' \phi_d + j\omega \mu_o \vec{j}_e^d \phi_d] dF' = 0 \quad (3.26)$$

$$\begin{aligned} -\hat{n} \times \frac{1}{4\pi} \int_{F_e} [-\vec{j}_e^d \times \nabla' \phi_d - \frac{q_m^d}{\mu_o} \nabla' \phi_d + j\omega \epsilon_d \eta_o \vec{j}_m^d \phi_d] dF' = 0 \\ (3.27) \end{aligned}$$

It can be shown that equalities exist between the following limiting surface integrals and their corresponding regular surface integrals

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$$\hat{n} \times \int_{F_i} q_e^d \nabla' \phi_f dF' = \hat{n} \times \int_F q_e^d \nabla' \phi_f dF'$$

$$\hat{n} \times \int_{F_e} q_e^d \nabla' \phi_d dF' = \hat{n} \times \int_F q_e^d \nabla' \phi_d dF'.$$

Subtracting the second equation from the first equation, we find that

$$\hat{n} \times \int_{F_i} q_e^d \nabla' \phi_f dF' - \hat{n} \times \int_{F_e} q_e^d \nabla' \phi_d dF' =$$

$$\hat{n} \times \int_F q_e^d \nabla' (\phi_f - \phi_d) dF' = -\hat{n} \times \nabla \int_F q_e^d (\phi_f - \phi_d) dF'.$$

However, since $q_e^d = \frac{j}{\omega} \nabla' \cdot \vec{j}_e^d$, the integral on the right-hand side of the above equation can be rewritten as

$$\int_F q_e^d (\phi_f - \phi_d) dF' = \frac{j}{\omega} \int_F (\nabla' \cdot \vec{j}_e^d) (\phi_f - \phi_d) dF'.$$

Applying theorem 1 to the right-hand side of the above equation, we can rewrite it as

$$\begin{aligned} \int_F q_e^d (\phi_f - \phi_d) dF' &= -\frac{j}{\omega} \int_F \vec{j}_e^d \cdot \nabla' (\phi_f - \phi_d) dF' \\ &= -\frac{j}{\omega} \int_F \vec{j}_e^d \cdot \nabla' (\phi_f - \phi_d) dF'. \end{aligned}$$

It should be mentioned that the last equality is true, because $\nabla' = \nabla - \hat{n} \frac{\partial}{\partial n}$, and $\vec{j}_e^d$ has no normal component. Then finally, we have

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$$\begin{aligned}
& \hat{n} \times \int_{F_i} q_e^d \nabla' \phi_f dF' - \hat{n} \times \int_{F_e} q_e^d \nabla' \phi_d dF' \\
& = -\hat{n} \times \nabla \left[-\frac{j}{\omega} \int_F j_e^d \cdot \nabla' (\phi_f - \phi_d) dF' \right] \\
& = \frac{j}{\omega} \hat{n} \times \int_F \nabla [j_e^d \cdot \nabla' (\phi_f - \phi_d)] dF'.
\end{aligned}$$

Based on the following vector identity

$$\nabla (\vec{A} \cdot \vec{B}) = (\vec{A} \cdot \nabla) \vec{B} + (\vec{B} \cdot \nabla) \vec{A} + \vec{A} \times (\nabla \times \vec{B}) + \vec{B} \times (\nabla \times \vec{A}),$$

it is easy to show that

$$\begin{aligned}
& \hat{n} \times \int_{F_i} q_e^d \nabla' \phi_f dF' - \hat{n} \times \int_{F_e} q_e^d \nabla' \phi_d dF' \\
& = -\frac{j}{\omega} \hat{n} \times \int_F [(j_e^d \cdot \nabla') \nabla' (\phi_f - \phi_d)] dF'. \quad (3.28)
\end{aligned}$$

Now, subtracting eq. (3.26) from eq. (3.24) with the latter multiplied by ϵ_0 and the former multiplied by ϵ_d , the result is

$$\begin{aligned}
& \epsilon_0 \hat{n} \times \frac{1}{4\pi} \int_{F_i} [\eta_o j_m^d \times \nabla' \phi_f - \frac{q_e^d}{\epsilon_0} \nabla' \phi_f + j\omega \mu_o j_o^d \phi_f] dF' - \\
& \epsilon_d \hat{n} \times \frac{1}{4\pi} \int_{F_e} [\eta_o j_m^d \times \nabla' \phi_d - \frac{q_e^d}{\epsilon_d} \nabla' \phi_d + j\omega \mu_o j_o^d \phi_d] dF' \\
& = -\epsilon_0 \eta_o j_o^d
\end{aligned}$$

or

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& \frac{1}{4\pi} \{ [\hat{n} \times \int_{F_i} \epsilon_o \eta_o \vec{j}_m^d \times \nabla' \phi_f \, dF' - \hat{n} \times \int_{F_e} \epsilon_d \eta_o \vec{j}_m^d \times \nabla' \phi_d \, dF'] \\
& - [\hat{n} \times \int_{F_i} q_e^d \nabla' \phi_f \, dF' - \hat{n} \times \int_{F_e} q_e^d \nabla' \phi_d \, dF'] + \\
& [\hat{n} \times \int_{F_i} j_{\omega\mu} \epsilon_o \vec{j}_e^d \phi_f \, dF' - \hat{n} \times \int_{F_e} j_{\omega\mu} \epsilon_d \vec{j}_e^d \phi_d \, dF'] \} \\
& = -\epsilon_o \eta_o \vec{j}_m^i .
\end{aligned}$$

Based on eq. (3.28) and theorem 2, the above equation becomes

$$\begin{aligned}
\vec{j}_m^d &= \frac{2}{\kappa^2 + 1} \vec{j}_m^i - \frac{1}{2\pi(\kappa^2 + 1)} \hat{n} \times \int_F [jk_o \vec{j}_e^d (\kappa^2 \phi_d - \phi_f) + \\
& \vec{j}_m^d \times \nabla' (\kappa^2 \phi_d - \phi_f) + \frac{j}{k_o} (\vec{j}_e^d \cdot \nabla') \nabla' (\phi_d - \phi_f)] \, dF' .
\end{aligned}$$

Similarly, the other surface integral equation can be derived from eqs. (3.25) and (3.27),

$$\begin{aligned}
\vec{j}_e^d &= \vec{j}_e^i + \frac{1}{4\pi} \hat{n} \times \int_F [jk_o \vec{j}_m^d (\kappa^2 \phi_d - \phi_f) - \vec{j}_e^d \times \nabla' (\phi_d - \phi_f) + \\
& \frac{j}{k_o} (\vec{j}_m^d \cdot \nabla') \nabla' (\phi_d - \phi_f)] \, dF'
\end{aligned}$$

where $\kappa = \sqrt{\epsilon_d / \epsilon_o}$ is the refractive index of the body. If we normalize all the quantities concerning the dimension of length by $1/k_o$, the above two coupled surface integral equations become (these two equations in similar forms have also been given by Müller [13] and Morita [19] through different derivations)

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$$\vec{j}_m^d = \frac{2}{\kappa^2 + 1} \vec{j}_m^i - \frac{1}{2\pi(\kappa^2 + 1)} \hat{n} \times \int_F [j(\kappa^2 \phi_d - \phi_f) \vec{j}_e^d + \vec{j}_m^d \times \nabla'(\kappa^2 \phi_d - \phi_f) + j(\vec{j}_e^d \cdot \nabla') \nabla'(\phi_d - \phi_f)] dF' \quad (3.29)$$

and

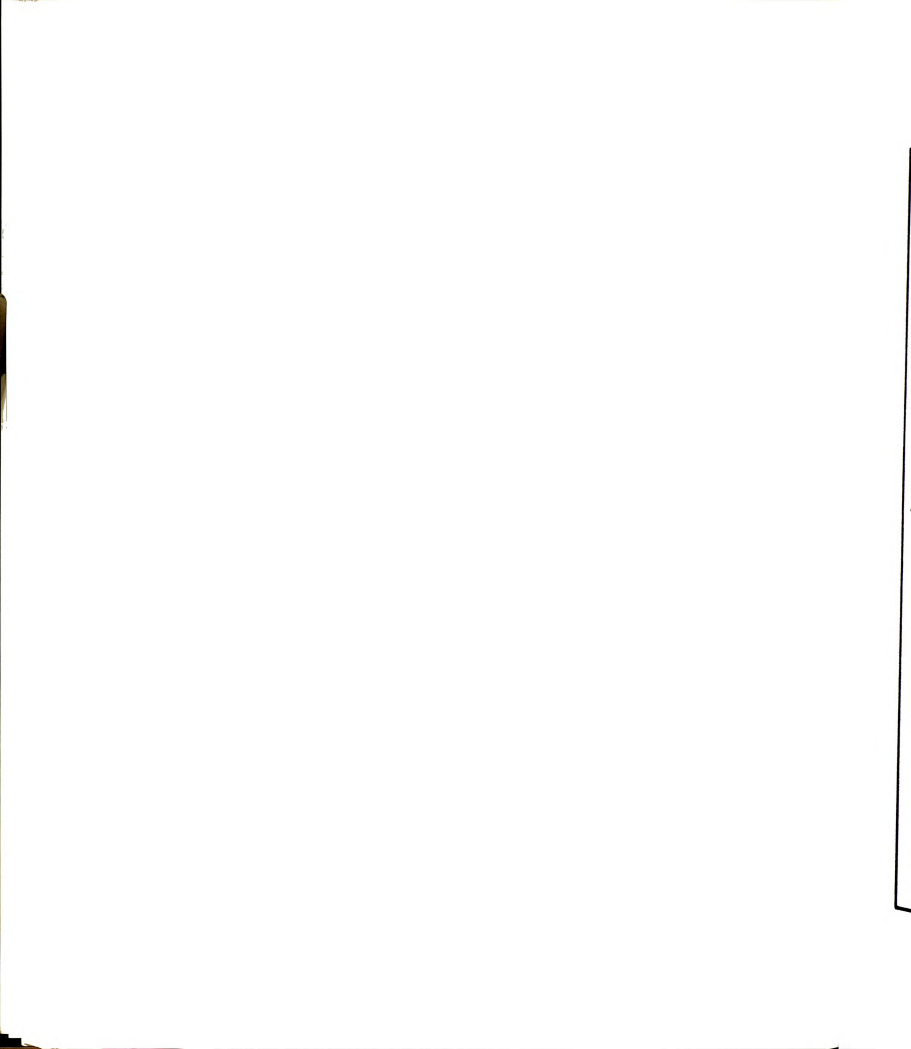
$$\vec{j}_e^d = \vec{j}_e^i + \frac{1}{4\pi} \hat{n} \times \int_F [j(\kappa^2 \phi_d - \phi_f) \vec{j}_m^d - \vec{j}_e^d \times \nabla'(\phi_d - \phi_f) + j(\vec{j}_m^d \cdot \nabla') \nabla'(\phi_d - \phi_f)] dF' . \quad (3.30)$$

The induced equivalent surface currents ($\vec{j}_e^d$, $\vec{j}_m^d$) (or the induced surface field ($\vec{E}^d$, $\vec{H}^d$)) can be calculated based on the above two coupled, surface integral equations. In these two equations, due to normalization with respect to $1/k_0$, the Green's functions ϕ_d and ϕ_f have the forms of

$$\phi_d = \frac{e^{-j\kappa|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} \quad \text{and} \quad \phi_f = \frac{e^{-j|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} .$$

3.5. The Special Case of an Infinite Plane Interface

In this section, we apply the surface integral equations to the special case of an infinite plane interface between two half spaces, a free-space in the region $z < 0$ with permeability μ_0 and permittivity ϵ_0 and a finite conducting medium in the region $z > 0$ with permeability μ_0 and complex permittivity ϵ_d , as shown in Figure 3.8. We assume that the incident plane wave, which propagates in



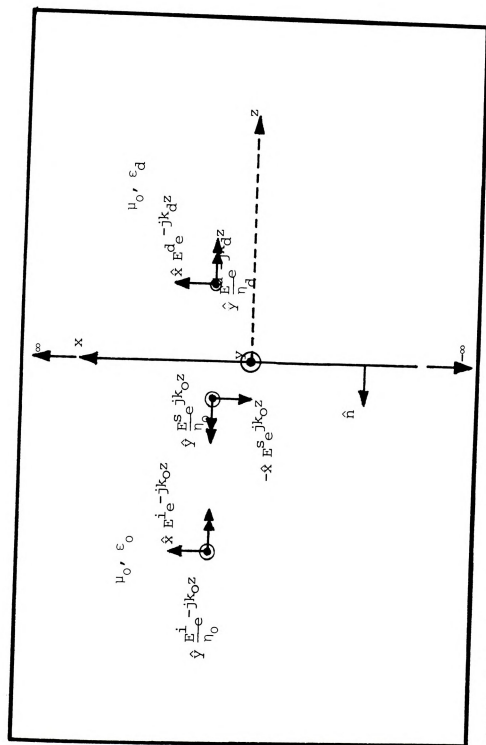


Figure 3.8. An infinite interface between free space and a finite conducting medium with permeability μ_0 and complex permittivity ϵ_d is impinged by a plane wave propagating in the $+z$ direction.

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the +z direction with its electric field $\vec{E}^i$ polarized in the +x direction and its magnetic field $\vec{H}^i$ polarized in the +y direction, is normally impinging on the interface, and produces a reflected field ($\vec{E}^s, \vec{H}^s$) and a transmitted field ($\vec{E}^d, \vec{H}^d$) in the free space region and the conducting medium region, respectively. Intuitively, these fields can be expressed as

$$\left. \begin{aligned} \vec{E}^i &= \hat{x} E^i e^{-jk_o z} \\ \vec{H}^i &= \hat{y} H^i e^{-jk_o z} = \hat{y} \frac{E^i}{\eta_o} e^{-jk_o z} , \\ \vec{E}^s &= -\hat{x} E^s e^{jk_o z} \\ \vec{H}^s &= \hat{y} \frac{E^s}{\eta_o} e^{jk_o z} \end{aligned} \right\} \text{ in the region } z < 0$$

and

$$\left. \begin{aligned} \vec{E}^d &= \hat{x} E^d e^{-jk_d z} \\ \vec{H}^d &= \hat{y} H^d e^{-jk_d z} = \hat{y} \frac{E^d}{\eta_d} e^{-jk_d z} \end{aligned} \right\} \text{ in the region } z > 0$$

where $\eta_d = \sqrt{\mu_o / \epsilon_d}$ is the complex wave impedance in the conducting medium. By solving a simple boundary value problem, it is easy to show that, on the interface, the transmitted field (E^d, H^d) is related to the incident field (E^i, H^i) by

$$\begin{aligned} E^d &= \frac{2}{1 + \kappa} E^i \\ H^d &= \frac{2\kappa}{1 + \kappa} H^i. \end{aligned} \quad (3.31)$$

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It would be interesting to show that the results of eq. (3.31) can be deduced from the rather complicated surface integral equations which have been derived in the previous section. To do that, let's start from the following definitions of the various equivalent surface current densities.

$$\left. \begin{aligned} \vec{j}_e^i &\equiv \hat{n} \times \vec{H}^i = -\hat{z} \times \hat{y} H^i = \hat{x} \frac{E^i}{\eta_0} \\ \vec{j}_m^i &\equiv -\hat{n} \times \frac{\vec{E}^i}{\eta_0} = \hat{z} \times \hat{x} \frac{E^i}{\eta_0} = \hat{y} \frac{E^i}{\eta_0} \end{aligned} \right\} \text{ at } z = 0$$

also

$$\begin{aligned} \vec{j}_e^d &\equiv \hat{n} \times \vec{H}^d = \hat{x} j_e^d \\ \vec{j}_m^d &\equiv -\hat{n} \times \frac{\vec{E}^d}{\eta_0} = \hat{y} j_m^d \end{aligned}$$

where $\hat{n}$ is the unit normal vector on the interface pointing in the $-z$ direction.

Now, substituting the above expressions of surface current densities into eq. (3.29), we may show that the last two terms of the surface integration contribute nothing to the total surface integral. We should consider these two terms one by one as below.

I) Let us consider the term $\hat{n} \times \int_F \vec{j}_m^d \times \nabla' (\kappa^2 \phi_d - \phi_F) dF'$ first.

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It is easy to show that

$$\hat{n} \times \int_F \vec{j}_m^d \times \nabla' (\kappa^2 \phi_d - \phi_f) dF' = (-\hat{z}) \times \int_F \hat{y} \times \vec{j}_m^d \times \nabla' (\kappa^2 \phi_d - \phi_f) dF' = 0 \quad (3.32)$$

is true, since the term $\nabla' (\kappa^2 \phi_d - \phi_f)$ has only two components, x and y components but not z component (nothing is z-dependent in this problem), both of which vanish after cross-producting with unit vectors $\hat{y}$ and $\hat{z}$ consecutively.

II) Secondly, consider the term of $\hat{n} \times \int_F (\vec{j}_e^d \cdot \nabla') \nabla' (\phi_d - \phi_f) dF'$.

Since the surface current $\vec{j}_e^d$ is x-directed, we may express the first part of the integrand as

$$(\vec{j}_e^d \cdot \nabla') = j_{ex}^d \frac{\partial}{\partial x'}.$$

Because everything is z-independent, we may write the second part of the integrand as

$$\nabla' (\phi_d - \phi_f) = \hat{x} \frac{\partial}{\partial x'} (\phi_d - \phi_f) + \hat{y} \frac{\partial}{\partial y'} (\phi_d - \phi_f).$$

Combining the above two equations, we have the following expression for the integrand:

$$(\vec{j}_e^d \cdot \nabla') \nabla' (\phi_d - \phi_f) = \hat{x} j_{ex}^d \frac{\partial^2}{\partial x'^2} (\phi_d - \phi_f) + \hat{y} j_{ex}^d \frac{\partial^2}{\partial x' \partial y'} (\phi_d - \phi_f).$$

Hence,

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Hence, the integral considered can be rewritten as

$$\begin{aligned} \hat{n} \times \int_F (\hat{j}_e^d \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' &= \hat{n} \times \int_F \hat{x} j_{ex}^d \frac{\partial^2}{\partial x'^2} \\ &(\phi_d - \phi_f) dF' + \hat{n} \times \int_F \hat{y} j_{ey}^d \frac{\partial^2}{\partial x' \partial y'} (\phi_d - \phi_f) dF'. \end{aligned} \quad (3.33)$$

If we choose the field point to be at origin; i.e., $\vec{r} = 0$, we may write the Green's functions ϕ_d and ϕ_f as

$$\phi_d = \frac{e^{-jk_r r'}}{r'}, \quad \text{and} \quad \phi_f = \frac{e^{-jr'}}{r'},$$

where $r' = |\vec{r}'| = \sqrt{x'^2 + y'^2}$. With some straightforward algebraic manipulations, it can be shown that

$$\begin{aligned} \frac{\partial^2}{\partial x'^2} (\phi_d - \phi_f) &= [j \frac{e^{-jr'} - \kappa e^{-jk_r r'}}{r'^2} + \frac{e^{-jr'} - e^{-jk_r r'}}{r'^3}] + \\ \cos^2 \theta &[\frac{e^{-jr'} - \kappa^2 e^{-jk_r r'}}{r'} - 3j \frac{e^{-jr'} - \kappa e^{-jk_r r'}}{r'^2} \\ &- 3 \frac{e^{-jr'} - e^{-jk_r r'}}{r'^3}] \end{aligned}$$

where θ is the angle between the source-point position vector $\vec{r}'$ and the x axis, as shown in Figure 3.9. Similarly, it can be shown that

$$\frac{\partial^2}{\partial x' \partial y'} (\phi_d - \phi_f) = \sin \theta \cos \theta [\frac{e^{-jr'} - \kappa^2 e^{-jk_r r'}}{r'}]$$

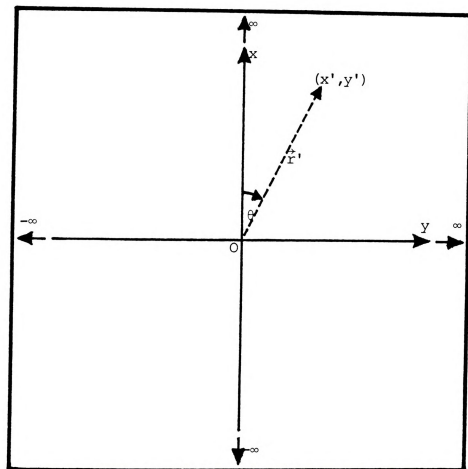


Figure 3.9. The coordinate system utilized in the infinite interface problem.

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$$- 3j \frac{e^{-jr'} - \kappa e^{-jk r'}}{r'^2} - 3 \frac{e^{-jr'} - e^{-jk r'}}{r'^3}] .$$

In the above derivation, the field point has been arbitrarily chosen at the origin. This assumption doesn't destroy the generality of this treatment, since the interface is infinitely large, and the field induced on it is uniformly distributed over it.

Now, we must show that both of the integrals on the right-hand side of eq. (3.33) vanish. Let us consider the last term of that equation first.

$$\int_F \frac{\partial^2}{\partial x' \partial y'} (\phi_d - \phi_f) dF' = \int_0^{2\pi} \sin \theta \cos \theta d\theta \cdot \int_0^\infty [e^{-jr'} - \kappa^2 e^{-jk r'} - 3j \frac{e^{-jr'} - \kappa e^{-jk r'}}{r'} - 3 \frac{e^{-jr'} - e^{-jk r'}}{r'^2}] dr'$$

It is clear that the first integral on the right-hand side of the above equation is equal to zero. Hence we have

$$\int_F \frac{\partial^2}{\partial x' \partial y'} (\phi_d - \phi_f) dF' = 0. \quad (3.34)$$

For the first term on the right-hand side of eq.

(3.33), we may write that

$$\int_F \frac{\partial^2}{\partial x'^2} (\phi_d - \phi_f) dF' = 2\pi \int_0^\infty [j \frac{e^{-jr'} - \kappa e^{-jk r'}}{r'} +$$

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$$\begin{aligned}
& \frac{e^{-jr'} - e^{-jkr'}}{r'^2} \bigg] dr' + \int_0^{2\pi} \cos^2 \theta \, d\theta \cdot \int_0^\infty [e^{-jr'} - \kappa^2 e^{-jkr'} \\
& - 3j \frac{e^{-jr'} - \kappa e^{-jkr'}}{r'} - 3 \frac{e^{-jr'} - e^{-jkr'}}{r'^2}] dr' \\
& = -j\pi \int_0^\infty \frac{e^{-jr'} - \kappa e^{-jkr'}}{r'} dr' - \pi \int_0^\infty \frac{e^{-jr'} - e^{-jkr'}}{r'^2} dr' \\
& + j\kappa\pi - j\pi = \pi \int_0^\infty \frac{d}{dr'} \left(\frac{e^{-jr'} - e^{-jkr'}}{r'} \right) dr' + j\kappa\pi - j\pi \\
& = \pi \lim_{r' \rightarrow 0} \left[\frac{e^{-jkr'} - e^{-jr'}}{r'} \right] + j\kappa\pi - j\pi = 0. \quad (3.35)
\end{aligned}$$

Note that for the last step in the above derivation, the radiation condition is applied.

After substitution of eqs. (3.34) and (3.35) into eq. (3.33), the result is

$$\hat{n} \times \int_F (\vec{j}_e^d \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' = 0. \quad (3.36)$$

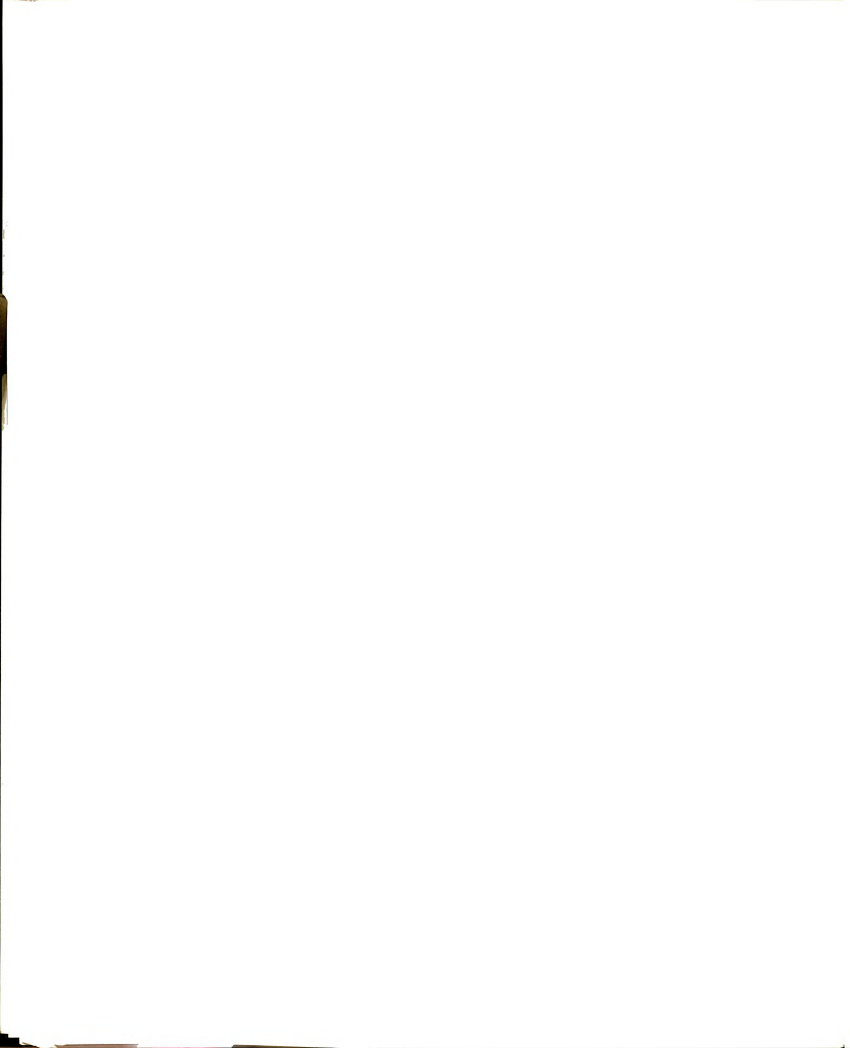
Then, substituting eqs. (3.32) and (3.36) into eq. (3.29), we obtain

$$j_m^d = \left(\frac{2}{\kappa^2 + 1} \right) \frac{E^i}{\eta_0} + \frac{j}{2\pi(\kappa^2 + 1)} j_e^d \int_F (\kappa^2 \phi_d - \phi_f) dF'. \quad (3.37)$$

The integral term is easy to evaluate. We write down the result directly.

$$\int_F (\kappa^2 \phi_d - \phi_f) dF' = j 2\pi(1 - \kappa)$$

Hence, eq. (3.37) becomes



$$j_m^d = \left(\frac{2}{\kappa+1}\right) \frac{E^i}{\eta_0} + \frac{\kappa-1}{\kappa+1} j_e^d. \quad (3.38)$$

Following a similar procedure, we can obtain another equation from eq. (3.30), which is

$$j_e^d = \frac{E^i}{\eta_0} + \frac{\kappa-1}{2} j_m^d. \quad (3.39)$$

Equations (3.29) and (3.30) have been simplified to become two linear equations with unknowns j_e^d and j_m^d . The solutions of these two unknowns are

$$j_e^d = \left(\frac{2\kappa}{\kappa+1}\right) \frac{E^i}{\eta_0}$$

and

$$j_m^d = \left(\frac{2}{\kappa+1}\right) \frac{E^i}{\eta_0},$$

which also can be written as

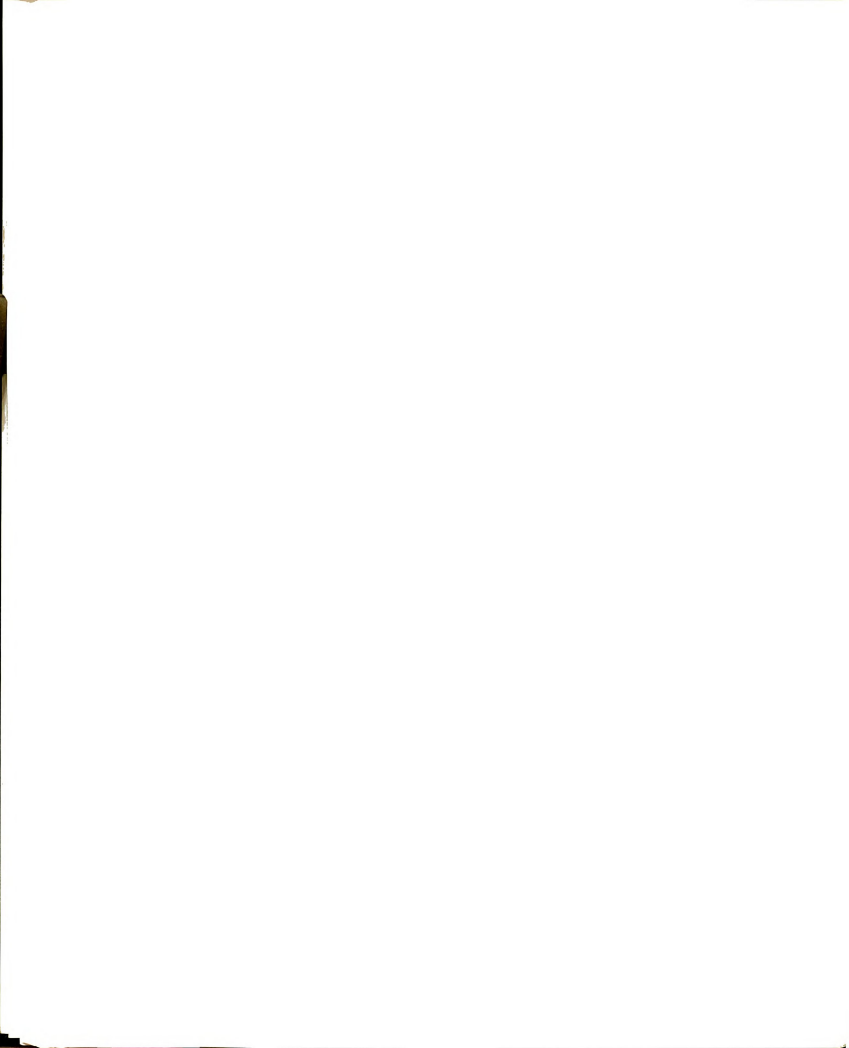
$$E^d = \left(\frac{2}{\kappa+1}\right) E^i$$

and

$$H^d = \left(\frac{2\kappa}{\kappa+1}\right) H^i$$

which are exactly the same as eq. (3.31).

We have shown that the coupled, surface integral equations can predict correct results for the simplified special case of an infinite two-dimensional interface between a finite conducting medium and free space. This



example analytically proves the validity of the surface integral equations. In the following sections, we will develop a numerical technique for solving these surface integral equations for the cases of arbitrarily shaped, finite conducting bodies; after that, some numerical examples will be presented to discuss the accuracy of this numerical technique.

3.6. Review of Moment-Method

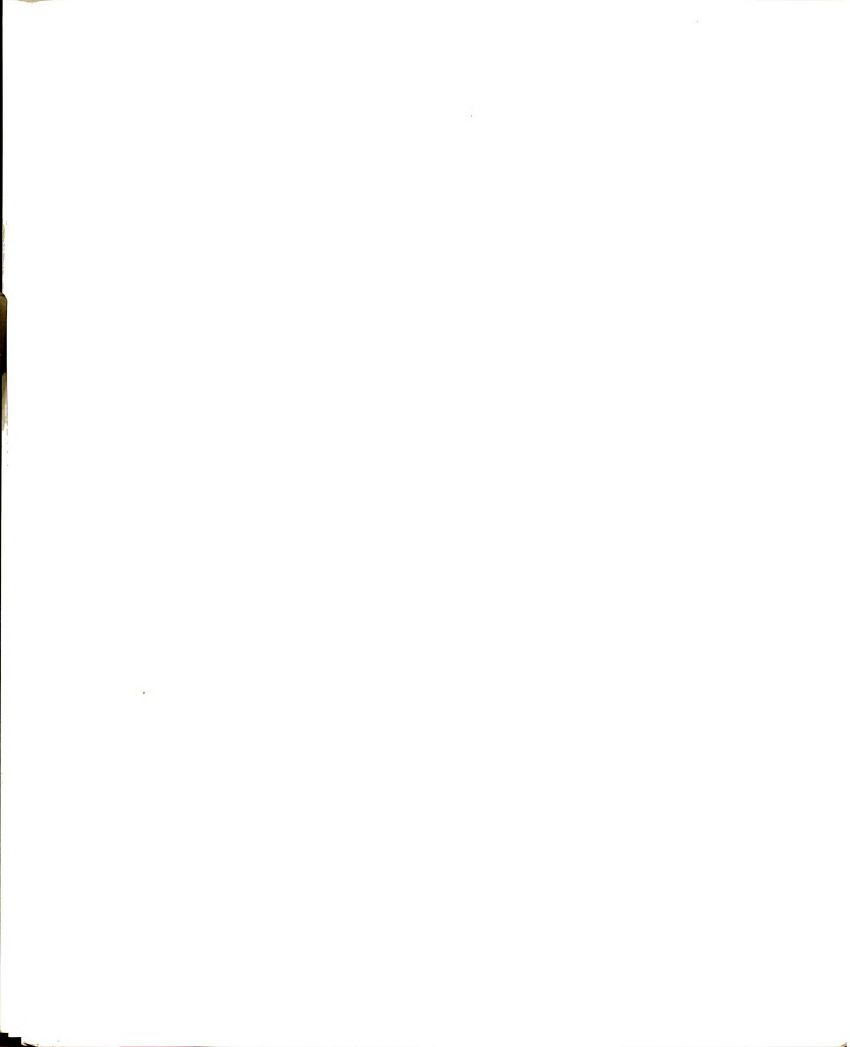
Before actually solving the surface integral equations, we should briefly review the basic steps involved in numerical implementation of such solutions, or the so-called Method of Moments (MOM) [17]. The idea of moment-method solutions is that the linear operator equation

$$L(f) = g \quad (3.40)$$

where L is a linear operator, g is some known driving function, and f is to be determined, may be solved by first expanding f in a series of basis functions f_n with unknown weighting coefficients α_n ,

$$f = \sum_n \alpha_n f_n \quad (3.41)$$

Next, a set of weighting or testing functions w_m is chosen and an inner product (denoted by symbol $\langle \rangle$) taken in eq. (3.40). After substitution of eq. (3.41) into eq. (3.40), the result is



$$\sum_n \alpha_n \langle w_m, L f_n \rangle = \langle w_m, g \rangle \quad (3.42)$$

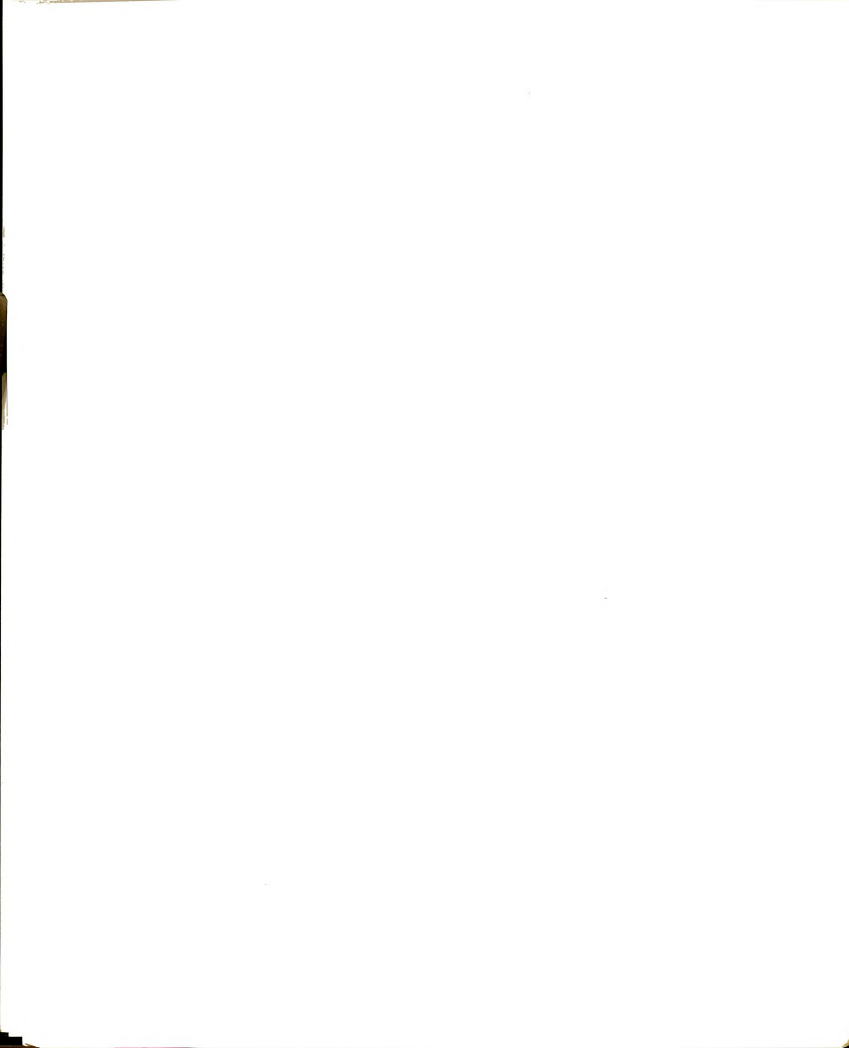
which may be written in matrix form as

$$[l_{mn}] [\alpha_n] = [g_m].$$

As illustrated in Figure 3.10, the idea when applied to the cubic box assumes that the box surface may be divided into square patches over each of which both components of surface current are assumed constant with unknown amplitudes (subdomain flat-top pulse-expansion functions). The testing functions chosen in this problem are δ functions at the patch centers (collocation or point matching). Thus, by successively enforcing the surface integral equations to be satisfied at these match points, a system of linear algebraic equations is generated. Matrix inversion then yields the solution for the unknown coefficients of the current samples.

3.7. The Numerical Technique

The validity of eqs. (3.29) and (3.30) has been checked in Section 3.5 through the consideration of a special case of a two-dimensional, infinite-plane interface between a finite conducting medium and free space. Since it is our goal to quantify the EM fields induced on the surfaces of arbitrarily shaped, finite conducting bodies, such as human bodies, it seems more realistic to consider a simulated human body with shape as shown in



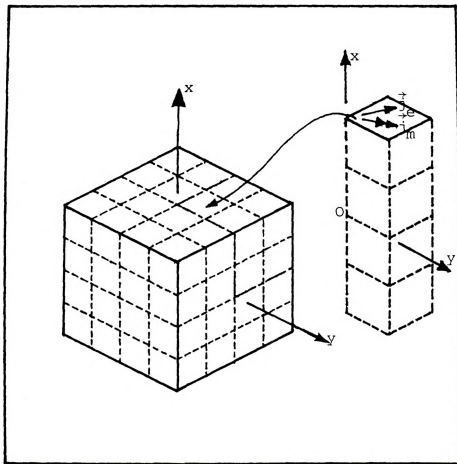


Figure 3.10. Division of the surface of a cubic body into square surface cells for moment-method solution using two-dimensional flat pulses as expansion functions and δ functions as testing functions.



Figure 3.11. We assume that the body region G , being characterized by permeability μ_0 and complex permittivity ϵ_d and enclosed by surface F , is illuminated by an incident plane wave which is propagating in the $+z$ direction with its electric field $\vec{E}^i$ polarized in the $+x$ direction. The body is immersed in free space as shown in Figure 3.11.

It is easy to see from Figure 3.11 that, for a simulated body, the body surface F consists of a number of plane surface elements. If the body is described in a rectangular coordinate system, which is the one used in Figure 3.11, the problem can be simplified by constructing the body surface with surface elements having normal vectors pointing in one of the six possible along-axis directions, i.e., the $\pm x$, the $\pm y$, and the $\pm z$ directions. For the purpose of clarity, we will name these surface elements according to the directions of their unit normal vectors. For example, the top surface element of the body head shown in Figure 3.11 will be considered as one of the $+x$ surface elements, since its unit normal vector is pointing in the $+x$ direction. It should be emphasized that a closed surface can consist of several surface elements with unit normal vectors pointing in the same direction, as can be seen from Figure 3.11.

Before going any further, we should mention that we choose a plane wave as the incident wave, because it is a realistic approximation when the body is distant from a



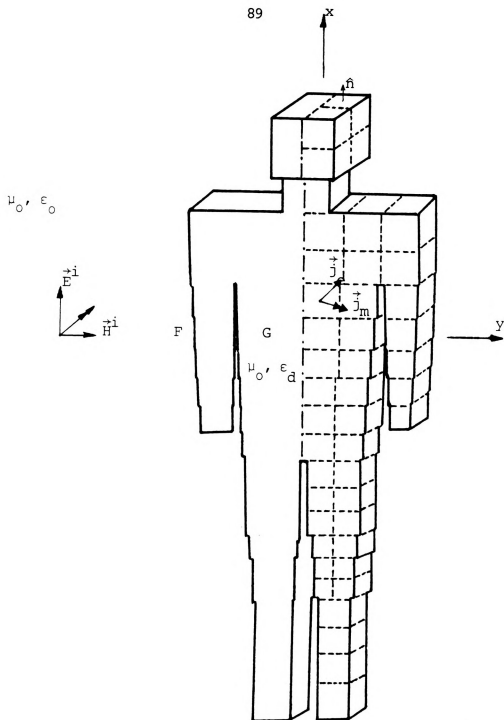
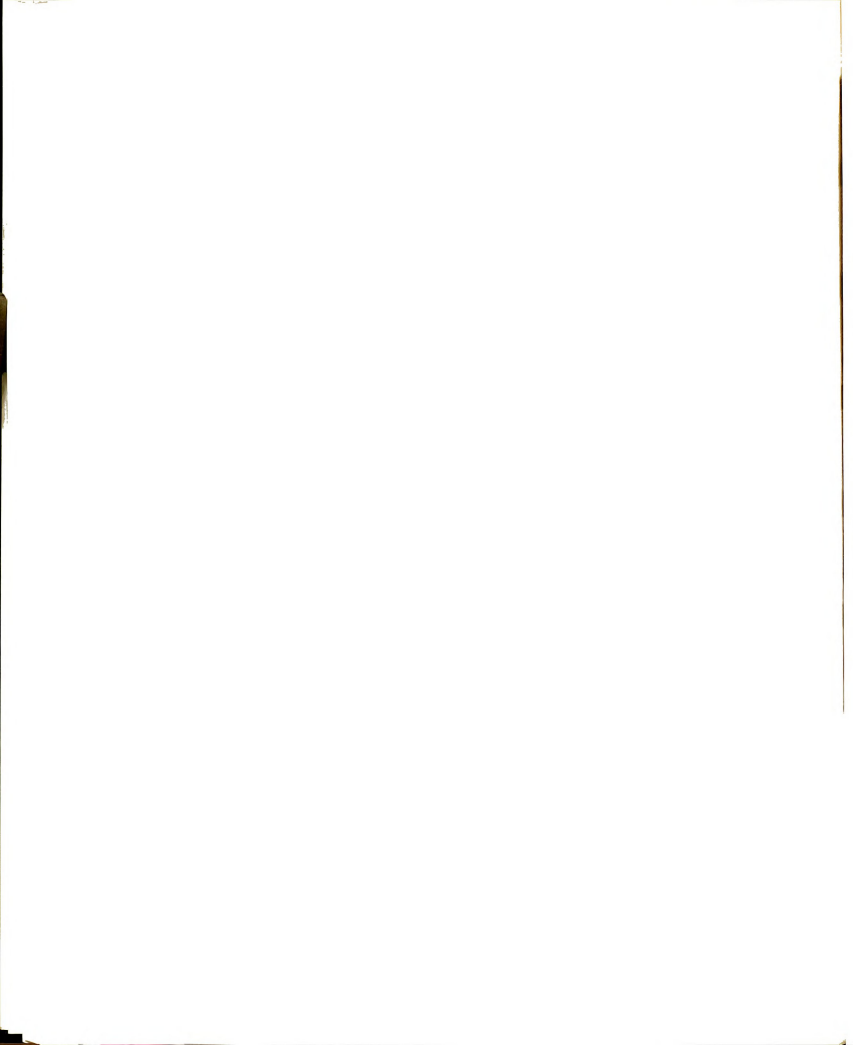


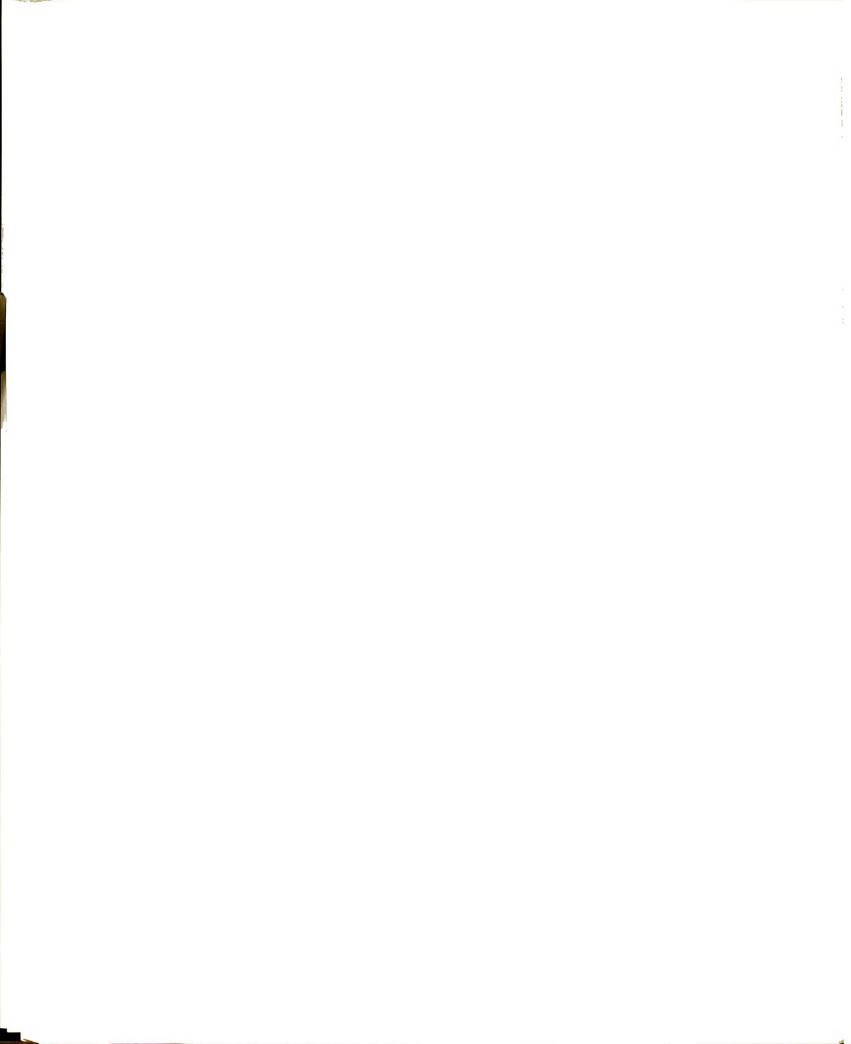
Figure 3.11. A simulated human body is irradiated by an incident plane wave propagating in the $+z$ direction with a x -polarized electric field. The body is characterized by complex permittivity ϵ_d , and its surface is subdivided into a number of square surface cells. The induced equivalent surface currents ($\vec{J}_e$, $\vec{J}_m$) at different locations on the surface can be obtained by solving two coupled, surface integral equations.



radiating source, and it facilitates the analysis. Also, we choose the rectangular coordinate system since it is the most adaptable to arbitrarily shaped bodies. However, it is to be noted that the surface integral equations derived previously are so general that they can be applied to any form of incident field, as well as any orthogonal coordinate system.

Now, let us consider Figure 3.11 again, where each surface element of F has been subdivided into several square subareas named surface cells. Each surface cell can have any arbitrary dimension which depends on the accuracy desired. Furthermore, the center point of each surface cell is considered as a reference point which could be either a field point or a source point or both. The EM field induced inside each surface cell is assumed to be uniform. We will use the notation F_x to represent the combined surface of all the $+x$ surface elements, and use the symbol N_x to denote the total number of surface cells included in F_x . Similar definitions will be applied to the notations, $F_y, N_y; F_z, N_z; F_{-x}, N_{-x}; F_{-y}, N_{-y};$ and F_{-z}, N_{-z} , etc. In other words, we may say that the body surface F has been divided into totally N_{cell} surface cells, where N_{cell} is the sum of $N_x, N_y, N_z, N_{-x}, N_{-y},$ and N_{-z} .

To solve the surface integral equations numerically by use of the moment-method, we must decompose the integral



equations first and then point-match them at the reference point of each surface cell. By doing this, a system of simultaneous, scalar, linear algebraic equations will be obtained. After transforming these linear equations into a matrix form, it can be solved by conventional matrix inversion techniques. To describe the procedures in more detail, let's start with the following definitions:

m: the numbering index of the surface cells in each subsurface ($F_x, F_y, F_z, F_{-x}, F_{-y}, F_{-z}$), when center points of the surface cells are considered as field points.

n: the numbering index of the surface cells in each subsurface, when center points of the surface cells are considered as source points.

A^n : the area of the nth surface cell.

$\vec{r}^m$: the position vector locating the mth field point, where $\vec{r}^m = (x^m, y^m, z^m)$.

$\vec{r}^n$: the position vector locating the nth source point, where $\vec{r}^n = (x^n, y^n, z^n)$.

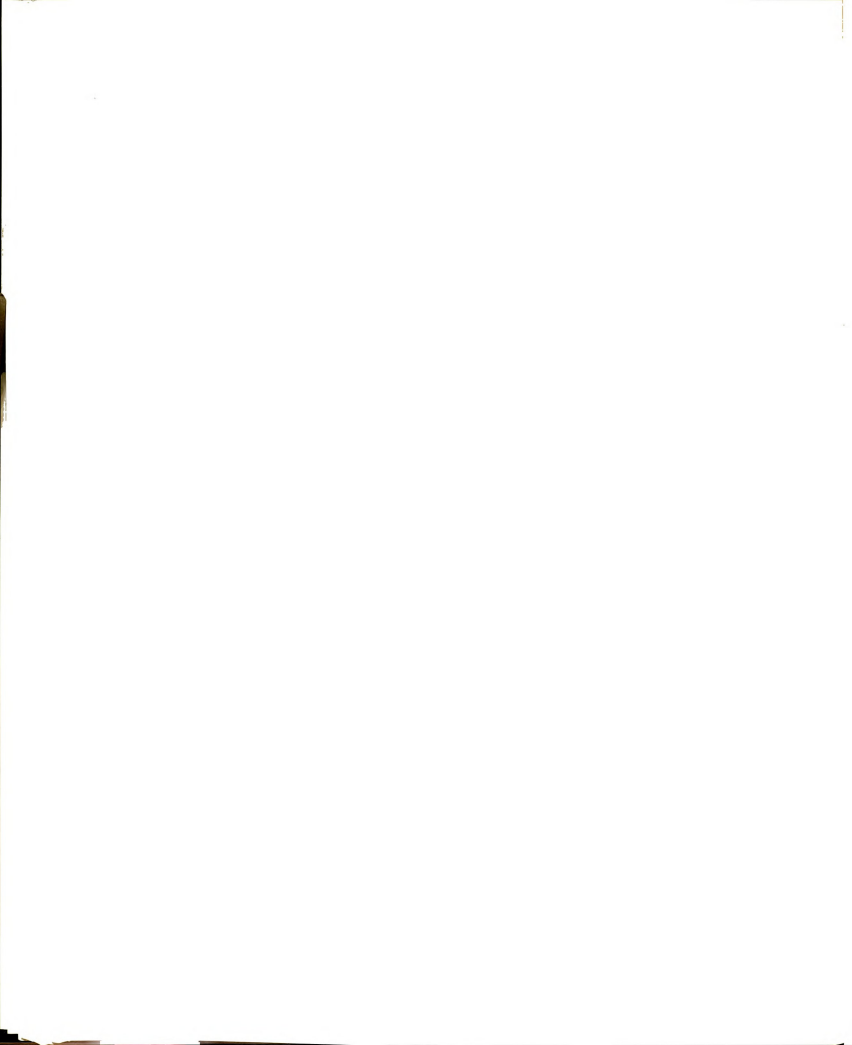
Also,

$$r^{nm} \equiv |\vec{r}^n - \vec{r}^m|, \quad x^{nm} \equiv x^n - x^m, \quad y^{nm} \equiv y^n - y^m,$$

$$z^{nm} \equiv z^n - z^m.$$

and

$a^m = \sqrt{A^m/\pi}$: the equivalent radius of the mth surface cell.



We also like to introduce the following associated Green's functions; the reasons behind these definitions will be clear when the surface integral equations are decomposed:

$$\begin{aligned}
 G_1^{nm} &= j \frac{e^{-jkr^{nm}} - \kappa^2 e^{-jk\kappa r^{nm}}}{r^{nm2}} + \frac{e^{-jkr^{nm}} - e^{-jk\kappa r^{nm}}}{r^{nm3}} \\
 G_2^{nm} &= \frac{e^{-jkr^{nm}} - \kappa^2 e^{-jk\kappa r^{nm}}}{r^{nm3}} - 3j \frac{e^{-jkr^{nm}} - \kappa^2 e^{-jk\kappa r^{nm}}}{r^{nm4}} - 3 \frac{e^{-jkr^{nm}} - e^{-jk\kappa r^{nm}}}{r^{nm5}} \\
 G_3^{nm} &= j \frac{e^{-jkr^{nm}} - \kappa^3 e^{-jk\kappa r^{nm}}}{r^{nm2}} + \frac{e^{-jkr^{nm}} - \kappa^2 e^{-jk\kappa r^{nm}}}{r^{nm3}} \\
 G_4^{nm} &= - \frac{e^{-jkr^{nm}} - \kappa^2 e^{-jk\kappa r^{nm}}}{r^{nm}}
 \end{aligned} \quad \left. \vphantom{\begin{aligned} G_1^{nm} \\ G_2^{nm} \\ G_3^{nm} \\ G_4^{nm} \end{aligned}} \right\} (3.43)$$

$$H_1^m = \pi (j e^{-ja^m} - j \kappa e^{-j\kappa a^m} + \frac{e^{-ja^m} - e^{-j\kappa a^m}}{a^m})$$

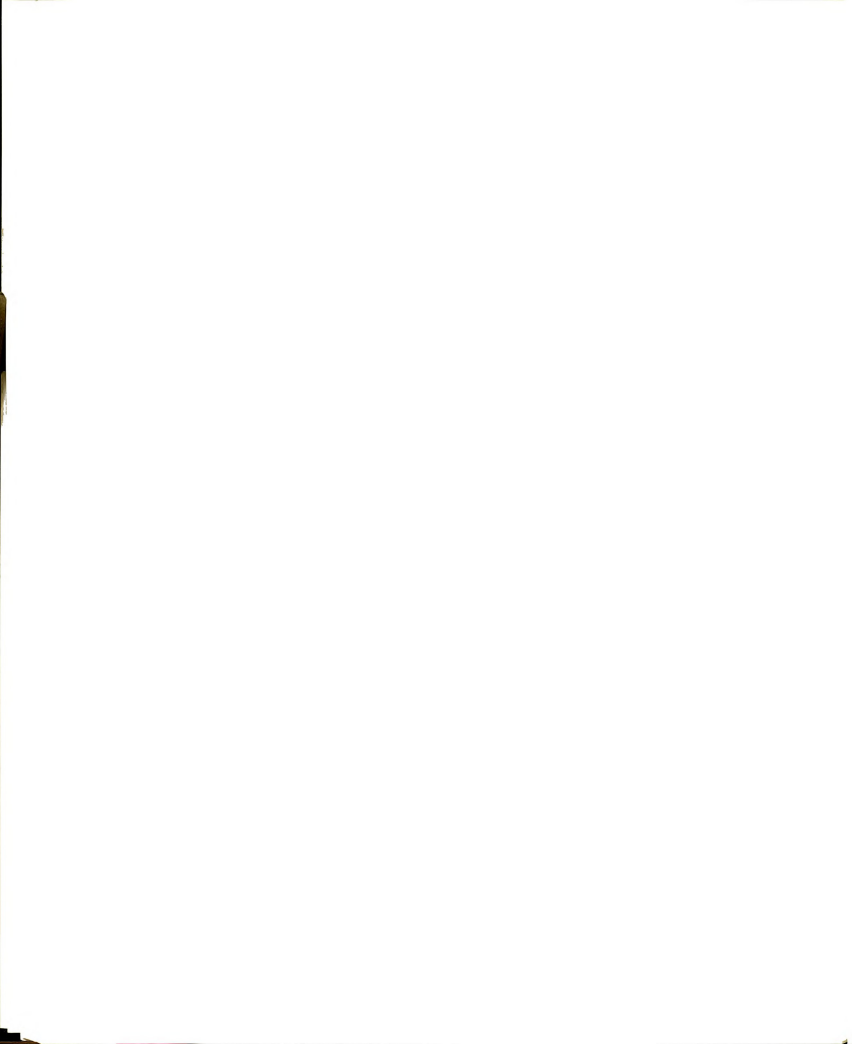
and

$$H_2^m = 2\pi (j \kappa e^{-j\kappa a^m} - j e^{-ja^m} - j \kappa + j).$$

Now, we are in a position to decompose the coupled, surface integral equations, i.e., eqs. (3.29) and (3.30). For convenience, we repeat them here,

$$\vec{J}_m = \frac{2}{\kappa^2 + 1} \vec{J}_m^i - \frac{1}{2\pi(\kappa^2 + 1)} \hat{n} \times \int_F [j(\kappa^2 \phi_d - \phi_f) \vec{J}_e + \vec{J}_m \times$$

$$\nabla'(\kappa^2 \phi_d - \phi_f) + j(\vec{J}_e \cdot \nabla') \nabla'(\phi_d - \phi_f)] dF' \quad (3.29')$$



$$\begin{aligned} \vec{j}_e = & \vec{j}_e^i + \frac{1}{4\pi} \hat{n} \times \int_F [j(\kappa^2 \phi_d - \phi_f) \vec{j}_m - \vec{j}_e \times \nabla'(\phi_d - \phi_f) \\ & + j(\vec{j}_m \cdot \nabla') \nabla'(\phi_d - \phi_f)] dF'. \end{aligned} \quad (3.30')$$

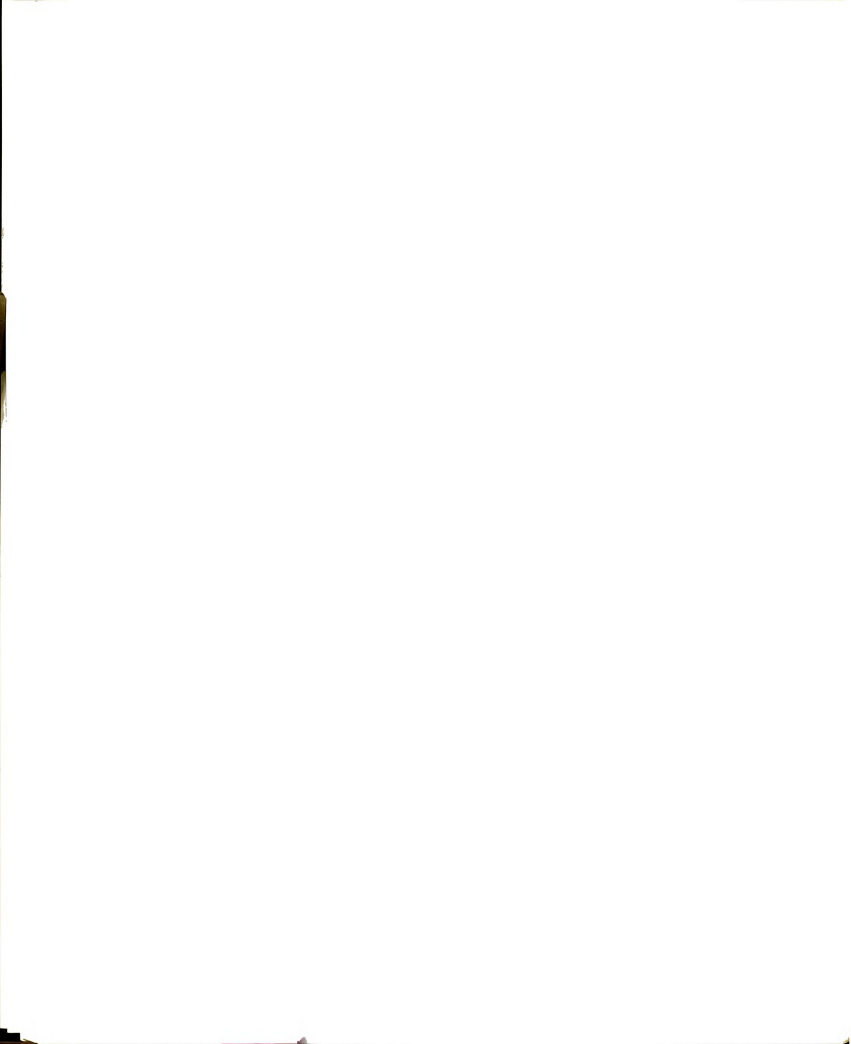
Note that, without any ambiguity, the superscript "d" for the induced equivalent surface current densities has been omitted in the above equations.

Let us consider eq. (3.29') first. To numerically implement the various surface integrals of this equation, we have to consider each possible case about the field point locations.

I) First of all, we consider the situation when the field point is located in one of the +x surface elements. Which means that the unit normal vector at field point, $\hat{n}$, is the same as the x-direction unit vector $\hat{x}$. For this case, we know that $1 \leq m \leq N_x$. Based on the previous definitions, it is easy to see that the last integral of eq. (3.29') can be written as

$$\begin{aligned} \hat{n} \times \int_F (\vec{j}_e \cdot \nabla') \nabla'(\phi_d - \phi_f) dF' = & \hat{x} \times \{ [\int_{F_x} + \int_{F_y} + \int_{F_z} + \\ & \int_{F_{-x}} + \int_{F_{-y}} + \int_{F_{-z}}] (\vec{j}_e \cdot \nabla') \nabla'(\phi_d - \phi_f) dF' \} \end{aligned} \quad (3.44)$$

where we have decomposed the total surface integral $\int_F$ into six different subintegrals, one for each subsurface. For example, $\int_{F_x}$ represents the integration of the integrand over all the surface elements having $\hat{x}$ as the unit normal



vector. The remaining five subintegrals in eq. (3.44) have similar meanings. Now, as an example, let's carry out the first subintegral of eq. (3.44).

$$\begin{aligned}
 \hat{x} \times \int_{F_x} (\hat{j}_e \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' &= \hat{x} \times \int_{F_x} [(j_{ey} \hat{y} + j_{ez} \hat{z}) \cdot (\hat{x} \frac{\partial}{\partial x'} + \\
 &\hat{y} \frac{\partial}{\partial y'} + \hat{z} \frac{\partial}{\partial z'})] [(\hat{x} \frac{\partial}{\partial x'} + \hat{y} \frac{\partial}{\partial y'} + \hat{z} \frac{\partial}{\partial z'}) (\phi_d - \phi_f)] dF' \\
 &= \hat{x} \times \int_{F_x} [j_{ey} \frac{\partial^2}{\partial x' \partial y'} + j_{ez} \frac{\partial^2}{\partial x' \partial z'} + \hat{y} (j_{ey} \frac{\partial^2}{\partial y'^2} + j_{ez} \frac{\partial^2}{\partial y' \partial z'}) + \\
 &\hat{z} (j_{ey} \frac{\partial^2}{\partial y' \partial z'} + j_{ez} \frac{\partial^2}{\partial z'^2})] (\phi_d - \phi_f) dF' \\
 &= \hat{x} \times \int_{F_x} [\hat{y} (j_{ey} \frac{\partial^2}{\partial y'^2} + j_{ez} \frac{\partial^2}{\partial y' \partial z'}) + \hat{z} (j_{ey} \frac{\partial^2}{\partial y' \partial z'} + j_{ez} \frac{\partial^2}{\partial z'^2})] \\
 &(\phi_d - \phi_f) dF'
 \end{aligned}$$

Since $\phi_d = \frac{e^{-jk\underline{r}}}{\underline{r}}$ and $\phi_f = \frac{e^{-jk\underline{r}}}{\underline{r}}$ where $\underline{r} = |\underline{r} - \underline{r}'|$, we

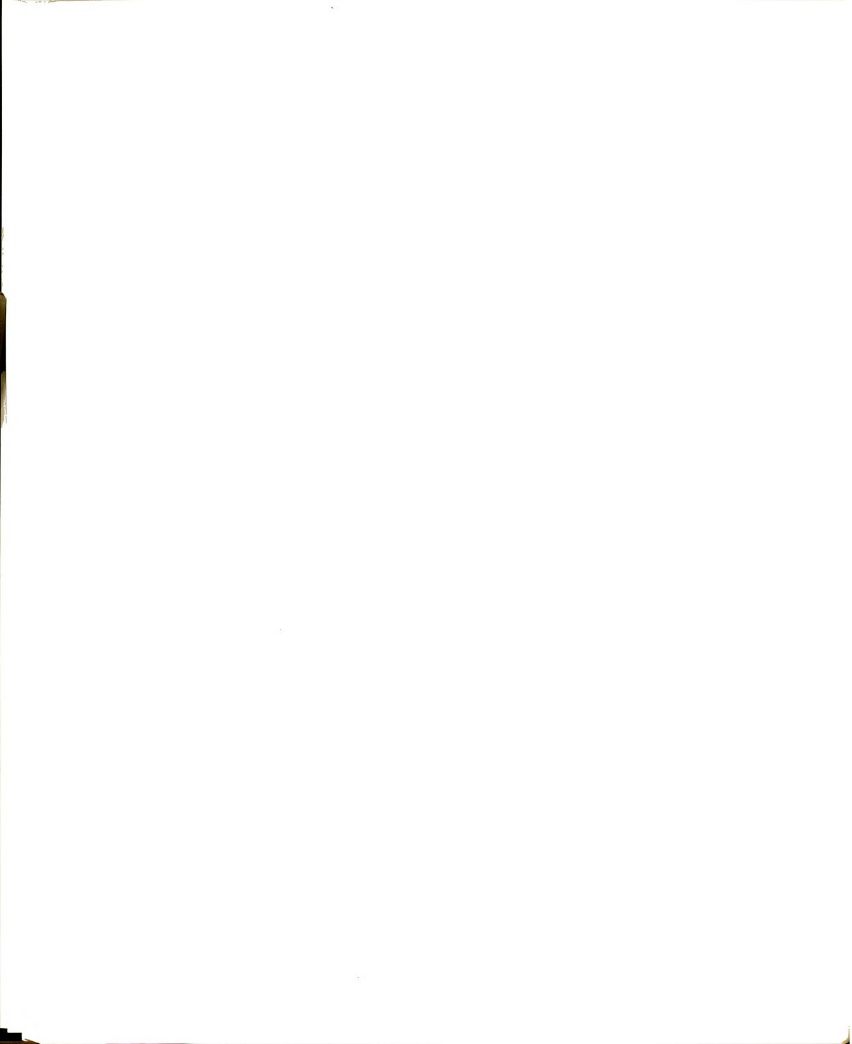
may show that

$$\frac{\partial^2}{\partial y'^2} (\phi_d - \phi_f) = G_1 + (y' - y)^2 G_2$$

$$\frac{\partial^2}{\partial z'^2} (\phi_d - \phi_f) = G_1 + (z' - z)^2 G_2$$

and

$$\frac{\partial^2}{\partial y' \partial z'} (\phi_d - \phi_f) = (y' - y)(z' - z) G_2$$

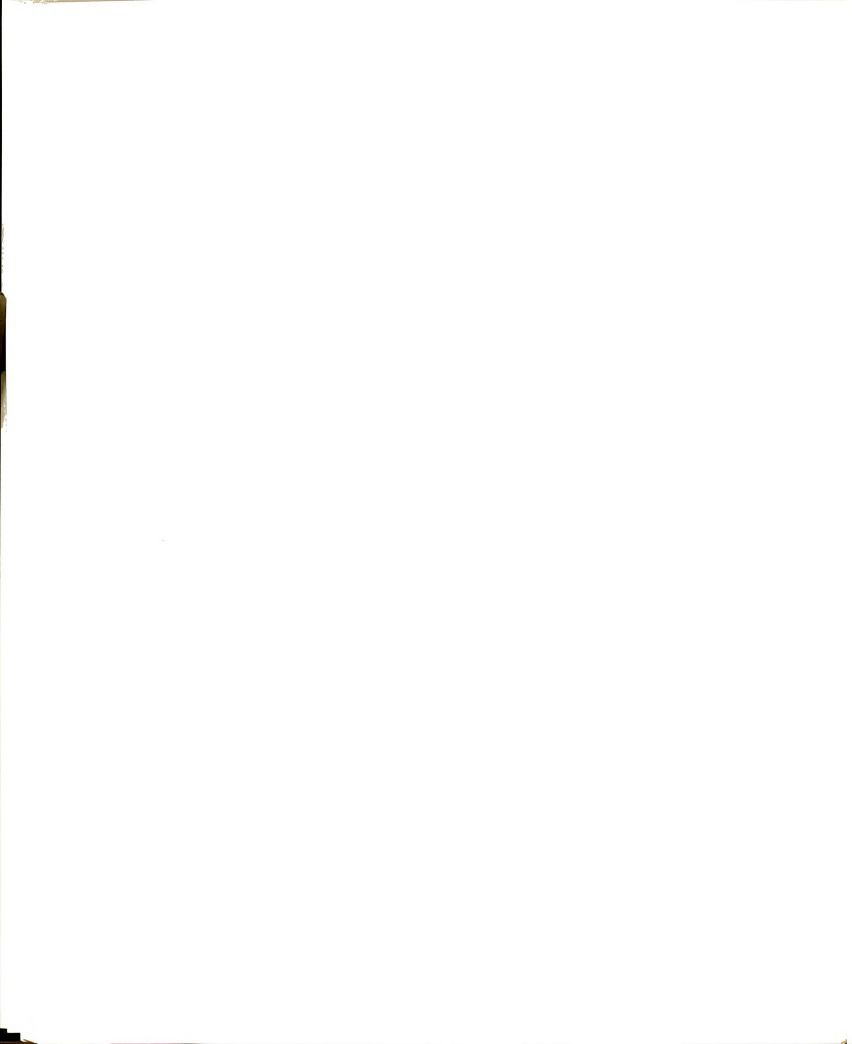


where G_1 and G_2 have the same expressions as that of G_1^{nm} and G_2^{nm} defined in eq. (3.43) except r^{nm} being replaced by $\underline{r}$. Hence we have

$$\begin{aligned} \hat{x} \times \int_{F_x} (\vec{j}_e \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' = & -\hat{y} \int_{F_x} \{ j_{ey} (y' - y) (z' - z) G_2 + \\ & j_{ez} [G_1 + (z' - z)^2 G_2] \} dF' + \hat{z} \int_{F_x} \{ j_{ey} [G_1 + (y' - y)^2 G_2] + \\ & j_{ez} (y' - y) (z' - z) G_2 \} dF'. \end{aligned}$$

The integrals on the right-hand side of the above equation can be numerically implemented by summing up the products of integrands and areas of the corresponding surface cells with a special consideration being paid to the singular terms which occur when the source point coincides with the field point. This process is completely based on the moment-method technique with two-dimensional pulse functions as the basis functions, and δ functions as the test-functions. So we may rewrite the above equation as

$$\begin{aligned} \hat{x} \times \int_{F_x} (\vec{j}_e \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' = & -\hat{y} \sum_{n=1}^{N_x} [j_{ey}^n A^n y^{nm} z^{nm} G_2^{nm} + \\ & j_{ez}^n A^n (G_1^{nm} + z^{nm} z^{nm} G_2^{nm})] + \hat{z} \sum_{n=1}^{N_x} [j_{ey}^n A^n (G_1^{nm} + y^{nm} y^{nm} G_2^{nm}) + \\ & j_{ez}^n A^n (y' - y) (z' - z) G_2^{nm}] \end{aligned}$$



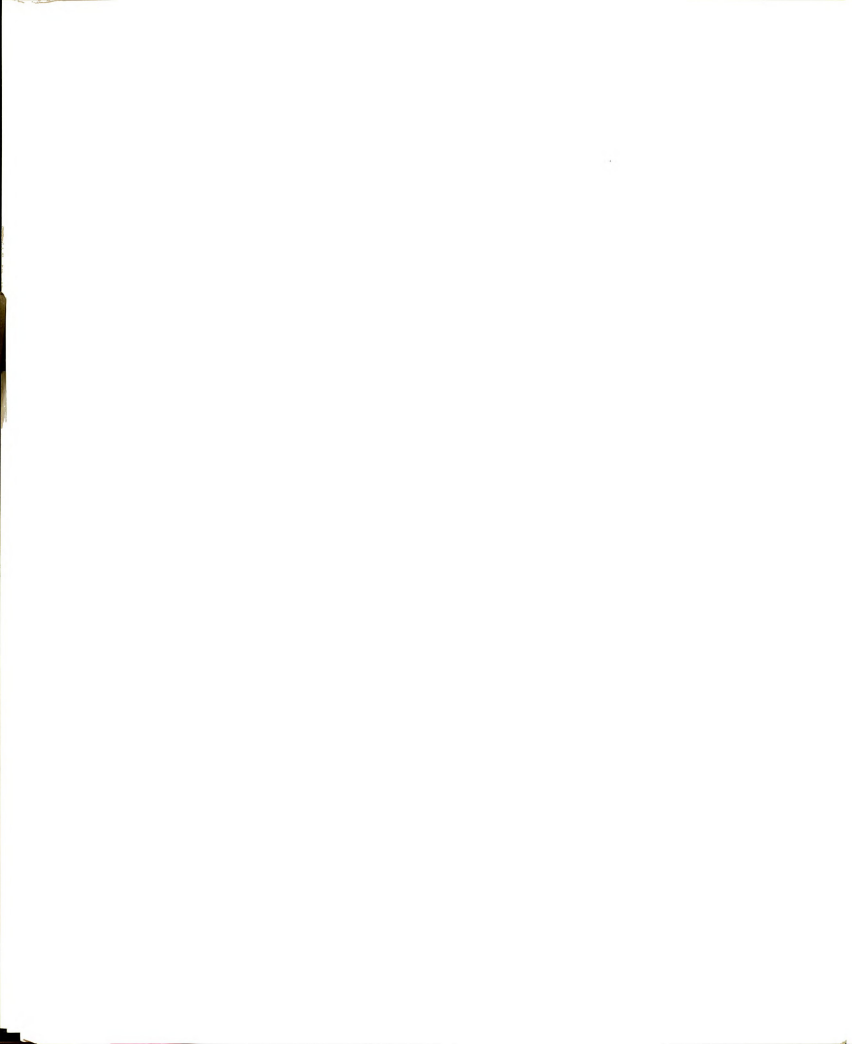
$$\begin{aligned}
& j_{ez}^n A^n y^{nm} z^{nm} G_2^{nm} - \hat{y} \int_{A^m} \{ j_{ey}^m (y' - y) (z' - z) G_2 + j_{ez}^m [G_1 + \\
& (z' - z)^2 G_2] \} dF' + \hat{z} \int_{A^m} \{ j_{ey}^m [G_1 + (y' - y)^2 G_2] + \\
& j_{ez}^m (y' - y) (z' - z) G_2 \} dF', \quad 1 \leq m \leq N_x
\end{aligned} \tag{3.45}$$

where j_{ey}^n , j_{ez}^n , j_{ey}^m , and j_{ez}^m are the unknown expansion coefficients of the equivalent surface current densities. Note that the above equation is satisfied at every field point location on the subsurface F_x . The notation $\sum_{n=1}^{N_x}$ implies a summation with respect to all the source points on F_x excluding the one coinciding with the field point where the singularities of Green's functions occur. The contributions of the singular terms are accounted for by the last two terms in eq. (3.45), where the notation $\int_{A^m}$ means the integration of integrand over a circle centered at the field point with area A^m . Now, let's investigate these singular terms more specifically. It can be shown that

$$\begin{aligned}
\int_{A^m} (y' - y) (z' - z) G_2 dF' &= \int_0^{A^m} \int_0^{2\pi} \underline{r}^2 \cos \theta \sin \theta G_2 d\theta \underline{r} d\underline{r} \\
&= \underbrace{\int_0^{2\pi} \cos \theta \sin \theta d\theta}_0 \int_0^{A^m} \underline{r}^3 G_2 d\underline{r} = 0,
\end{aligned} \tag{3.46}$$

also,

$$\int_{A^m} [G_1 + (y' - y)^2 G_2] dF' = \int_0^{2\pi} \int_0^{A^m} [G_1 + \underline{r}^2 \cos^2 \theta G_2] \underline{r} d\underline{r} d\theta$$



and

$$\int_A^m [G_1 + (z' - z)^2 G_2] dF' = \int_0^{2\pi} \int_0^{\alpha_m} [G_1 + \underline{r}^2 \sin^2 \theta G_2] \underline{r} d\underline{r} d\theta.$$

After substituting the expressions of G_1 and G_2 into the above equations, we have

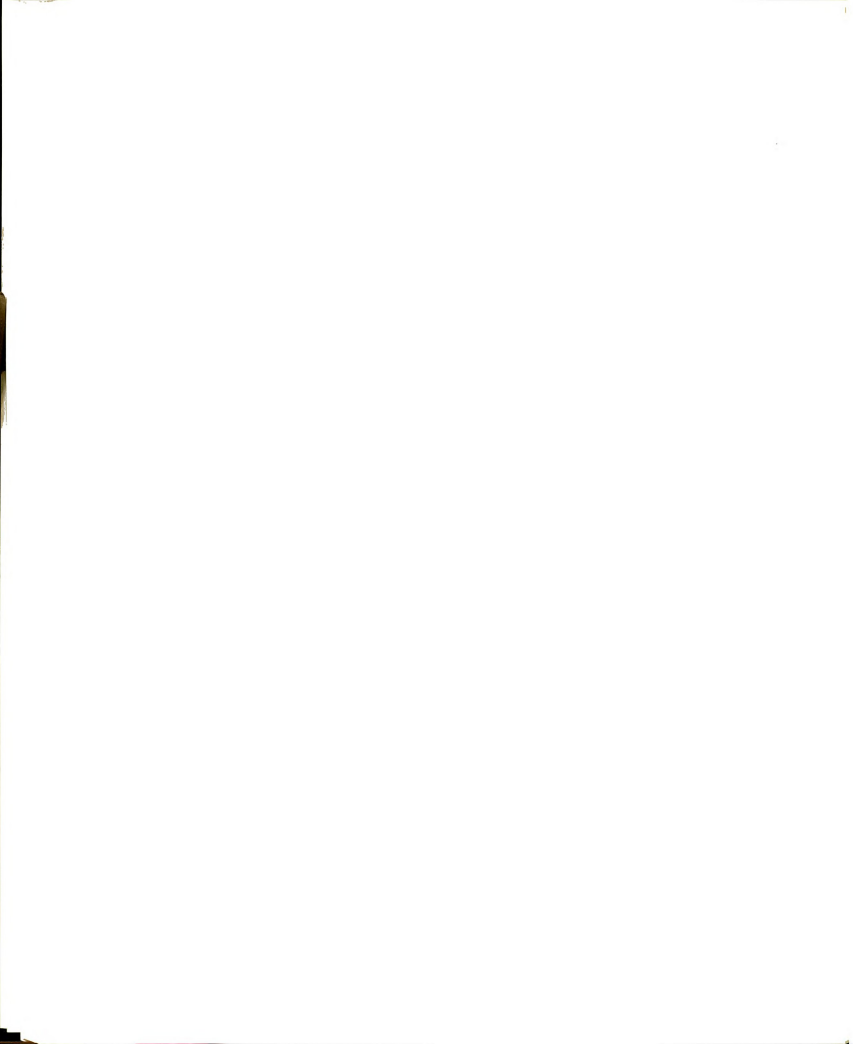
$$\int_A^m [G_1 + (y' - y)^2 G_2] dF' = \int_A^m [G_1 + (z' - z)^2 G_2] dF' = H_1^m, \quad (3.47)$$

where H_1^m is defined in eq. (3.43). Then, substituting eqs. (3.46) and (3.47) into eq. (3.45), finally we obtain

$$\begin{aligned} \hat{x} \times \int_{F_x} (\vec{j}_e \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' = & -\hat{y} \left\{ \sum_{\substack{n=1 \\ n \neq m}}^{N_x} A^n [j_{ey}^n y^{nm} z^{nm} G_2^{nm} + \right. \\ & \left. j_{ez}^n (G_1^{nm} + z^{nm} \cdot z^{nm} G_2^{nm})] + j_{ez}^m H_1^m \right\} + \\ & \hat{z} \left\{ \sum_{\substack{n=1 \\ n \neq m}}^{N_x} A^n [j_{ey}^n (G_1^{nm} + y^{nm} \cdot y^{nm} G_2^{nm}) + \right. \\ & \left. j_{ez}^n y^{nm} z^{nm} G_2^{nm}] + j_{ey}^m H_1^m \right\}, \text{ where } 1 \leq m \leq N_x. \end{aligned} \quad (3.48)$$

This concludes the evaluation of the first subintegral in eq. (3.44). The remaining five subintegrals in that equation can be carried out similarly. Combining these results with eq. (3.44), the latter will become

$$\hat{x} \times \int_F (\vec{j}_e \cdot \nabla') \nabla' (\phi_d - \phi_f) dF' = -\hat{y} \left\{ \left(\sum_{\substack{n=1 \\ n \neq m}}^{N_x} + \sum_{n=1}^{N-x} \right) A^n [j_{ey}^n y^{nm} z^{nm} G_2^{nm} + \right.$$

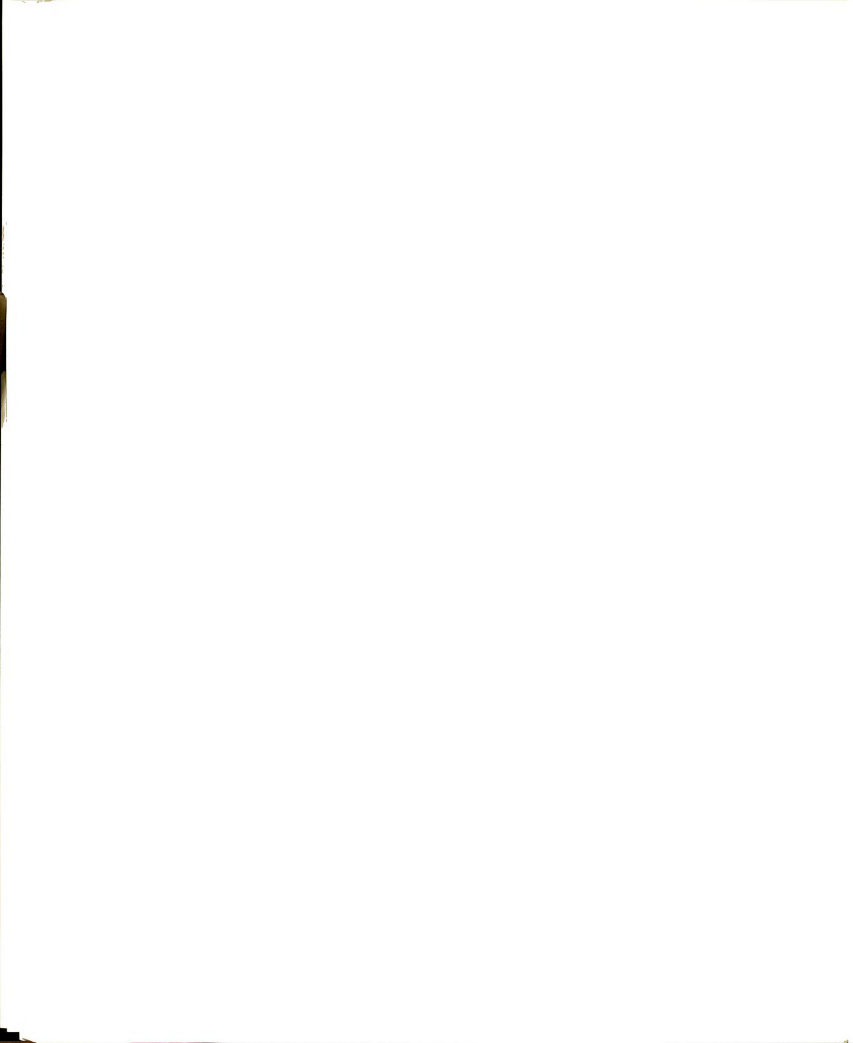


$$\begin{aligned}
& j_{ez}^n (G_1^{nm} + z^{nm} \cdot z^{nm} G_2^{nm}) + \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N-y} \right) A^n [j_{ex}^n x^{nm} z^{nm} G_2^{nm} + \\
& j_{ez}^n (G_1^{nm} + z^{nm} \cdot z^{nm} G_2^{nm})] + \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N-z} \right) A^n [j_{ex}^n x^{nm} z^{nm} G_2^{nm} + \\
& j_{ey}^n y^{nm} z^{nm} G_2^{nm}] + j_{ez}^m H_1^m \} + \\
& \hat{z} \cdot \left(\sum_{\substack{n=1 \\ n \neq m}}^{N_x} + \sum_{n=1}^{N-x} \right) A^n [j_{ey}^n (G_1^{nm} + y^{nm} \cdot y^{nm} G_2^{nm}) + \\
& j_{ez}^n y^{nm} z^{nm} G_2^{nm}] + \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N-y} \right) A^n [j_{ex}^n x^{nm} y^{nm} G_2^{nm} + \\
& j_{ez}^n y^{nm} z^{nm} G_2^{nm}] + \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N-z} \right) A^n [j_{ex}^n x^{nm} y^{nm} G_2^{nm} + \\
& j_{ey}^n (G_1^{nm} + y^{nm} \cdot y^{nm} G_2^{nm})] + j_{ey}^m H_1^m \}, \text{ where } 1 \leq m \leq N_x.
\end{aligned}
\tag{3.49}$$

This concludes the investigation of the last integral on the right-hand side of eq. (3.29').

If we repeat the whole process for the second surface integration of eq. (3.29'), we will find that it can be decomposed into the following form:

$$\begin{aligned}
\hat{x} \times \int_F \hat{j}_m \times \nabla' (\kappa^2 \phi_d - \phi_f) dF' = \hat{y} \cdot \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N-x} \right) A^n x^{nm} G_3^{nm} j_{my}^n - \\
\left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N-y} \right) A^n y^{nm} G_3^{nm} j_{mx}^n - \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N-z} \right) A^n [j_{mx}^n y^{nm} G_3^{nm} -
\end{aligned}$$



$$\begin{aligned}
& j_{ny}^n x^{nm} G_3^{nm} \} + 2 \{ \left(\frac{N_x}{\Sigma} + \frac{N_x}{\Sigma} \right) A^n x^{nm} G_3^{nm} j_{nz}^n + \\
& \left(\frac{N_y}{\Sigma} + \frac{N_y}{\Sigma} \right) A^n j_{nz}^n x^{nm} G_3^{nm} - j_{mx}^n z^{nm} G_3^{nm} \} - \\
& \left(\frac{N_z}{\Sigma} + \frac{N_z}{\Sigma} \right) A^n z^{nm} G_3^{nm} j_{mx}^n \}, \text{ where } 1 \leq m \leq N_x, \quad (3.50)
\end{aligned}$$

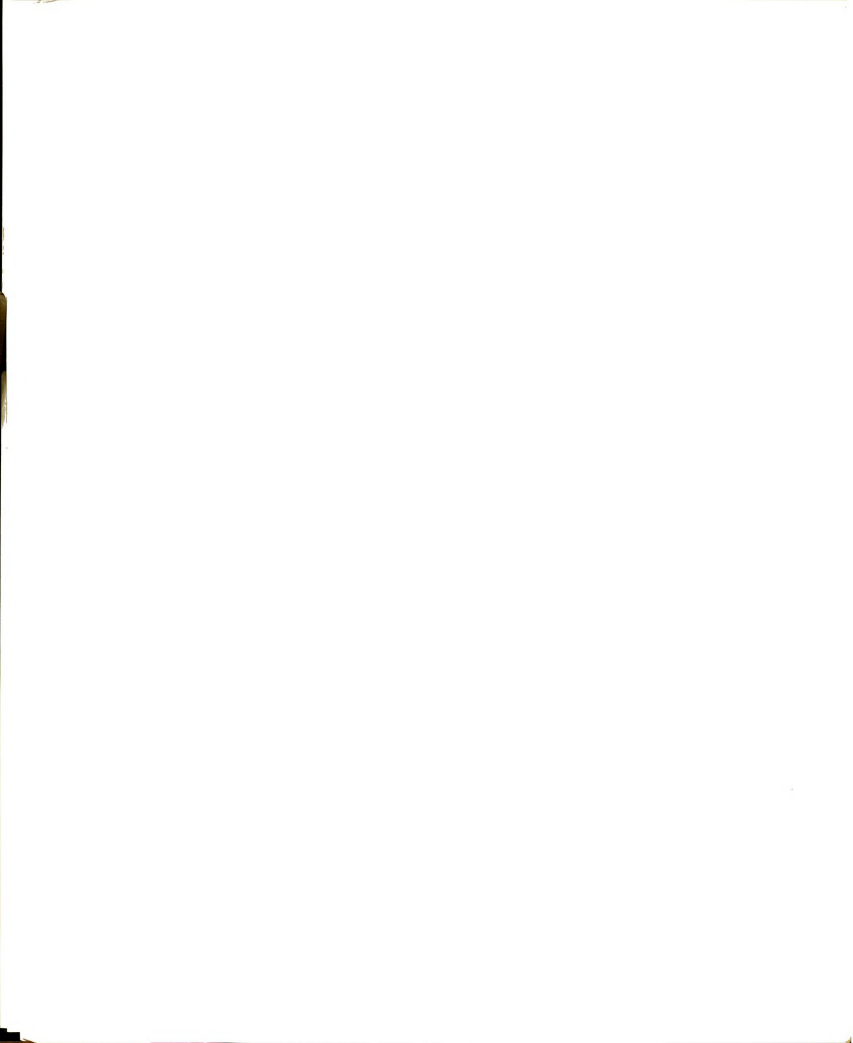
and G_3^{nm} is defined in eq. (3.43). It is noted that the singular terms have no contribution to eq. (3.50).

Similarly, the first integral term on the right-hand side of eq. (3.29') can be transformed to become

$$\begin{aligned}
& \hat{x} \times \hat{r}_F \cdot \hat{j}_E (\kappa^2 \phi_d - \phi_f) dF' = -\hat{y} \{ \left(\frac{N_x}{\Sigma} + \frac{N_x}{\Sigma} \right) A^n G_4^{nm} j_{ez}^n + \\
& \left(\frac{N_y}{\Sigma} + \frac{N_y}{\Sigma} \right) A^n G_4^{nm} j_{ez}^n + j_{ez}^m H_2^m \} + 2 \{ \left(\frac{N_x}{\Sigma} + \frac{N_x}{\Sigma} \right) A^n G_4^{nm} j_{ey}^n + \\
& \left(\frac{N_z}{\Sigma} + \frac{N_z}{\Sigma} \right) A^n G_4^{nm} j_{ey}^n + j_{ey}^m H_2^m \}, \text{ where } 1 \leq m \leq N_x, \quad (3.51)
\end{aligned}$$

and G_4^{nm} , H_2^m are defined in eq. (3.43).

If we enforce satisfaction of eq. (3.29') at each reference point on the subsurface F_x , with the help of eqs. (3.49), (3.50), and (3.51), we can obtain totally N_x vector equations, i.e., one equation for each matching point. Furthermore, it is possible to decompose these N_x vector equations into $2N_x$ scalar, linear algebraic equations since



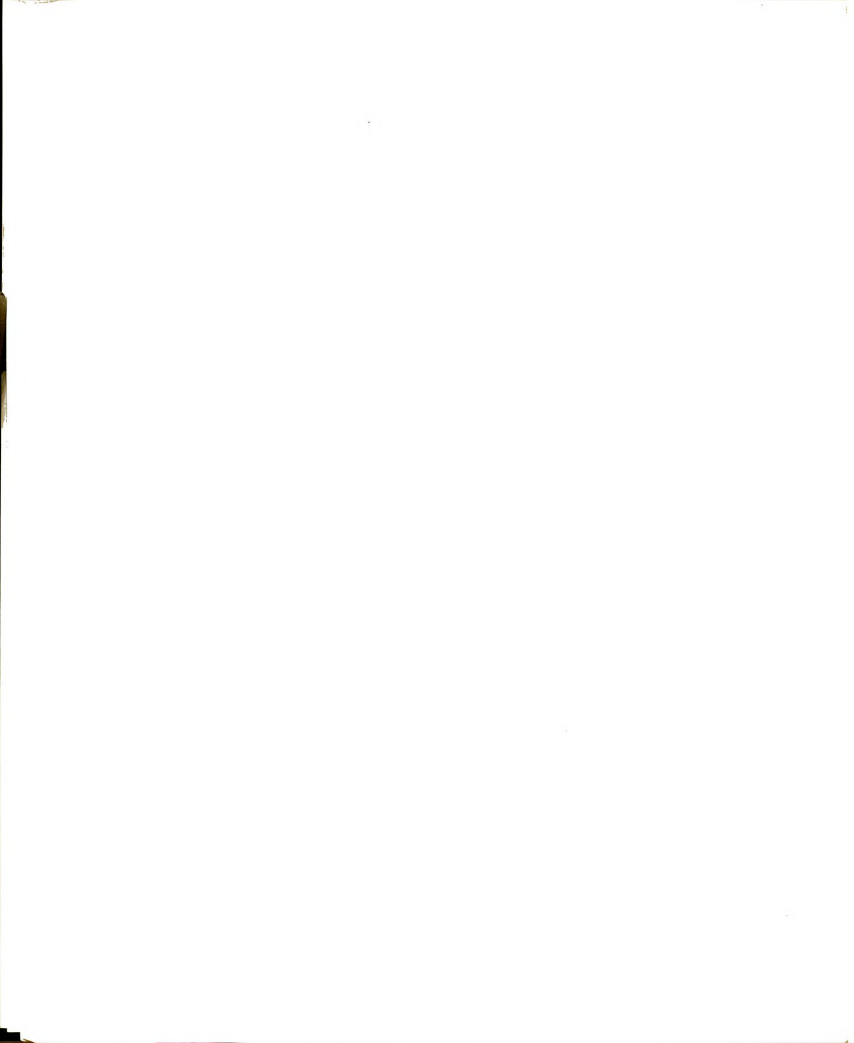
each surface current possesses two tangential components. Based on eqs. (3.29'), (3.49), (3.50), and (3.51), it can be shown that the following are those $2N_x$ linear equations:

$$\begin{aligned} & \left\{ \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N_y} \right) \cdot [-j_{ey}^n \psi_{1x}(n,m) - j_{ez}^n \psi_{2z}(n,m) + j_{my}^n \psi_{3x}(n,m)] - \right. \\ & \left. j_{ez}^m H(m) + j_{my}^m C \right\} - \left\{ \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N_z} \right) \cdot [j_{ex}^n \psi_{1y}(n,m) + j_{ez}^n \psi_{2z}(n,m) + \right. \\ & \left. j_{mx}^n \psi_{3y}(n,m)] \right\} - \left\{ \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N_x} \right) \cdot [j_{ex}^n \psi_{1y}(n,m) + j_{ey}^n \psi_{1x}(n,m) + \right. \\ & \left. j_{mx}^n \psi_{3y}(n,m) - j_{my}^n \psi_{3x}(n,m)] \right\} = 4\pi j_{my}^m \quad (3.52) \end{aligned}$$

also,

$$\begin{aligned} & \left\{ \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N_z} \right) \cdot [j_{ey}^n \psi_{2y}(n,m) + j_{ez}^n \psi_{1x}(n,m) + j_{mz}^n \psi_{3x}(n,m)] + \right. \\ & \left. j_{ey}^m H(m) + j_{mz}^m C \right\} + \left\{ \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N_z} \right) \cdot [j_{ex}^n \psi_{1z}(n,m) + j_{ez}^n \psi_{1x}(n,m) - \right. \\ & \left. j_{mx}^n \psi_{3z}(n,m) + j_{mz}^n \psi_{3x}(n,m)] \right\} + \left\{ \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N_x} \right) \cdot [j_{ex}^n \psi_{1z}(n,m) + \right. \\ & \left. j_{ey}^n \psi_{2y}(n,m) - j_{mx}^n \psi_{3z}(n,m)] \right\} = 4\pi j_{mz}^m \quad (3.53) \end{aligned}$$

where $1 \leq m \leq N_x$, and the various ψ functions are defined as



$$\psi_{1x}(n,m) = j A^n y^{nm} z^{nm} G_2^{nm}$$

$$\psi_{1y}(n,m) = j A^n x^{nm} z^{nm} G_2^{nm}$$

$$\psi_{1z}(n,m) = j A^n x^{nm} y^{nm} G_2^{nm}$$

$$\psi_{2x}(n,m) = j A^n (G_1^{nm} + x^{nm} \cdot x^{nm} G_2^{nm} + G_4^{nm})$$

$$\psi_{2y}(n,m) = j A^n (G_1^{nm} + y^{nm} \cdot y^{nm} G_2^{nm} + G_4^{nm})$$

$$\psi_{2z}(n,m) = j A^n (G_1^{nm} + z^{nm} \cdot z^{nm} G_2^{nm} + G_4^{nm})$$

$$\psi_{3x}(n,m) = A^n x^{nm} G_3^{nm}$$

$$\psi_{3y}(n,m) = A^n y^{nm} G_3^{nm}$$

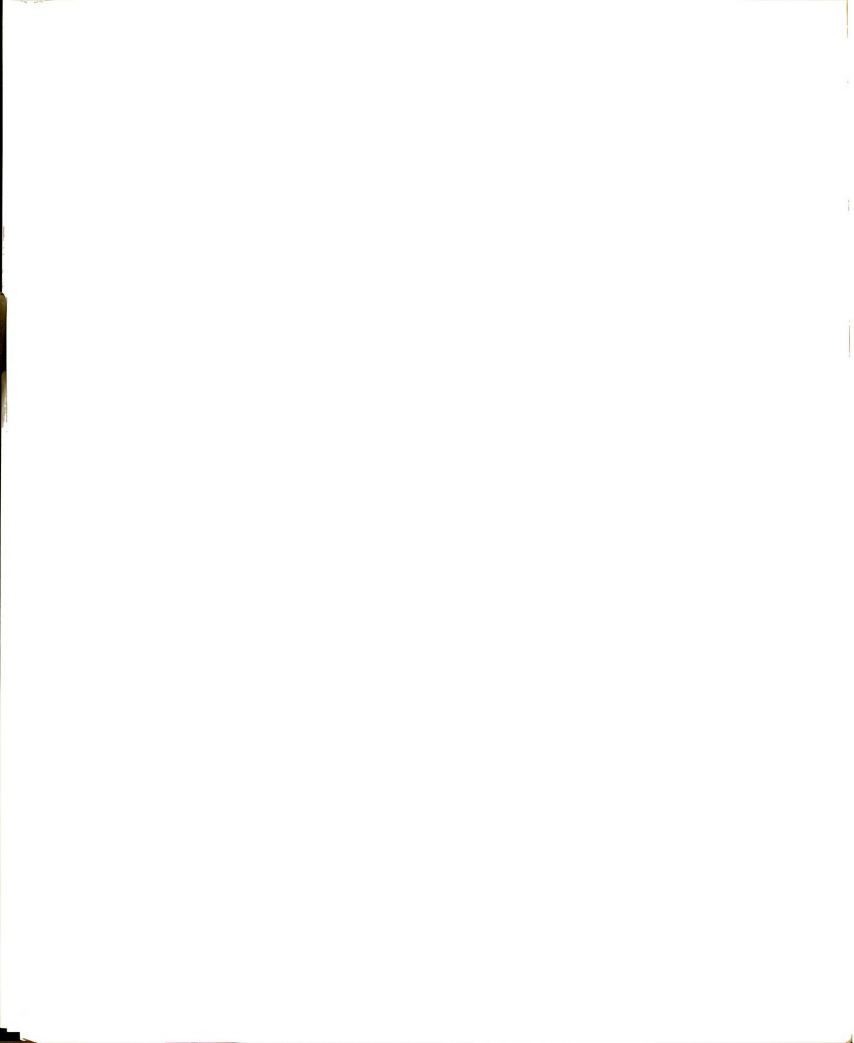
$$\psi_{3z}(n,m) = A^n z^{nm} G_3^{nm}$$

$$H(m) = j (H_1^m + H_2^m)$$

$$C = 2\pi(\kappa^2 + 1).$$

This completes the whole analysis for the case of $\hat{n} = \hat{x}$.

II) If the field point is located in one of the $+y$ -surface elements, i.e., when $\hat{n} = \hat{y}$, other $2N_y$ linear equations can be obtained by following a similar procedure as that described in case I). We will present these equations without repeating the details of the derivations:



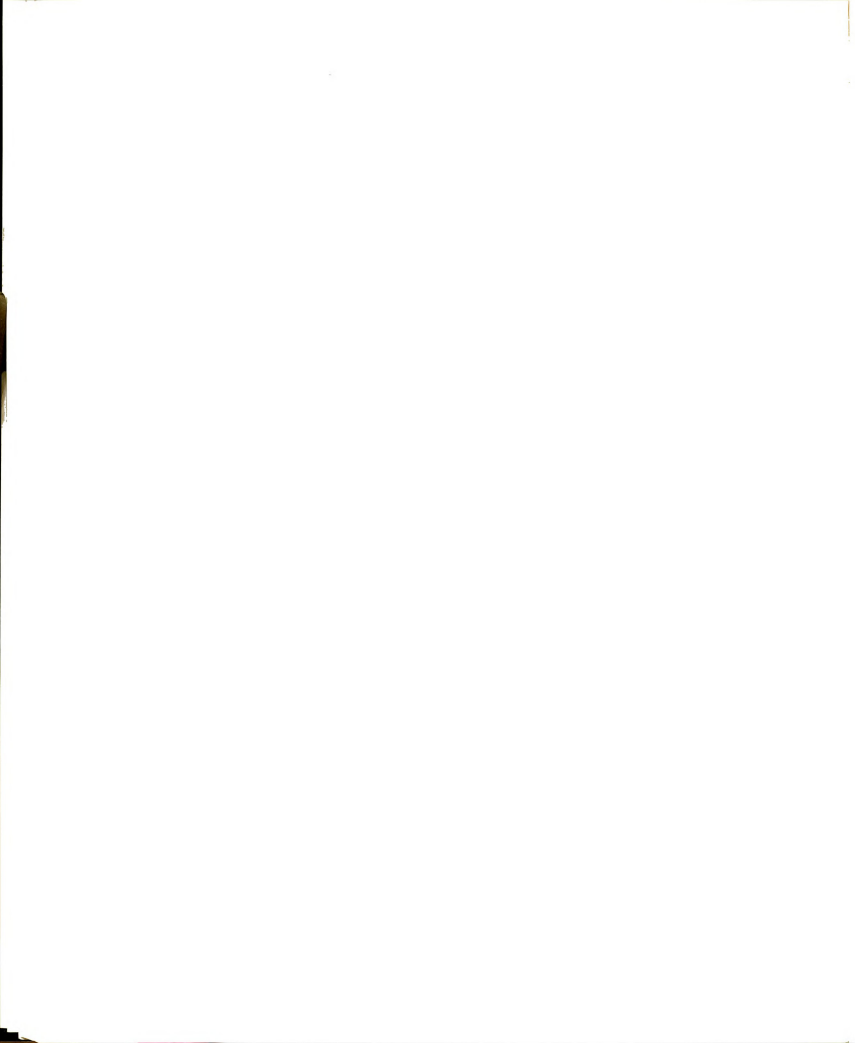
$$\begin{aligned}
& \left\{ \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}} \right) \cdot [j_{ey}^n \psi_{1x}(n, m) + j_{ez}^n \psi_{2z}(n, m) - j_{my}^n \psi_{3x}(n, m)] \right\} + \\
& \left\{ \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N_{-y}} \right) \cdot [j_{ex}^n \psi_{1y}(n, m) + j_{ez}^n \psi_{2z}(n, m) + j_{mx}^n \psi_{3y}(n, m)] \right\} + \\
& j_{ez}^m H(m) + j_{mx}^m C \} + \left\{ \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}} \right) \cdot [j_{ex}^n \psi_{1y}(n, m) + j_{ey}^n \psi_{1x}(n, m) + \right. \\
& \left. j_{mx}^n \psi_{3y}(n, m) - j_{my}^n \psi_{3x}(n, m)] \right\} = 4\pi j_{mx}^m i \quad (3.54)
\end{aligned}$$

also,

$$\begin{aligned}
& \left\{ \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}} \right) \cdot [-j_{ey}^n \psi_{1z}(n, m) - j_{ez}^n \psi_{1y}(n, m) - j_{my}^n \psi_{3z}(n, m) + \right. \\
& \left. j_{mz}^n \psi_{3y}(n, m)] \right\} - \left\{ \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N_{-y}} \right) \cdot [j_{ex}^n \psi_{2x}(n, m) + j_{ez}^n \psi_{1y}(n, m) - \right. \\
& \left. j_{mz}^n \psi_{3y}(n, m) - j_{ex}^m H(m) + j_{mz}^m C] - \left\{ \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}} \right) \cdot [j_{ex}^n \psi_{2x}(n, m) + \right. \right. \\
& \left. \left. j_{ey}^n \psi_{1z}(n, m) + j_{my}^n \psi_{3z}(n, m)] \right\} = 4\pi j_{mz}^m i \quad (3.55)
\end{aligned}$$

where $1 \leq m \leq N_y$.

III) Next, let's consider the case where the field point is located in one of the $+z$ -surface elements. As in the previous two cases, by use of the point-matching technique, we may obtain the following $2N_z$ scalar, linear equations:



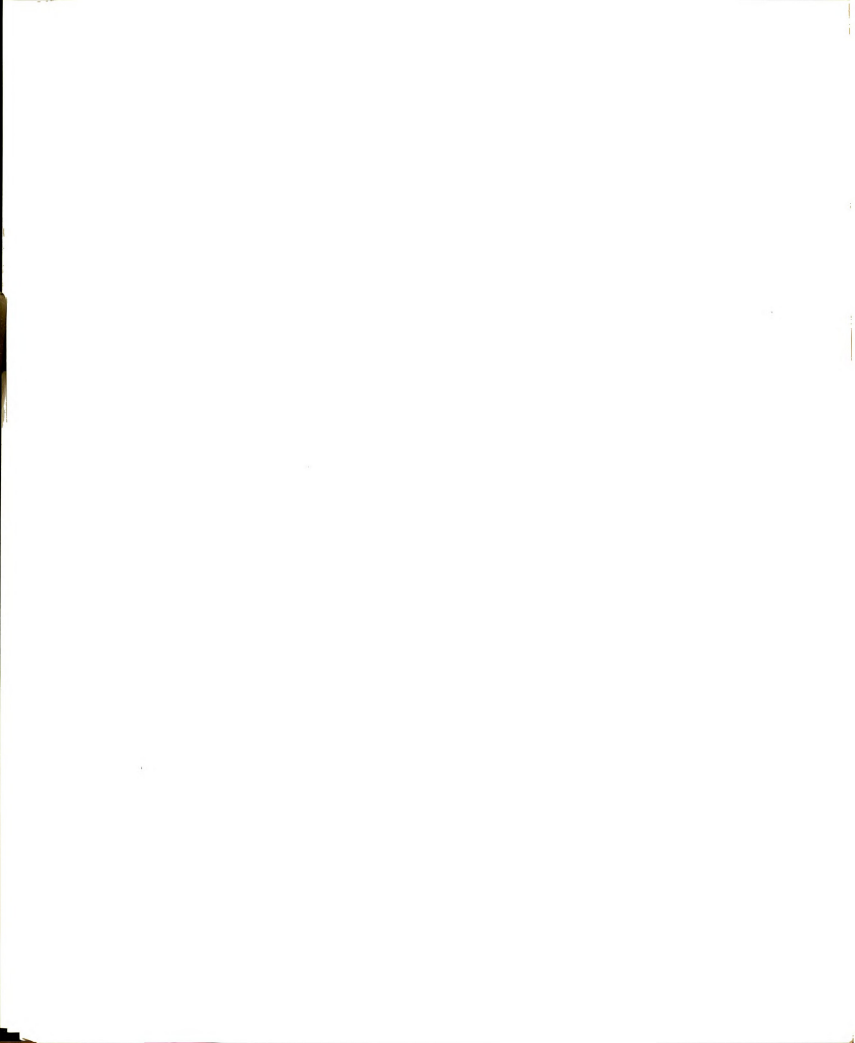
$$\begin{aligned}
& \{ (\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}}) \cdot [-j_{ey}^n \psi_{2y}(n,m) - j_{ez}^n \psi_{1x}(n,m) - j_{mz}^n \psi_{3x}(n,m)] - \\
& \{ (\sum_{n=1}^{N_y} + \sum_{n=1}^{N_{-y}}) \cdot [j_{ex}^n \psi_{1z}(n,m) + j_{ez}^n \psi_{1x}(n,m) - j_{mx}^n \psi_{3z}(n,m) + \\
& j_{mz}^n \psi_{3x}(n,m)] \} - \{ (\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}}) \cdot [j_{ex}^n \psi_{1z}(n,m) + j_{ey}^n \psi_{2y}(n,m) - \\
& j_{mx}^n \psi_{3z}(n,m)] - j_{ey}^m H(m) + j_{mx}^m C \} = 4\pi j_{mx}^m i \quad (3.56)
\end{aligned}$$

also,

$$\begin{aligned}
& \{ (\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}}) \cdot [j_{ey}^n \psi_{1z}(n,m) + j_{ez}^n \psi_{1y}(n,m) + j_{my}^n \psi_{3z}(n,m) - \\
& j_{mz}^n \psi_{3y}(n,m)] \} + \{ (\sum_{n=1}^{N_y} + \sum_{n=1}^{N_{-y}}) \cdot [j_{ex}^n \psi_{2x}(n,m) + j_{ez}^n \psi_{1y}(n,m) - \\
& j_{mz}^n \psi_{3y}(n,m)] \} + \{ (\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}}) \cdot [j_{ex}^n \psi_{2x}(n,m) + j_{ey}^n \psi_{1z}(n,m) + \\
& j_{my}^n \psi_{3z}(n,m)] + j_{ex}^m H(m) + j_{my}^m C \} = 4\pi j_{my}^m i \quad (3.57)
\end{aligned}$$

where $1 \leq m \leq N_z$.

Similarly, we can obtain $2N_{-x}$, $2N_{-y}$, and $2N_{-z}$ linear equations by enforcing eq. (3.29') at various reference points on the subsurfaces F_{-x} , F_{-y} , and F_{-z} , respectively. We list all these equations as follows:



$$\text{IV) } \hat{n} = -\hat{x}:$$

$$\begin{aligned} & \{ (\sum_{n=1}^{N_x} + \sum_{\substack{n=1 \\ n \neq m}}^{N_x}) \cdot [-j_{ey}^n \psi_{1x}(n, m) - j_{ez}^n \psi_{2z}(n, m) + j_{my}^n \psi_{3x}(n, m)] - j_{ez}^m H(m) - \\ & j_{my}^m C \} - \{ (\sum_{n=1}^{N_y} + \sum_{n=1}^{N_y}) \cdot [j_{ex}^n \psi_{1y}(n, m) + j_{ez}^n \psi_{2z}(n, m) + \\ & j_{mx}^n \psi_{3y}(n, m)] \} - \{ (\sum_{n=1}^{N_z} + \sum_{n=1}^{N_z}) \cdot [j_{ex}^n \psi_{1y}(n, m) + j_{ey}^n \psi_{1x}(n, m) + \\ & j_{nx}^n \psi_{3y}(n, m) - j_{my}^n \psi_{3x}(n, m)] \} = -4\pi j_{my}^m i \end{aligned} \quad (3.58)$$

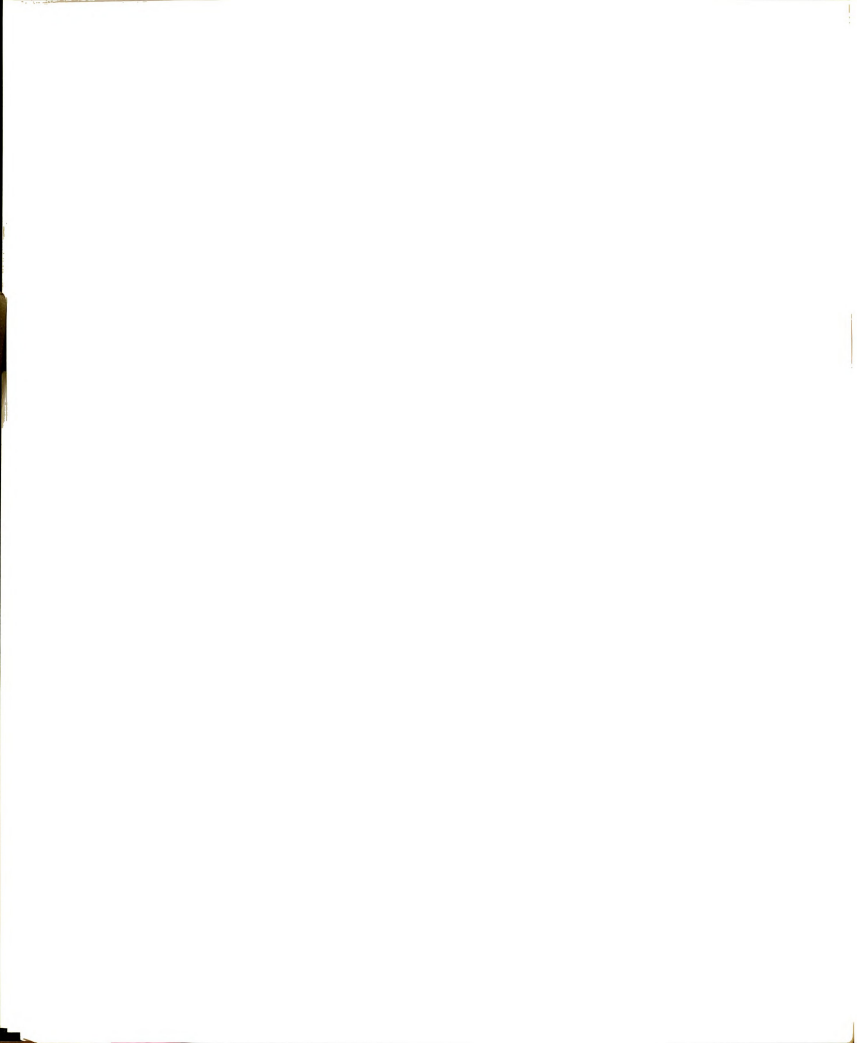
also,

$$\begin{aligned} & \{ (\sum_{n=1}^{N_x} + \sum_{\substack{n=1 \\ n \neq m}}^{N_x}) \cdot [j_{ey}^n \psi_{2y}(n, m) + j_{ez}^n \psi_{1x}(n, m) + j_{mz}^n \psi_{3x}(n, m)] + j_{ey}^m H(m) - \\ & j_{mz}^m C \} + \{ (\sum_{n=1}^{N_y} + \sum_{n=1}^{N_y}) \cdot [j_{ex}^n \psi_{1z}(n, m) + j_{ez}^n \psi_{1x}(n, m) - \\ & j_{mx}^n \psi_{3z}(n, m) + j_{mz}^n \psi_{3x}(n, m)] \} + \{ (\sum_{n=1}^{N_z} + \sum_{n=1}^{N_z}) \cdot [j_{ex}^n \psi_{1z}(n, m) + \\ & j_{ey}^n \psi_{2y}(n, m) - j_{mx}^n \psi_{3z}(n, m)] \} = -4\pi j_{mz}^m i \end{aligned} \quad (3.59)$$

where $1 \leq m \leq N_x$.

$$\text{V) } \hat{n} = -\hat{y}:$$

$$\{ (\sum_{n=1}^{N_x} + \sum_{n=1}^{N_x}) \cdot [j_{ey}^n \psi_{1x}(n, m) + j_{ez}^n \psi_{2z}(n, m) - j_{my}^n \psi_{3x}(n, m)] \} +$$



$$\begin{aligned}
& \left\{ \left(\sum_{n=1}^{N_y} + \sum_{\substack{n=1 \\ n \neq m}}^{N_{-y}} \right) \cdot [j_{ex}^n \psi_{1y}(n, m) + j_{ez}^n \psi_{2z}(n, m) + j_{mx}^n \psi_{3y}(n, m)] + \right. \\
& j_{ez}^m H(m) - j_{mx}^m C \} + \left\{ \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}} \right) \cdot [j_{ex}^n \psi_{1y}(n, m) + j_{ey}^n \psi_{1x}(n, m) + \right. \\
& j_{mx}^n \psi_{3y}(n, m) - j_{my}^n \psi_{3x}(n, m)] \} = -4\pi j_{mx}^m i \quad (3.60)
\end{aligned}$$

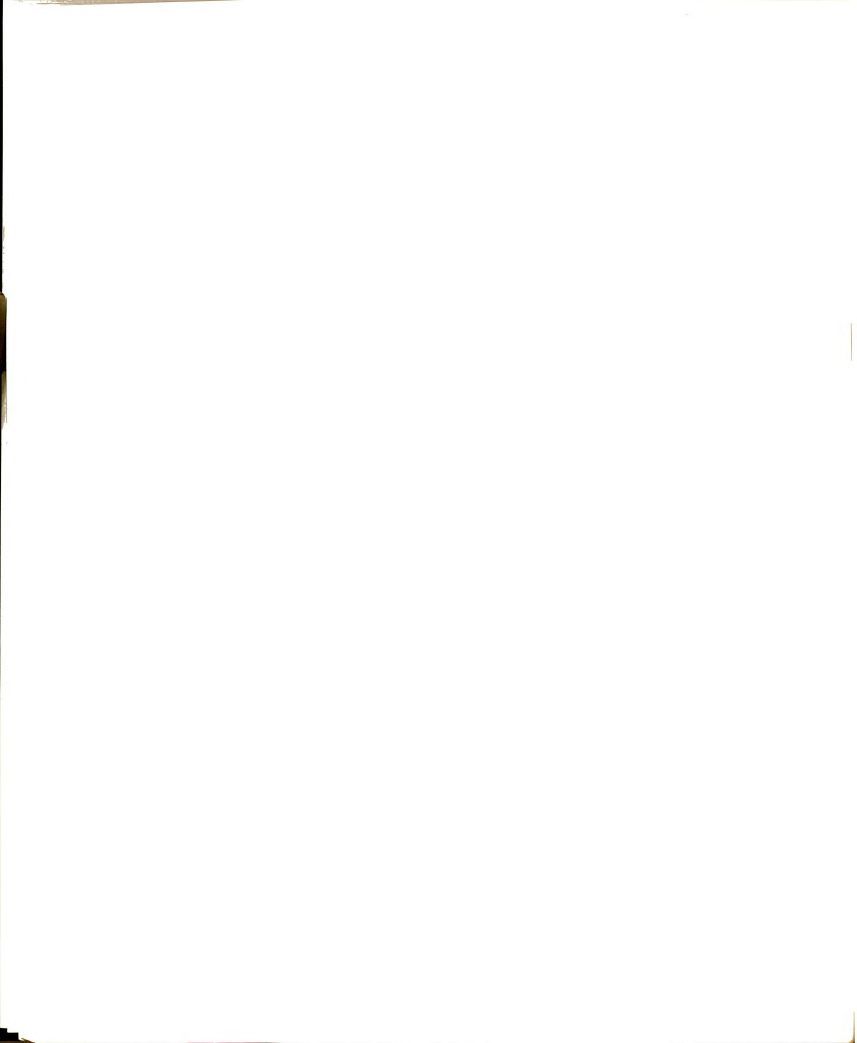
also,

$$\begin{aligned}
& \left\{ \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}} \right) \cdot [-j_{ey}^n \psi_{1z}(n, m) - j_{ez}^n \psi_{1y}(n, m) - j_{my}^n \psi_{3z}(n, m) + \right. \\
& j_{mz}^n \psi_{3y}(n, m)] \} - \left\{ \left(\sum_{n=1}^{N_y} + \sum_{\substack{n=1 \\ n \neq m}}^{N_{-y}} \right) \cdot [j_{ex}^n \psi_{2x}(n, m) + j_{ez}^n \psi_{1y}(n, m) - \right. \\
& j_{mz}^n \psi_{3y}(n, m) - j_{ex}^m H(m) - j_{mz}^m C] - \left\{ \left(\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}} \right) \cdot [j_{ex}^n \psi_{2x}(n, m) + \right. \\
& j_{ey}^n \psi_{1z}(n, m) + j_{my}^n \psi_{3z}(n, m)] \} \} = -4\pi j_{mz}^m i \quad (3.61)
\end{aligned}$$

where $1 \leq m \leq N_{-y}$.

VI) $\hat{n} = -\hat{z}$:

$$\begin{aligned}
& \left\{ \left(\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}} \right) \cdot [-j_{ey}^n \psi_{2y}(n, m) - j_{ez}^n \psi_{1x}(n, m) - j_{mz}^n \psi_{3x}(n, m)] \} - \right. \\
& \left\{ \left(\sum_{n=1}^{N_y} + \sum_{n=1}^{N_{-y}} \right) \cdot [j_{ex}^n \psi_{1z}(n, m) + j_{ez}^n \psi_{1x}(n, m) - j_{mx}^n \psi_{3z}(n, m) + \right. \\
& j_{mz}^n \psi_{3x}(n, m)] \} - \left\{ \left(\sum_{n=1}^{N_z} + \sum_{\substack{n=1 \\ n \neq m}}^{N_{-z}} \right) \cdot [j_{ex}^n \psi_{1z}(n, m) + j_{ey}^n \psi_{2y}(n, m) - \right.
\end{aligned}$$



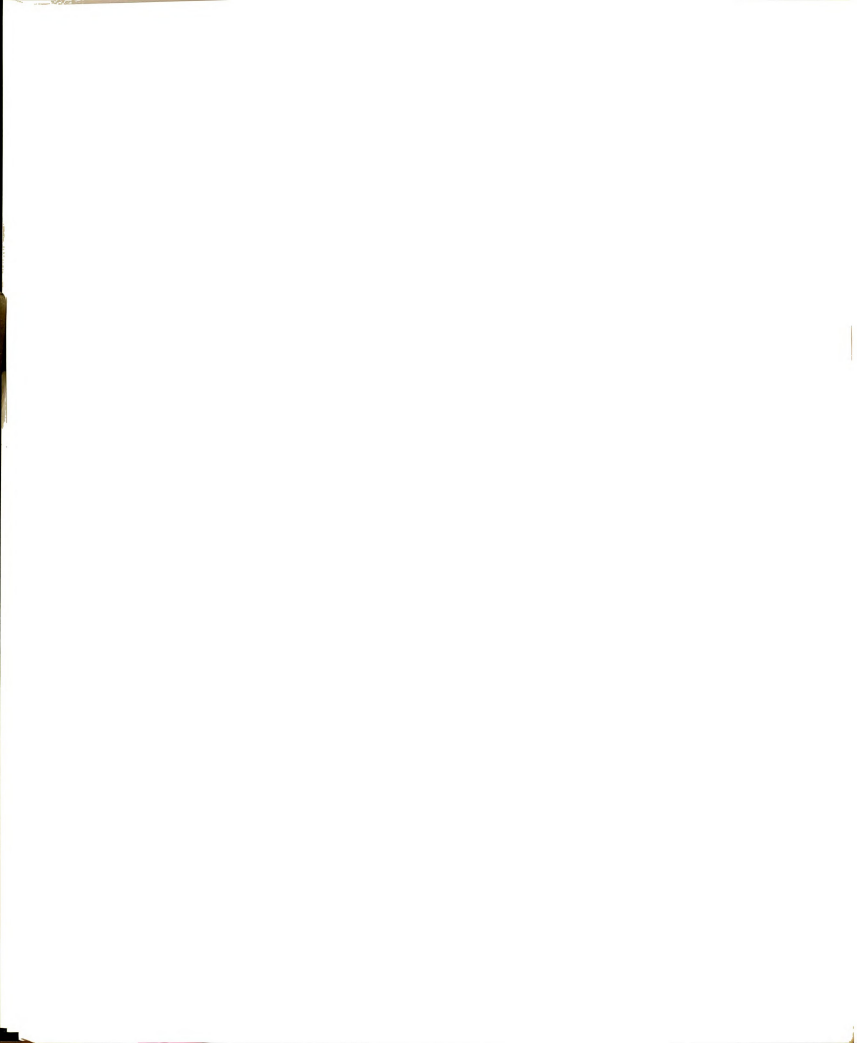
$$j_{mx}^n \psi_{3z}(n,m) - j_{ey}^m H(m) - j_{mx}^m C \} = -4\pi j_{mx}^m i \quad (3.62)$$

also,

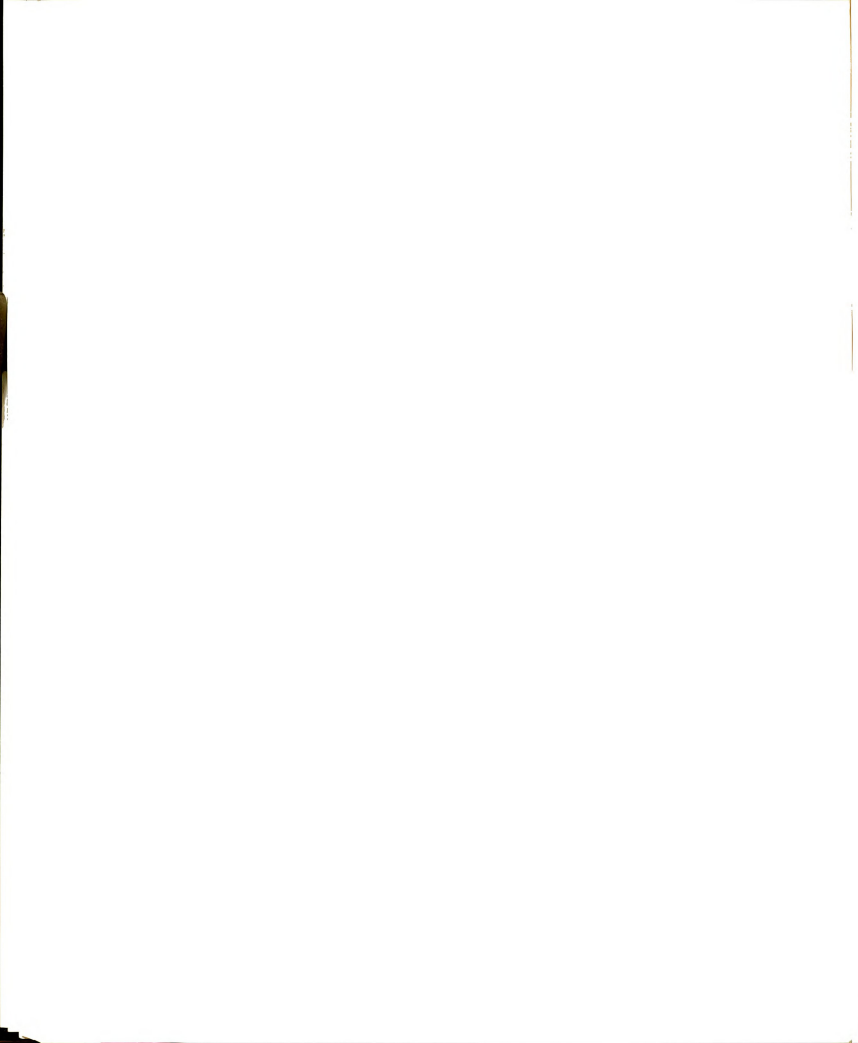
$$\begin{aligned} & \{ (\sum_{n=1}^{N_x} + \sum_{n=1}^{N_{-x}}) \cdot [j_{ey}^n \psi_{1z}(n,m) + j_{ez}^n \psi_{1y}(n,m) + j_{my}^n \psi_{3z}(n,m) - \\ & j_{mz}^n \psi_{3y}(n,m)] \} + \{ (\sum_{n=1}^{N_y} + \sum_{n=1}^{N_{-y}}) \cdot [j_{ex}^n \psi_{2x}(n,m) + j_{ez}^n \psi_{1y}(n,m) - \\ & j_{mz}^n \psi_{3y}(n,m)] \} + \{ (\sum_{n=1}^{N_z} + \sum_{n=1}^{N_{-z}}) \cdot [j_{ex}^n \psi_{2x}(n,m) + j_{ey}^n \psi_{1z}(n,m) + \\ & j_{my}^n \psi_{3z}(n,m)] + j_{ex}^m H(m) - j_{my}^m C \} = -4\pi j_{my}^m i \quad (3.63) \end{aligned}$$

where $1 \leq m \leq N_{-z}$.

Up to this point, we have obtained totally $2N_{\text{cell}}$ ($N_{\text{cell}} = N_x + N_y + N_z + N_{-x} + N_{-y} + N_{-z}$) scalar, linear algebraic equations based on eq. (3.29'). Repeating the same procedure for eq. (3.30'), we can obtain another $2N_{\text{cell}}$ linear equations. We therefore have totally $4N_{\text{cell}}$ equations, which are adequate to solve for the surface unknowns. For the purpose of saving space, we will not list the other $2N_{\text{cell}}$ equations; however, since eqs. (3.29') and (3.30') have very similar forms, they can be carefully deduced from eqs. (3.52)-(3.63). The combination of these $4N_{\text{cell}}$ linear equations can be transformed into a matrix form as shown in eq. (3.64), which then can be solved by a conventional

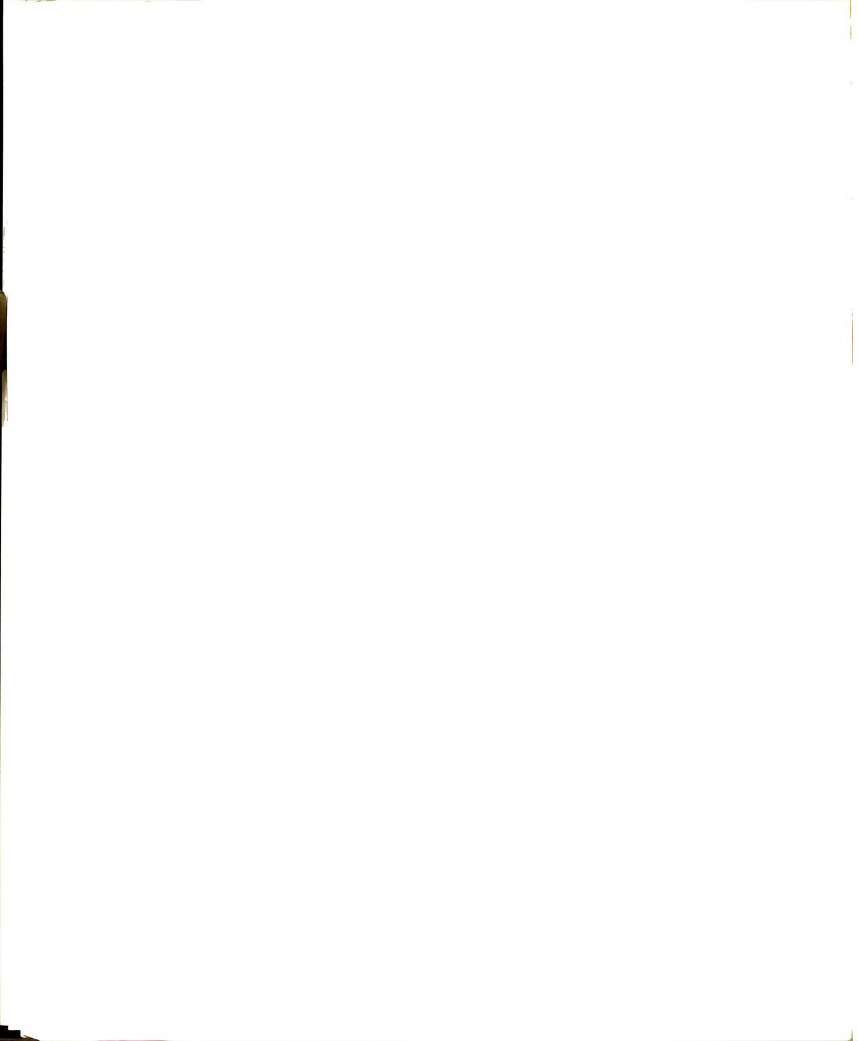


[illegible]



matrix inversion technique. The matrix elements in eq. (3.64) may be obtained from the above linear equations.

It should be mentioned that the matrix $[M]$ shown in eq. (3.64) possesses a four-fold symmetry, which is easy to see from the coefficients of the linear equations. These symmetry properties of the equations have been found very useful for the computer implementation of the numerical method. The initial version of the computer program is implemented directly based on eq. (3.64). Unfortunately, this program has only limited usefulness due to the prohibitive requirement of the matrix size needed to adequately sample current variations. However, if the incident plane wave is decomposed into four basic modes (cosine and sine variations of E and H fields), an eight-fold symmetrical property can be found for each type of surface current. These symmetrical properties are shown in Figure 3.12. A final version of the computer program, with the above symmetry conditions imposed, is then developed. By use of this computer program, it is possible to reduce the matrix size by a factor of 8 when an eight-fold symmetrical body is considered since it is sufficient to solve for the surface currents induced in only $1/8$ of the biological body. The induced surface currents for the rest of the body can be readily obtained by intuition. Of course, since the original incident plane wave has been



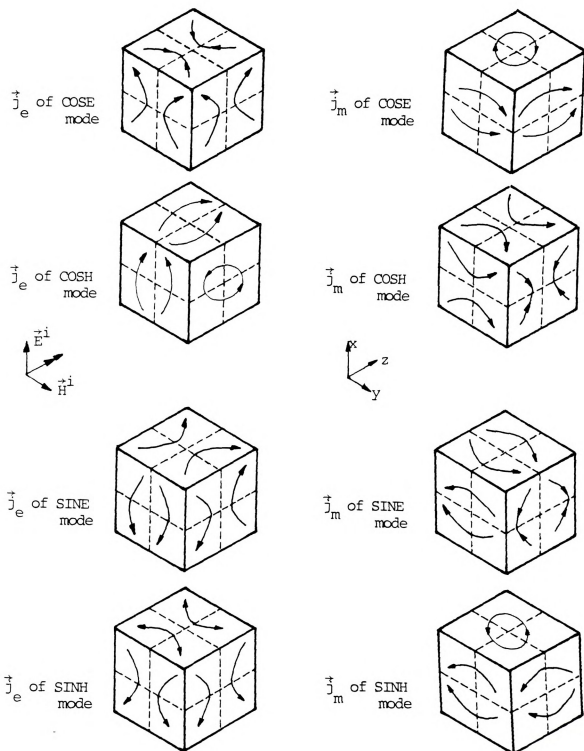
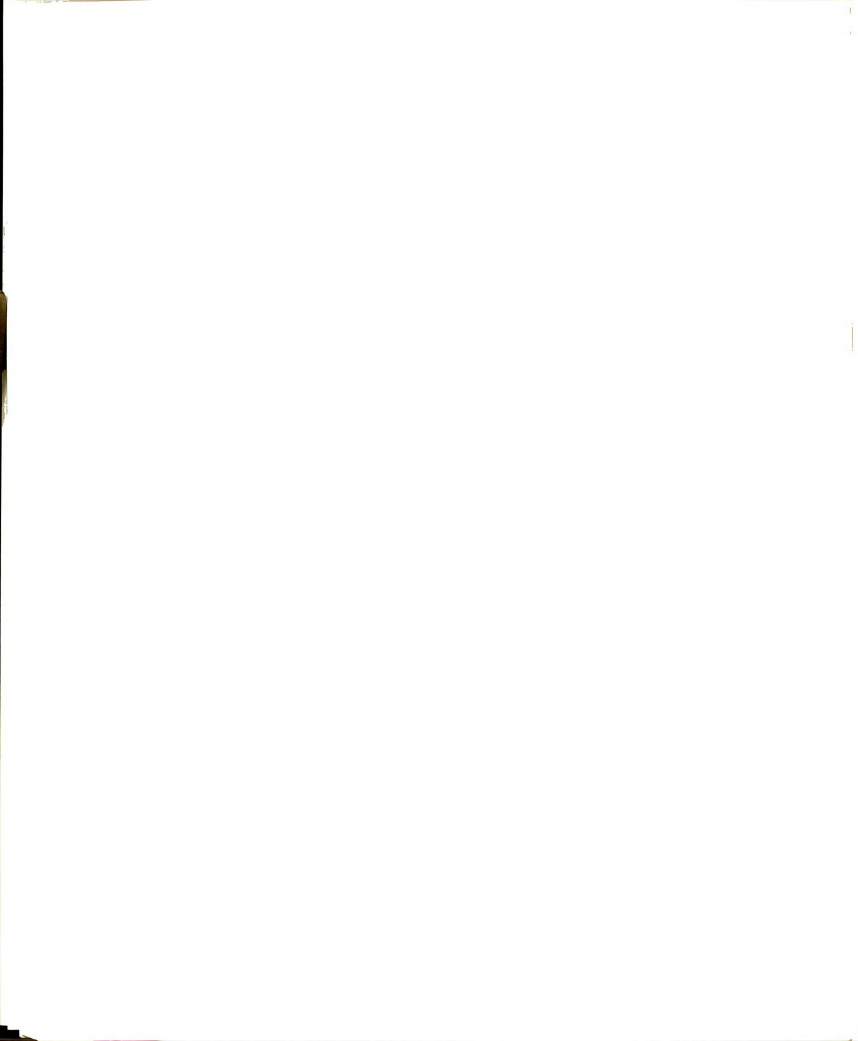


Figure 3.12. The symmetrical properties of various modes of the induced equivalent surface currents ($\vec{j}_e, \vec{j}_m$) on the surface of a cubic body.

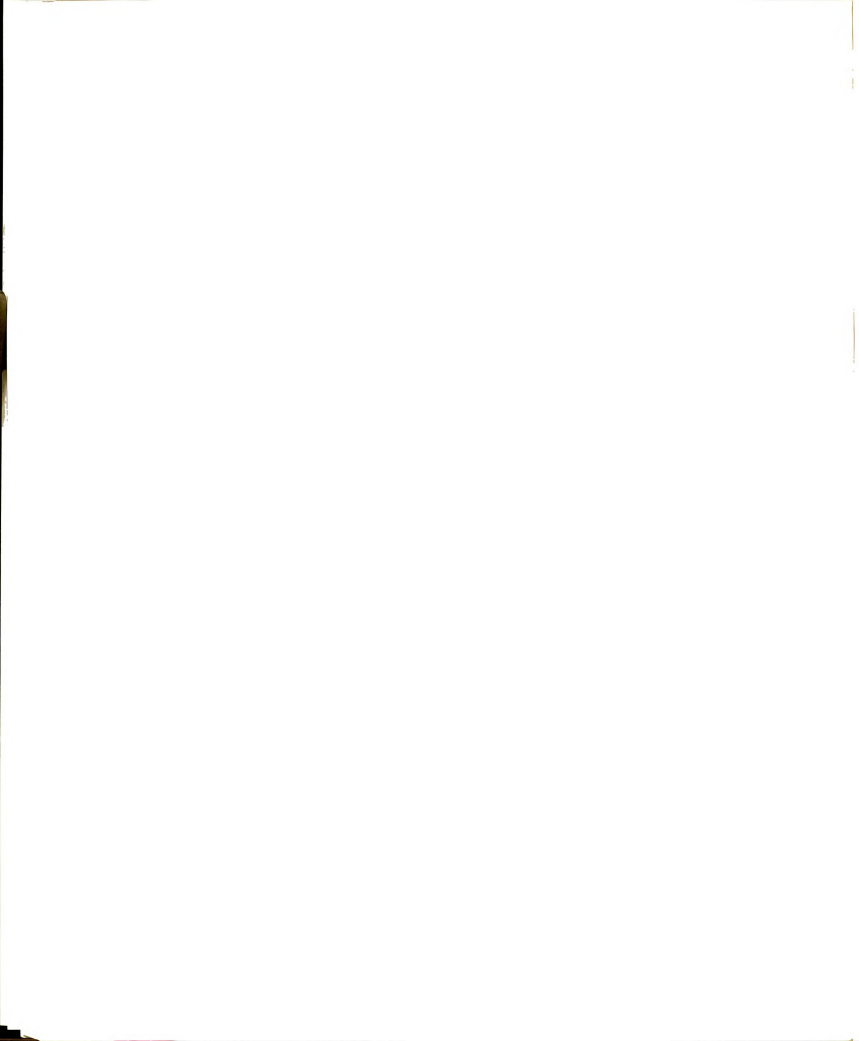


separated into four different modes, the final solutions must be the combinations of results obtained for each basic mode. More details about the computer programs will be discussed in Chapter IV.

3.8. Numerical Results

The numerical technique developed in the previous section has been applied to solve for the EM fields induced on the surfaces of several finite conducting bodies. These results of induced surface fields obtained by the surface integral equation method (SIEM), accompanied with the results of internal fields induced in the same bodies obtained by the volume integral equation method (VIEM) [18], will be presented in this section. It should be mentioned that although SIEM gives solutions for both the induced electric field and induced magnetic field, for the purpose of simplifying the presentation we will only emphasize the induced electric fields, especially the component which predominates.

Figure 3.13 shows the vertical component of the induced electric field, E_x , in a muscle layer of 6cm x 6cm x 1 cm irradiated by an EM wave of 100 MHz with a vertical incident electric field of 1 V/m at side-on incidence. The conductivity and dielectric constant of the body are assumed to be 0.889 S/m and 71.7. The top figure of Figure 3.13 shows the results for E_x (magnitude



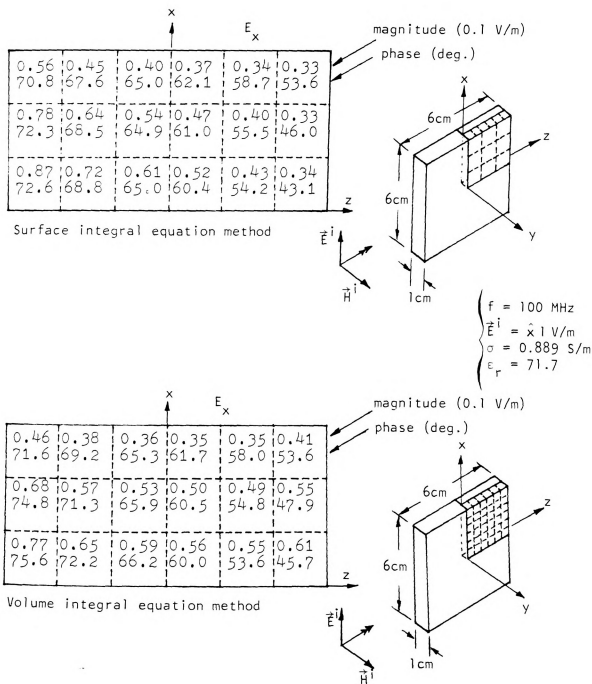
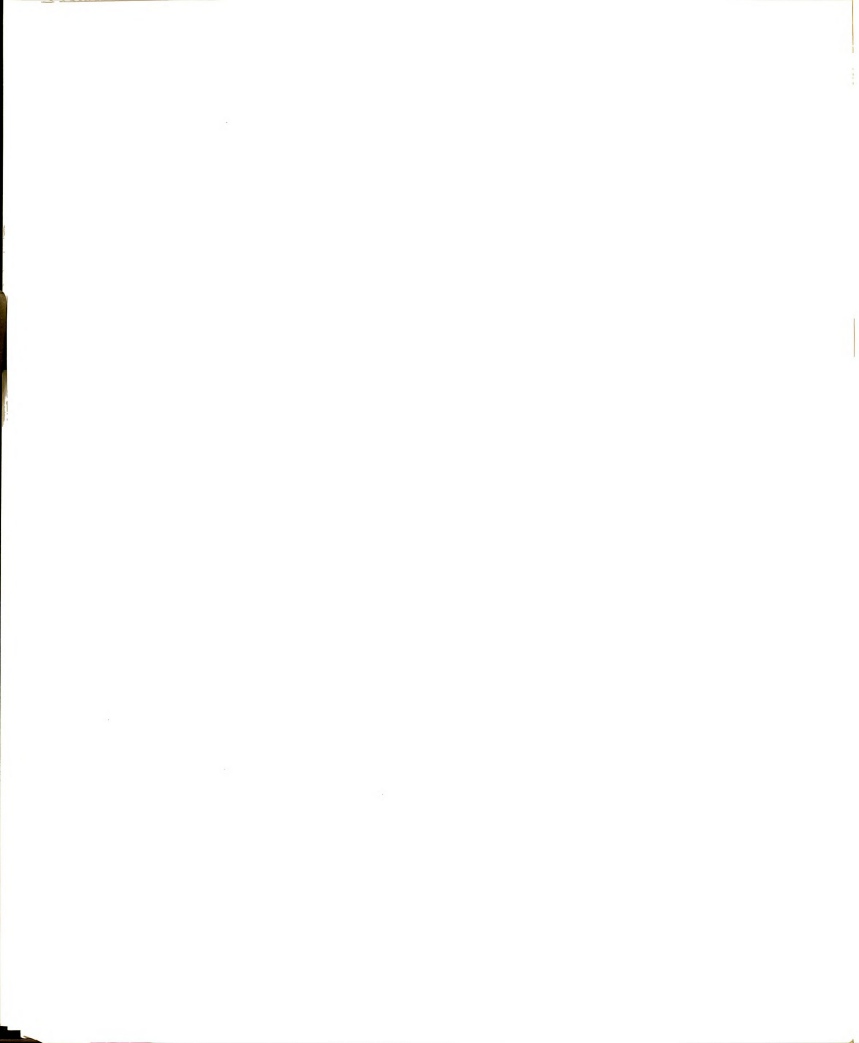
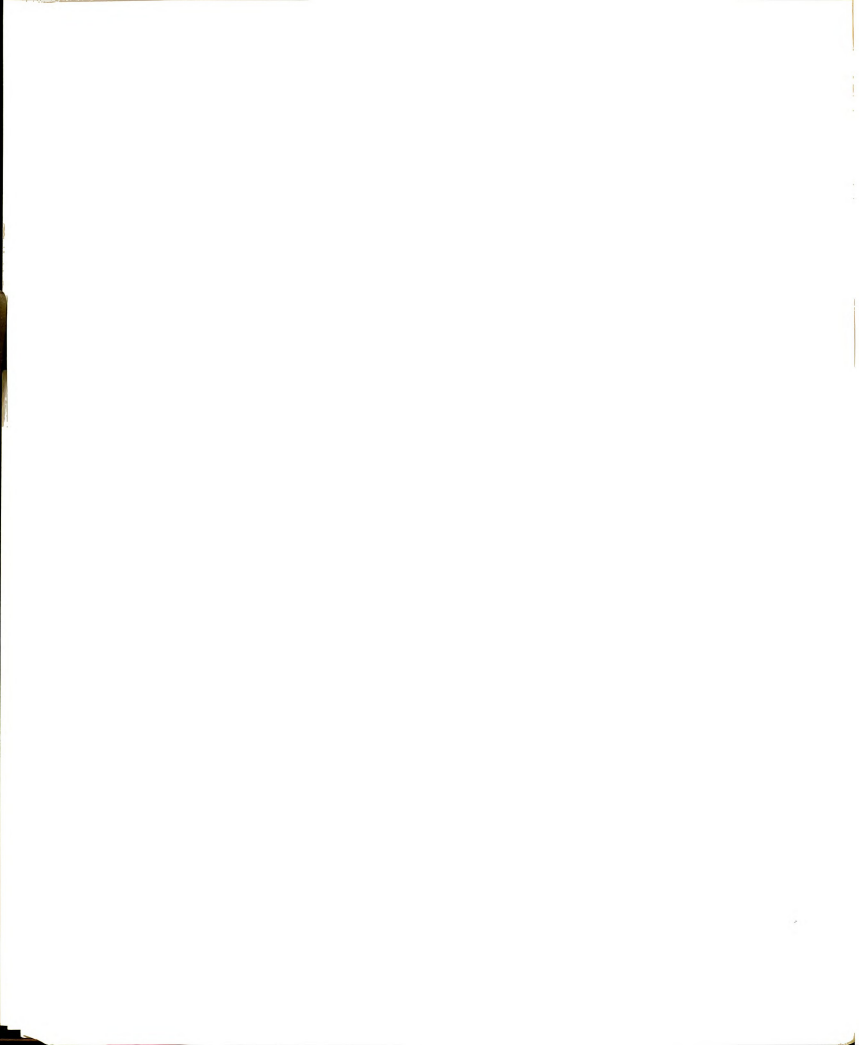


Figure 3.13. The x-components of the induced electric fields determined based on the surface integral equation method and the volume integral equation method.



and phase angle) on the surface of the body obtained with the surface integral equation method. In this numerical calculation, $1/8$ of the body surface is divided into 21 subareas of two different sizes. The bottom figure of Figure 3.13 shows the results for E_x (magnitude and phase angle) at the centers of the first layer cells, or at $y = 0.25$ cm plane, obtained with the tensor integral equation method. In this numerical calculation, $1/8$ of the body is divided into 36 volume cells. Comparison of these two sets of results, obtained with a surface integral equation method and a volume integral equation method, shows a qualitatively good agreement. This comparison is possible because the body is electrically thin in the y -direction, and the induced EM field should remain quite uniform in that direction. The disagreement between the two sets of results occurs mainly over the vertical edges of the body where the circulatory magnetic mode of the induced electric field has a significant contribution; the volume integral equation method usually gives poor results for that circulatory magnetic mode of the induced electric field.

Figure 3.14 shows the horizontal component of the induced electric field, E_z , in the same body under the same irradiation as the case of Figure 3.13. E_z near the central portion of the body, or the region near the x -axis, can be considered as mainly consisting of the



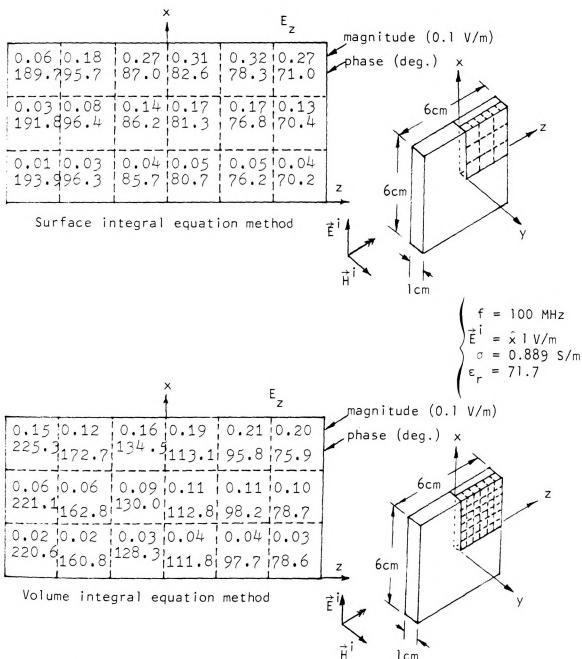
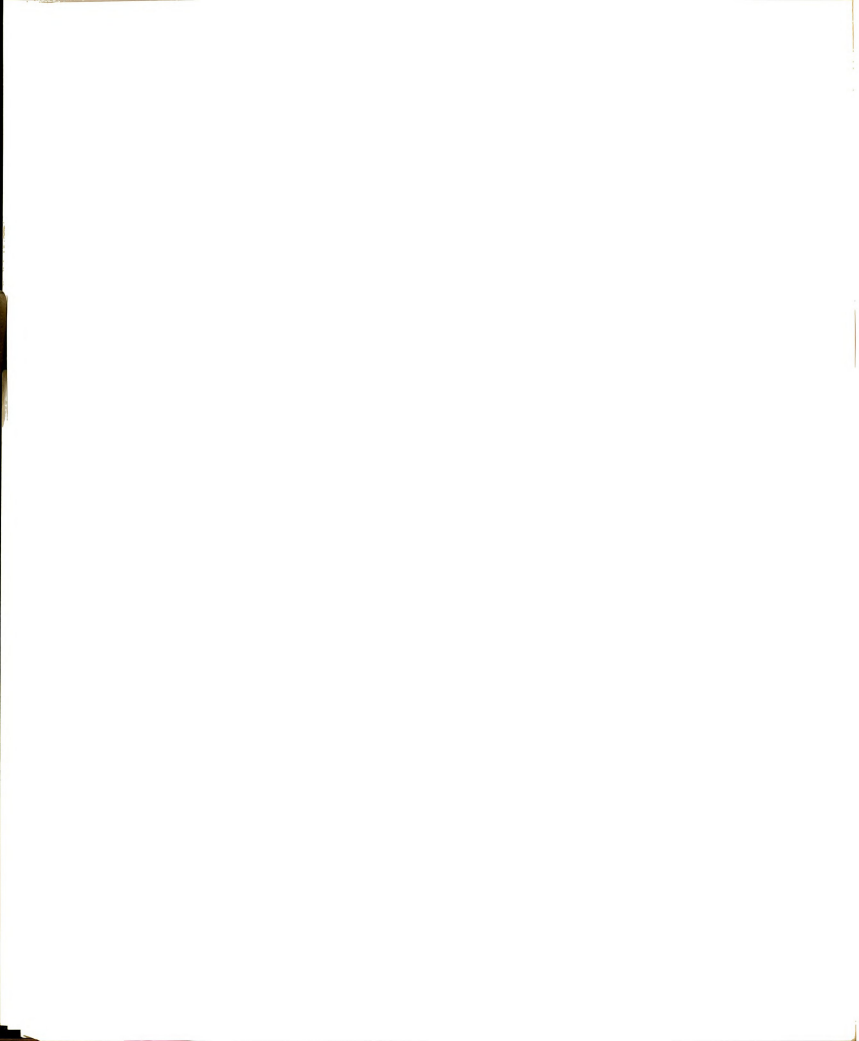


Figure 3.14. The z -components of the induced electric fields determined based on the surface integral equation method and the volume integral equation method.



circulatory magnetic mode of the induced electric field. For this component of induced electric field, only a fair agreement was obtained between the two sets of results obtained with the surface integral equation method and the volume integral equation method. The main reason for discrepancy is the inefficiency of a volume integral equation method in quantifying the circulatory magnetic mode of the induced electric field. The results obtained with the surface integral equation method appear to be more accurate for this case, because this set of results agrees with the results calculated by a method of quantifying the eddy current induced by a RF magnetic field as discussed in Chapter II.

As the second numerical example, we consider a rectangular body with a height of 180 cm, a width of 30 cm, and a thickness of 7.5 cm illuminated by a plane wave as shown in Figure 3.15. The incident wave is assumed to be x-polarized with a frequency of 10 MHz and an electric field intensity of 1 V/m, and the body is assumed to have a conductivity (σ) of 0.625 S/m and a dielectric constant (ϵ_r) of 160. For the purpose of comparison, we plot the amplitude and phase distributions of the x-component of induced electric field in the top and bottom figures of Figure 3.15, where the results of both SIEM and VIEM are included. In the calculation by SIEM, 1/8 of the body surface is divided into 52 surface cells of two different



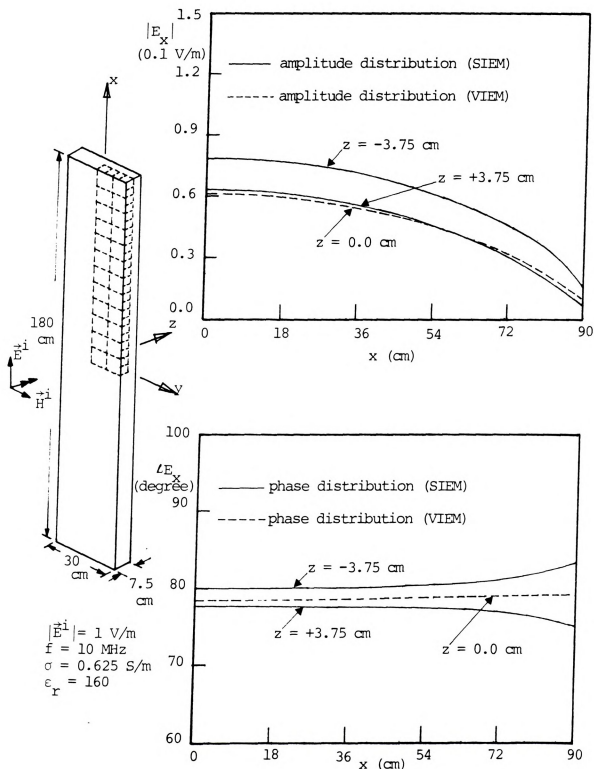
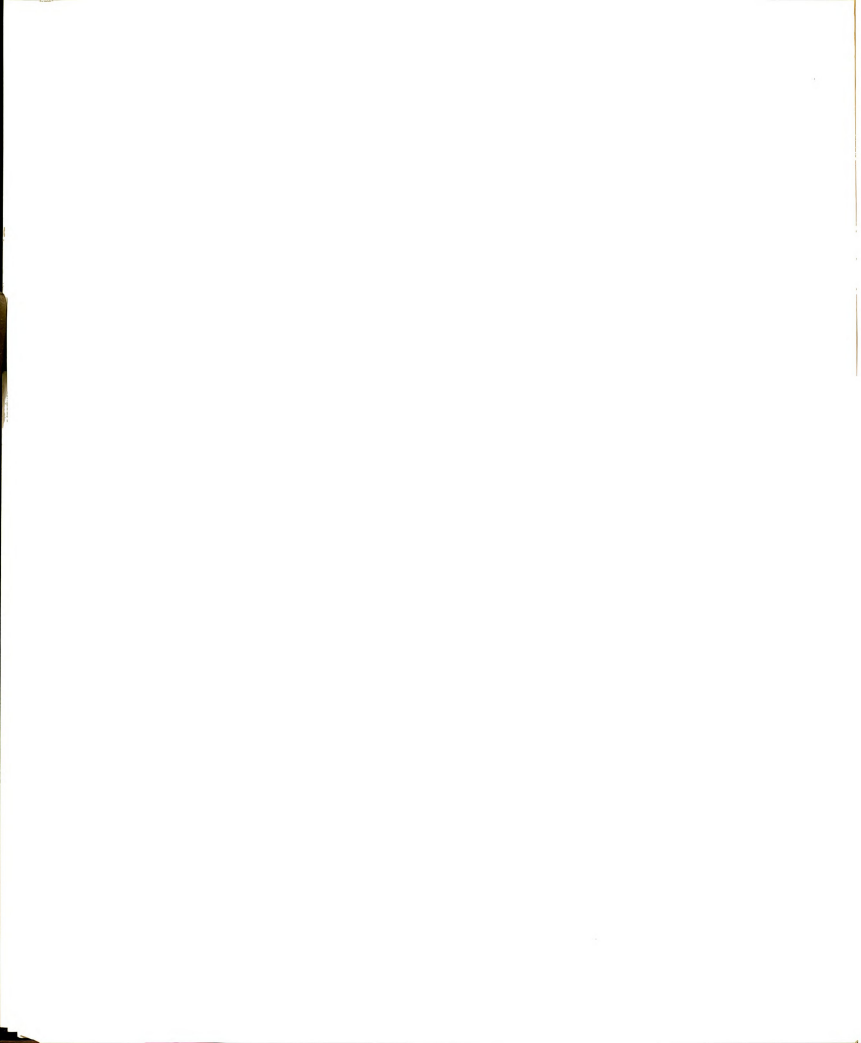
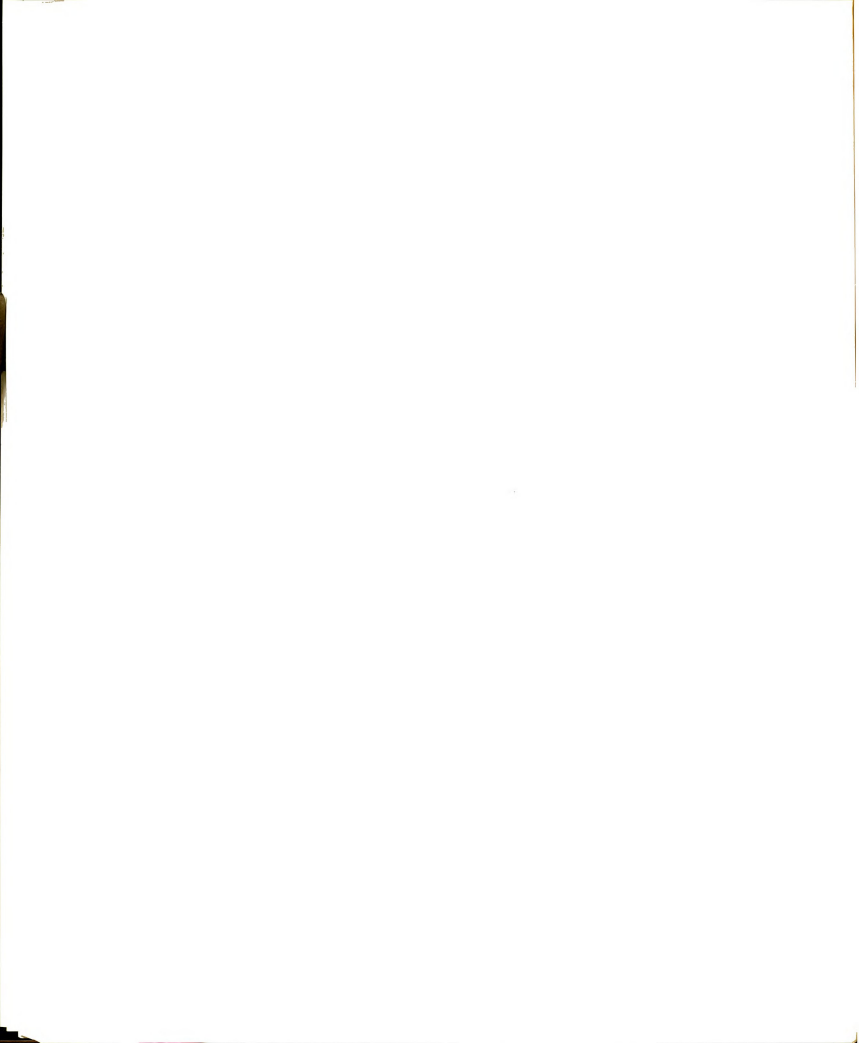


Figure 3.15. Amplitude and phase distributions of the x -component of induced electric field along the x direction in various layers at $y = 3.75$ cm.

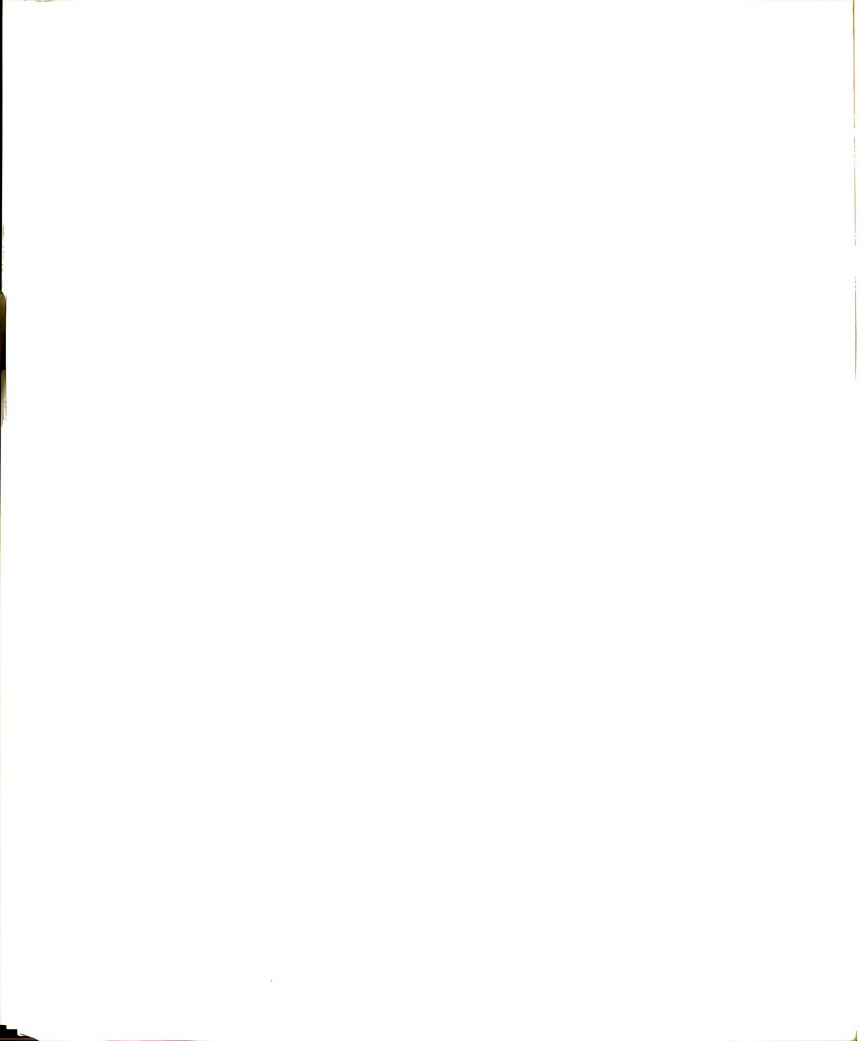


sizes. The EM fields induced at center points of the surface cells are numerically evaluated. In the top figure of Figure 3.15, we plot, in solid lines, the amplitude distributions of the x-component of electric field, E_x , induced at $y = 3.75$ cm on the front surface and the back surface as functions of the x-coordinate. The corresponding phase distributions are plotted in the bottom figure of Figure 3.15. For the numerical calculation by VIEM, 1/4 of the body is divided into 24 volume cells. The internal electric fields induced at center points of the volume cells are then numerically calculated. In Figure 3.15, we plot, in dashed lines, the amplitude and phase distributions of the x-component of internal electric field induced at $y = 3.75$ cm and $z = 0.0$ cm as functions of the x-coordinate. The comparison between these two sets of results, the results from SIEM and that from VIEM, shows a reasonably good agreement. However, we observe that the amplitude of the surface field induced on the front surface is almost 25% larger than that induced on the back surface, while the latter is very close to the results of VIEM as shown in Figure 3.15. Based on this observation, we would tend to conclude that the SIEM doesn't predict accurate results for the field induced on the front surface since we expect that the induced field is quite uniform along the z-direction. But, if we consider an extreme case with the assumption that the body



is extended to infinity in both transverse directions, we find that the surface field induced on the front surface of this infinite plane slab is about 40% larger than that induced on the back surface (based on solution of plane-wave reflection and transmission in multiple-layer region). This special case gives some qualitative explanations regarding the difference between the surface fields of the front surface and the back surface. Finally, we should emphasize that the comparison of the phase distributions obtained by SIEM and VIEM shows a very good agreement, especially in the central portion of the body.

Figure 3.16 shows the amplitude and phase distributions of the x-component of induced electric field along the x-direction in various layers at $y = 11.25$ cm for the same body illuminated by the same incident wave. The results shown in Figure 3.16 are very similar to those of Figure 3.15. This is understandable since for this case, the x-component of the induced electric field should be rather insensitive to the location along the y-direction. In Figure 3.16, the results obtained by SIEM are plotted in solid lines, while those obtained by VIEM are plotted in dashed lines. As we can see from the top figure of Figure 3.16, the amplitude distributions of the x-component of induced electric field decrease monotonically along the x-direction. On the other hand, we observe from the bottom figure of Figure 3.16 that the phase distributions are



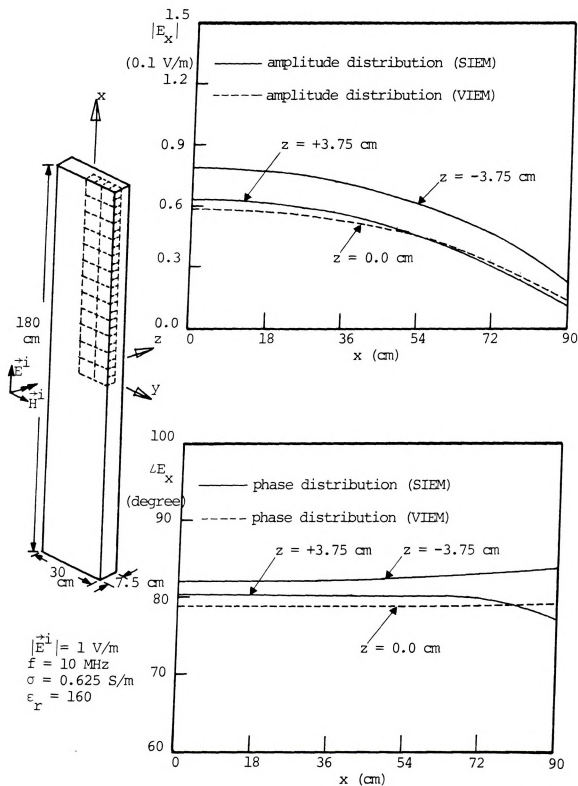
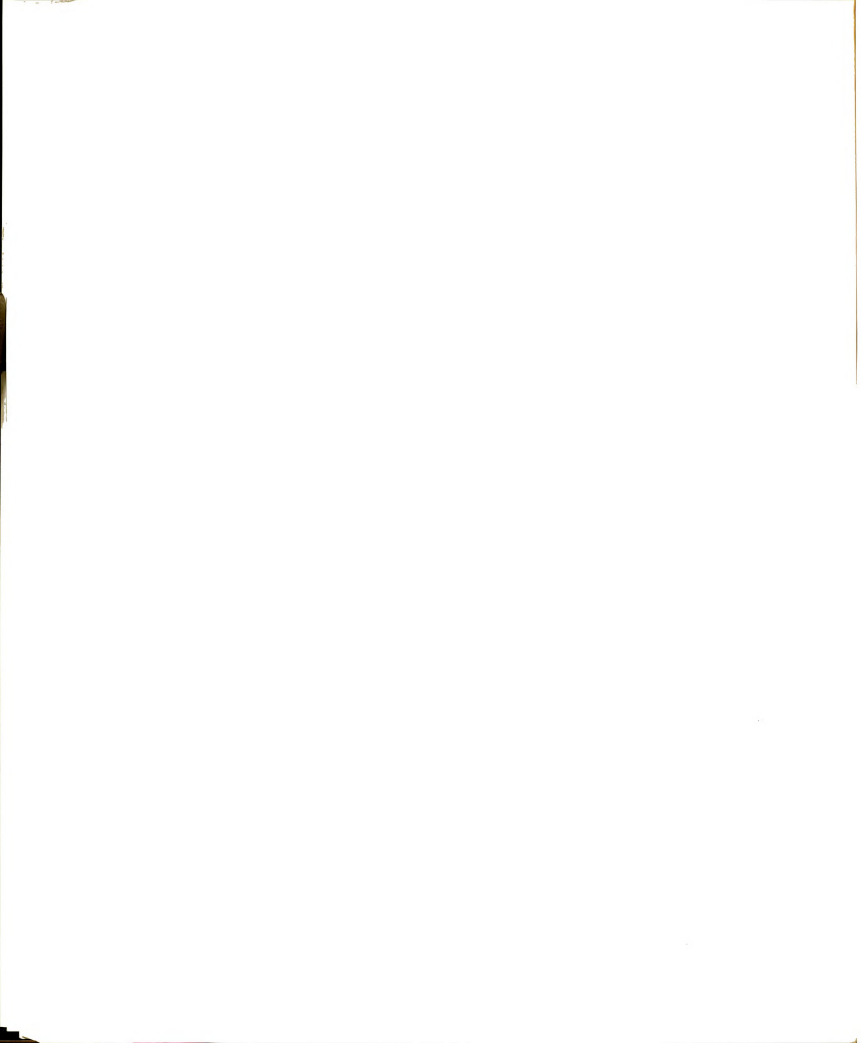


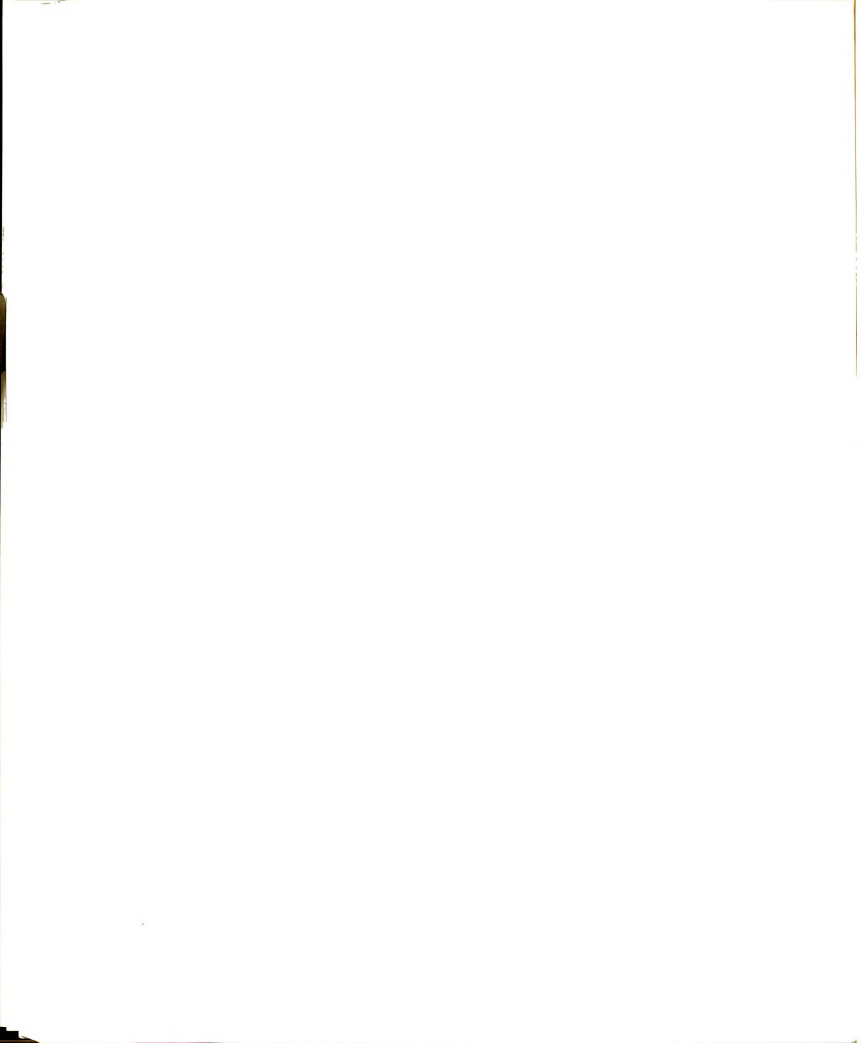
Figure 3.16. Amplitude and phase distributions of the x-component of induced electric field along the x direction in various layers at $y = 11.25$ cm.



nearly constant along the x-direction. Again, as in the case of Figure 3.15, the comparison between SIEM and VIEM shows reasonably good agreement, especially for the phase distributions.

The amplitude and phase distributions of the y-component of induced electric field along the x-direction in various layers at $y = 11.25$ cm and $y = 3.75$ cm are plotted in Figure 3.17. It is shown in Figure 3.17 that the field distributions at $y = 11.25$ cm (top figure) are very similar to those at $y = 3.75$ cm (bottom figure), except that the amplitude distributions of the former are three times as large as that of the latter. This is reasonable since the y-component of induced electric field must depend strongly on the locations along the y-direction. Figure 3.17 shows that the amplitude distributions of the y-component of induced electric field increase monotonically along the x-direction, while the phase distributions are nearly constant along the x-direction. It is to be noted that the y-component of induced electric field on the front surface and that on the back surface are exactly the same. This is quite different from what was observed in the previous two figures. Finally, it is clear from Figure 3.17 that the comparison between SIEM and VIEM demonstrates very good agreement for this case.

As the last example, we consider a more complicated geometry as shown in Figures 3.18 and 3.19. The body



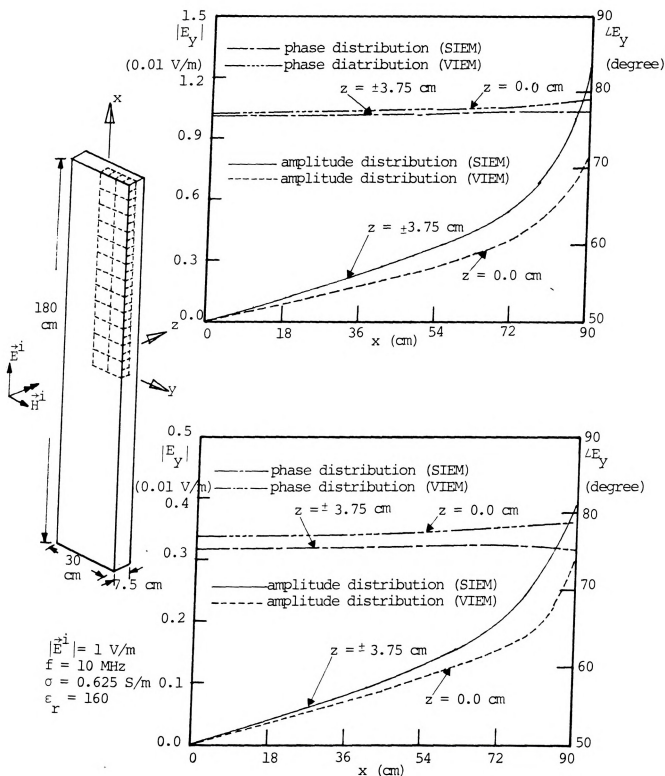
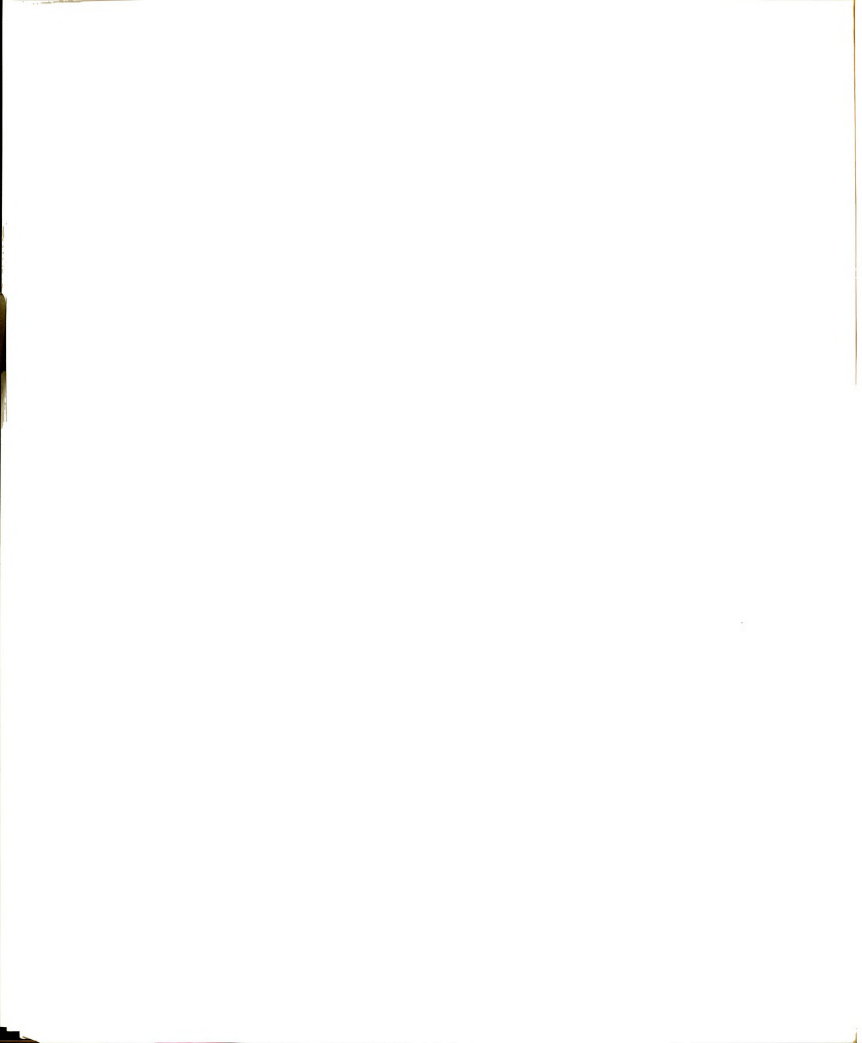
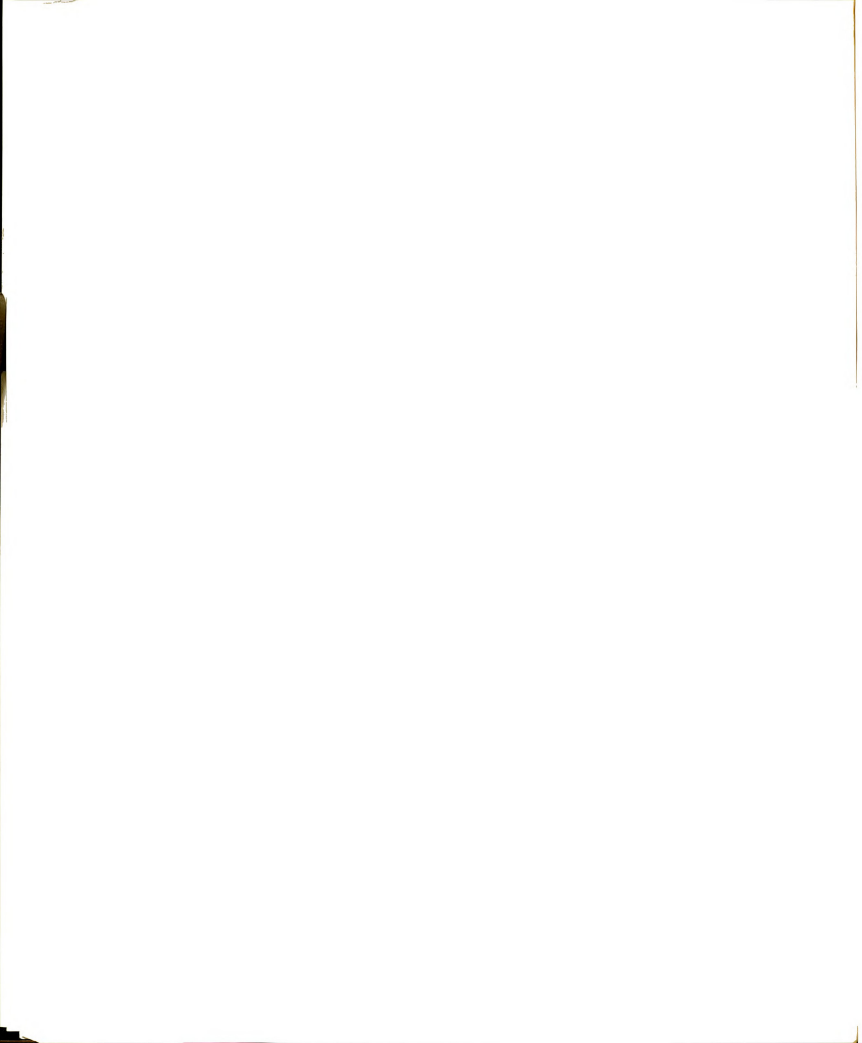


Figure 3.17. Amplitude and phase distributions of the y-component of induced electric field along the x direction in various layers at $y = 11.25 \text{ cm}$ (top figure) and $y = 3.75 \text{ cm}$ (bottom figure).



considered has a shape similar to that of a human body with the following dimensions: 70 cm in height, 20 cm in maximum width, and 10 cm in thickness. This situation can be considered as a simulation to the case of a two-year-old child innocently exposed to a high-intensity EM wave. For the purpose of reducing the number of unknowns to keep the computational cost low, we still assume that the frequency is as low as 10 MHz. The electrical properties of the body are assumed to be $\sigma = 0.625$ S/m and $\epsilon_r = 160$. In the calculation by SIEM, 1/4 of the body surface is divided into 29 surface cells of two different sizes. The amplitude distribution of the x-component of induced electric field is shown in Figure 3.18 for both the front surface and the back surface of the body. It is observed from Figure 3.18 that the electric field induced on the front surface is larger than that induced on the back surface. In Figure 3.19, the phase distribution of the x-component of induced electric field is shown. Unlike the amplitude distribution, the phase distribution is quite uniform over the entire body surface.



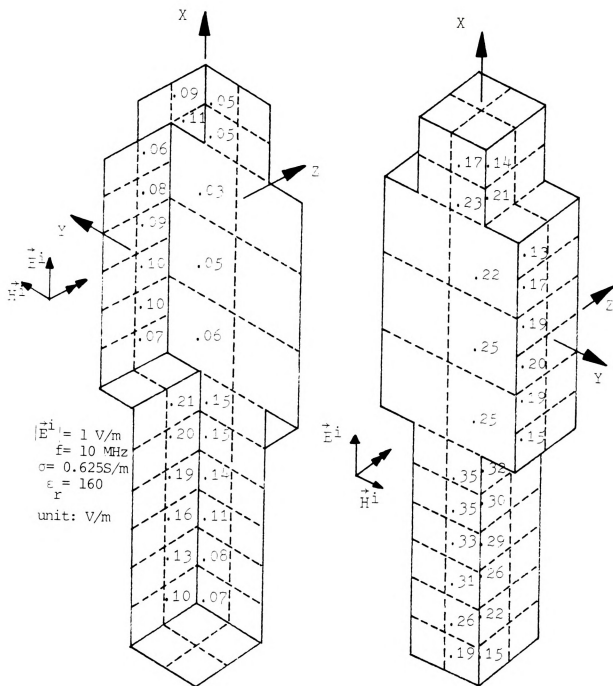
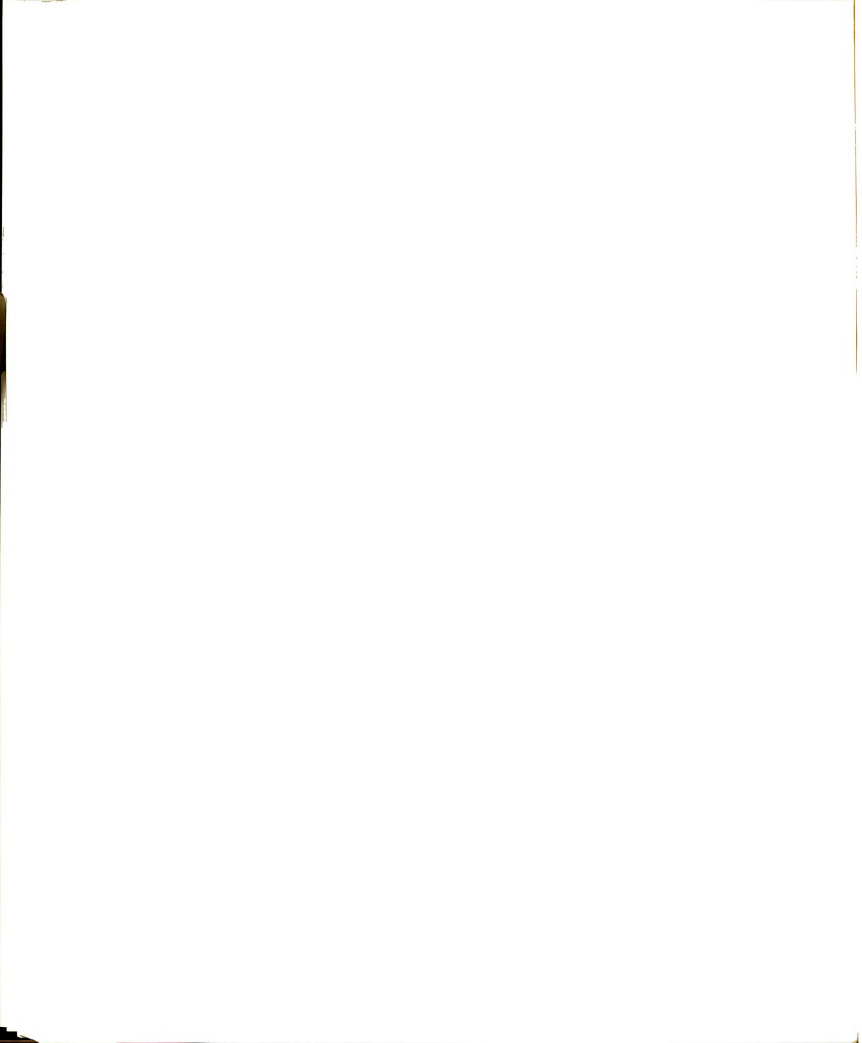


Figure 3.18. Amplitude distribution of the x-component of induced electric field on the surface of a simulated human body.



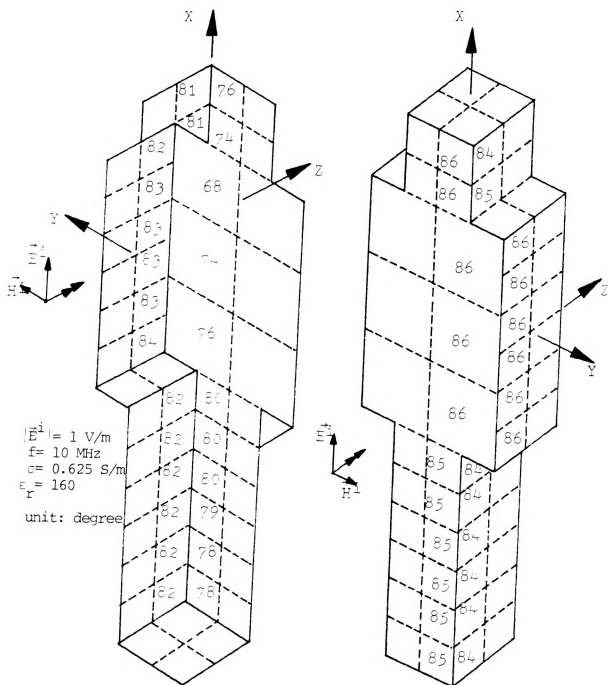
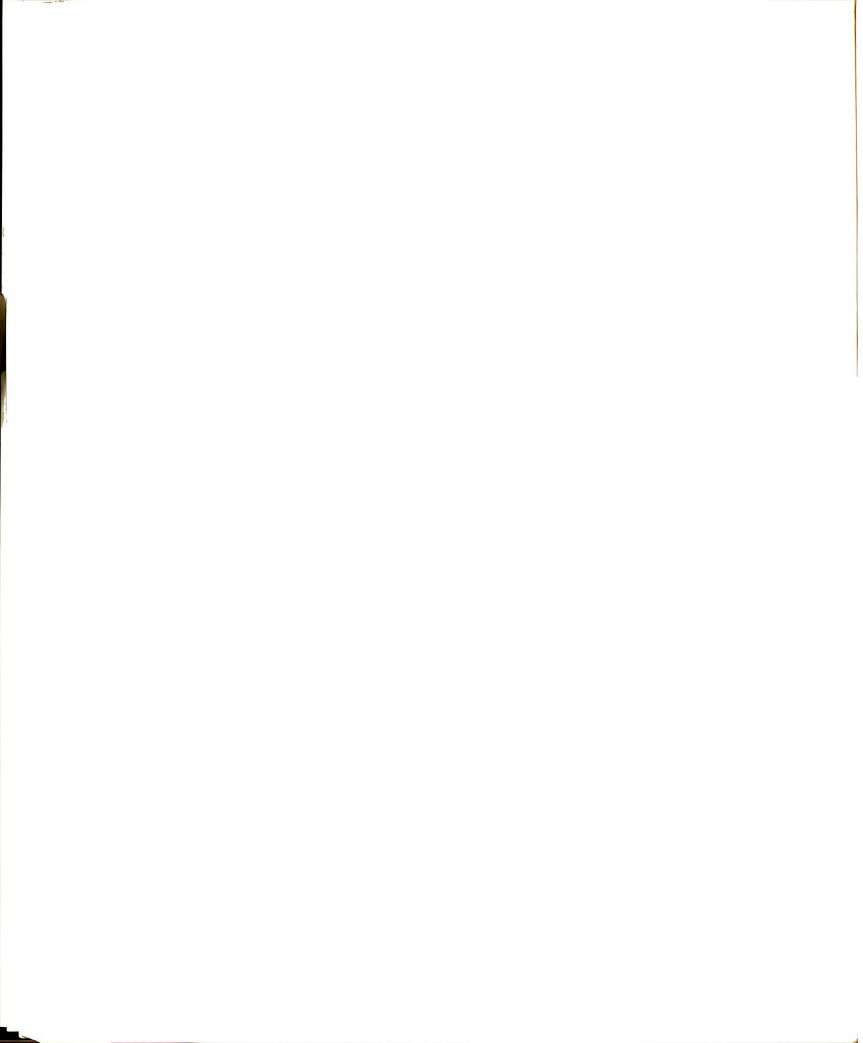


Figure 3.19. Phase distribution of the x-component of induced electric field on the surface of a simulated human body.



CHAPTER IV

PART 1

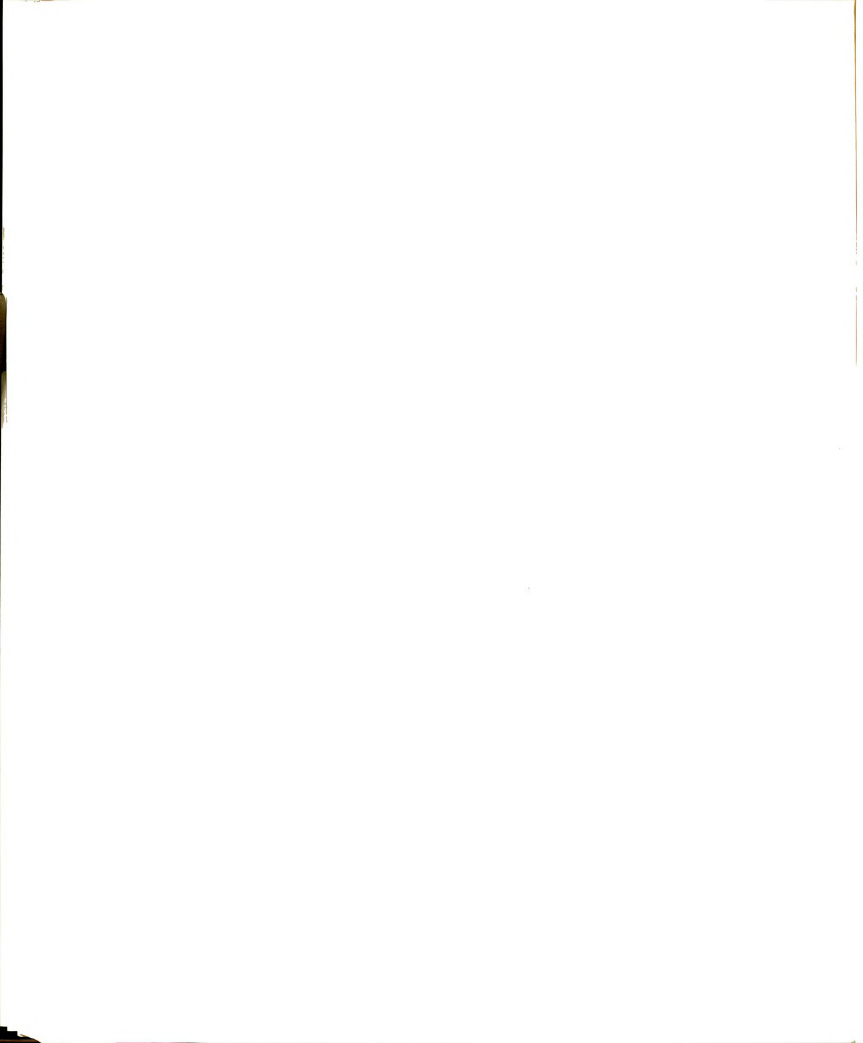
A USER'S GUIDE TO COMPUTER PROGRAM FOR INDUCED EDDY CURRENT INSIDE A FINITE CONDUCTING BODY WITH ROTATIONAL SYMMETRY

The first part of Chapter IV explains the computer program used to evaluate the eddy current induced at various locations within a finite conducting body with rotational symmetry. The theory behind this numerical technique has been presented in Chapter II. In addition to the listing of the program, an example will be worked out to help readers understand the sequential order of data files and the sample print out.

4.1. Formulation of the Problem

The program "EDDY" can handle two different cases:

(1) the case of a body immersed in a uniform magnetic field and (2) the case of a body irradiated by a uniform magnetic beam. The first step in utilizing this program is to divide a body with rotational symmetry into N circular rings with various dimensions and radii. Inside each circular ring the induced electric field (induced electric field and induced eddy current are interchangeable, since



they are only related by the conductivity) and the electrical properties are assumed to be uniformly distributed; however, they may be changed from ring to ring.

For the first case considered, the impressed magnetic field can be written as

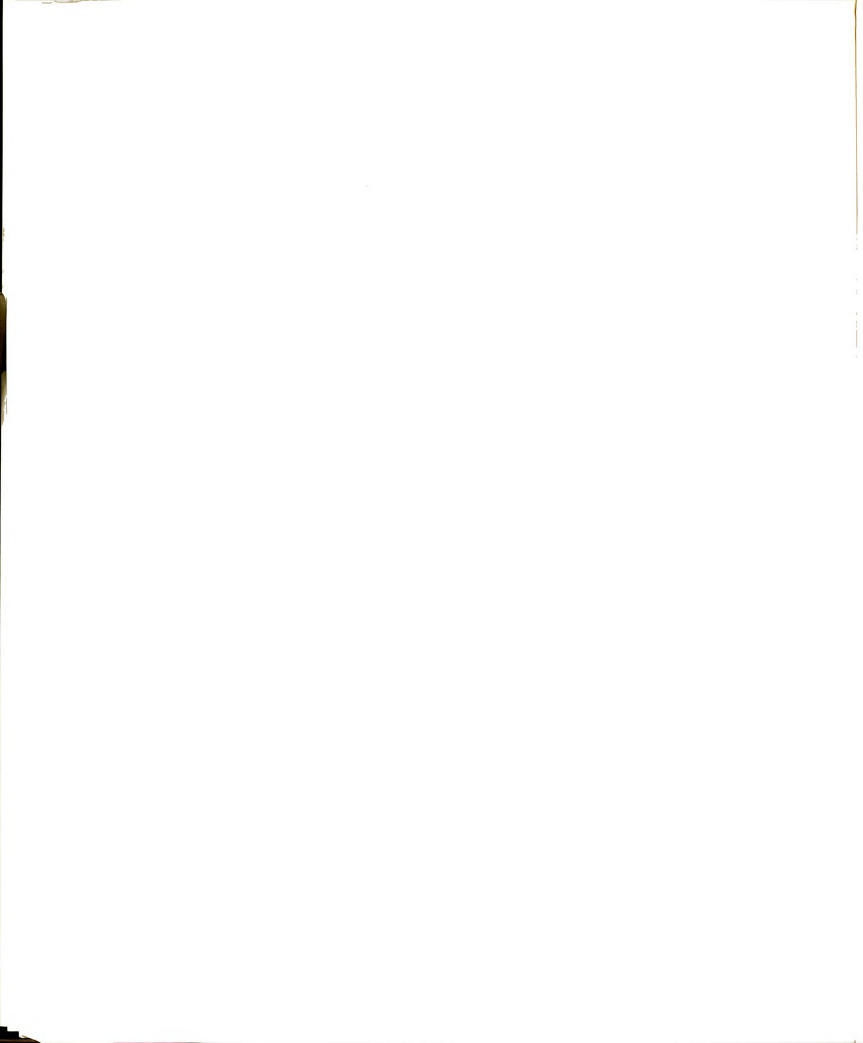
$$\vec{H}^i(r) = \hat{z}H^i.$$

Similarly, for the second case, it can be expressed as

$$\begin{aligned}\vec{H}^i(r) &= \hat{z}H^i & \text{for } 0 \leq r < b \\ &= 0 & \text{for } b \leq r.\end{aligned}$$

Here the time-varying factor $\exp(j\omega t)$ has been suppressed. Usually, we assume that the impressed magnetic field has unit intensity, $H^i = 1$ A/m. The body considered has a symmetrical axis in the $+z$ direction, as shown in Figure 4.1. For convenience, we also assume that the cross-section of each ring is a square which is then approximated by a circle in the integration. Since the induced electric fields have only one component (in the ϕ direction), we have N unknowns in total.

The second step in the numerical formulation of the problem is the specification of the location of each ring, its physical dimensions and its electrical parameters. As mentioned earlier, this program has been written for the case of uniformly impressed magnetic field or for the case of uniform magnetic beam. However, other types of



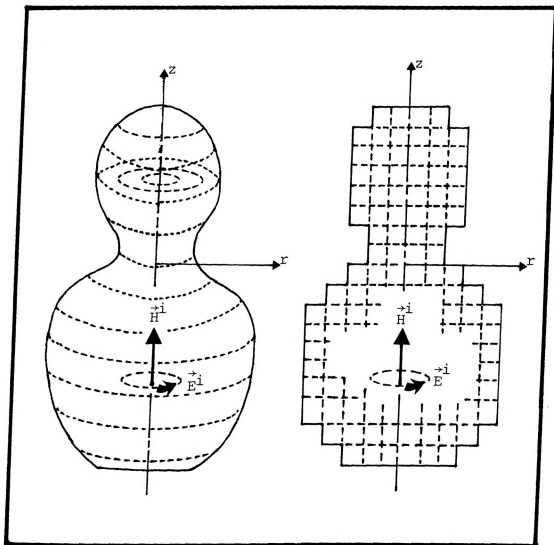
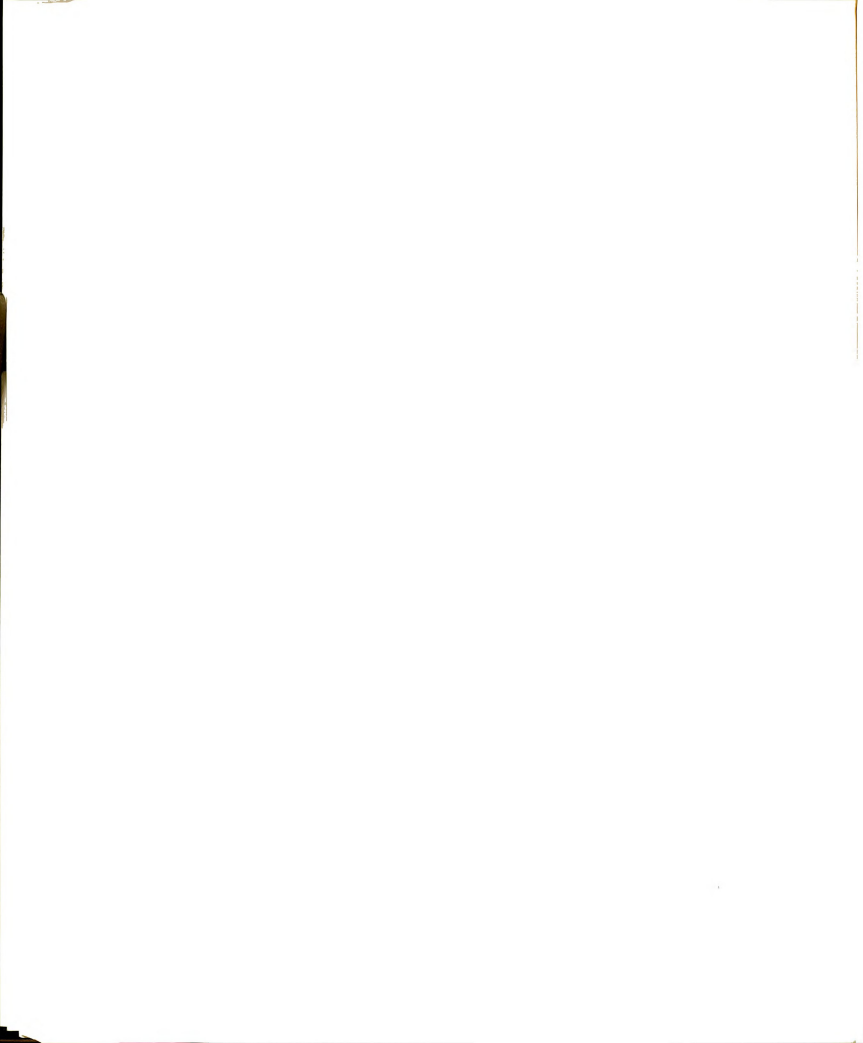


Figure 4.1. A rotationally symmetrical body with the symmetrical axis in the $+z$ direction is irradiated by a magnetic field with unit intensity. The body is subdivided into a number of circular rings of various radii.



impressed magnetic fields may be used with few changes in the main program.

4.2. Description of Computer Program

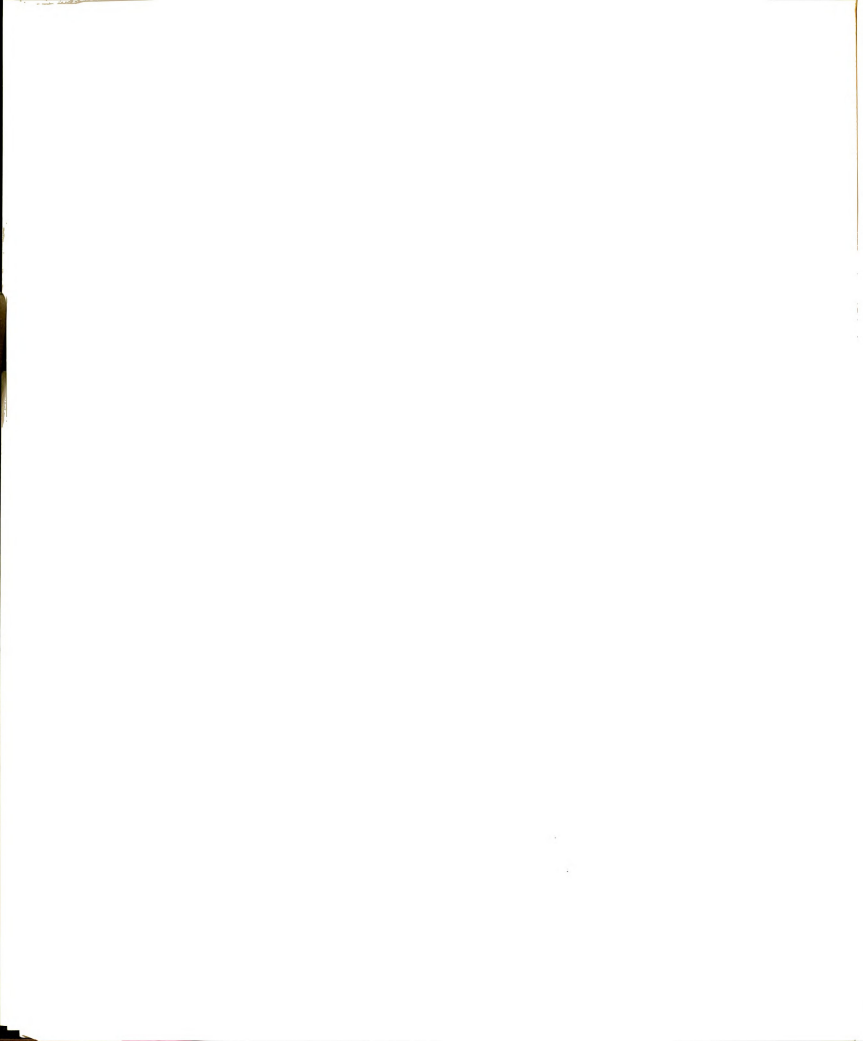
The program "EDDY" is coded in FORTRAN. It used the following complex functions and subroutines:

"LEEMAT"--is a subprogram which calculates the elements of $[M]$ matrix based on eqs. (2.16) and (2.17).

"FMM C," "FMMS," "FMNC," "FMNS"--are function subprograms which determine the integrands of integrals K_{mn} , and K_{mm} as expressed in eqs. (2.6) and (2.8). Note that we decompose each integrand into real and imaginary parts. Since these four functions are used as calling arguments for function subprogram "DCADRE," they must be declared external in subroutine "LEEMAT" which calls subprogram "DCADRE" for numerical integration.

"DCADRE"--is an IMSL routine which integrates a function by using technique of Cautious Adaptive Romberg Extrapolation.

"CMATP"--is a subprogram which calculates the induced electric field in each ring by solving $N \times N$ matrix. It is actually a Gauss-Seidel



method for solving a system of N equations in N unknowns.

"MAGPHA"--this subprogram, as the name suggests, determines magnitude and phase of the induced electric field in each ring.

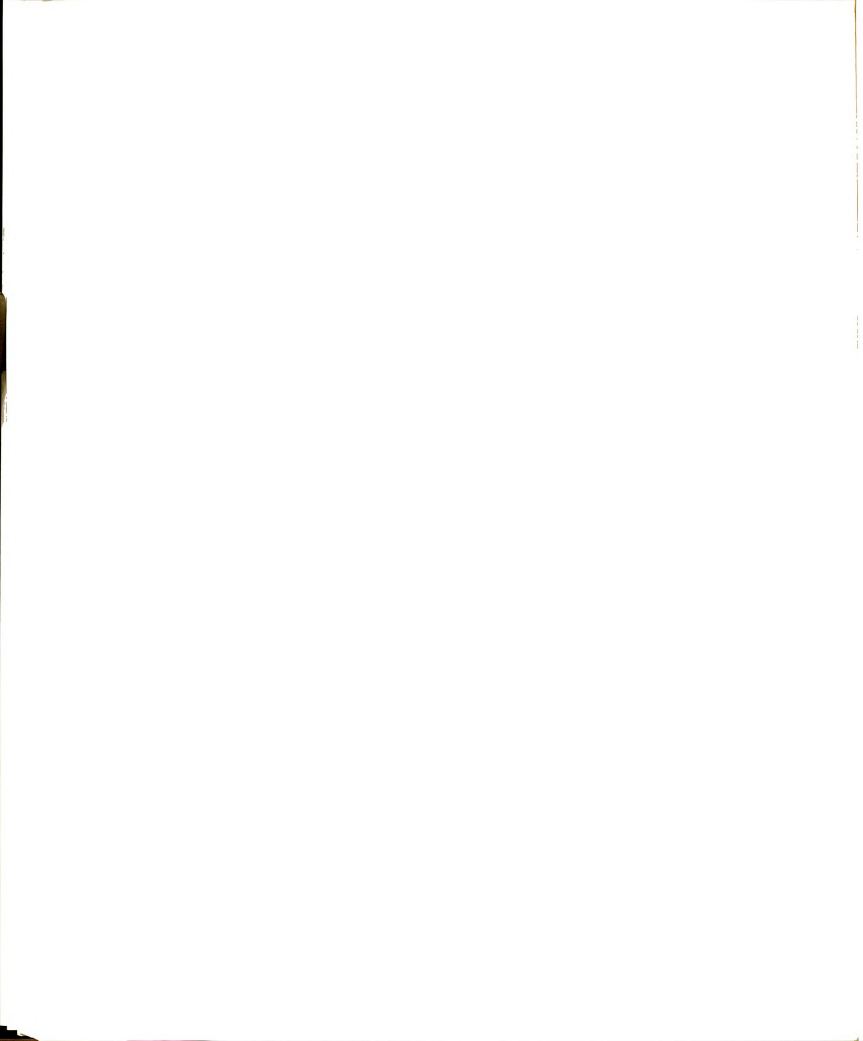
"MCURVE," "PCURVE"--these two subroutines plot the magnitude and phase distributions of the induced electric fields as functions of radial distance.

A listing of the program "EDDY" is given at the end of this chapter.

4.3. Data Structure and Input Variables

Figure 4.2 shows a typical example of a body with rotational symmetry we refer to from time to time. This body is divided into six circular rings. The origin of the coordinate system can be arbitrarily chosen at any point along the central axis of the body; then the location of each ring is determined with respect to this origin. In this example, it is assumed that all the rings have the same cross-sectional area and electrical parameters.

The sequential structure of the input data files, the format specifications and the symbolic names of the variables appearing on each file are outlined in Table 4.1. There are totally three input files; each of them consists



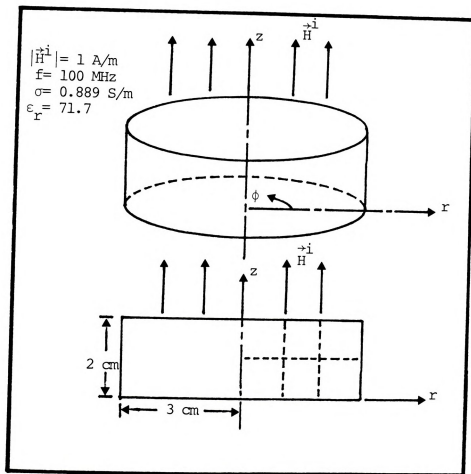


Figure 4.2. A cylindrical biological body with a diameter of 6 cm and a height of 2 cm is subdivided into 6 circular rings with a cross-section of 1 cm x 1 cm for each ring. The body is assumed to have a conductivity of 0.889 S/m and a dielectric constant of 71.7, and is immersed in a 100 MHz uniform magnetic field with an intensity of 1 A/m.

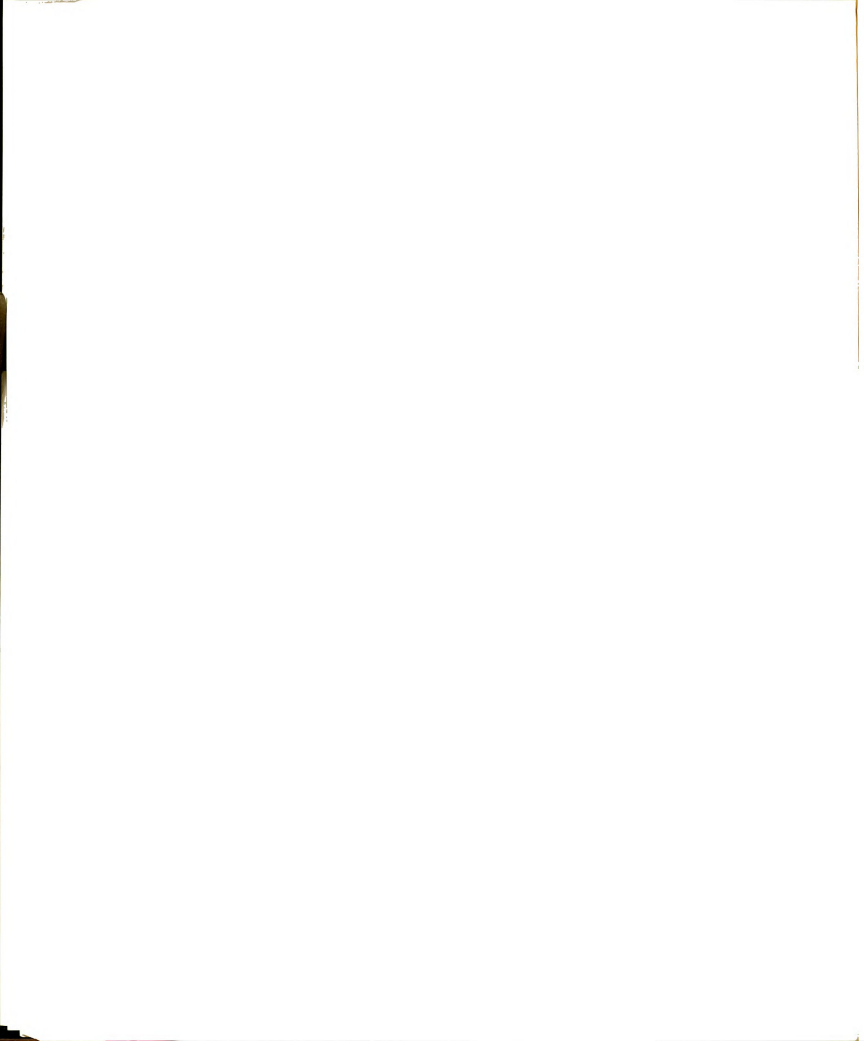
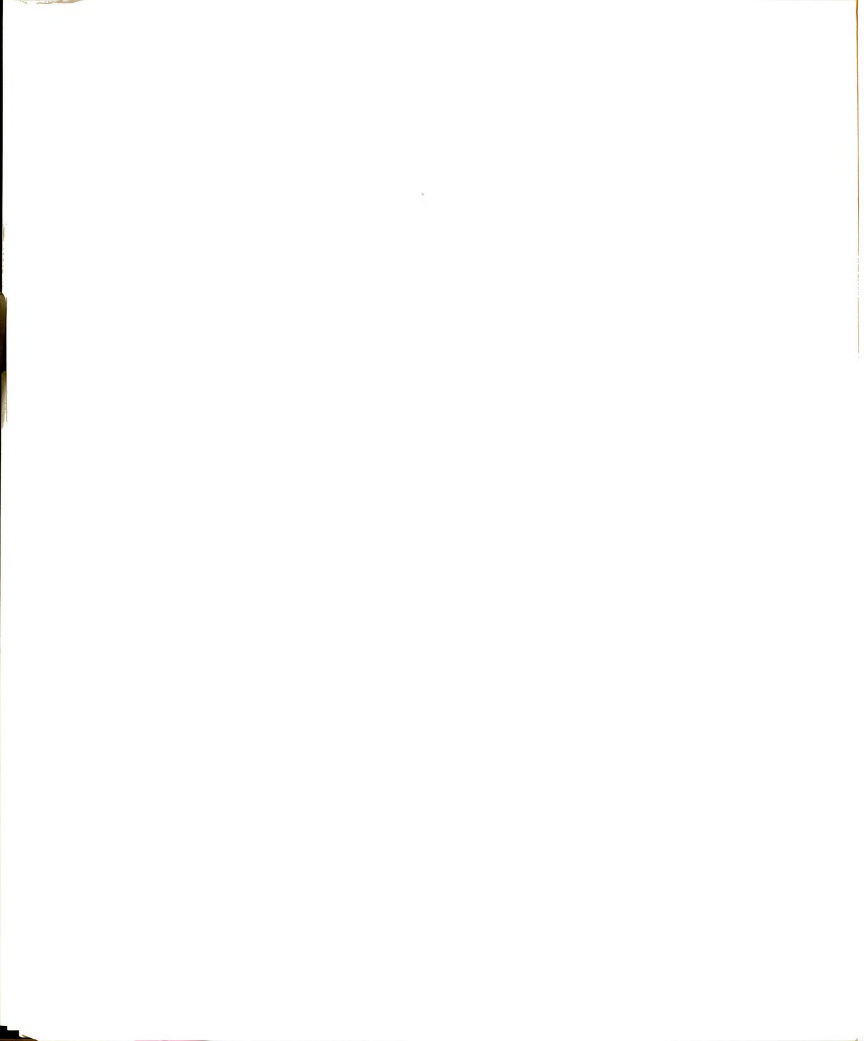


Table 4.1. The symbolic names of input variables and corresponding specifications for the data files used in data structure for the program "EDDY."

File No.	Card No.	Symbolic Name	Columns	Format
1	1	N	1-5	I5
		FREINM	6-15	F10.5
		NX	16-20	I5
		NXB	21-25	I5
2	1	NPAR	1-5	I5
		AERR	6-15	F10.5
		RERR	16-25	F10.5
3	1-N	XEND	1-12	F12.5
		ZEND	13-24	F12.5
		XAA	25-36	F12.5
		ZBB	37-48	F12.5
		REP	49-60	F12.5
		SIG	61-73	E13.6



of at least one data card. Specifically, the first two data files consist of only one data card each, while the third file consists of N data cards. The information on each data file is explained as below:

First Data File

This data file defines four variables with symbolic names N, FREINM, NX, and NXB. Where N is the total number of circular rings inside the body. NX is the number of rings per layer, and NXB is the number of rings per layer within the range of the magnetic beam. Finally FREINM is simply the frequency of the impressed magnetic field in MHz.

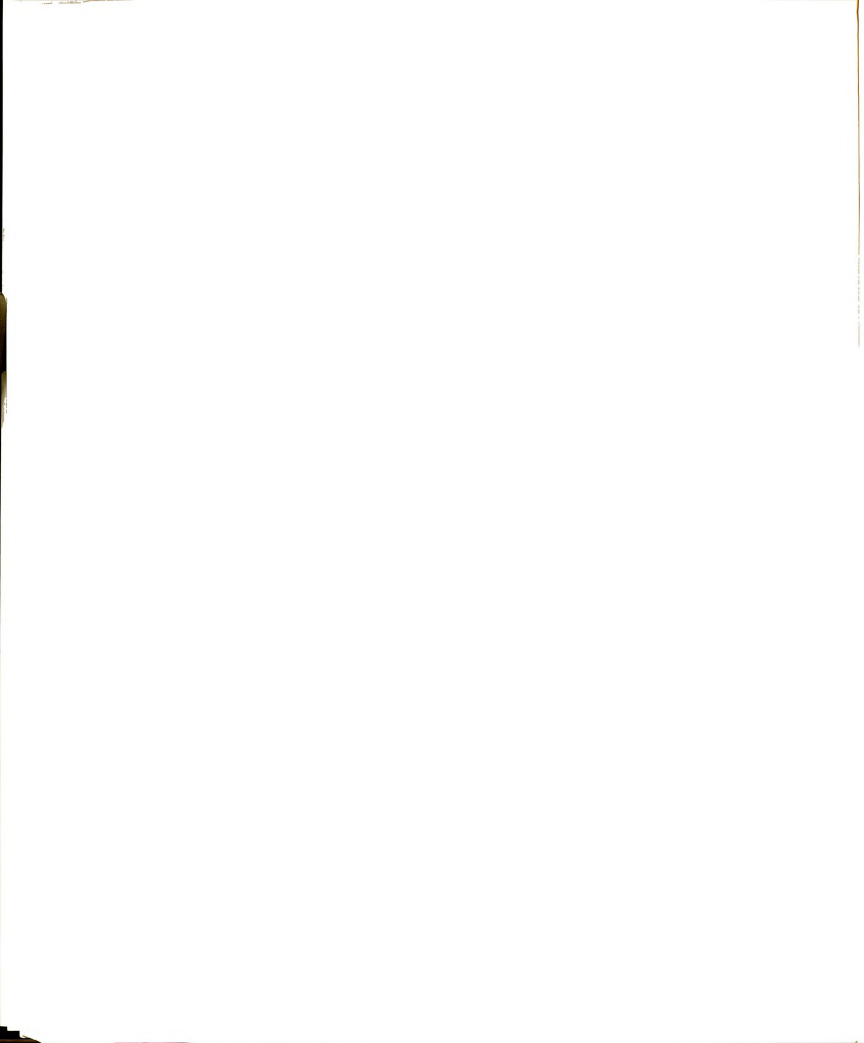
Second Data File

This file consists of only one data card and determines the following variables:

"NPAR"--for the integrals K_{mn} and K_{mm} , the integration limits 0° and 180° will be divided into "NPAR" equal-size subintervals. The purpose of this partitioning is to save the computational cost.

"AERR"--desired minimum absolute error for the numerical integration.

"RERR"--desired minimum relative error for the numerical integration.



Third Data File

This data file consists of N data cards, one for each circular ring. This set of data cards helps simulate the finite conducting body being considered. Each data card contains the following information:

"XEND," and "ZEND"--these codes correspond to the maximum boundaries of a ring cross-section in the x-, and the z-directions with reference to the origin. This information is supplied by the user in centimeters.

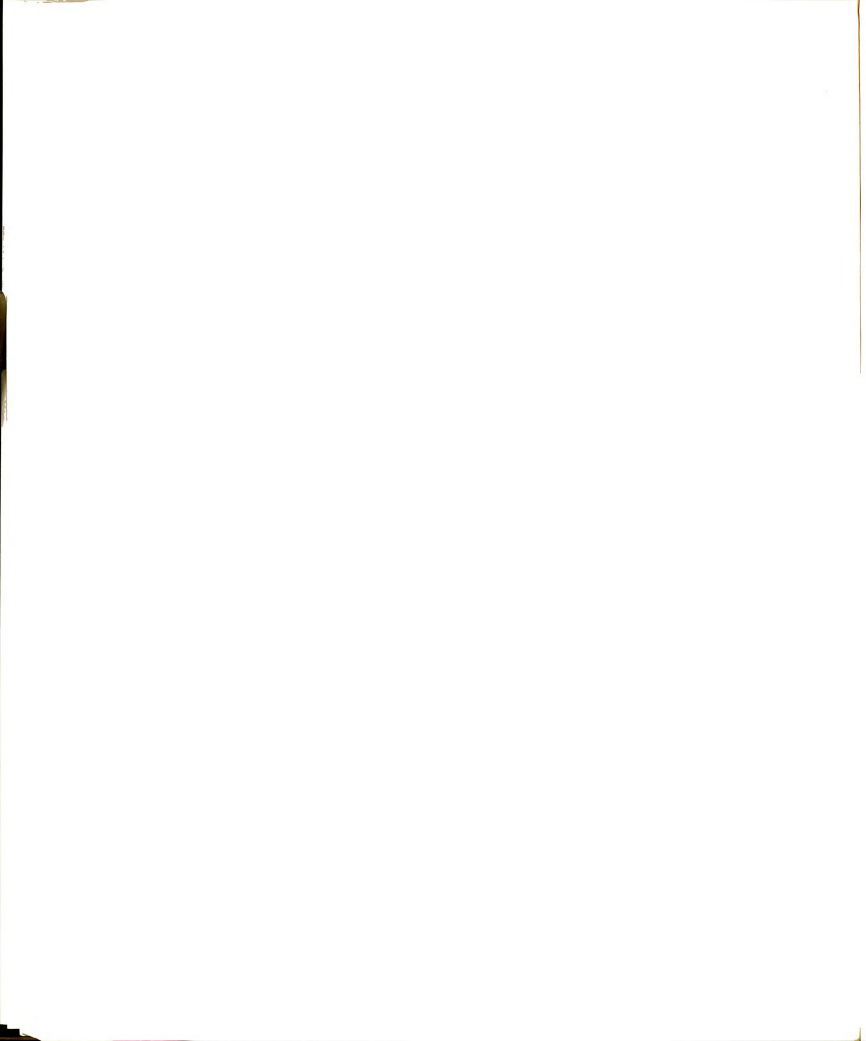
"XAA," and "ZBB"--are the symbolic names for the dimensions of the ring cross-section in the x-, and the z-directions. Since each ring cross-section is a square, the column for "XAA," and "ZBB" will contain the same information.

"REP," and "SIG"--are the codes for dielectric constant and conductivity (S/m) of each circular ring.

This completes the structure of the input data files needed to specify all the necessary information for the quantification of induced electric fields inside a finite conducting body with rotational symmetry. An example is worked out in the next section.

4.4. An Example to Use the Program

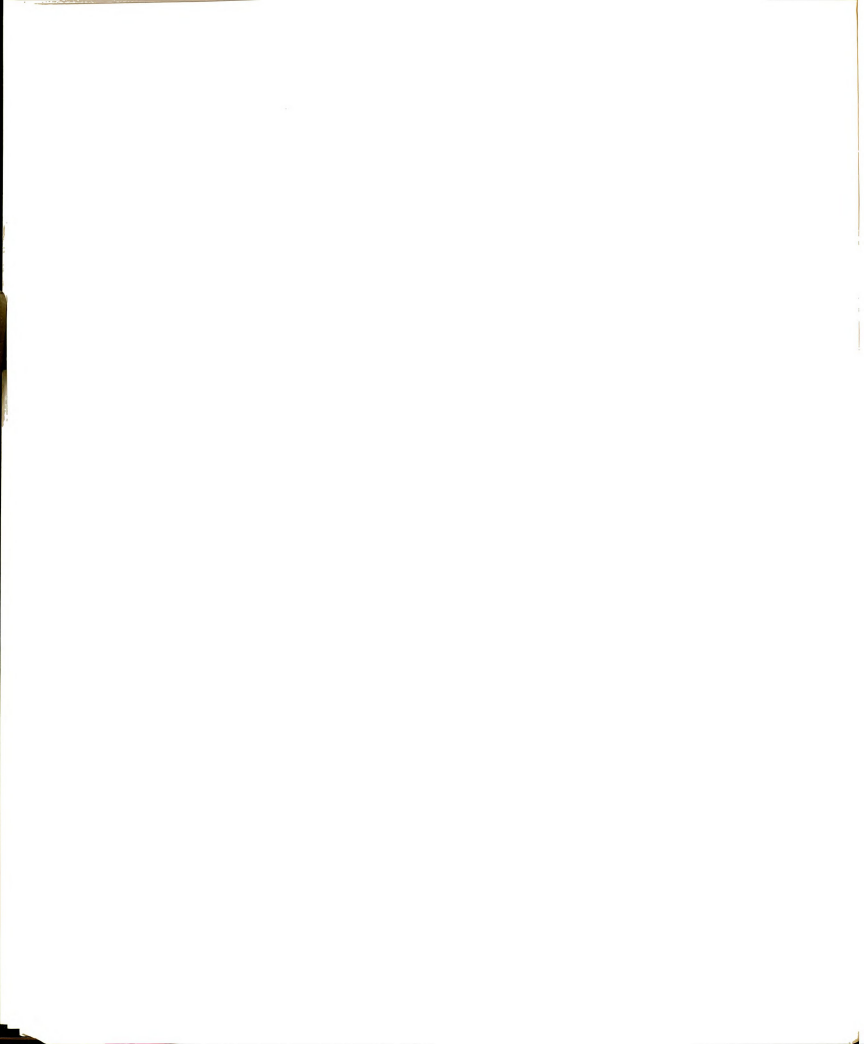
As an example, let us try to determine the electric field induced inside a circular cylinder, as shown in Figure 4.2. We assume that the body has radius 3 cm and



height 2 cm, and is immersed in a uniform magnetic field with an intensity of 1 A/m. Let us further assume that the frequency of the impressed magnetic field is 100 MHz, and the electrical parameters of the conducting body are $\epsilon_r = 71.7$, $\sigma = 0.889$ S/m, and with a ring cross-section of 1 cm x 1 cm. With the aids of Section 4.3 and Table 4.1, the sequential order of the input data files is as follows:

<u>File No.</u>	<u>Information on the File</u>						
1	6	100.0			3	3	
2	18	0.0			0.01		
3.1	1.0	1.0	1.0	1.0	71.7	0.88900E+00	
3.2	2.0	1.0	1.0	1.0	71.7	0.88900E+00	
3.3	3.0	1.0	1.0	1.0	71.7	0.88900E+00	
3.4	1.0	2.0	1.0	1.0	71.7	0.88900E+00	
3.5	2.0	2.0	1.0	1.0	71.7	0.88900E+00	
3.6	3.0	2.0	1.0	1.0	71.7	0.88900E+00	

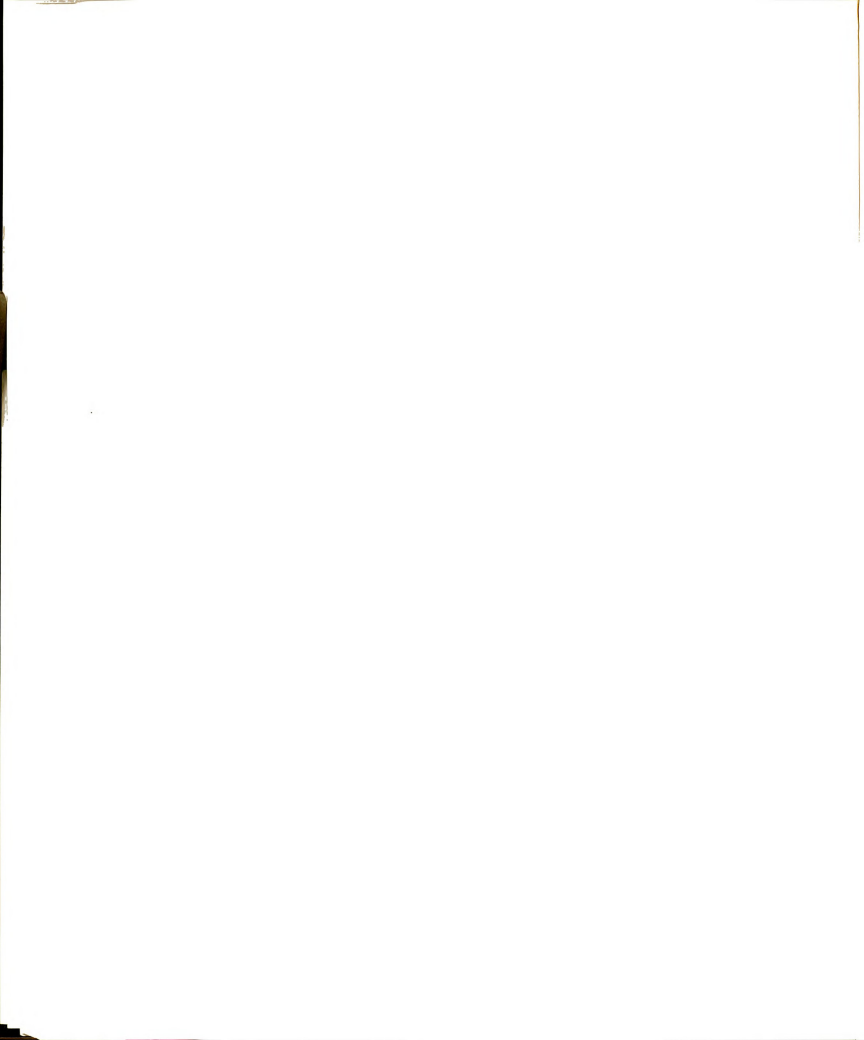
Now suppose that the program "EDDY" and the input data files are both stored on the magnetic disc under the permanent file names of "EDDYCURRENTW" and "EDDYCURRENTDATA," respectively. In this case, the list of the commands needed for the execution of the program when submitted through a terminal interactively is as follows:



<u>Statement No.</u>	<u>Its Purpose</u>	<u>Information in the Statement</u>
1.	Job request	*JOB CARD*, CM170000, RGL,T600,JC1000.
2.	Dispose to engineering terminal	DISPOSE,**,V.
3.	Identification name	HAL,BANNER,LEE, EDDY,CURRENT.
4.	Calling subroutine DCADRE	HAL,LGO=DCADRE, UERTST.
5.	Calling the PF	ATTACH,EWFIL, EDDYCURRENTW.
6.	Calling the EDITOR	EDITOR.
7.	Compile the program	FTN,I=Z,R,T.
8.	Return the EWFIL	RETURN,EWFIL.
9.	Calling the data	ATTACH,EWFIL, EDDYCURRENTDATA.
10.	Calling the EDITOR	EDITOR.
11.	Execute the program	LGO,W.
12.	End of Record	*EOR
13.	Change EWFIL to scope coded format	SAVE,Z.
14.	End of Record	*EOR
15.	Change EWFIL to scope coded format	SAVE,W,NS.
16.	End of file	*EOF

4.5. Printed Output

Now we would like to explain briefly the various sections of the output data files.



First, the coordinates of the maximum boundary limits on the x-z plane for each ring cross-section are listed, accompanied with the dimension and the cross-sectional area for each ring.

Then, the second output file contains the internally calculated coordinates for the central location of each ring cross-section, as well as the dielectric constant and the conductivity.

The third output file includes the information of the associated impressed electric field due to the impressed magnetic field of unit intensity. Also printed out is type of magnetic field applied.

The fourth output file has the information of the real and imaginary parts of the induced electric field inside each ring. Also included are the parameters used for numerical integrations, and the frequency of the impressed magnetic field.

The last output file is the induced electric field in magnitude and phase form, where the phase angles are in degrees.

An example print out of this program is shown in Table 4.2.

4.6. Listing of the Program

A FORTRAN listing of the program "EDDY" begins on page 175. Generally speaking, when compiled, the program "EDDY" requires about 46500B storage units of core memory.

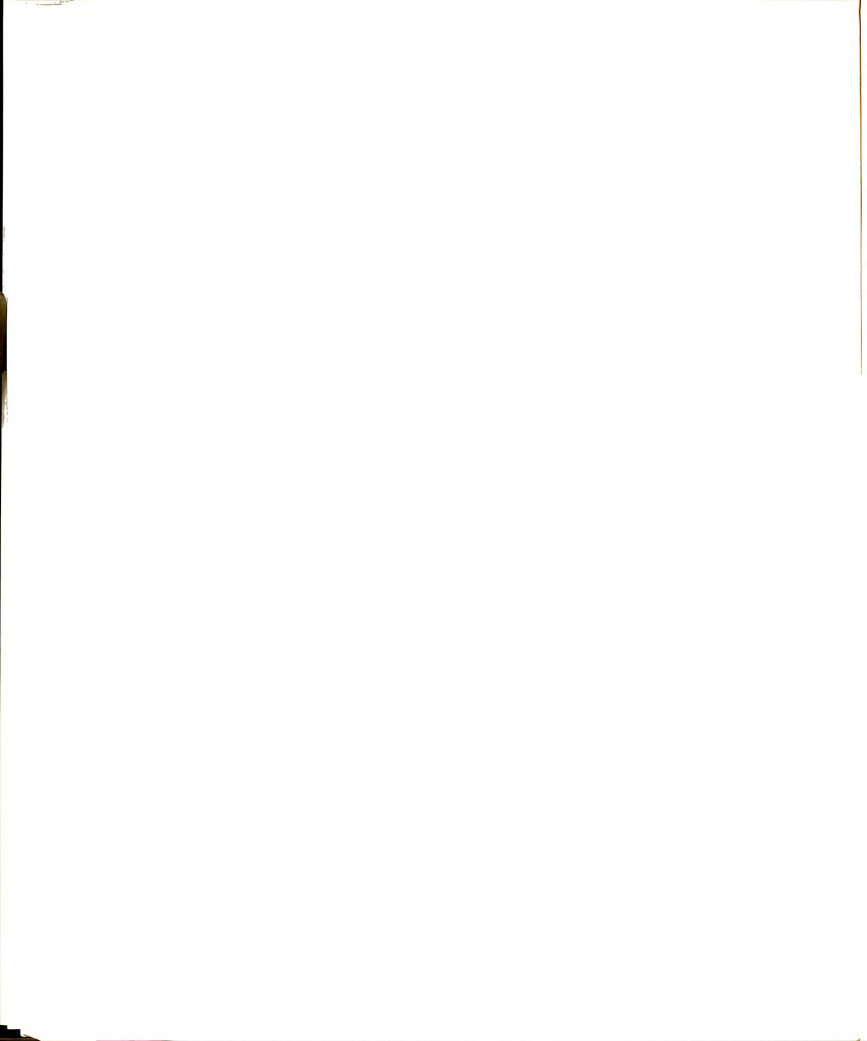


Table 4.2. Printed output of program "EDDY."

THE EXTENDING FUNCTIONS WILL AS CRO-SEC AREA OF EACH BITS ARE GIVEN BELOW IN WRTTMS					
N	XI*10	Z*10 ¹⁰	X-DIMENSION	Z-DIMENSION	CRO-SEC AREA
1	.11111111-01	.11111111-01	.11111111-01	.11111111-01	.11111111-01
2	.20000001-01	.10000001-01	.11111111-01	.11111111-01	.11111111-01
3	.31111111-01	.11111111-01	.11111111-01	.11111111-01	.11111111-01
4	.40000001-01	.11111111-01	.11111111-01	.11111111-01	.11111111-01
5	.21111111-01	.21111111-01	.11111111-01	.11111111-01	.11111111-01
6	.30000001-01	.21111111-01	.11111111-01	.11111111-01	.11111111-01

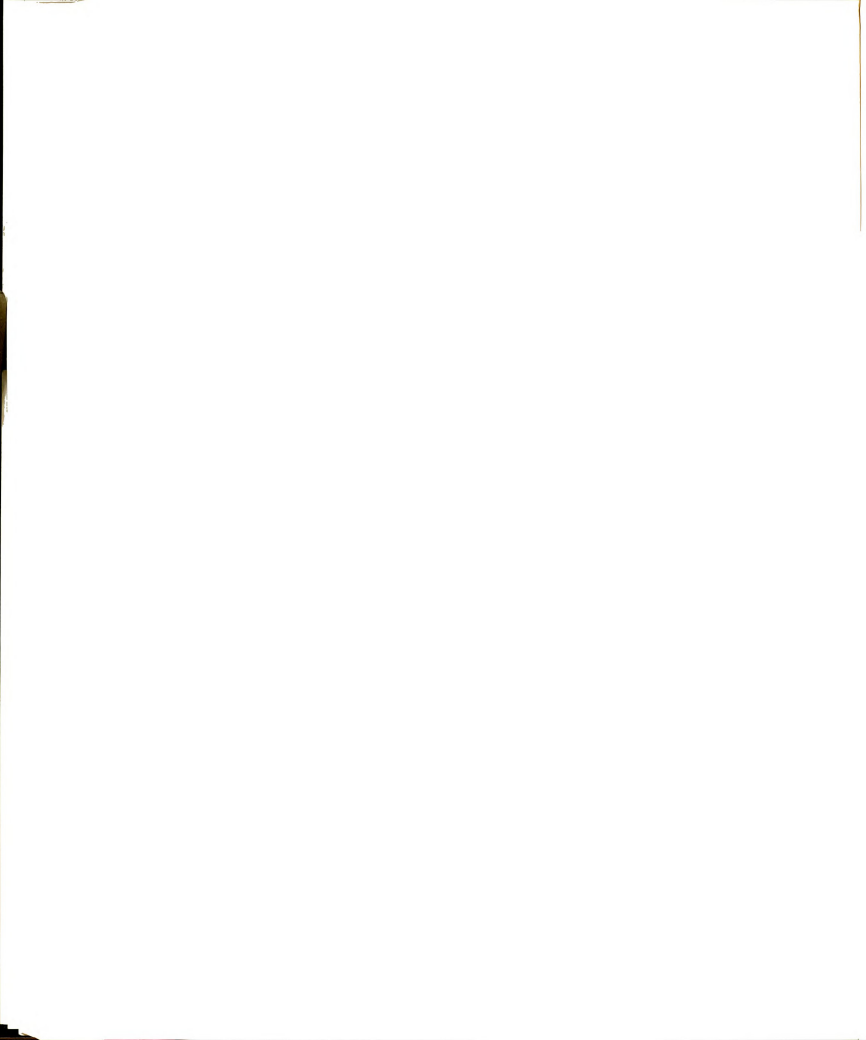


Table 4.2. (Cont'd.)

THE CENTRAL POINT LOCATION AS WELL AS ELECTRICAL PROPERTIES OF EACH SITE ARE GIVEN BELOW				
N	X (CM)	Y (CM)	DEPTH (CM)	CONDUCTIVITY (MH/CM)
1	500000-02	500000-02	217600-02	800000000
2	1510000-11	500000-02	217600-02	800000000
3	250000-11	500000-02	217600-02	800000000
4	500000-02	150000-01	217600-02	800000000
5	150000-11	150000-01	217600-02	800000000
6	250000-11	150000-01	217600-02	800000000

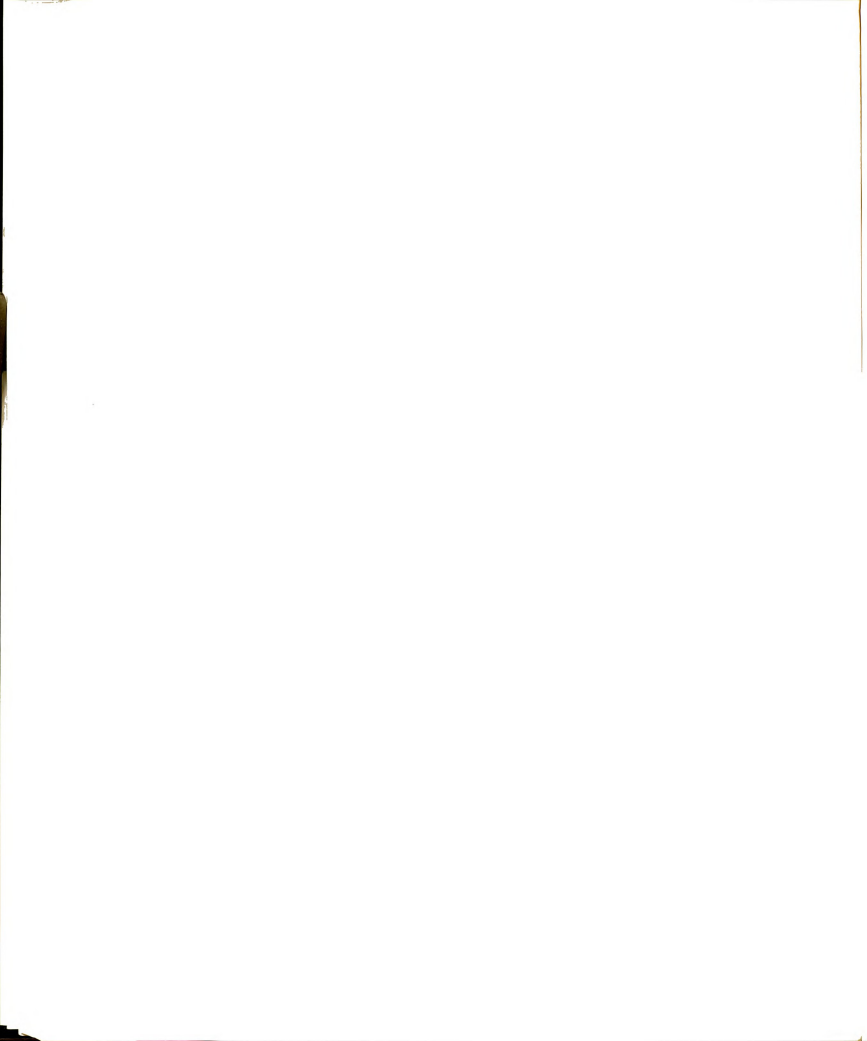


Table 4.2. (Cont'd.)

THE IMPRESS ELECTRIC FIELD DUE TO DIPOLE AND QUADRUPOLE FIELDS WITH UNIT INCIDENT ELECTRIC FIELD E^0 AS FOLLOWS			
N	IMPRESS ELECTRIC FIELD (GAUSS/UNIT)	REAL	IMAG
1	0.	0.	-0.53556E-02
2	0.	0.	-0.157187E-01
3	0.	0.	-0.341791E-01
4	0.	0.	-0.524906E-02
5	0.	0.	-0.157187E-01
6	0.	0.	-0.261938E-01

THE BODY IS IMMERSED IN AN UNIFORM ELECTRIC FIELD

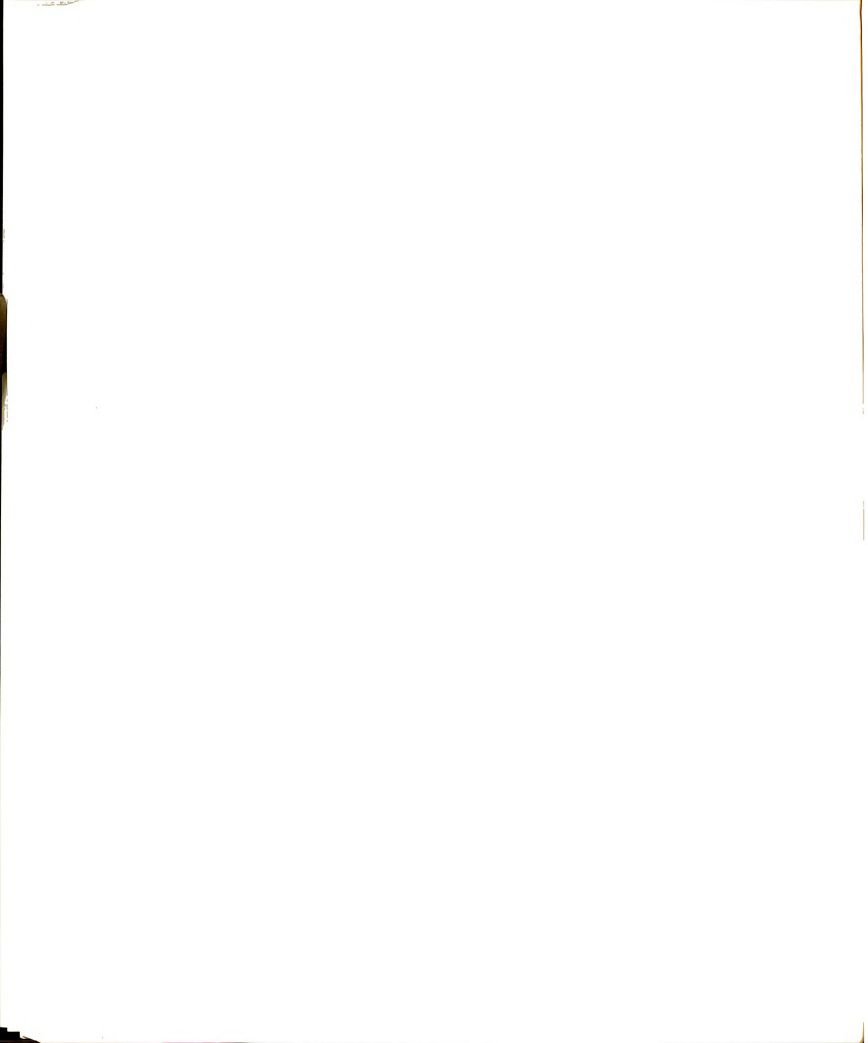


Table 4.2. (Cont'd.)

THE INTEGRATIONS INVOLVED WITHIN				THE ESTIMATED VALUES AT CENTRAL OF EACH SQUARE ARE GIVEN BELOW	
THE FREQUENCY BEING USED IS 10^{-10} HERTZ, ZERO TO BE BEING DIVIDED INTO 1000 INTERVALS, DIRECT NUMERICAL INTEGRATIONS				AND/OR	
THE INTEGRATIONS INVOLVED WITHIN				RELATIVE ERROR	
N	INDUCED-POL (VOLTS/UNIT)	INDUCED-POL (VOLTS/UNIT)	INDUCED-POL (VOLTS/UNIT)	INDUCED-POL (VOLTS/UNIT)	RELATIVE ERROR
1	-1.78467E-05	-1.78467E-05	-1.78467E-05	-1.78467E-05	-1.78467E-05
2	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02
3	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02
4	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02
5	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02
6	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02	-1.17806E-02

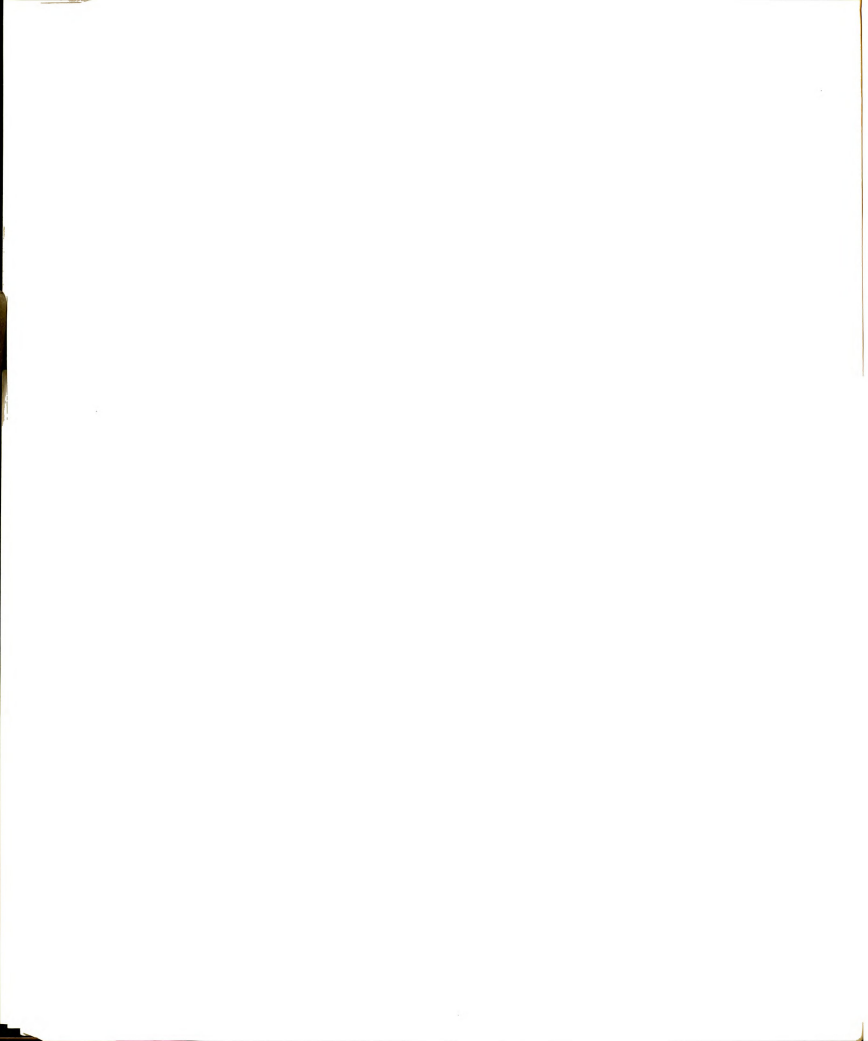
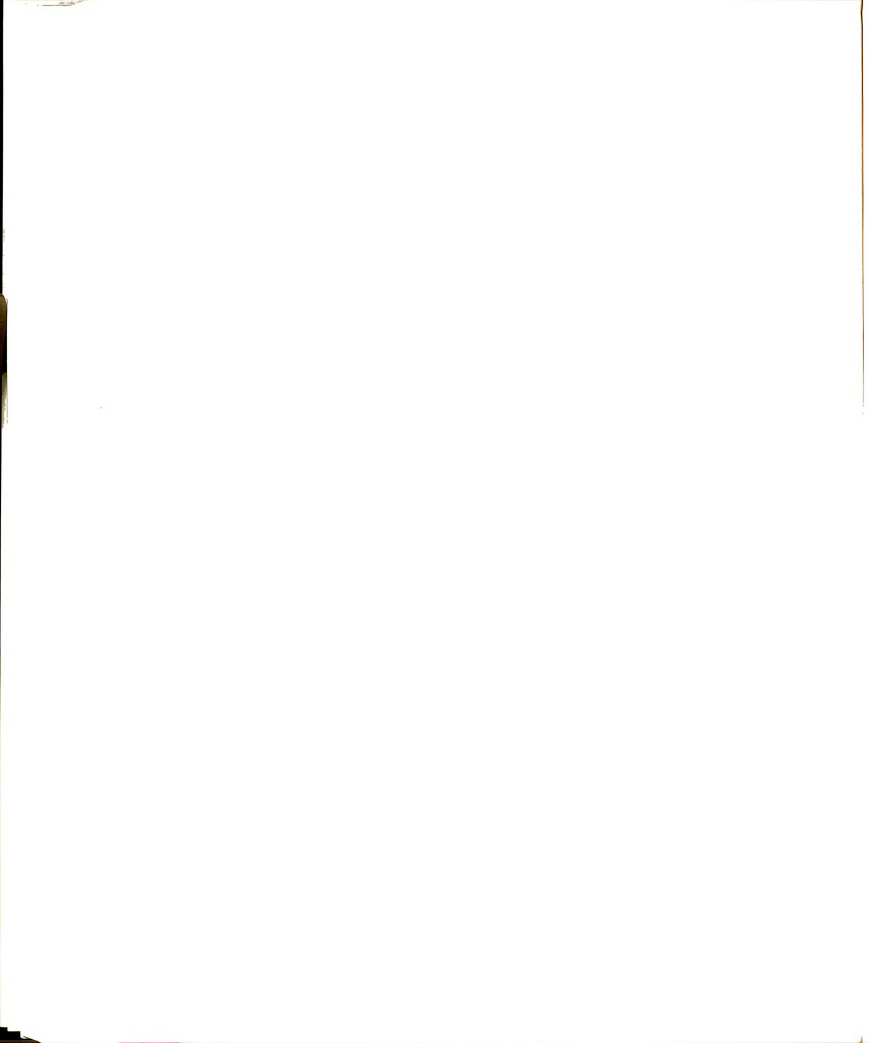


Table 4.2. (Cont'd.)

THE NEGATIVES AND POSITS OF PRODUCTS ARE GIVEN BELOW		
N	PRODUCTS	PRODUCTS (IN DEGREES)
1	2.536291×10^{-1}	-2.942191×10^{-2}
2	4.613461×10^{-1}	-2.757041×10^{-2}
3	2.66035×10^{-1}	-2.252401×10^{-2}
4	2.294291×10^{-1}	-2.34119×10^{-3}
5	2.111361×10^{-1}	-2.333781×10^{-3}
6	2.66035×10^{-1}	-2.252401×10^{-2}



PART 2

A USER'S GUIDE TO COMPUTER PROGRAM FOR INDUCED EM FIELD ON THE SURFACE OF A FINITE CONDUCTING BODY WITH ARBITRARY SHAPE

The purpose of Part 2 is to explain the computer program used for quantifying the EM field induced on the surface of a finite conducting body when illuminated by an incident plane wave. The theoretical derivation and the development of this numerical technique have been discussed in Chapter III. Besides explaining the usage of this program, we will also present an example accompanied with a sample print out for the purpose of better understanding.

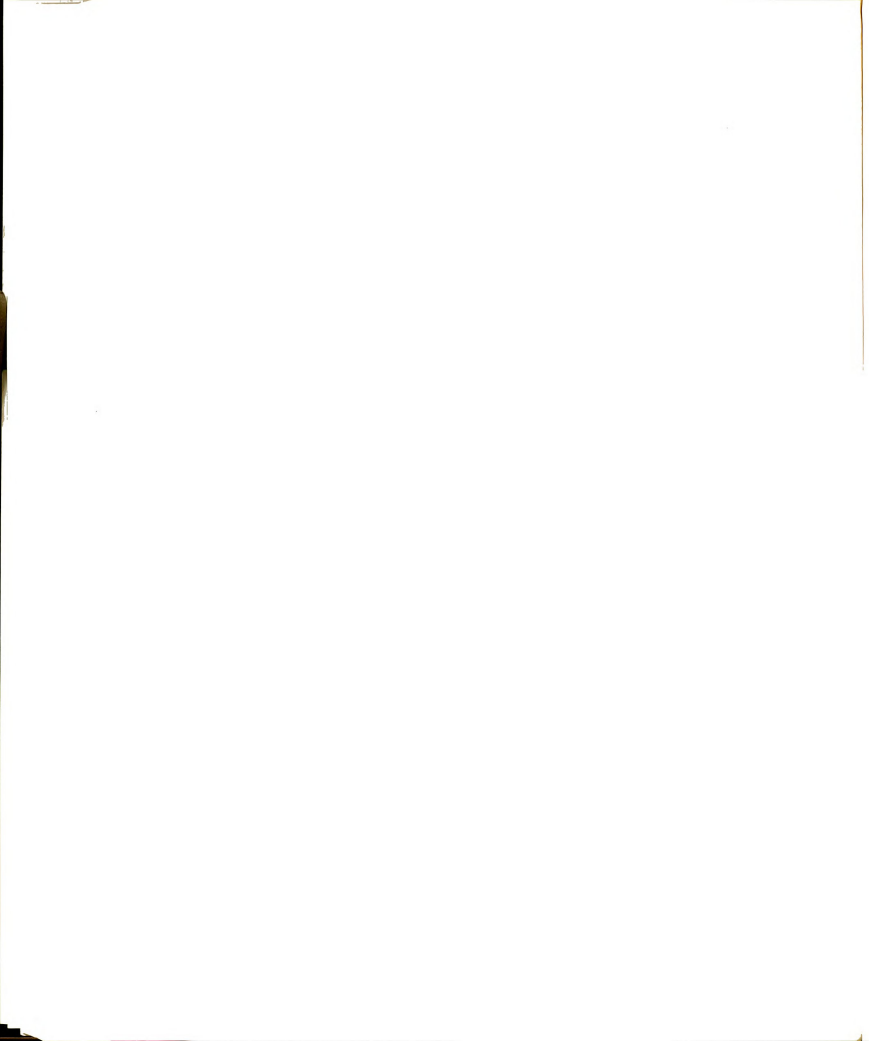
4.7. Formulation of the Problem

We are considering a finite conducting body illuminated by an incident wave which, for the purpose of simplicity, has been assumed to be a plane wave. Mathematically, the incident plane wave can be expressed in the following form:

$$\vec{E}^i(z) = \hat{x} e^{-jz} = \hat{x} [\cos(z) - j \sin(z)] \quad (4.1)$$

$$\vec{H}^i(z) = \hat{y} \frac{1}{\eta_0} e^{-jz} = \hat{y} \frac{1}{\eta_0} [\cos(z) - j \sin(z)]$$

We like to mention a few points concerning the above expressions:



- 1) The wave is propagating in the +z direction, and has a time factor $\exp(j\omega t)$ which is suppressed.
- 2) All quantities with the dimension of length are normalized by $1/k_0$.
- 3) The incident electric field has unit intensity, 1 V/m.
- 4) η_0 is the intrinsic wave impedance of free space.

The first step in utilizing this program is the numerical formulation of the problem; hence, the first thing to do is to visualize the shape, dimensions and the orientation of the conducting body with respect to the incident electromagnetic wave. The incident wave may illuminate the body either at normal incidence or at end-on incidence. To begin with, the body surface is divided into NCELL surface cells with NX, NY, NZ, NNX, NNY, and NNZ cells facing the positive-x, positive-y, positive-z, negative-x, negative-y, and negative-z directions, respectively. Each surface cell is a square with suitable dimension in order to obtain optimum results.

The maximum number of unknowns that can be handled is about 150. Since there are four unknowns for each surface cell, the maximum number of surface cells, without any simplifications, can not be greater than 40. This imposes a restraint on the physical size of the body. But, due to the symmetrical properties which exist in most of the bodies we are considering, it seems adequate to

apply these properties and reduce the number of unknowns by a factor of 4 or 8. It must be mentioned here that in order to apply symmetrical conditions the incident wave has to be decomposed into four different modes named COSE, COSH, SINE, and SINH, respectively, as can be understood from eq. (4.1). The above-mentioned symmetry conditions are justifiable since most of the bodies considered (either biological bodies or other finite conducting materials encountered in engineering) possess four or eight similar-looking segments and hereafter called quadrants, such that it is sufficient merely to calculate the induced electromagnetic field at each surface cell of the first quadrant and then convert them into the induced field for the rest of the body.

After determining the symmetry conditions, the next step in the numerical formulation of the problem is the specification of the location of each surface cell, its physical dimensions, and the electrical parameters of the body. The size of each cell can vary but the electrical properties are assumed to be uniform throughout the body. Note that the central location of each surface cell is predetermined by the user; then the incident field intensities are automatically evaluated for each cell.

4.8. Description of Computer Program

This program is also coded in standard FORTRAN and can be compiled on either FTN or MNF compilers. The program is symbolically named as "SURFLDS" with input and output formats on any undefined logic units in conjunction with four magnetic tapes, "TAPE1," TAPE2," "TAPE3," and "TAPE4." "TAPE1"- "TAPE4" are the names of the local files which are used to temporarily store the computed results due to COSE, COSH, SINE, and SINH components of incident field.

Program "SURFLDS" makes use of the following sub-programs:

"MNSB"--is a subroutine mainly used to print out the echo information which includes all the messages just read in from the input data files by the main program.

"MATRI," "LEE"--are subroutines which generate the matrix of the linear equations set. The matrix elements are related to the associated Green's functions evaluated at various locations.

"ELEMZ," "ELEMZ," and "ELEMZ"--are subroutines which evaluate the matrix elements when source points are in the $\pm x$ -, the $\pm y$ -, and the $\pm z$ -directed surfaces, respectively. In these

subroutines, the symmetry conditions have been imposed.

"EXIM," "CEIM," "CHIM," "SEIM," and "SHIM"--are subroutines which calculate the incident electric and magnetic fields for EXPZ, COSE, COSH, SINE, and SINH modes, respectively. In these subroutines, the incident electric field is polarized in the +x direction and has unit intensity, 1 V/m.

"T"--is a subroutine which converts the equivalent surface currents to surface field quantities.

"PRNT"--this subroutine, as the name suggests, is a program used to generate the print out.

"COEFX," "COEFY," "COEFZ," "FCNXX"--"FCNZZ," and "CEFFX," "--"CEFZZ," etc.--are subprograms which evaluate the elements of the Tensor Green's function.

"CMATP"--is the program being used to solve a system of N equations in N unknowns. It is actually a Gauss-Seidel method of numerical technique.

A listing of the program "SURFLDS," including all the subprograms called, is given at the end of this chapter. In the next section, the structure of the input data files as well as the associated input variables are explained in detail.

4.9. Data Structure and Input Variables

Figure 4.3 shows a sample body with finite conductivity. With origin of the coordinate system being chosen at the center of the body, we divide it into eight different sections which are called quadrants. The numbering system used is also shown in Figure 4.3. Note that the symmetry conditions exist if the physical dimensions of each surface cell in the first quadrant are the same as their counterparts in other quadrants, which is the case for this sample body because we assume that all the surface cells have the same physical dimensions. Based on these existing symmetry conditions, we only intend to solve for the induced electromagnetic fields at various surface cells of the first quadrant and then convert them into the fields induced at the surface cells on the rest of the body. Now we like to introduce the input data files before we go any further in determining the induced surface field. First of all, the sequential structure of the data files, the format specifications and the symbolic names of the variables appearing on each file are outlined in Table 4.3.

There are totally five input data files. Only the fourth data file contains NCELL data cards, while the rest of the data files contain one data card each. The input variables associated with these input data files are discussed in detail as below:

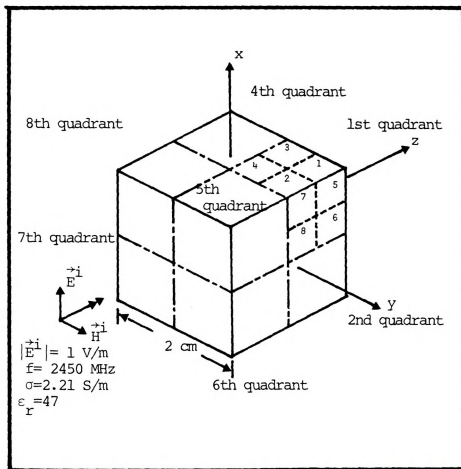


Figure 4.3. A cubic biological body ($\sigma = 2.21 \text{ S/m}$, $\epsilon_r = 47$) with dimensions of $2 \text{ cm} \times 2 \text{ cm} \times 2 \text{ cm}$ is irradiated by an incident plane wave which has a frequency of 2.45 GHz and a unit intensity, x-polarized electric field. The body surface is subdivided into 95 surface cells with dimensions of $0.5 \text{ cm} \times 0.5 \text{ cm}$.

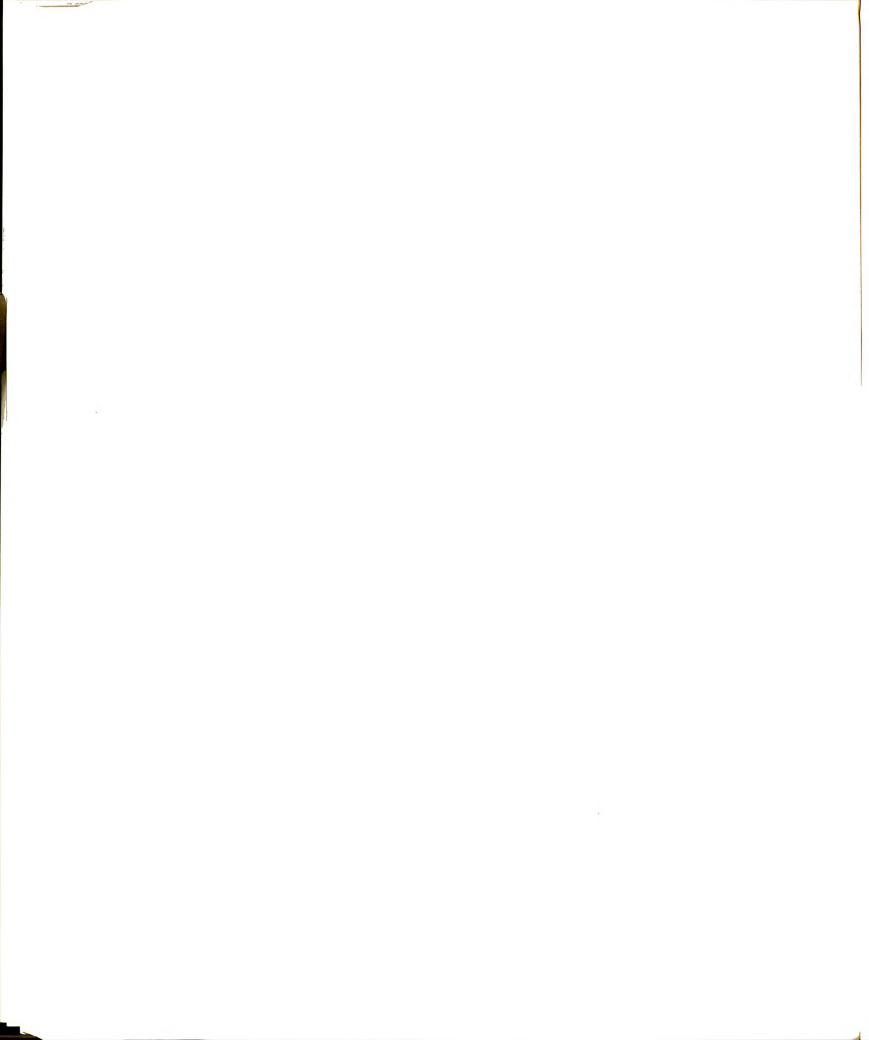
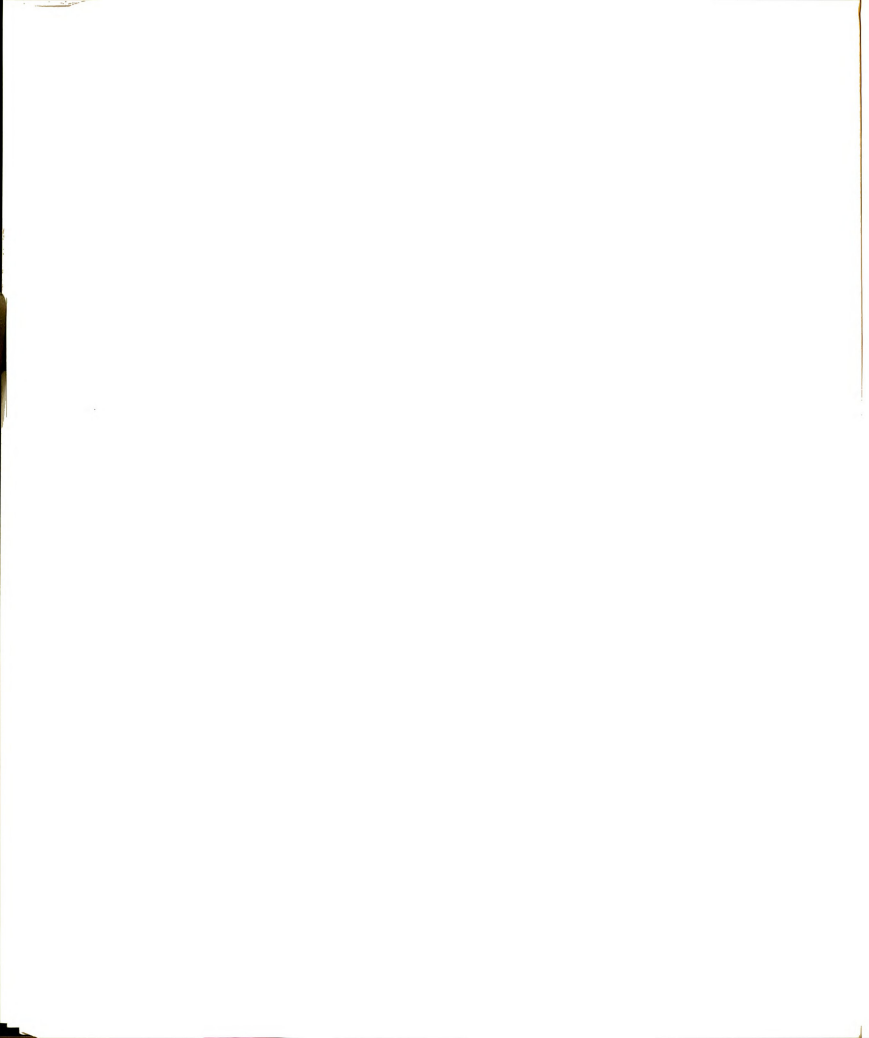


Table 4.3. The symbolic names of input variables and corresponding specifications for the data files used in data structure for the program "SURFLDS."

File No.	Card No.	Symbolic Name	Columns	Format
1	1	NDIV	1-2	I2
2	1	IQ	1-4	I4
3	1	NX	1-3	I3
		NY	4-6	I3
		NZ	7-9	I3
		NNX	10-12	I3
		NNY	13-15	I3
		NNZ	16-18	I3
4	1-NCELL	X	1-10	F10.5
		Y	11-20	F10.5
		Z	21-30	F10.5
		DCELL	31-40	F10.5
5	1	RLEP	1-13	E13.6
		SIG	14-26	E13.6
		FMEG	27-39	E13.6
		MODE	40-43	A4



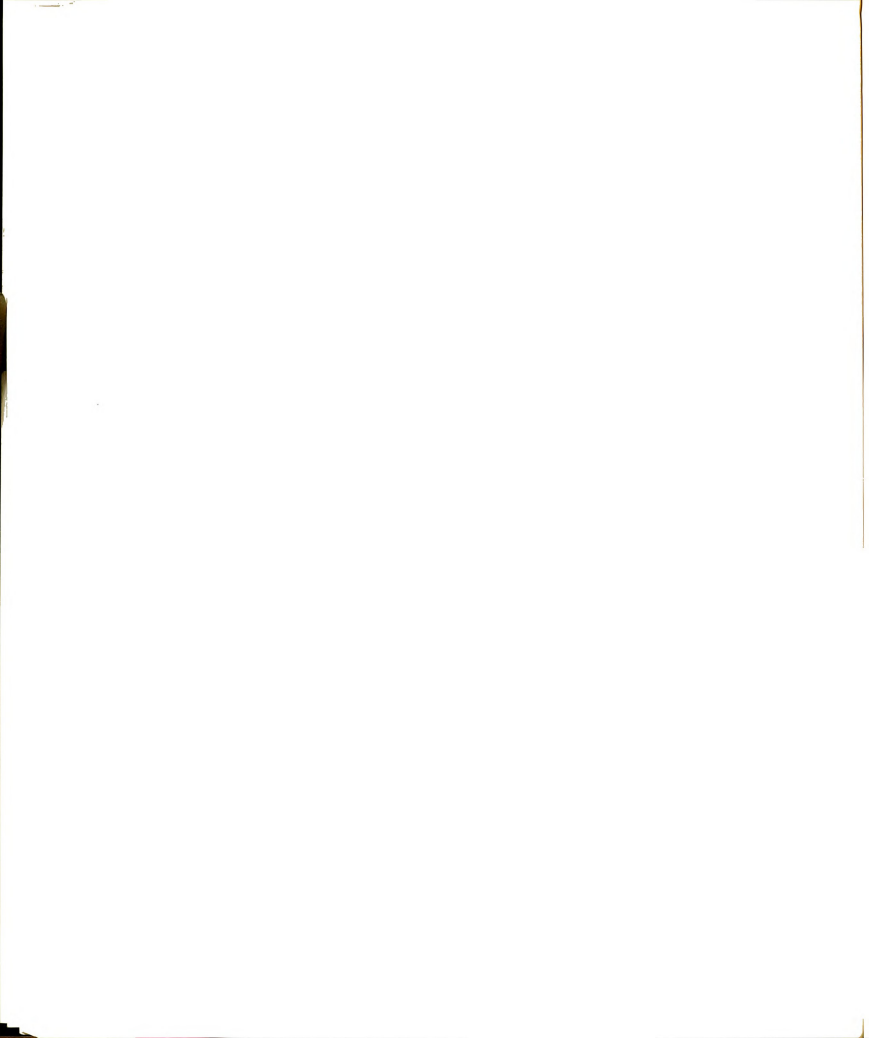
First Data File

This file has only one data card with only one variable on it; namely, NDIV. NDIV is the number of subdivisions each side of the square surface cell will be divided into when the numerical integration over the surface cell is performed. The purpose of subdividing the surface cell is to improve the accuracy of the results. But the computational cost increases substantially when the value of NDIV goes up. Usually, we find that NDIV with value 3 gives excellent results and still keeps the required CPU time reasonably low. For the case of NDIV=3, each surface cell is subdivided into nine subsurface cells when numerical integration is carried out by the computer.

Second Data File

This data file has only one data card which defines a variable with symbolic name IQ. IQ is a four-digit code of the symmetrical property used. There are four different types of symmetrical property that can be handled by this program, as described in the following:

- IQ = 0008 This is for the case of an eight-quadrants symmetrical body.
- IQ = 1234 This is used when the body has symmetry with respect to the y-z plane and the x-z plane.



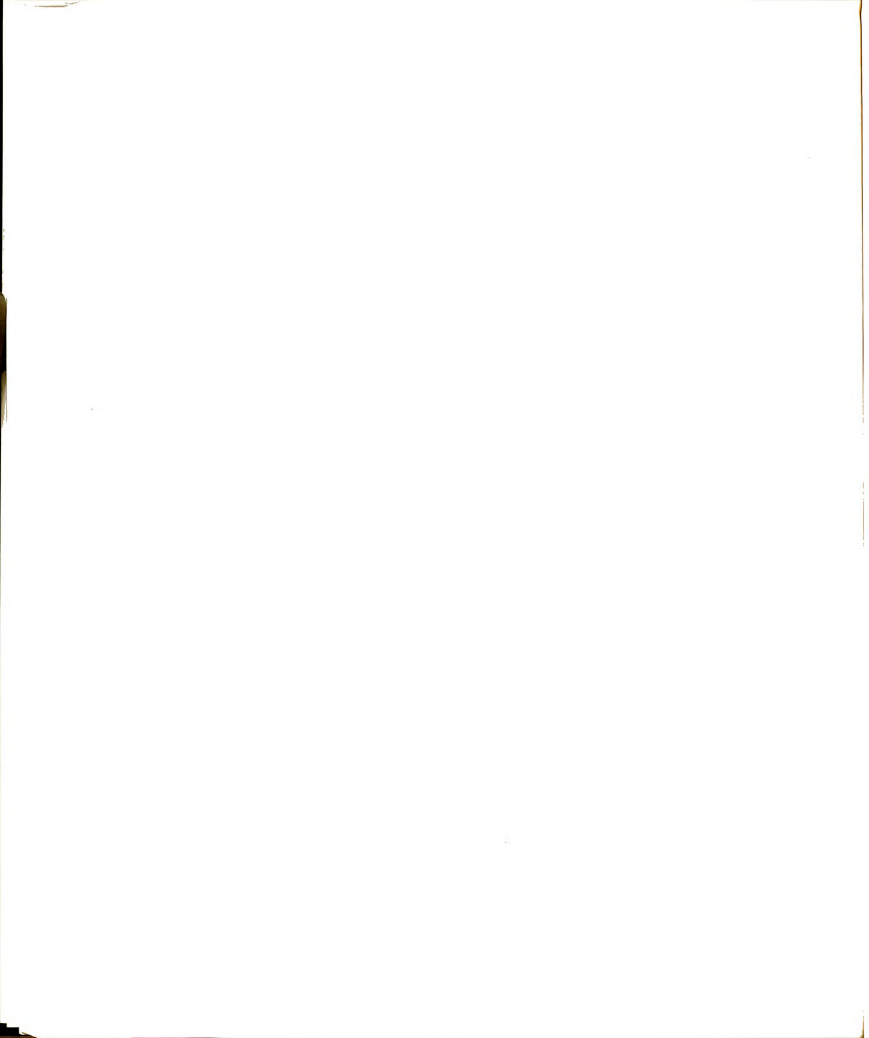
- IQ = 1458 This is used when the body considered has symmetry with respect to the x-y plane and the x-z plane.
- IQ = 1256 This is used when symmetry properties with respect to the x-y plane and the y-z plane exist.

Third Data File

This file has only one data card which defines six integer variables, NX, NY, NZ, NNX, NNY, and NNZ. These six integers are the numbers of surface cells on the +x, the +y, the +z, the -x, the -y, and the -z-directed surfaces of the first quadrant, respectively. The total number of surface cells of the first quadrant, NCELL, is obtained by summing up all these six integers.

Fourth Data File

This data file consists of NCELL data cards; i.e., one data card for each surface cell. On each data card, it defines four real variables, X, Y, Z, and DCELL. Where X, Y, Z are the coordinates of the central location for each corresponding surface cell with reference to the origin, while DCELL is the dimension of the corresponding surface cell. Note that these quantities are in centimeters.



Fifth Data File

This data file consists of one data card with four variables defined on it. The first three variables, RLEP, SIG, and FMEG, are specified under E13.6 format. RLEP is the symbolic name for dielectric constant, and SIG is that for conductivity in S/m, while FMEG stands for frequency in terms of MHz. Finally, the last variable of this last data file is symbolically named MODE, which is read in under format A4 and may have one of the following values:

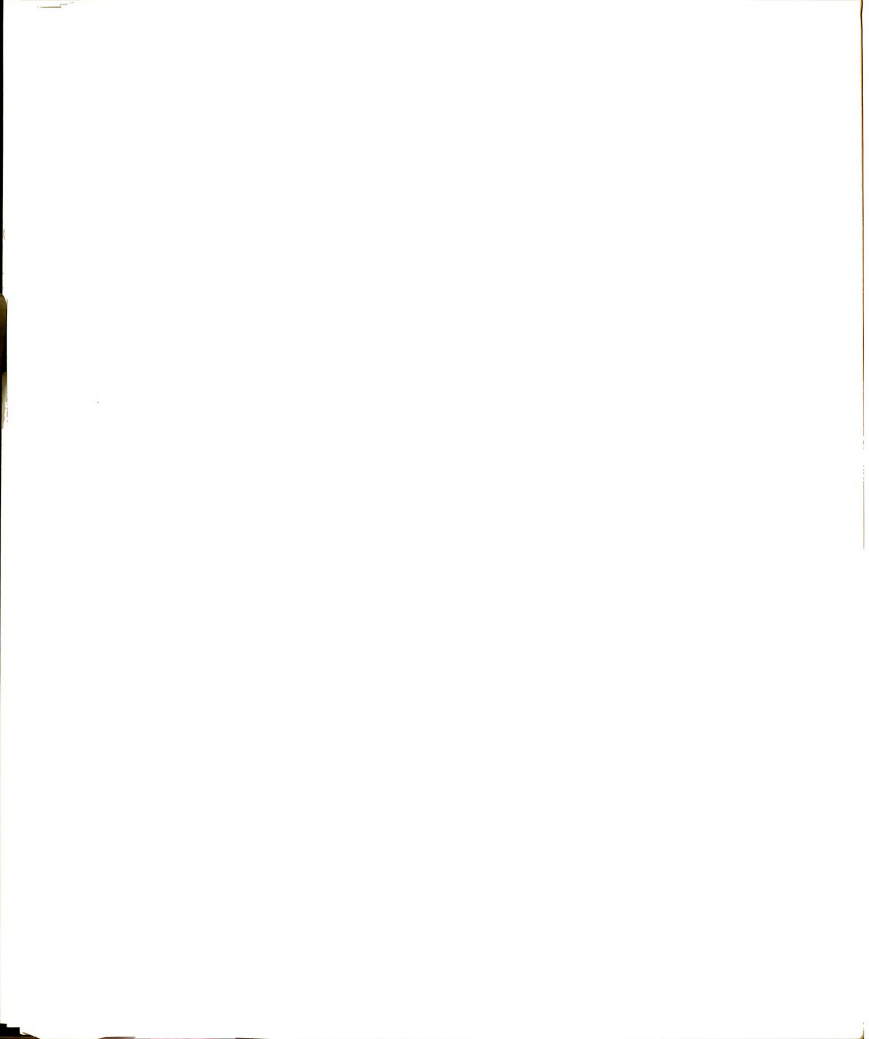
"EXPZ"--is a four-letter code for plane wave which has complex exponential variation in z direction.

"COSE"--is the code for incident field which has only the x-component of E field with cosine variation in the z direction.

"COSH"--is the code for incident field which has only the y-component of H field with cosine variation in the z direction.

Similarly, "SINE" ("SINH")--is the code for incident field which has only the x-component (the y-component) of E (H) field with sine variation in the z direction.

Obviously, the effect of EXPZ is equivalent to the combined results due to COSE, COSH, SINE, and SINH. As we have mentioned, the reason of dividing a plane wave into

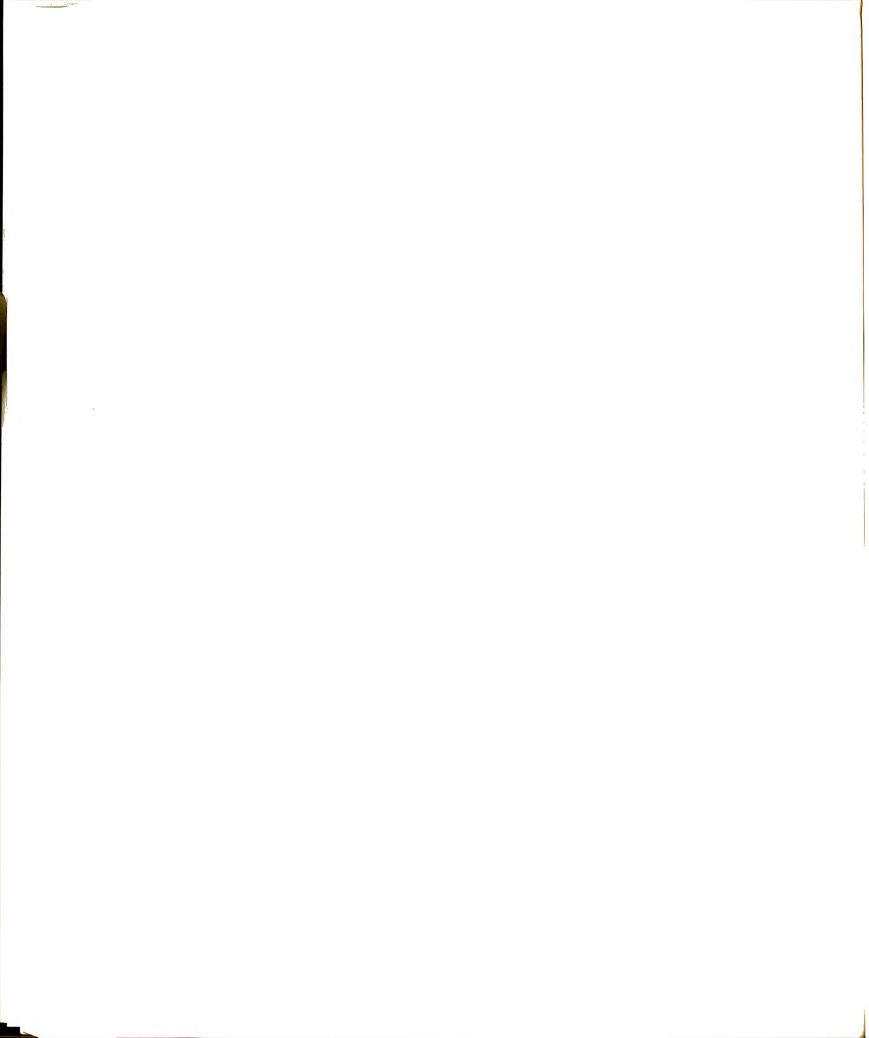


four different modes is for the purpose of applying symmetry conditions to a symmetrical body.

We have explained all the details about the input data files in this section. Now we are ready to use "SURFLDS" to solve for the EM field induced on the surface of an arbitrary body. An example is worked out in the next section.

4.10. An Example to Use the Program

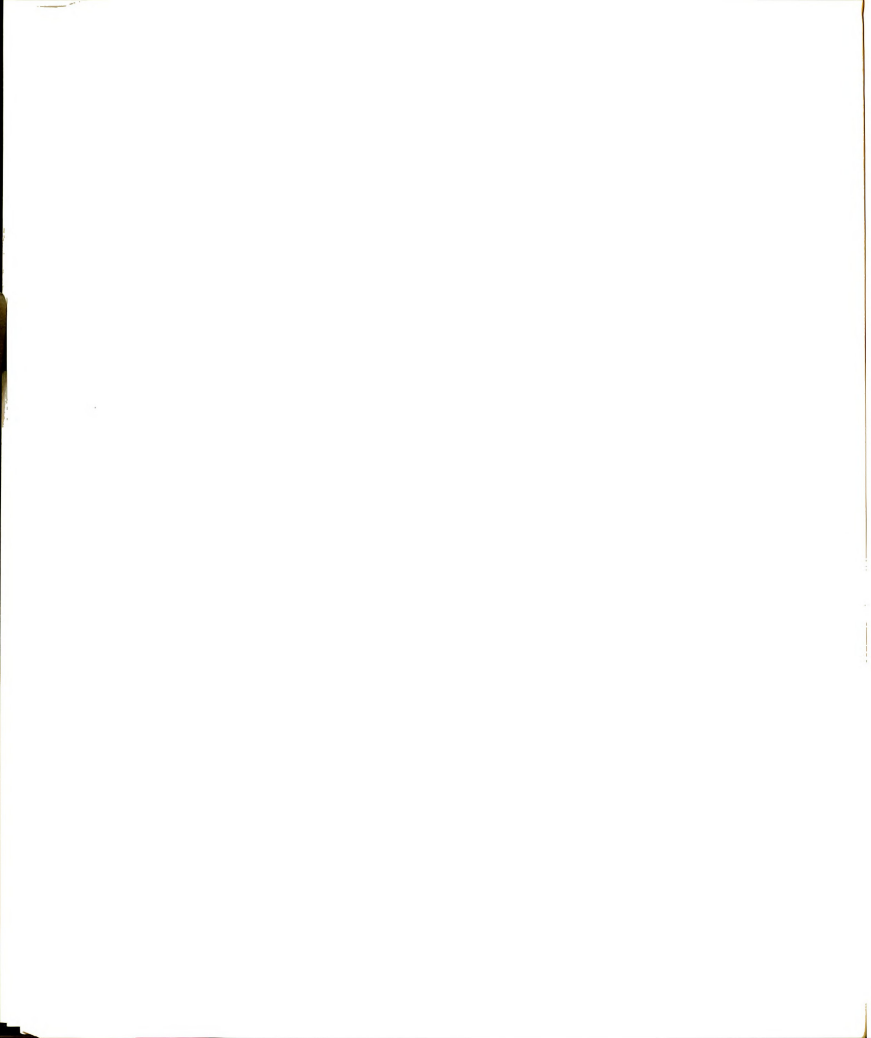
As an example about the usage of the program "SURFLDS," let us again consider the case of a cube illuminated by a plane wave as shown in Figure 4.3. We assume that the cube has dimensions 2 cm x 2 cm x 2 cm, and is constructed with a material with the following electrical parameters at the frequency of 2.45 GHz: $\sigma = 2.21 \text{ S/m}$, $\epsilon_r = 47$. The cube has eight-fold symmetry, such that it can be divided into eight different quadrants. Although the same problem can be solved by dividing the body into four quadrants, for the purpose of demonstrating the logic of decomposing the plane wave into four separate modes, we apply eight-quadrant symmetry conditions to this body. We further divide each surface of the first quadrant into four surface cells, such that there are totally 12 surface cells (or totally 48 surface unknowns) to be considered, where each of these surface cells has dimensions of 0.5 cm x 0.5 cm.



As mentioned previously, since the incident plane wave is decomposed into COSE, COSH, SINE, and SINH modes, we have to run the program in four parts--one for each constituting mode of the incident wave. The results due to each incident mode will be buffered into a logical unit with one of the following logical names, "TAPE1," "TAPE2," "TAPE3," and "TAPE4." And then the results contained within these units will be appropriately combined in order to obtain the true EM fields induced on the surfaces of the first and the fifth quadrants.

With the help of Section 4.9 and Table 4.3, we can easily list the input data files for the case of COSE incident mode as follows:

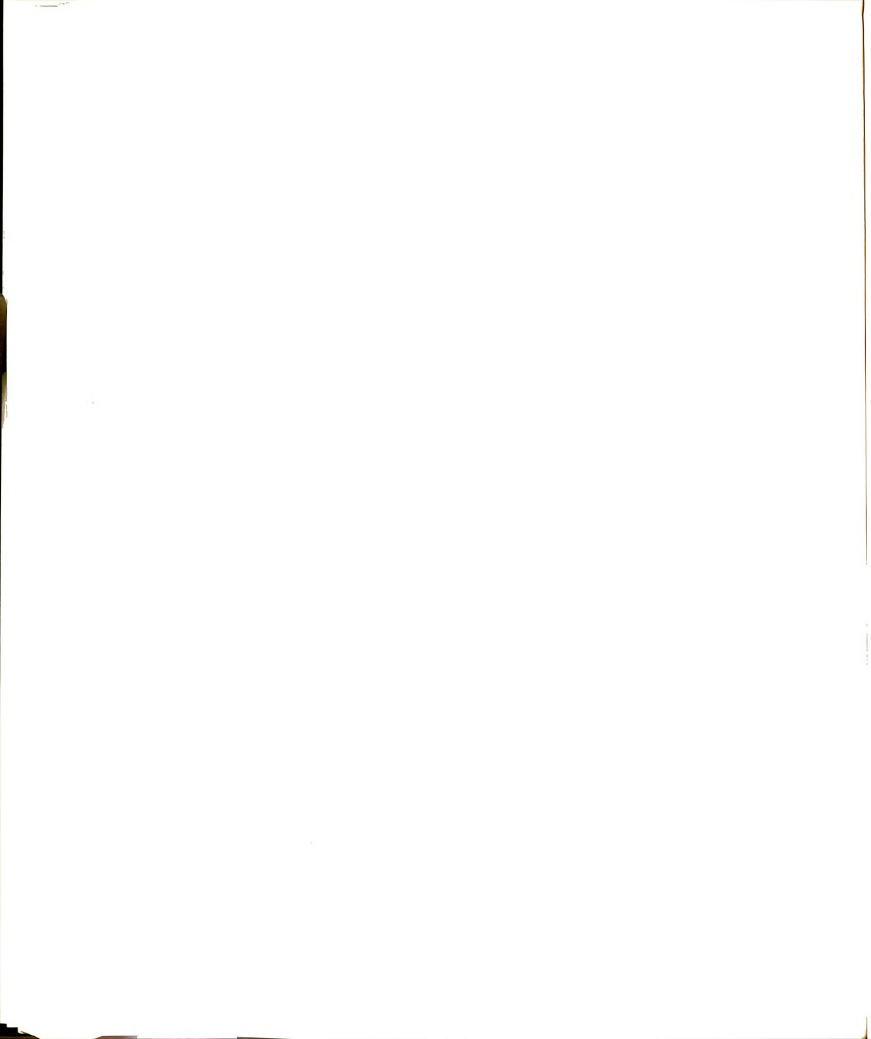
<u>File No.</u>	<u>Information on the File</u>					
1	03					
2	0008					
3	004	004	004	000	000	000
4.1	1.0	0.75	0.75	0.5		
4.2	1.0	0.75	0.25	0.5		
4.3	1.0	0.25	0.75	0.5		
4.4	1.0	0.25	0.25	0.5		
4.5	0.75	1.0	0.75	0.5		
4.6	0.25	1.0	0.75	0.5		
4.7	0.75	1.0	0.25	0.5		
4.8	0.25	1.0	0.25	0.5		



4.9	0.75	0.75	1.0	0.5
4.10	0.25	0.75	1.0	0.5
4.11	0.75	0.25	1.0	0.5
4.12	0.25	0.25	1.0	0.5
5	47.000000E+00	2.210000E+00	2.450000E+03	COSE

Now we assume that both the program "SURLDS" and the input data files are already stored in different permanent files with "SURFACEEW" and "SURFACEDATACOSE" as permanent file names, respectively. In this case, the commands needed in order to submit the job through the interactive terminals should be as follows:

<u>Statement No.</u>	<u>Information in the Statement</u>
1.	*JOB CARD*, RGL, JC1000, T600, CM170000.
2.	DISPOSE, **, V.
3.	HAL, BANNER, LEE, SURFLDS, 8SYMM, COSE.
4.	ATTACH, EWFILE, SURFACEEW.
5.	EDITOR.
6.	ATTACH, A, SURFACEB.
7.	FTN, I=Z, R, T.
8.	COPYL, A, LGO, B.
9.	PURGE, A.
10.	CATALOG, B, SURFACEB, RP=999.
11.	RETURN, EWFILE.
12.	ATTACH, EWFILE, SURFACEDATACOSE.
13.	EDITOR.

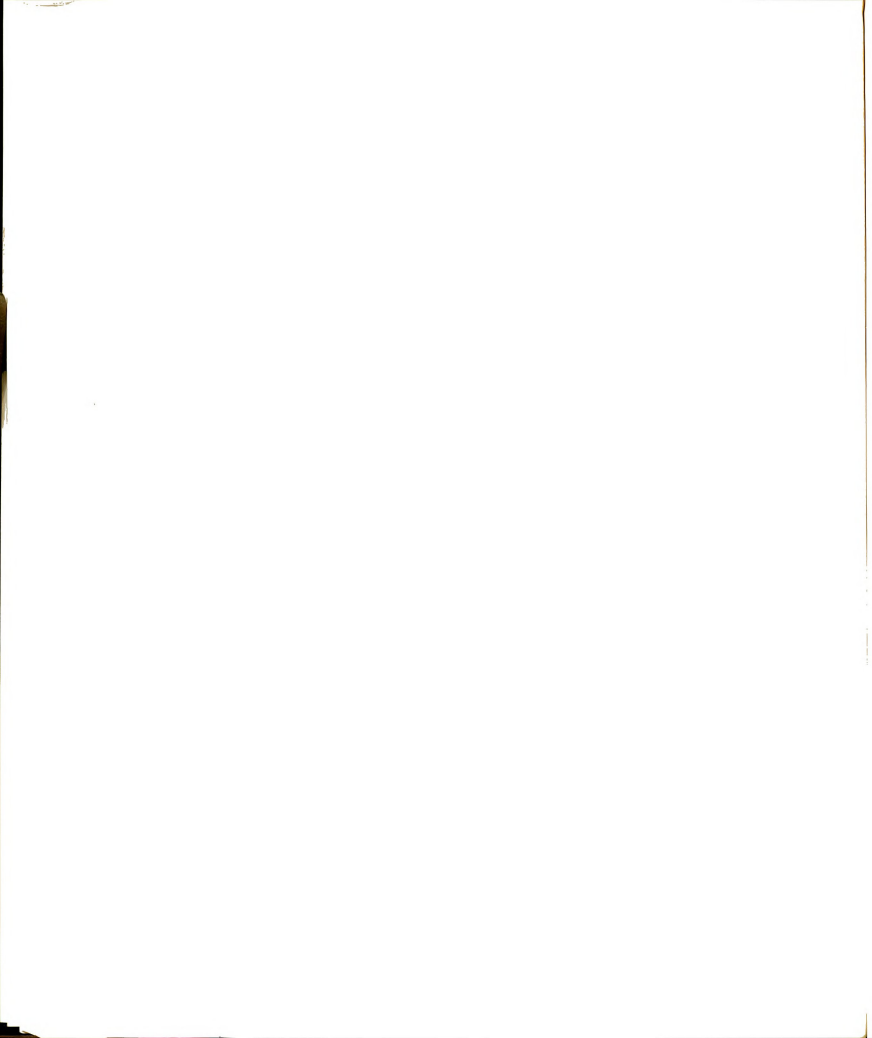


14.	B,W.
15.	CATALOG,TAPE1,SURFACEDUETOCOSE.
16.	*EOR
17.	SAVE,Z,100-780.
18.	*EOR
19.	SAVE,W,NS.
20.	*EOF

There are several points we like to mention about the above job:

- 1) We assume that the program "SURFLDS" has been compiled, and the compiled binary codes are already stored in a file with permanent file name "SURFACEB."
- 2) We don't have to recompile the whole program each time we use it. If the array dimensions have been changed, the only section that needs to be recompiled is the main program.
- 3) The computed induced surface fields are buffered out to TAPE1, and then cataloged as a permanent file named "SURFACEDUETOCOSE."

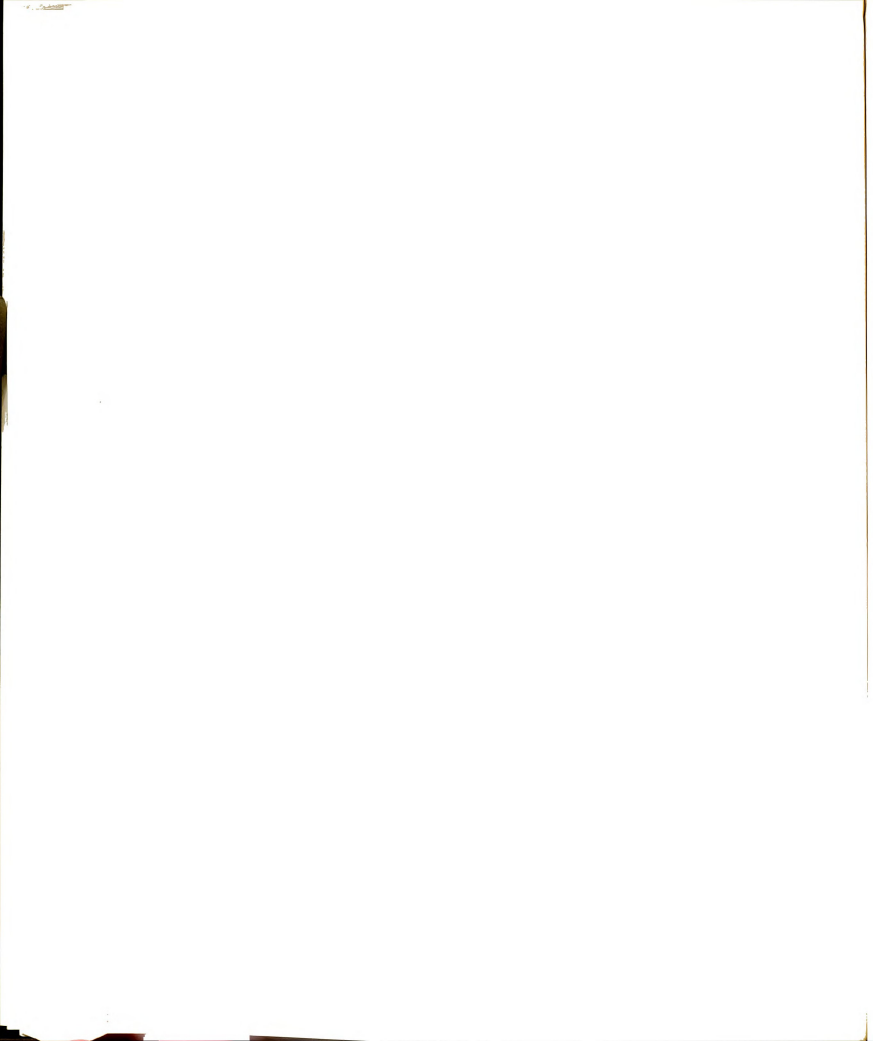
So the results due to COSE mode of incident field have been obtained and stored into "SURFACEDUETOCOSE." To obtain the results due to the other three incident modes, the same program has to be rerun three more times with the only change in the last data file where the correct code



for the incident mode must be used. For example, when COSH mode is considered, the program can be executed without any recompilation by the use of the following commands:

<u>Statement No.</u>	<u>Information in the Statement</u>
1.	*JOB CARD*,RG1,JC1000,T600,CM170000.
2.	DIPOSE,**,V.
3.	HAL,BANNER,LEE,SURFLDS,8SYMM,COSH.
4.	ATTACH,A,SURFACEB.
5.	ATTACH,EWFILE,SURFACEDATACOSH.
6.	EDITOR.
7.	A,W.
8.	CATALOG,TAPE2,SURFACEDUETOCOSH.
9.	*EOR
10.	SAVE,W,NS.
11.	*EOF

After obtaining results for all four different modes and creating four different permanent files for them separately, we are now in a position to combine these results in such a way so as to yield the total induced surface fields on the body surfaces. Another program called "COMBINE" has been written for this purpose. Program "COMBINE" has output including the magnitude and phase information of the total surface fields in quadrants 1 and 5. The surface fields on the rest of the body, either

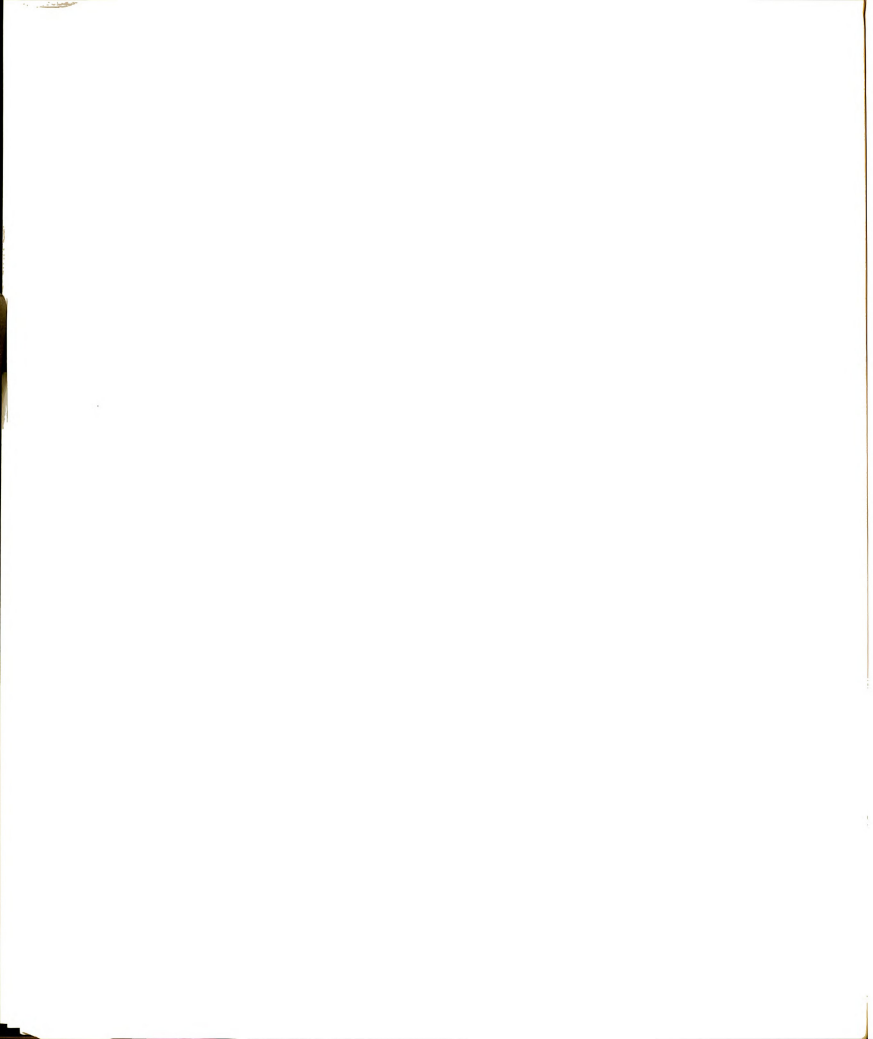


magnitude or phase, can be deduced from the information of the first or the fifth quadrant by intuition.

The data files for "COMBINE" are the same as those for program "SURFLDS," except that the variable MODE of the last data file for "SURFLDS" is redundant here. Hence the input data files for "SURFLDS" can be used for "COMBINE" without any modifications.

The commands needed in order to submit the program "COMBINE" interactively, assuming that both the program and the data files have been cataloged into permanent files, will be as follows:

<u>Statement No.</u>	<u>Information in the Statement</u>
1.	*JOB CARD*, RGL, JC500, T100.
2.	DISPOSE, **, V.
3.	HAL, BANNER, LEE, SURFLDS, 8SYMM, COMBINE.
4.	ATTACH, EWFILE, COMBINEEW.
5.	EDITOR.
6.	ATTACH, A, COMBINEB.
7.	FTN, I=Z, R, T.
8.	COPYL, A, LGO, B.
9.	PURGE, A.
10.	CATALOG, B, COMBINEB, RP=999.
11.	RETURN, EWFILE.
12.	ATTACH, TAPE1, SURFACEDUETOCOSE.
13.	ATTACH, TAPE2, SURFACEDUETOCOSH.

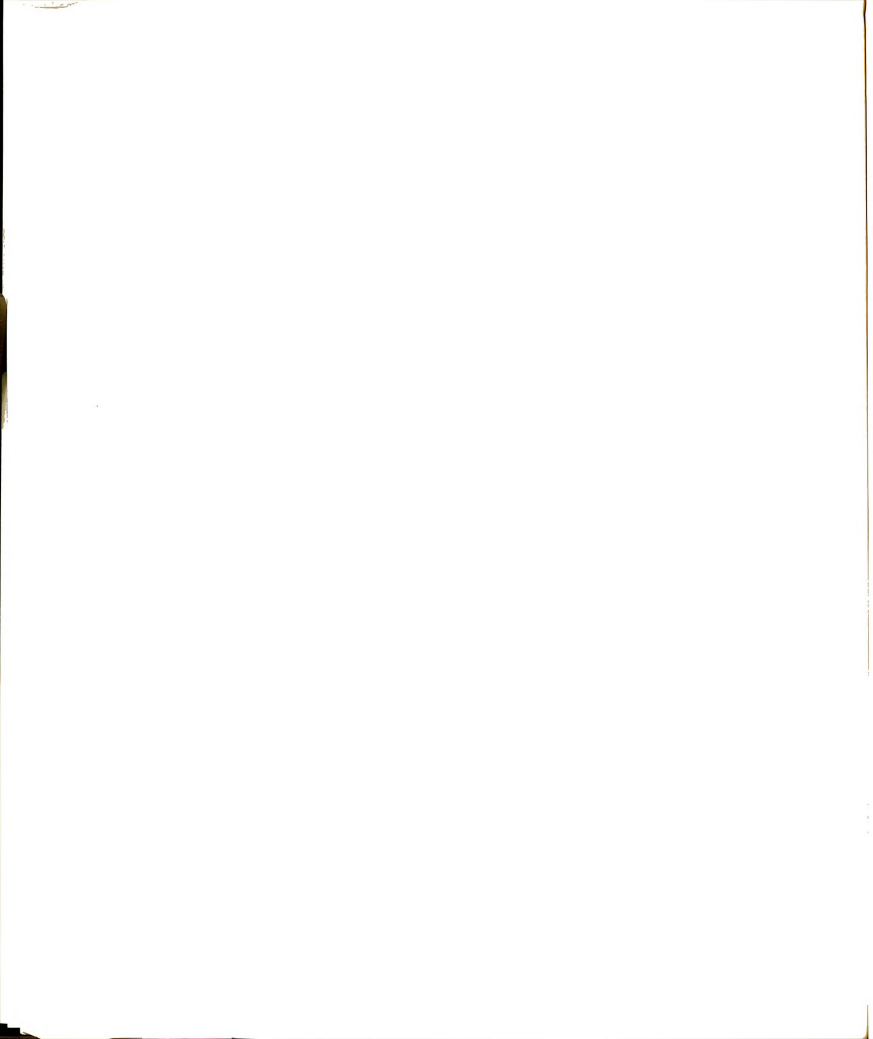


14. ATTACH,TAPE3,SURFACEDUETOSINE.
 15. ATTACH,TAPE4,SURFACEDUETOSINH.
 16. ATTACH,EWFIL, SURFACEDATACOSE.
 17. EDITOR.
 18. B,W.
 19. RETURN,EWFIL.
 20. PURGE,TAPE1.
 21. PURGE,TAPE2.
 22. PURGE,TAPE3.
 23. PURGE,TAPE4.
 24. *EOR
 25. SAVE,Z,100-1820.
 26. *EOR
 27. SAVE,W,NS.
 28. *EOF

Note that in the above job, we have assumed that program "COMBINE" has been cataloged as permanent file "COMBINEEW," and that the compiled program has been stored in permanent file named "COMBINEB." Furthermore, instead of creating a new input data file, the file "SURFACEDATACOSE" has been used.

4.11. Printed Output

The first output file of program "SURFLDS" mainly performs the echo checking of the input data. The location and dimension of each surface cell are listed first, then



the number of surface cells on surfaces of the first quadrant, the electrical parameters of the body, as well as the frequency and the type of mode of the incident field are all printed out as shown in Table 4.4.

The second output file is the listing of the real and imaginary parts of the incident equivalent surface currents which are internally determined for each surface cell. This file is shown in Table 4.5.

Then the computed induced surface currents at various surface locations are listed in the third output file where all the surface currents are expressed in terms of Amp/meter, as shown in Table 4.6.

The last output file consists of the information of the induced surface fields at various locations on the body surface as shown in Table 4.7. It is this file which will be buffered out to a magnetic tape with one of the following local file names, TAPE1, TAPE2, TAPE3, and TAPE4. Then this magnetic tape will be cataloged as a permanent file which has one of the following PFN's, "SURFACEDUETOSE," "SURFACEDUETOCOSH," "SURFACEDUETOSINE," and "SURFACEDUETOSINH."

Now we briefly discuss the output files of the program "COMBINE." The first file is the echo messages of the input variables. The next four files contain the values of the induced surface fields which are buffered in from TAPE1, TAPE2, TAPE3, TAPE4, respectively.

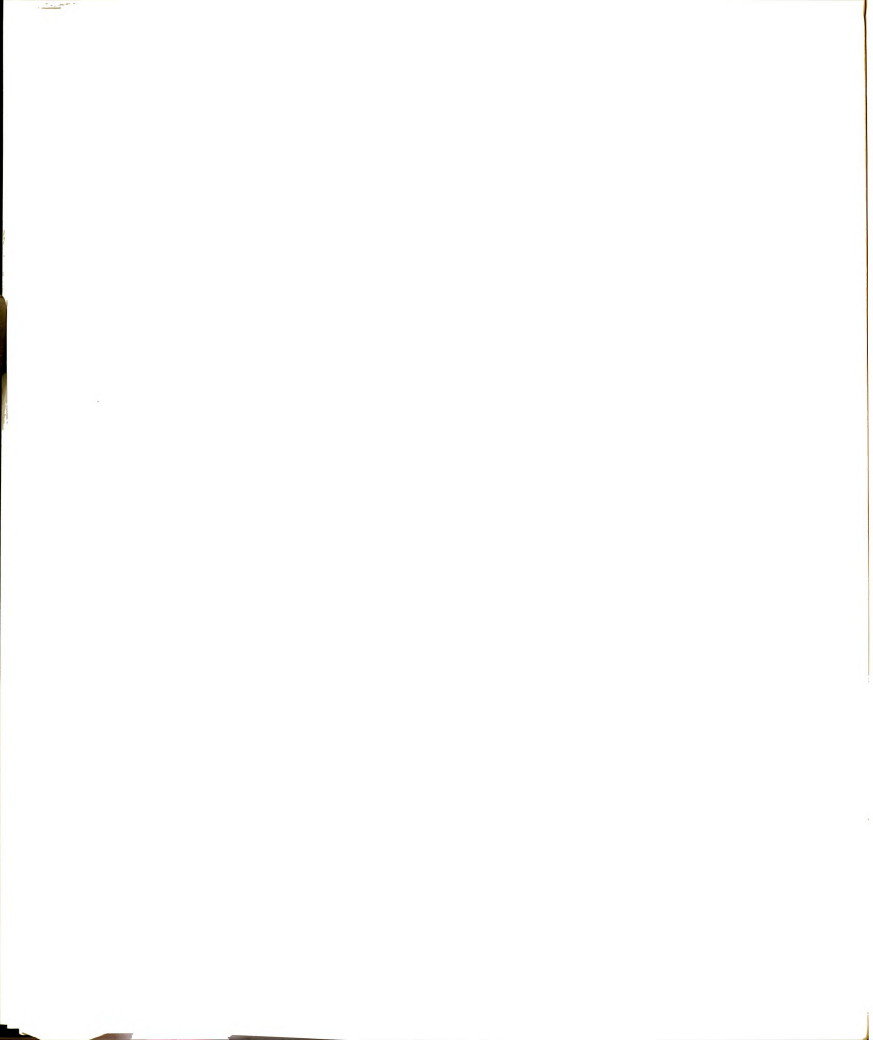


Table 4.4. First output file of program "SURFLDS."

THIS IS SOLUTION OF INDUCED FIELDS ON THE SURFACES OF AN N-D-S SYSTEM WITH AN ARBITRARY ELECTROSTATIC FIELD				
THE LOCATION AND DIMENSION OF EACH SURFACE-CELL ARE LISTED BELOW				
N	X (IN CM)	Y (IN CM)	Z (IN CM)	D (IN CM)
1	1.000E-01	.7500E-01	.7500E-01	.5000E-01
2	1.000E-01	.7500E-01	.2500E-01	.5000E-01
3	1.000E-01	.2500E-01	.7500E-01	.5000E-01
4	1.000E-01	.2500E-01	.2500E-01	.5000E-01
5	.7500E-01	1.0000E-01	.7500E-01	.5000E-01
6	.7500E-01	1.0000E-01	.2500E-01	.5000E-01
7	.7500E-01	.2500E-01	.7500E-01	.5000E-01
8	.7500E-01	.2500E-01	.2500E-01	.5000E-01
9	.7500E-01	.7500E-01	1.0000E-01	.5000E-01
10	.2500E-01	.7500E-01	1.0000E-01	.5000E-01
11	.2500E-01	.7500E-01	.2500E-01	.5000E-01
12	.2500E-01	.2500E-01	1.0000E-01	.5000E-01
THERE ARE TWO SURFACES OF THE FIRST QUADRANT				
EACH SURFACE CELL IS DIVIDED INTO 5 BY 5 SUB-SURFACE CELLS				
DIELECTRIC CONSTANT = 1.00000E+02				
CONDUCTIVITY = .001000E+02 OHM/CM				
FREQUENCY = 2400000000.0 HZ				
TYPE OF IMPRESSED FIELDING				

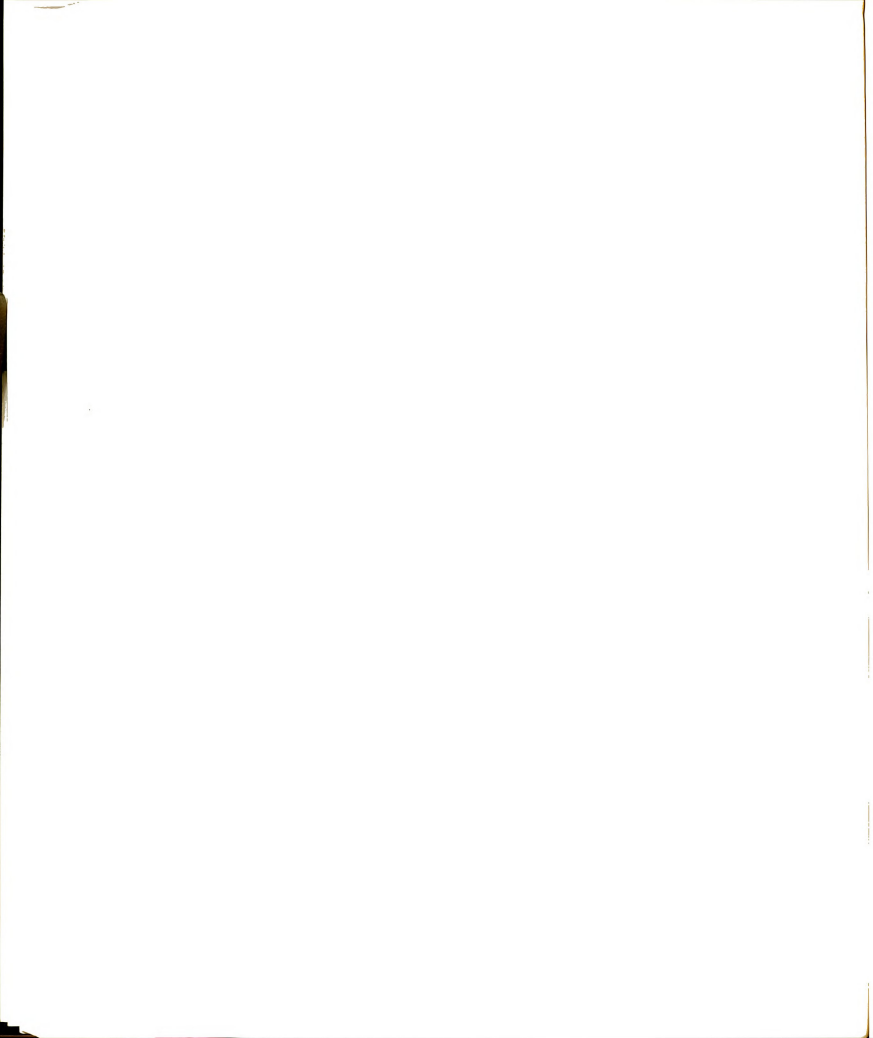


Table 4.5. Second output file of program "SURFLDS."

THE INMR CONT ARE LISTED AS FOLLOWS (ANYWHERE THE COMMENTS ON IN FIELDS WILL BE LIKE THIS: FIELDS)									
V	REAL	IMAG	REAL	IMAG	REAL	IMAG	REAL	IMAG	REAL
ON THE TOP SURFACE WE HAVE THE FOLLOWS									
J/T ON -H2									
J/T ON H2									
ON THE REAR SURFACE WE HAVE THE FOLLOWS									
J/T ON -H2									
J/T ON H2									
ON THE REAR SURFACE WE HAVE THE FOLLOWS									
J/T ON -H2									
J/T ON H2									
ON THE REAR SURFACE WE HAVE THE FOLLOWS									
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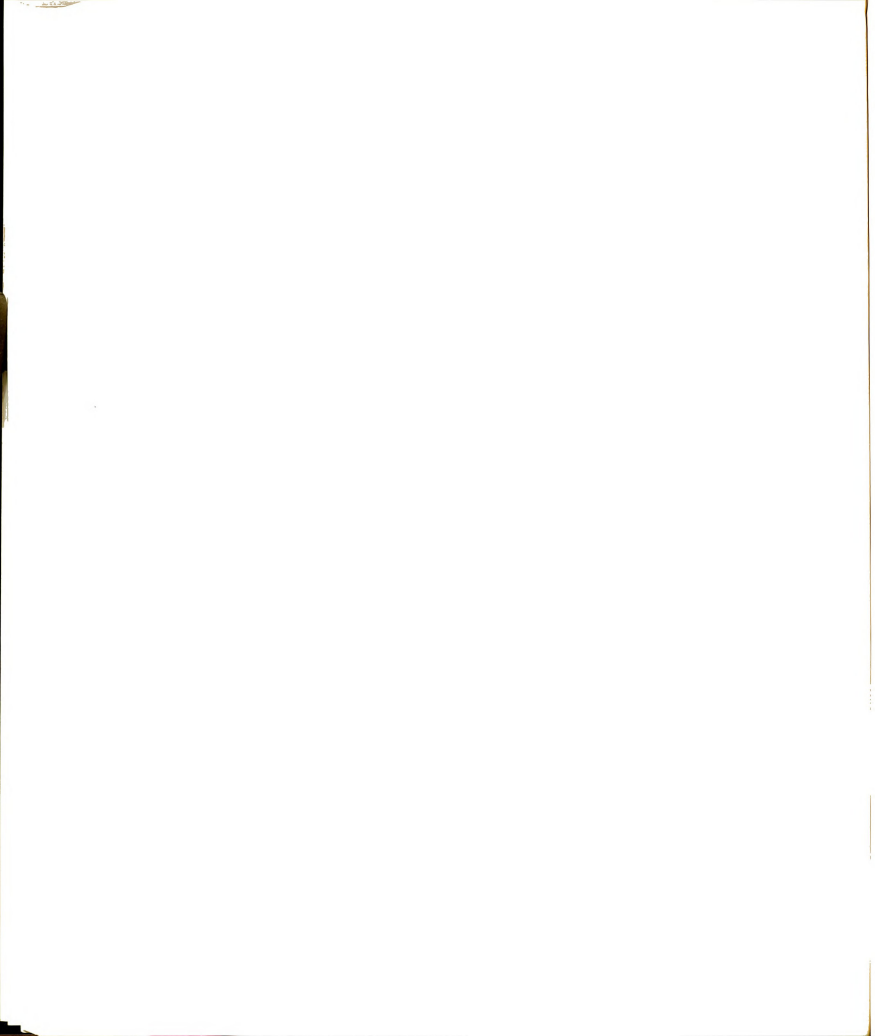
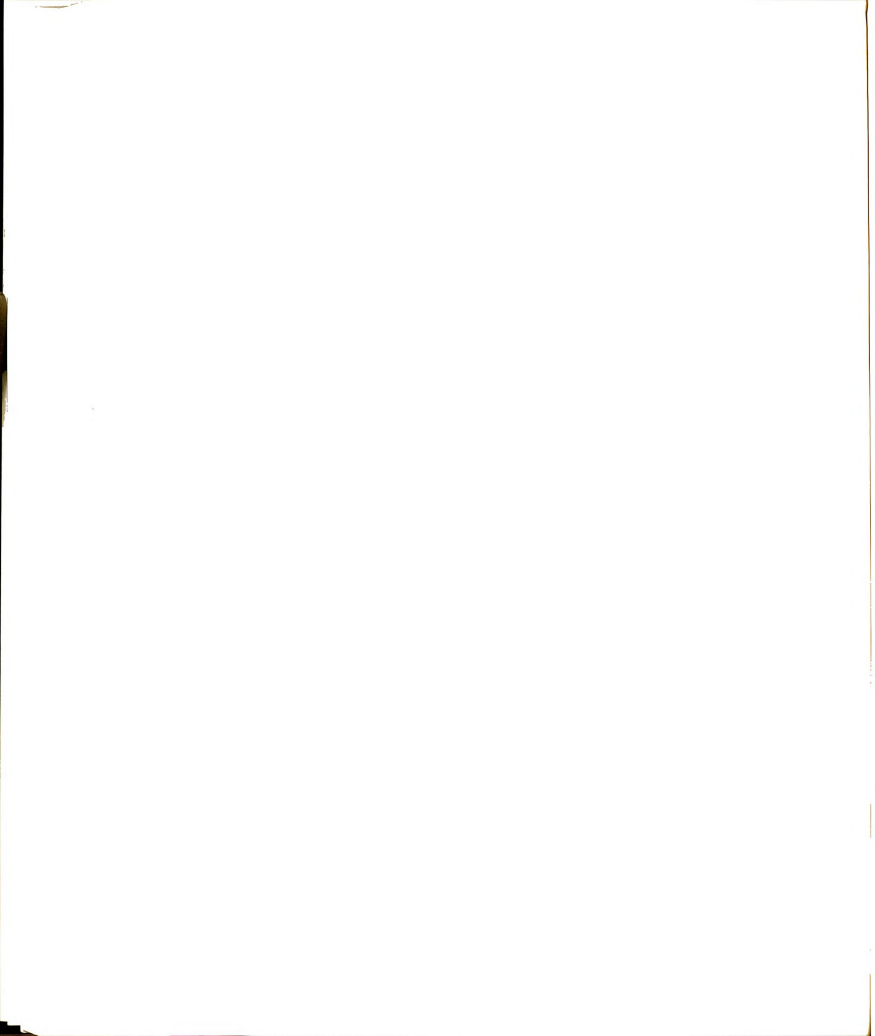
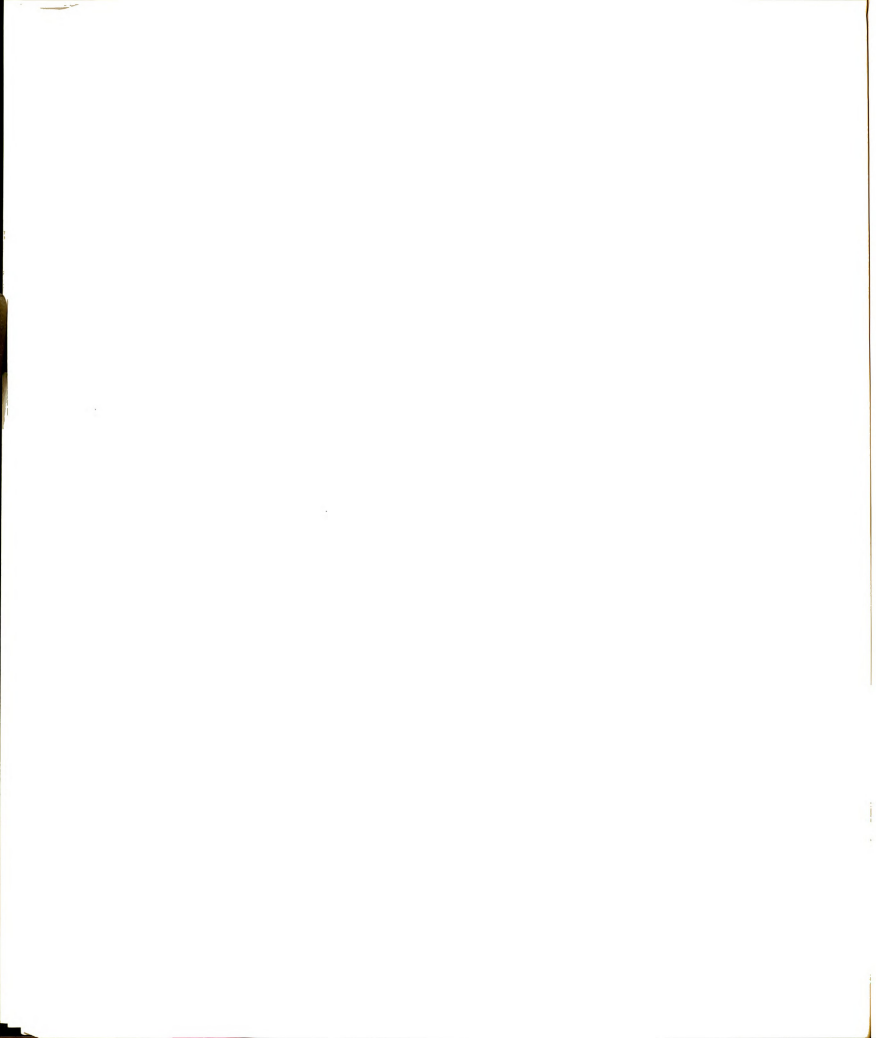


Table 4.6. Third output file of program "SURFLDS."

THE FOUR CASES ARE LISTED IN THE FOLLOWING ORDER: 1. THE FIRST TWO CASES ARE IN THE FIRST TWO COLUMNS, AND THE OTHER TWO ARE IN THE LAST TWO COLUMNS.									
ON THE TWO SURFACES WE HAVE THE FOLLOWING:									
JULY ON H ₂									
Y	REAL	IMAG	REAL	IMAG	REAL	IMAG	REAL	IMAG	TIME
JULY ON H ₂									
1	-.1585147E-03	-.4396270E-04	-.1519137E-03	-.5370923E-04	-.1526609E-04	-.1350853E-03	-.1526717E-04	-.165177E-03	JULY ON H ₂ 27/7/83
2	-.4871197E-03	.1058511E-03	-.1101313E-03	.4001190E-04	-.1376607E-03	-.6756441E-04	-.1278565E-03	.1171330E-04	
3	-.11091237E-03	.6057679E-04	-.1014091E-03	.1410531E-04	-.1331322E-03	-.1647127E-03	-.1117917E-03	.1738325E-04	
4	-.1144018E-03	.1654802E-03	-.1453131E-03	.1506277E-03	-.1643107E-03	-.6600131E-04	-.1516377E-03	.1318062E-04	
ON THE OTHER SURFACE WE HAVE THE FOLLOWING:									
JULY ON H ₂									
5	-.1015631E-02	.6468800E-03	-.2950573E-03	-.1757577E-03	-.1653116E-03	.1174401E-03	-.1339565E-03	.1317402E-03	JULY ON H ₂ 27/7/83
6	.1281115E-02	.8770807E-03	-.1316637E-03	-.1992313E-03	-.1571300E-03	.2894345E-04	-.1571823E-03	.2317305E-03	
7	.1359570E-02	-.7179767E-04	-.1713131E-03	-.1281837E-03	-.1713030E-03	.4673532E-04	-.1216547E-03	.1771551E-04	
8	.1171031E-02	-.592593E-03	.1356417E-03	-.1294483E-03	-.2008453E-04	.1319767E-04	-.1732670E-03	.1318062E-04	
ON THE OTHER SURFACE WE HAVE THE FOLLOWING:									
JULY ON H ₂									
9	.2996250E-03	.6465346E-03	-.1216417E-03	-.1281293E-03	-.1648360E-03	.1171700E-03	-.1339565E-03	.1317402E-03	JULY ON H ₂ 27/7/83
10	.1311321E-02	.8773551E-03	-.1654774E-03	-.1377104E-03	-.1559032E-03	.2894345E-04	-.1571823E-03	.2317305E-03	
11	.1316559E-02	-.7181826E-03	-.1818267E-03	-.1297240E-03	-.1736670E-03	.4673532E-04	-.1216547E-03	.1771551E-04	
12	.1407101E-02	-.651171E-03	.1297127E-03	-.1232560E-03	-.2073192E-04	.1317510E-04	-.1732670E-03	.1318062E-04	



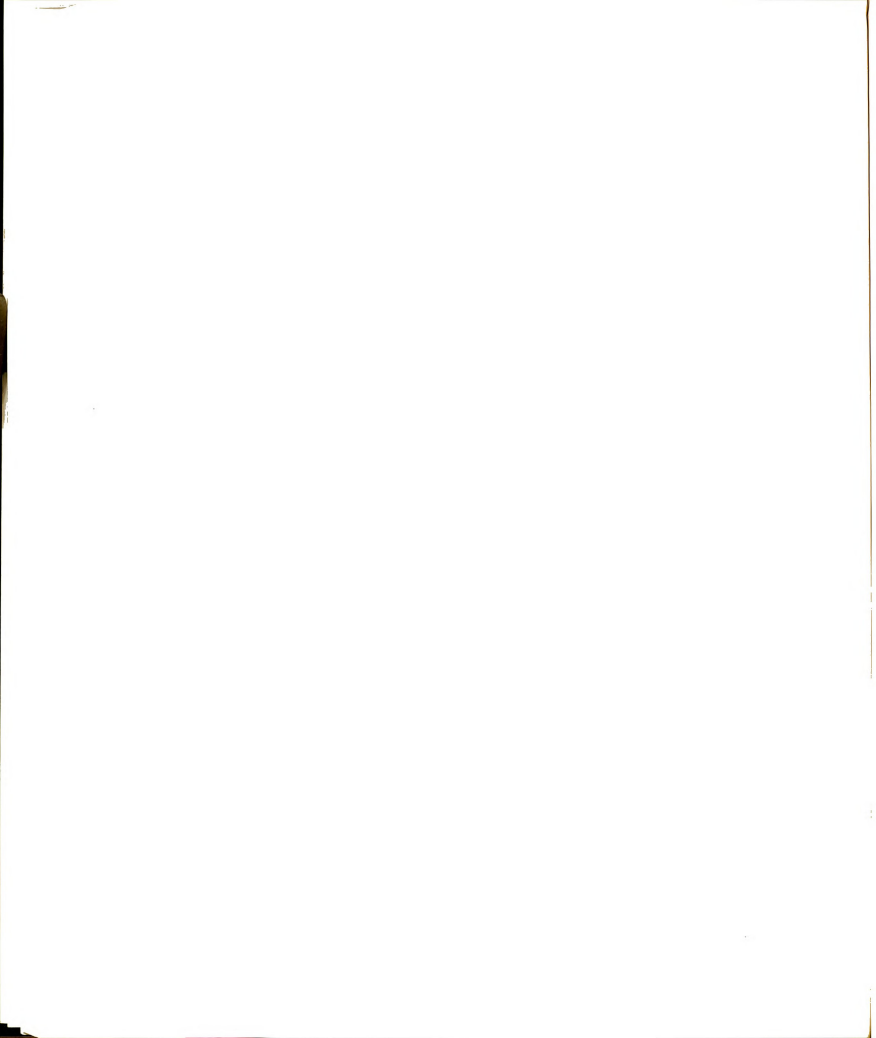


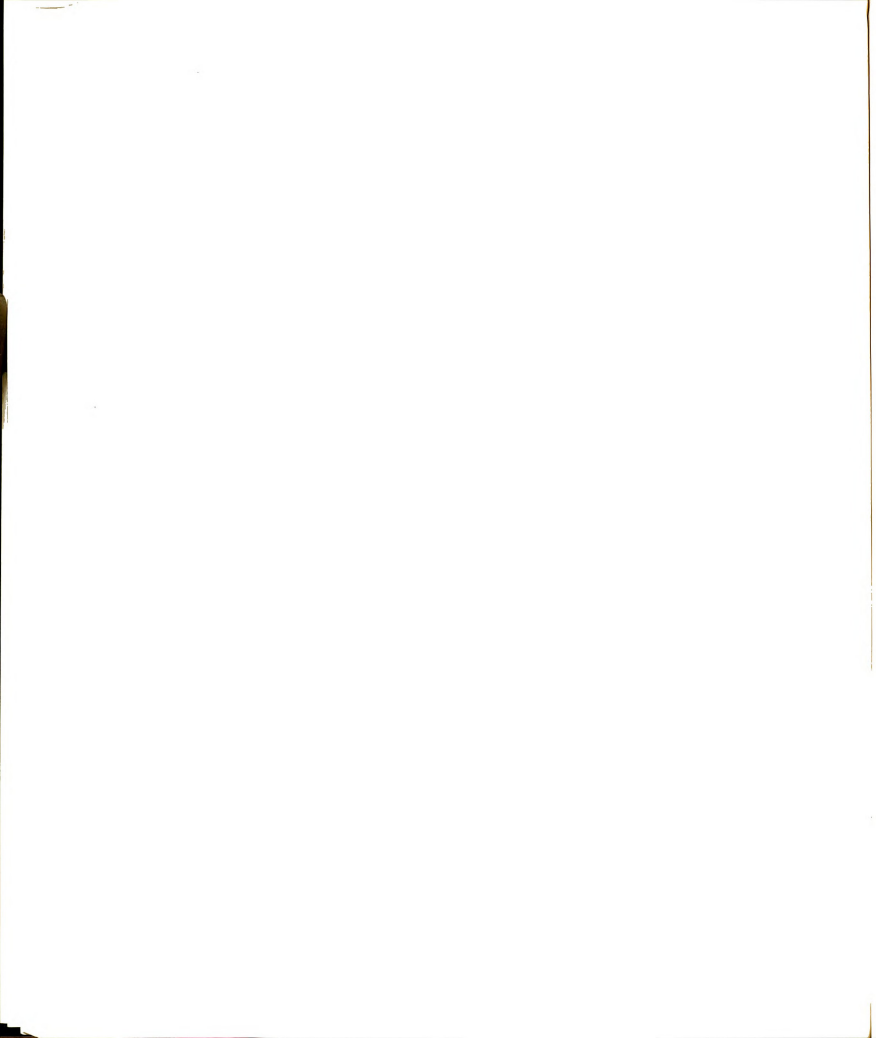
The sixth and the eighth output files list the total surface fields induced by plane wave on the surfaces of the first quadrant in real-imaginary and magnitude-phase forms, respectively.

Finally, the seventh and the last output files list the total surface fields for the fifth quadrant in real-imaginary and magnitude-phase forms, respectively. The above output files of the program "COMBINE" are shown in Table 4.8.

4.12. Listing of the Program

For the purpose of reference, the programs "SURFLDS" and "COMBINE" are listed on pages 183 -229. For about 150 surface unknowns, "SURFLDS" requires approximately 170000B words memory, while "COMBINE" requires about 30000B words. When the number of unknowns is more than 150, the memory required for "SURFLDS" exceeds that obtainable from CDC 6500 at MSU. To overcome this difficulty, the job may be submitted through Merit Network System for execution at other computing facilities. In this case, due to different operating systems involved, apparently different control commands must be employed. Nevertheless, since the program itself has been coded in standard FORTRAN, there is no modification needed except a subroutine named "BLOCK DATA" must be added into program "SURFLDS." The function





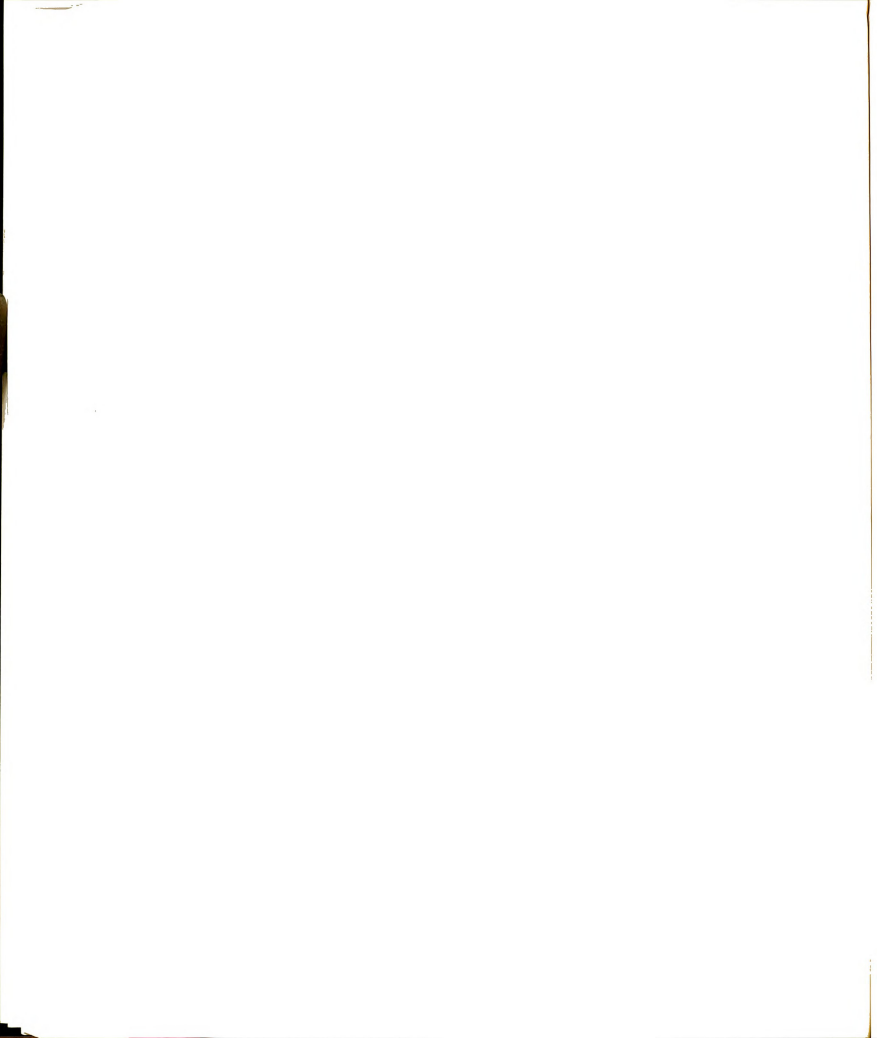


Table 4.8. (Cont'd.)

[illegible]

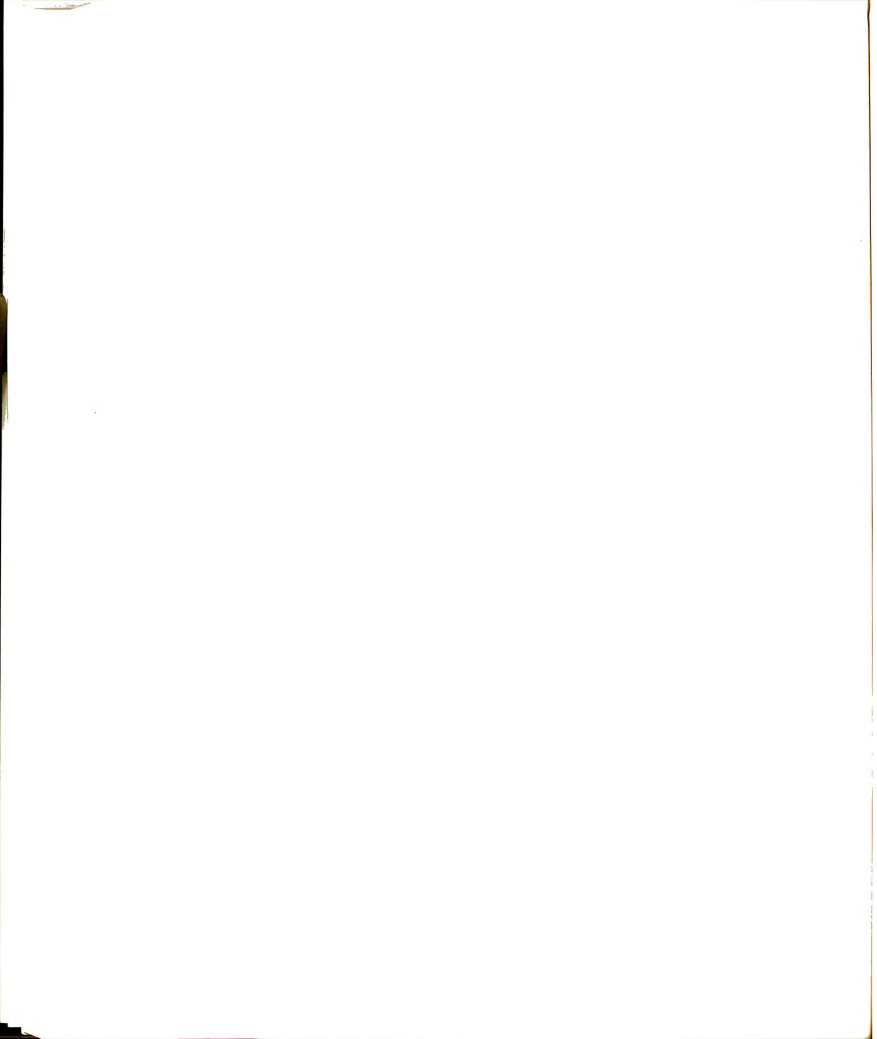


Table 4.8. (Cont'd.)

THE FILLS THINER BY STEP TWO OF IMPROVED FILLS IN THIS MOMENT ARE LISTED BELOW TOGETHER WITH THE FILLS WHICH ARE THINER									
N	REAL	1965	REAL	1966	REAL	1966	REAL	1967	1968
ON THE TOP SURFACE OF HART THE FILLS									
12									
ON THE RIGHT SURFACE OF HART THE FILLS									
17									
1	(-2879497-05	-2328801-75)	(-2853899-04	-2893355-04)	(-2859176-72	-2774825-73)	(-2835073-71	-2781315-01)	
2	(-2229737-05	-1727607-43)	(-2193497-13	-2093472-04)	(-2022347-01	-1645840-70)	(-2052712-00	-2033307-59)	
3	(-1318007-14	-9796431-15)	(-1127061-13	-2200564-04)	(-2077974-72	-2333420-73)	(-2012407-52	-2779435-03)	
4	(-5517018-05	-1532801-75)	(-2307534-03	-1335092-03)	(-1850222-71	-2600340-19)	(-2035304-72	-2300227-00)	
ON THE RIGHT SURFACE OF HART THE FILLS									
17									
5	(-2617737-03	-1862707-13)	(-2085777-04	-2136477-04)	(-2033427-72	-2410007-02)	(-2110377-11	-2781315-01)	
6	(-2870707-03	-2200737-53)	(-2220077-04	-2334707-04)	(-2295813-72	-2716797-53)	(-2002707-01	-2300007-01)	
7	(-1071607-13	-1100367-13)	(-2295217-04	-2078307-04)	(-2012417-00	-2684037-12)	(-2077974-72	-2300227-00)	
8	(-1234297-03	-1035207-13)	(-2000467-18	-2200007-04)	(-1674207-02	-2717177-53)	(-2064667-72	-2779435-03)	
ON THE RIGHT SURFACE OF HART THE FILLS									
17									
9	(-2273017-13	-1760007-04)	(-2007717-00	-2091707-04)	(-2122077-53	-2747707-10)	(-2100007-01	-2317177-53)	
10	(-2710107-13	-2180007-13)	(-2007717-00	-2091707-04)	(-2077974-72	-2410007-02)	(-2054777-11	-2779435-03)	
11	(-1000007-03	-2000207-13)	(-2122077-53	-2747707-10)	(-2077974-72	-2410007-02)	(-2054777-11	-2779435-03)	
12	(-2716007-13	-2180007-13)	(-2007717-00	-2091707-04)	(-2077974-72	-2410007-02)	(-2054777-11	-2779435-03)	

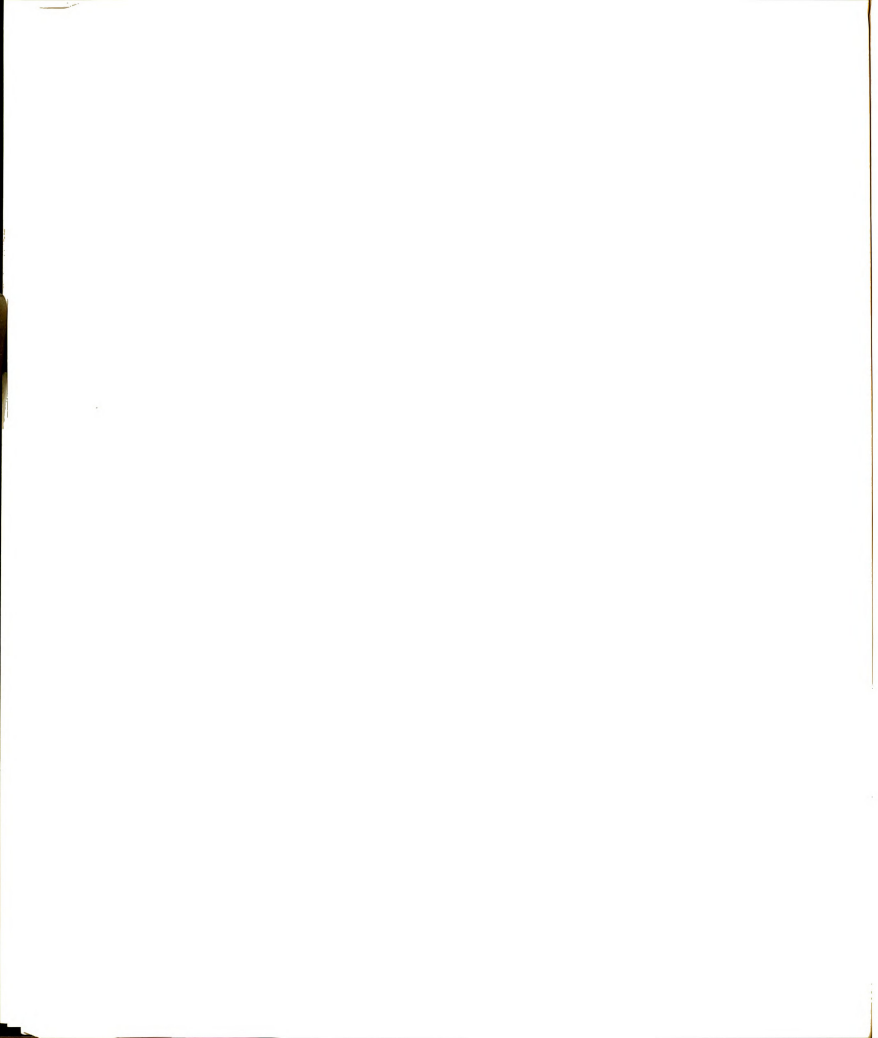
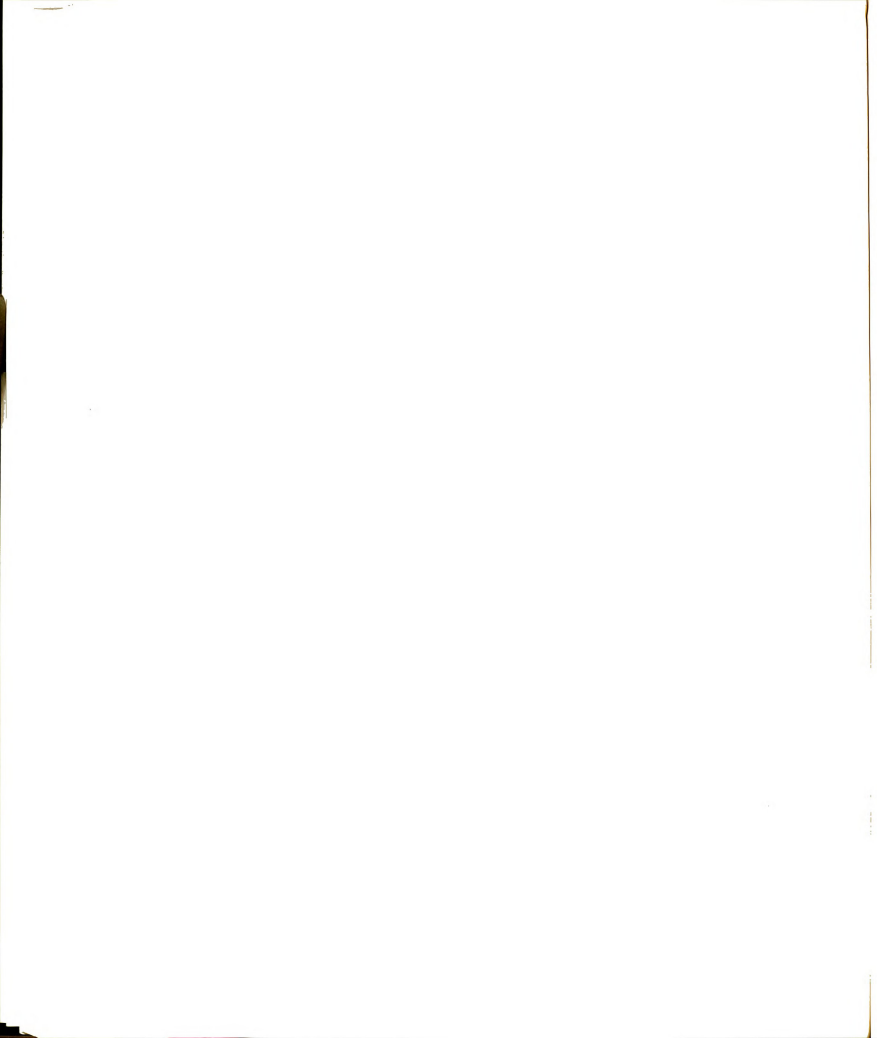
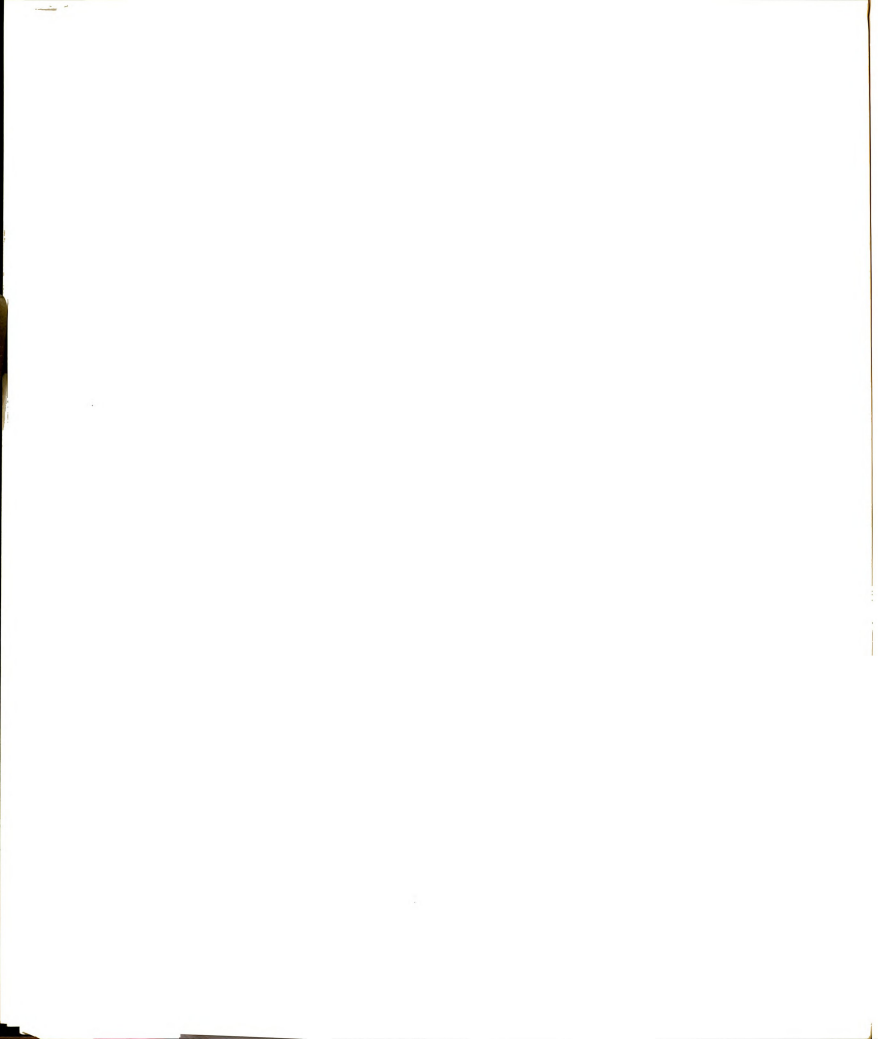
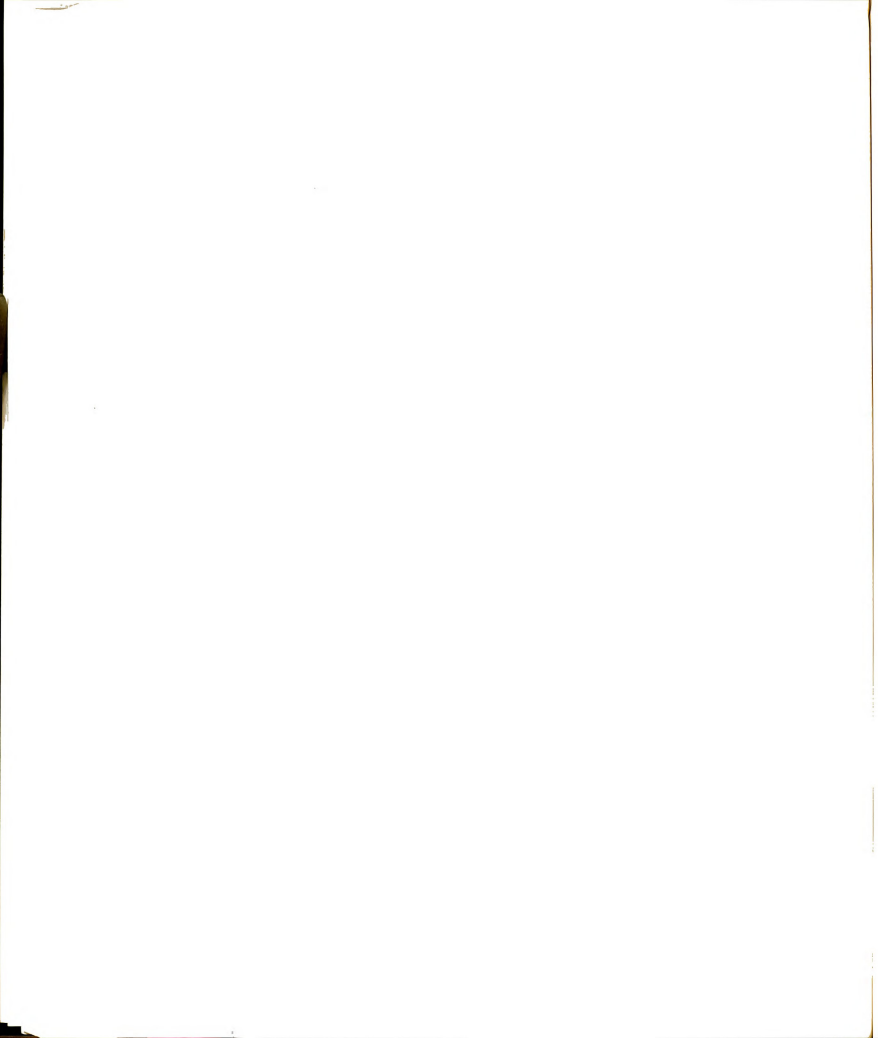


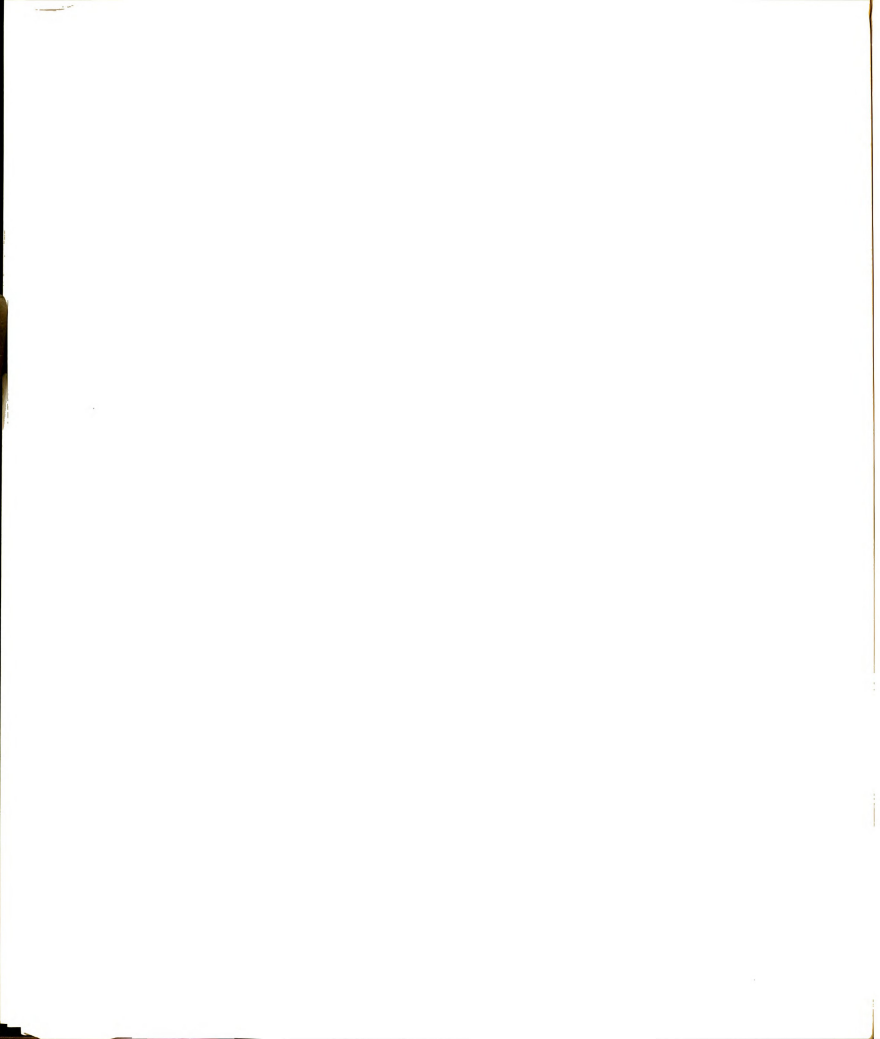
Table 4.8. (Cont'd.)

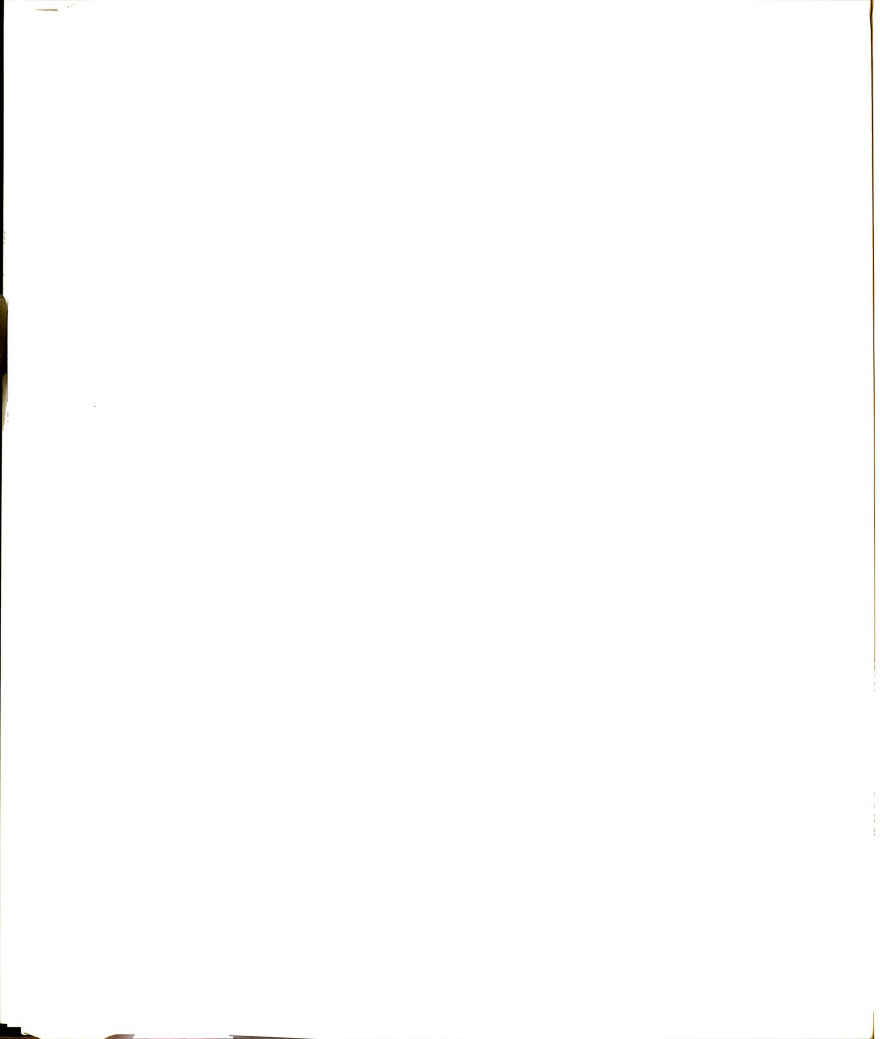
THE FILLS INDICED BY SIGN TYPE OF IMPROVED FIELD IN THE SUBSTANT ARE LISTED BELOW FOR EACH OF THE FOLLOWING SURFACES									
ON THE TOP SURFACE OF THE FOLLOWS									
W	REAL	1946	REAL	1946	REAL	1945			
ON THE REAR SURFACE OF THE FOLLOWS									
17	W	17	W	17	W	17			
1	-1592231-03	-2643841-03	-22355061-03	-15976001-03	-13063407-03	-7629171-03	-6481641-03	-5531111-03	
2	-50655001-03	-7892561-03	-7676961-03	-5816071-03	-5900037-03	-2603171-03	-11	-68633551-03	-27521171-03
3	-6906231-03	-2193907-03	-4541901-03	-1770011-03	-13264101-03	-1233407-03	-2780611-03	-4555511-03	
4	-1954347-03	-2160071-03	-1035521-03	-5902227-03	-5896637-03	-1751611-03	-1706641-03	-2793071-03	
ON THE REAR SURFACE OF THE FOLLOWS									
17	W	17	W	17	W	17			
5	-6659371-03	-1316071-03	-5084327-03	-4352317-03	-1790187-03	-6118561-03	-75733471-03	-1772921-03	
6	-1700111-03	-1612511-03	-5827961-03	-3556701-03	-1170671-03	-2805331-03	-11	-6553791-03	-27521171-03
7	-8569531-03	-1700111-03	-1270007-03	-3018061-03	-3405171-03	-5600111-03	-11	-6856611-03	-43194071-03
8	-1875411-03	-2720711-03	-5137011-03	-1473307-03	-5073307-03	-2750031-03	-11	-55864071-03	-2770071-03
ON THE REAR SURFACE OF THE FOLLOWS									
17	W	17	W	17	W	17			
9	-5536631-03	-7706371-03	-1677771-03	-1612211-03	-4512371-03	-4311671-03	-11	-1893211-03	-1894671-03
10	-1737001-03	-3700111-03	-5816111-03	-3557907-03	-7900111-03	-2737001-03	-11	-1270271-03	-2737001-03
11	-1798471-03	-1627707-03	-2800031-03	-3800331-03	-1299707-03	-5552111-03	-11	-5000221-03	-1717071-03
12	-6679001-03	-1716407-03	-1175111-03	-5651607-03	-1103307-03	-1800711-03	-11	-5073707-03	-2770071-03



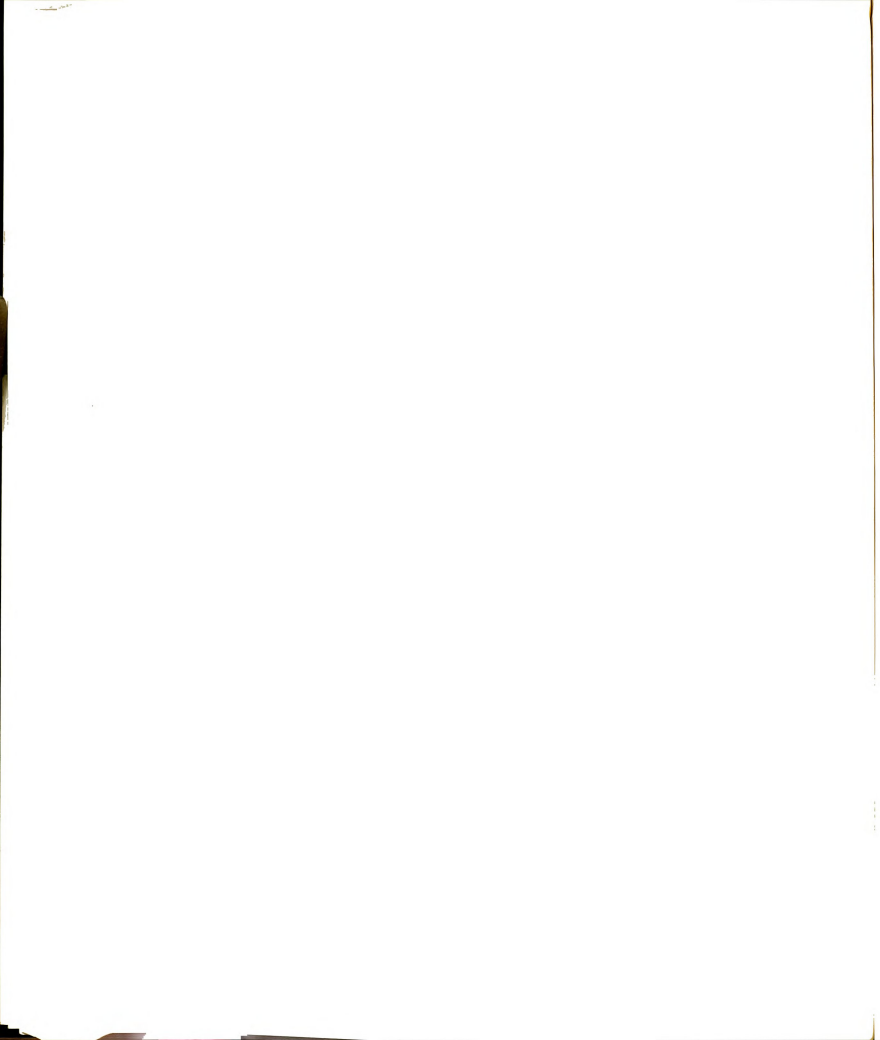


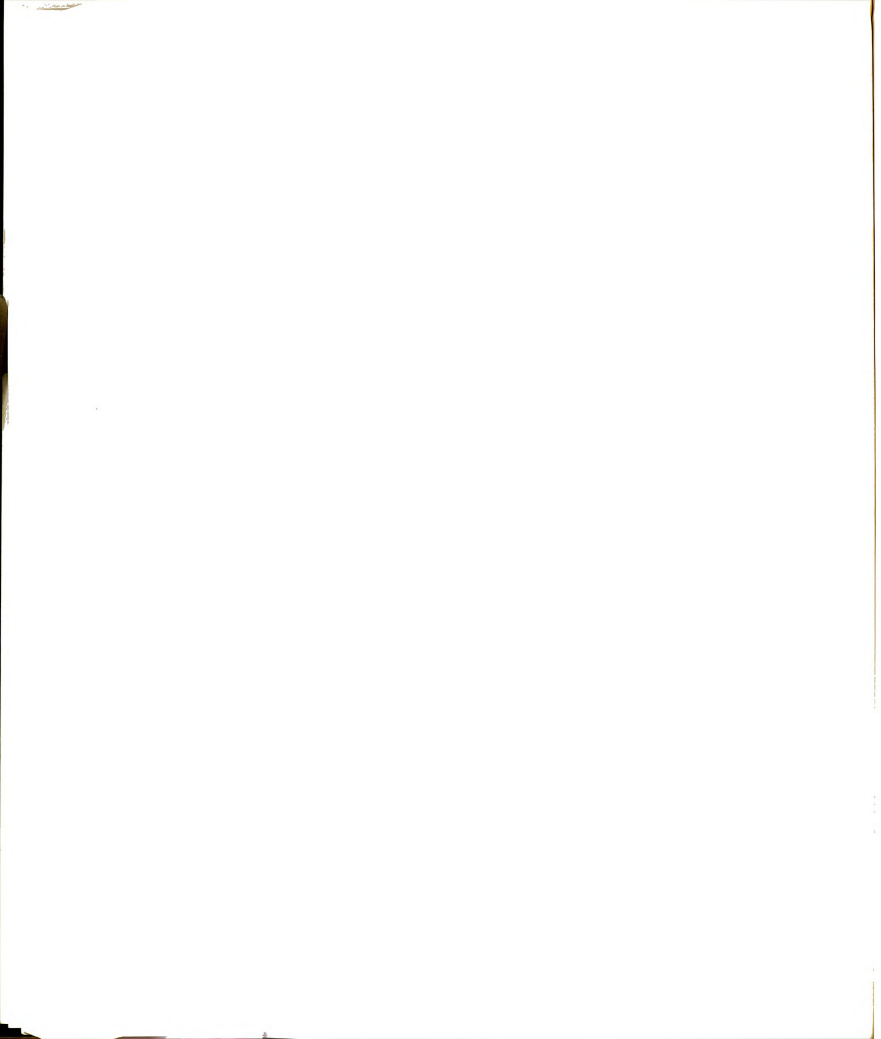






of "BLOCK DATA" can be found in most of the computer user's manuals.

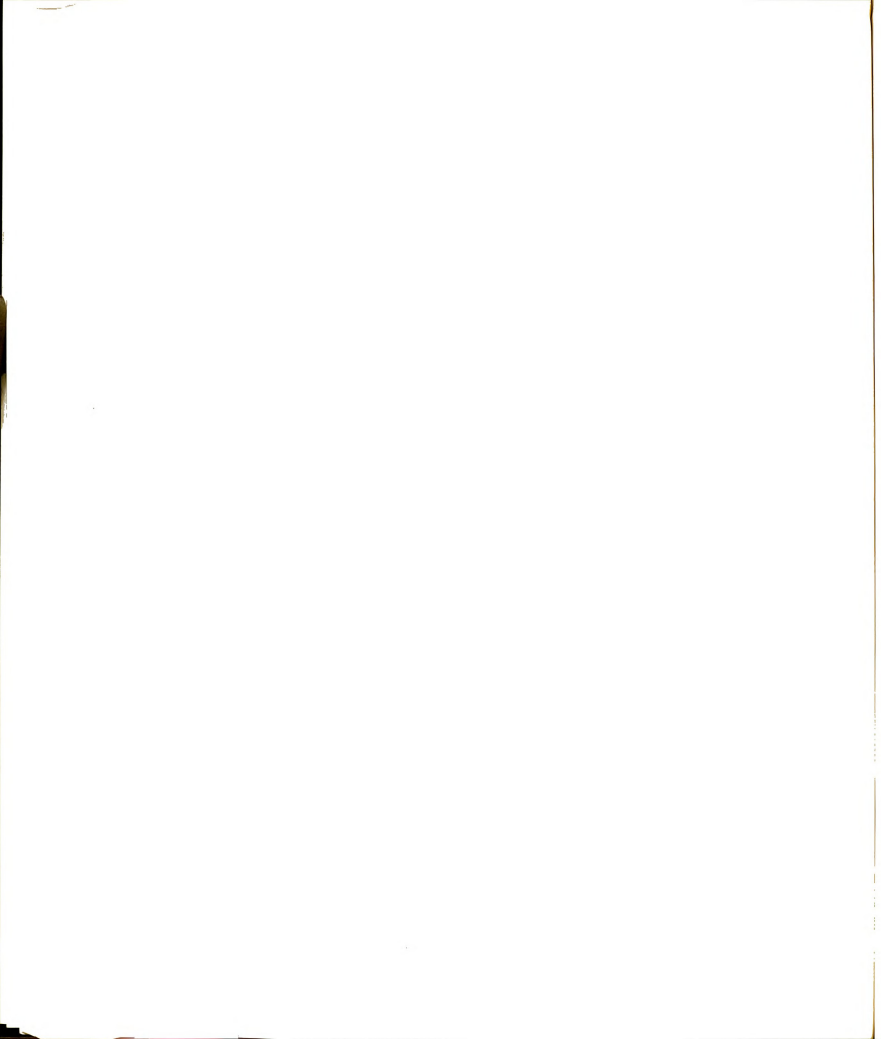




```

103  LEE(I,NP1)=EINC(I)
      PRINT 9, I, EINC(I)
      CONTINUE
      GO TO 51
50    I=1, NX
      EINC(I)=1.0*CHPLX(X(0,0,1.0)*KO*XCEN(I)/2.0
106  CONTINUE
      NXSP1=NXS+1
      DO 107 I=NXSP1,NX
107    EINC(I)=-1.0*CHPLX(X(0,0,1.0)*XEND(NXS)*XEND(NXS)*KO/(2.0*XCEN(I))
      NXSP1=NX
      N7M1=N7-1
      DO 108 J=1,N7M1
      DO 108 K=1,NX
      I=J*NX+K
108    EINC(I)=EINC(K)
      CONTINUE
      DO 109 I=1,N
      LEE(I,NP1)=EINC(I)
109  PRINT 9, I, EINC(I)
      CONTINUE
51    CALL OMTP(LEE,N,I,DET,I,OE-50)
      PRINT 10, FREINM, NPAR, AERR, RERR
      DO 104 I=1,N
      EIND(I)=LEE(I,NP1)
104  PRINT 11, I, LEE(I,NP1)
      CONTINUE
      PHA(EIND,N,MAG,PHA)
      PRINT 12, MAG, PHA
      DO 105 I=1,N
105  PRINT 11, I, MAG(I), PHA(I)
      CONTINUE
      CALL PCURVE(EINC,N,NX,XCEN,MAG)
      1 FORMAT(15,F10.5,2I5)
      2 FORMAT(15,2F10.5)
      3 FORMAT(15,2F10.5)
      4 FORMAT(15,2F10.5,2F10.5)
      5 FORMAT(15,2F10.5,2F10.5,2F10.5)
      6 FORMAT(15,2F10.5,2F10.5,2F10.5,2F10.5)
      7X,"X-DIMENSION",4X,"Z-DIMENSION",4X,"CRO-SEC AREA OF EA
      2CH RING ARE GIVEN BELOW IN METERS",10X,"XEND",11X,
      3"XEND",7X,"X-DIMENSION",4X,"Z-DIMENSION",4X,"CRO-SEC AREA")
      9 FORMAT(10,"BX,I3.5(2X,E13.6))
      10X,"THE CENTRAL POINT LOCATION AS WELL AS ELECTRICAL P
      PROPERTIES OF EACH RING ARE GIVEN BELOW"/"---",10X,"N",2X,"XCEN(
      SRG)",2X,"ZCEN (METERS)",2X,"DIELEC CONST.",2X,"CONDUCT.(MHO/M)".

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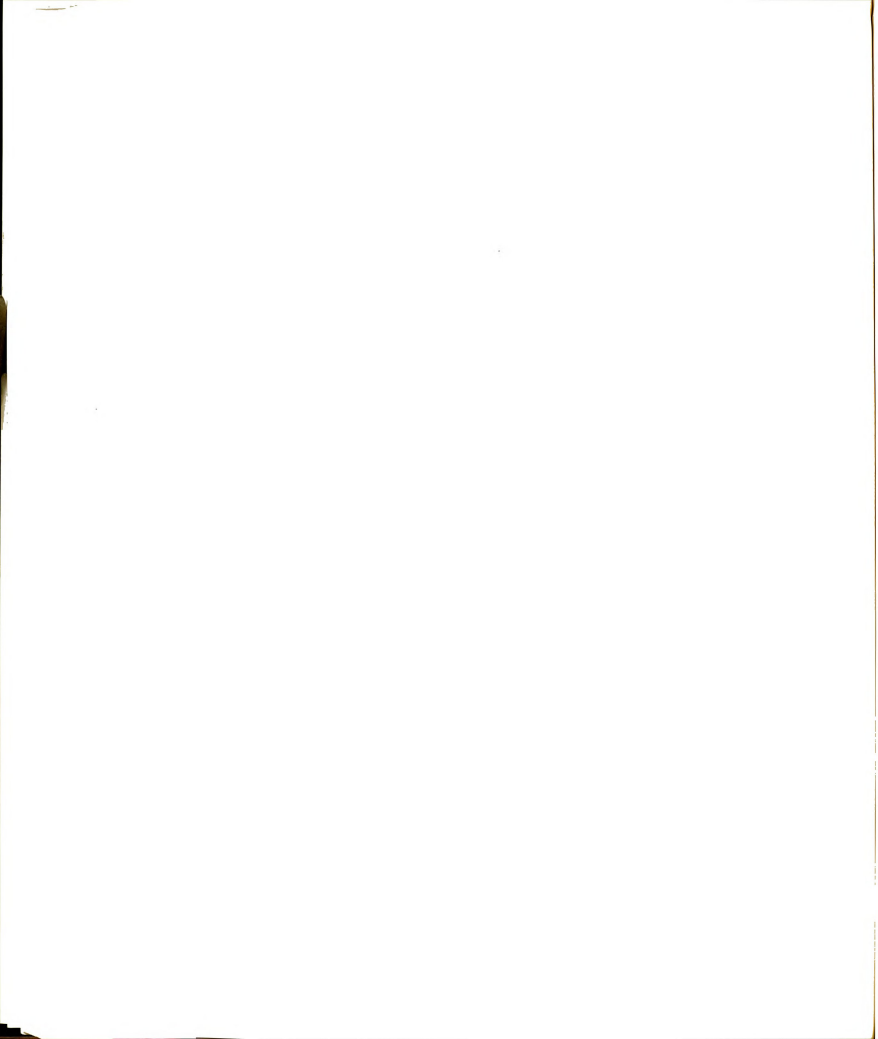
7 FORMAT("0", BX, I3.4(2X, E13.6))
8 FORMAT("1", IOX, "THE IMPRESS ELECTRIC FIELD DUE TO INCIDENT MAGNETIC
2C FIELD (WITH UNIT INCIDENT ELECTRIC FIELD) IS AS FOLLOWS", "0X",
3N, IOX, "IMPRESS ELECTRIC FIELD (VOLTS/METER)", "-.", "2X", "REAL", "2X",
4X, "IMAG")
9 FORMAT("0", BX, I3, I3X, E13.6, 4X, E13.6)
10 FORMAT("1", "THE INDUCED ELECTRIC FIELD IS IN AZIMUTHAL DIRECTION,
2THE ESTIMATED VALUES AT CENTRAL OF EACH SURFACE OF CYLINDER BELOW",
3"-.", "THE FREQUENCY BEING USED IS", F10.5, "MHZ, AND THE WAVELENGTH
4BEH INTO IS", SUBINTERVALS DURING NUMERICAL INTEGRATIONS, "1X", IOX,
52, "ANALOGUE TIME INVOLVED RETAIN", F10.5, "ABSOLUTE ERROR", "2X", IOX,
62, "ANALOGUE TIME INVOLVED RETAIN", F10.5, "ABSOLUTE ERROR", "2X", IOX,
71S/METER", "1X", "EINDUCED ELECTRIC VOLTS/METER", "1X")
11 FORMAT("0", BX, I3, I6X, E13.6, 23X, E13.6)
12 FORMAT("1", "THE MAGNITUDES AND PHASES OF INDUCED E FIELDS ARE GIVE
2N BELOW", "-.", "1X", "N", IOX, "EINDUCED-MAGN(VOLTS/METER)", IOX,
3N, "PHASES (IN DEGREES)", "-.", "1X", "N", IOX, "PHASES (IN DEGREES)",
43 FORMAT("0", "9X", "THE BODY IS IMMERSED IN AN UNIFORM MAGNETIC FIELD")
14 FORMAT("-.", "9X", "THE BODY IS ILLUMINATED BY A MAGNETIC BEAM")
STOP
END
*****
SUBROUTINE LEMMAT*****
SUBROUTINE LEMMAT(N, AN, FMNC, AERR, RERR, LEE)
EXTERNAL FMNC, FMFS, FMNC, AERR, RERR, LEE)
COMPLEX LEE(12, 13), KINTEG(12, 12), TAU(12)
REAL KO, MUO
COMMON/CONS1/FREQ, MUO, EPO, PI, OMEG, KO
COMMON/CONS2/VS(12), Z(12), A(12)
COMMON/PARA2/VS(12), Z(12), A(12)
COMMON/PARA3/XM
COMMON/PARA4/XM
COMMON/PARA5/XN, ZN, ZM
COMMON/PARA6/AREA(12)
DO 10 I=1, NP1
DO 10 J=1, NP1
LEE(I, J)=CMPLX(0.0, 0.0)
DEL T=PI/NUM
DO 10 K=1, NUM
TAU(K)=EIG(1)+CMPLX(0.0, 1.0)*OMEG*(EP(1)*EPO-EPO)
XM=A(I)
XN=X(1)
ZM=Z(1)
SUNC=0.0
SUNC=0.0
SUNC=0.0
SUNC=0.0
BETA=DEL
BETA=DEL
SUNC=SUNC+DCADRE(FMNC, ALPHA, BETA, AERR, RERR, ERROR, IER)
SUNC=SUNC+DCADRE(FMNC, ALPHA, BETA, AERR, RERR, ERROR, IER)

```

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C** SUBROUTINE LEENATV( A, Z, IZ, N, M, L, K, J, I)
      DIMENSION A(N,N), Z(N,N), IZ(N,N), L(M,M), K(M,M), J(M,M), I(M,M)

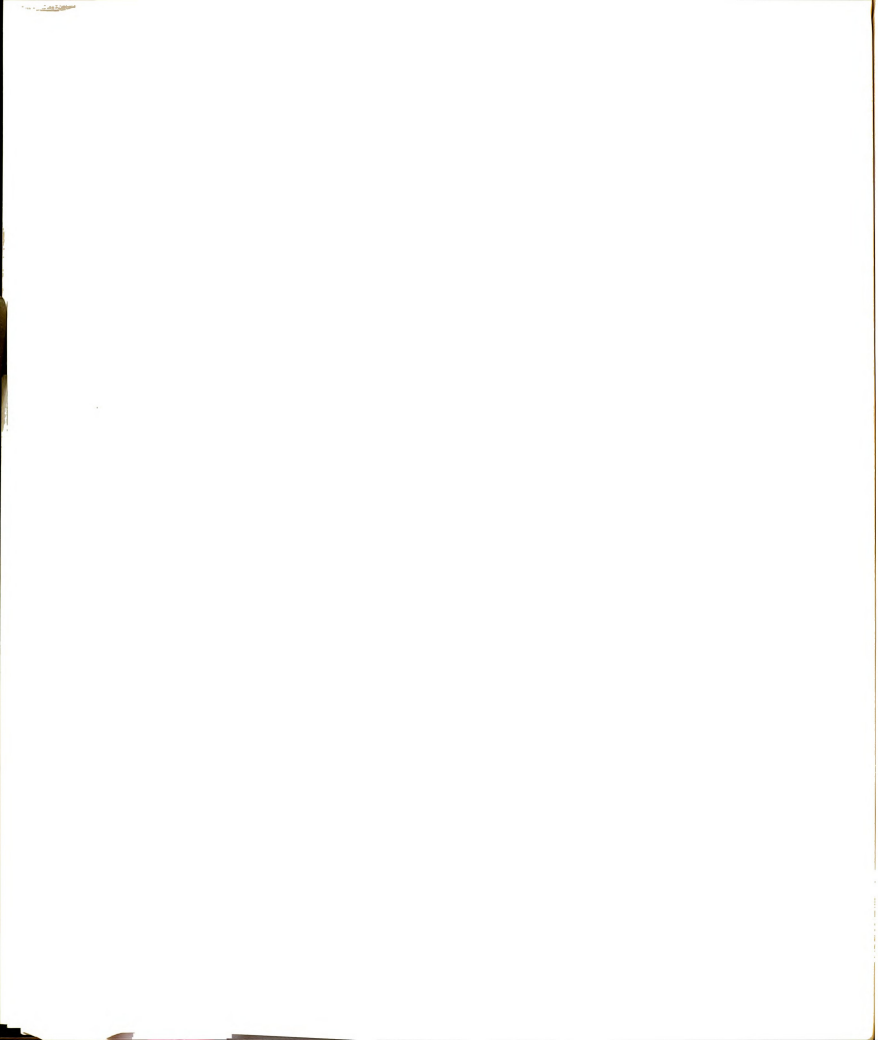
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      ALPHA=BETA
      BETA=BETA*DEL T
      IF (BETA.GT.PI+0.5*DEL T) GO TO 20
20  DO 10 I=1,N
      KINTEG(I,I)=(SUMC+CMPLX(0.0,1.0)*SUMS)*2.0*PI/AREA(I)
      IM1=I-1
      XM=X(I)
      ZM=Z(I)
      DO 30 J=1,IM1
      ZN=X(J)
      ZM=X(J)
      SUMC=0.0
      SUMS=0.0
      ALPHA=0.0
      BETA=DEL T
25  SUMC=SUMC+DCADRE(FMNC,ALPHA,BETA,AERR,ERR,ERROR,IER)
      SUMS=SUMS+DCADRE(FMNS,ALPHA,BETA,AERR,ERR,ERROR,IER)
      ALPHA=BETA
      BETA=BETA*DEL T
      IF (BETA.GT.PI+0.5*DEL T) GO TO 29
29  DO 10 I=1,N
      KINTEG(I,J)=SUMC+CMPLX(0.0,1.0)*SUMS
30  KINTEG(I,I)=KINTEG(I,J)
      DO 40 I=1,N
      DO 40 J=1,N
      GO TO 35
      LEE(I,J)=CMPLX(0.0,1.0)*FREQ*MUO*AREA(J)*X(J)*TAU(J)*KINTEG(I,J)
35  LEE(I,J)=1.0+CMPLX(0.0,1.0)*FREQ*MUO*AREA(J)*X(J)*TAU(J)*KINTEG(I,
40  J)
40  CONTINUE
      RETURN
      END
C*****FUNCTION FMNC*****
      FUNCTION FMNC(T)
      REAL KO,MUO
      COMMON/CONST/PI,OMEG,KO
      COMMON/CONST/ALPHA,XI
      COMMON/PARA4/AM
      ALPHAM=2.0*AM/SQRT(PI)
      F1=2.0*XM*XM*(1.0-COS(T))
      F2=SQRT(F1)
      FMNC=SQRT(1.0+ALPHAM*ALPHAM)-F2
      FMNC=OS(T)*COS(KO*F2)*ROOT
      RETURN
      END

```



```

C*****FUNCTION FMNS(T)*****
FUNCTION FMNS(T)
  REAL KO, MUO
  COMMON/CONS1/ FREQ, MUO, EPO, PI, OMEG, KO
  COMMON/PARA3/ XN
  COMMON/PARA4/ XH
  ALPHAM=2.0*PI/SQRT(PI)
  F1=2.0*X*M*X*(1.0-COS(T))
  F2=SQRT(F1)
  FMNS=SQRT(F1+ALPHAM*ALPHAM)-F2
  FMNS=COS(T)*SIN(KO*F2)*ROOT
  RETURN
END

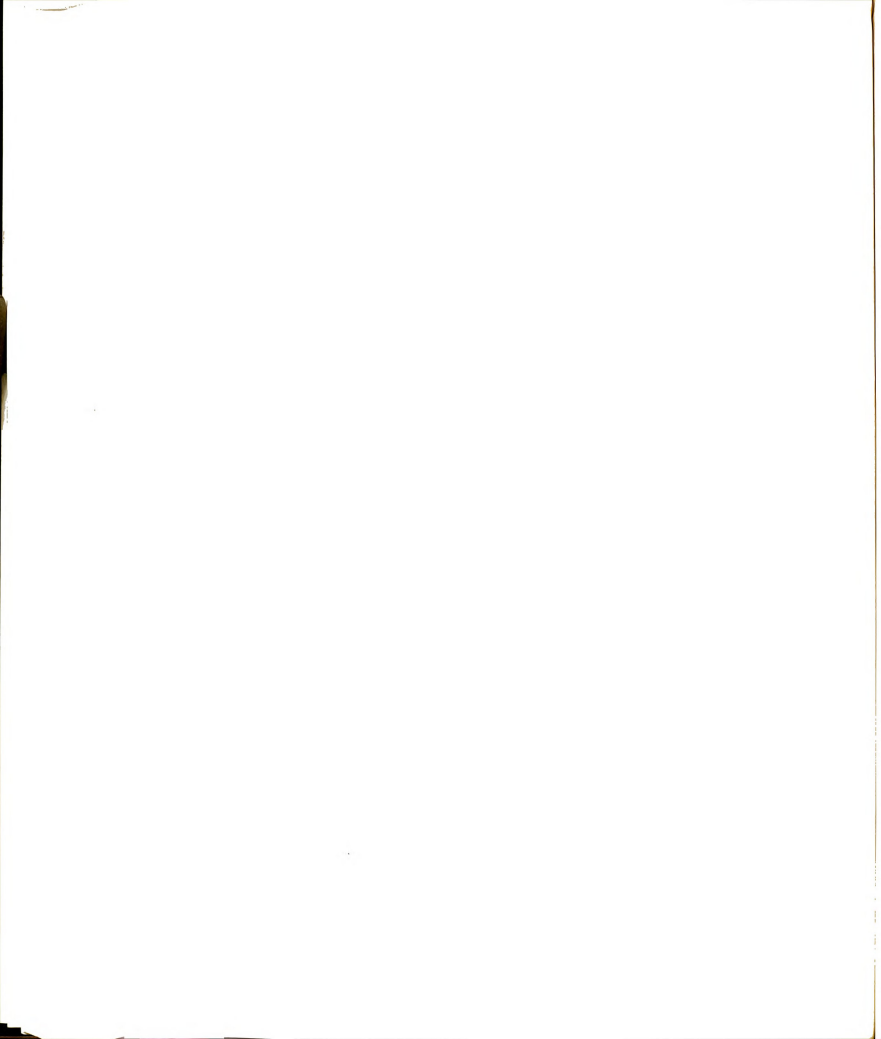
C*****FUNCTION FMNC(T)*****
FUNCTION FMNC(T)
  REAL KO, MUO
  COMMON/CONS1/ FREQ, MUO, EPO, PI, OMEG, KO
  COMMON/PARA3/ XN
  COMMON/PARA5/ XN, ZN, ZM
  ROOT=SQRT(XM*XN+XN*XN-2.0*X*M*X*(COS(T)+(ZN-ZN)*(ZM-ZN)))
  FMNC=COS(T)*COS(KO*ROOT)/ROOT
  RETURN
END

C*****FUNCTION FMNS(T)*****
FUNCTION FMNS(T)
  REAL KO, MUO
  COMMON/CONS1/ FREQ, MUO, EPO, PI, OMEG, KO
  COMMON/PARA3/ XN
  COMMON/PARA5/ XN, ZN, ZM
  ROOT=SQRT(XM*XN+XN*XN-2.0*X*M*X*(COS(T)+(ZN-ZN)*(ZM-ZN)))
  FMNS=-COS(T)*SIN(KO*ROOT)/ROOT
  RETURN
END

C*****SUBROUTINE CHATP(A,N,M,DEL,EP)*****
SUBROUTINE CHATP(A,N,M,DEL,EP)
  COMPLEX A,B,DET,CONST,S
  DIMENSION A(12,13)

  DET = CMPLX(1.0, 0.0)
  NPI = N + 1
  NPM = N + N
  NM1 = N - 1
  DO 4 J=1,NM1
    C = CABS(A(J,J))
    VP1 = J + 1
    DO 5 I = VP1,N
      D = CABS(A(I,J))

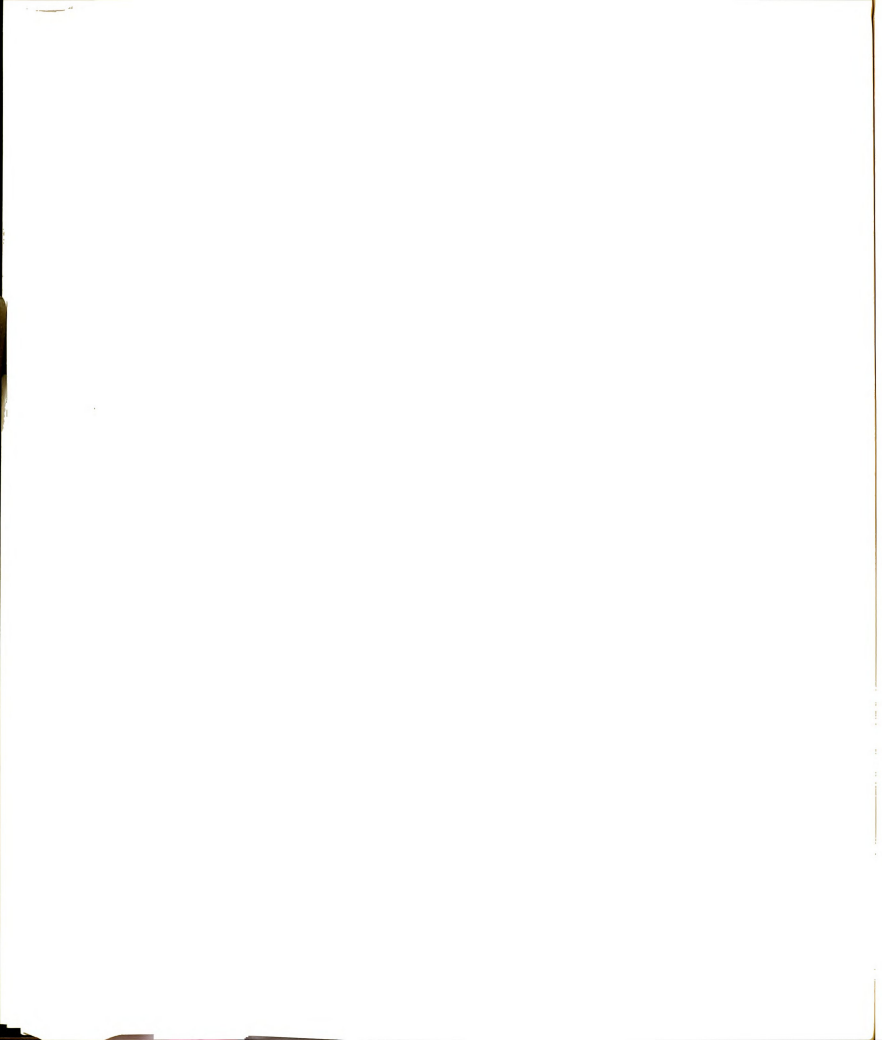
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6      IF ( C-D ) G, 5, 3
      DET = D-I
      DO 7 K = 1, NPH
      P = A(I,K) * A(J,K)
7      A(I,K) = A(J,K)
      C = D
5      CONTINUE
15     DO (CABS(A(J,I))-EP) 14, 15, 15
      DET = A(I,I) * A(J,I)
      CONST = A(I,I) / A(J,I)
4      DO 4 K = NPI, NPN
      A(I,K) = A(I,K) - CONST * A(J,K)
14     IF (CABS(A(N,N))-EP) 14, 18, 18
      DET = CMPLX(0, 0)
      PRINT 30
      RETURN
18     CONTINUE
      DO 11 I = 1, N
      KPI = K + I
      DO 12 L = NPI, NPN
      S = CMPLX(0, 0)
19     IF ( N - KPI ) 12, 19, 19
      DO 13 N = KPI, N
      S = S + A(K,I) * A(J,L)
13     A(K,L) = (A(K,L) - S) / A(K,K)
30     FORMAT(1X, 42H THE DETERMINANT OF THE SYSTEM EQUALS ZERO. /
      1X, 36H THE PROGRAM CANNOT HANDLE THIS CASE. //)
      RETURN
      END
C*****SUBROUTINE MAGPHA*****
SUBROUTINE MAGPHA(EIND, N, MAG, PHA)
      COMPLEX EIND(12)
      REAL MAG(12), PHA(12)
      DIMENSION I(1), NE(12), PHA(12)
      DO 10 I = 1, NE(12), PHA(12)
      MAG(I) = CABS(EIND(I))
      RE(I) = REAL(EIND(I))
      IM(I) = AIMAG(EIND(I))
      PHA(I) = ATAN2(IM(I), RE(I))
10     PHA(I) = PHA(I) * 57.29577951
      CONTINUE
      RETURN
      END

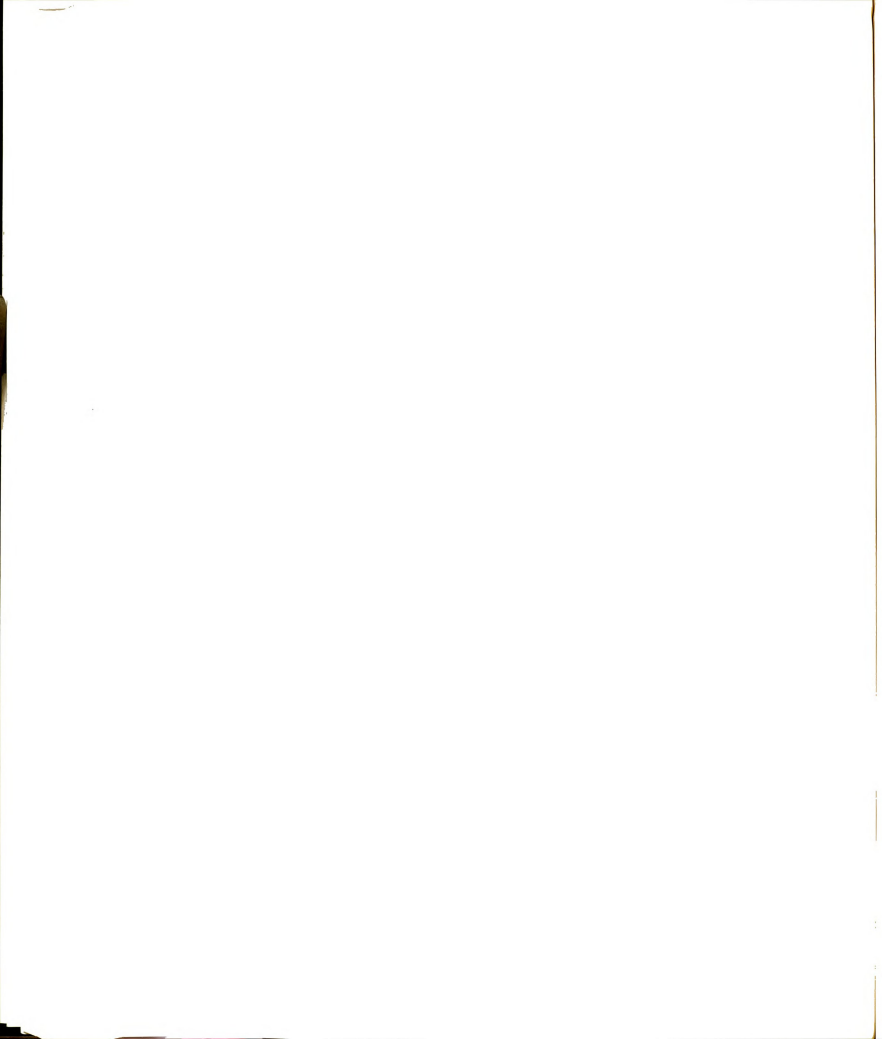
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C*****SUBROUTINE PCURVE*****
SUBROUTINE PCURVE(EINC, N, NX, XCEN, MAG)
COMPLEX EINC(30)
DELEN=MAG(30), MEINC(17)
CALL PLOT(XCEN(30), IBUF(385), 0)
DO 10 I=1, NX
  MEINC(I)=CABS(EINC(I))
  X(I)=XCEN(I)
10 CONTINUE
  N=NEN/2
  LEE=JEN*2
  IF(LEE.NE. NZ) JEN=JEN+1
  DO 30 J=1, JEN
    DO 20 K=1, NX
      TMA(K, J)=MAG(L)
20 CONTINUE
  DO 40 J=1, JEN
    CALL SCALE(X, S, 0, NX, I)
    CALL MSCALE(TMAG(1, I), 0, NX, I)
40 CONTINUE
    DO 50 JJ=1, JEN
      TMA(NX+2, JJ)=MEINC(NX+1)
      TMA(NX+2, JJ)=MEINC(NX+2)
50 CONTINUE
    CALL AXIS(0, 0, 0, 11HX IN METERS, -11.5, 0, 0, X(NX+1), X(NX+2))
    PWEINC(NX+1), MEINC(NX+2))
    CALL LK(X, MEINC, NX, 1, 1, 3)
    DO 60 LK=1, JEN
      CALL LINE(X, TMA(1, KK), NX, 1, 0, 0)
60 CONTINUE
    CALL PLOT(12, 0, 0, 0, 999)
  END
END
C*****SUBROUTINE PCURVE*****
SUBROUTINE PCURVE(N, NX, XCEN, PHA)
DIMENSION PHA(30), XCEN(30), IBUF(385), X(17), TPHA(17, 3)
CALL PLOT(XCEN(30), IBUF(385), 0)
X(1)=XCEN(1)
10 CONTINUE

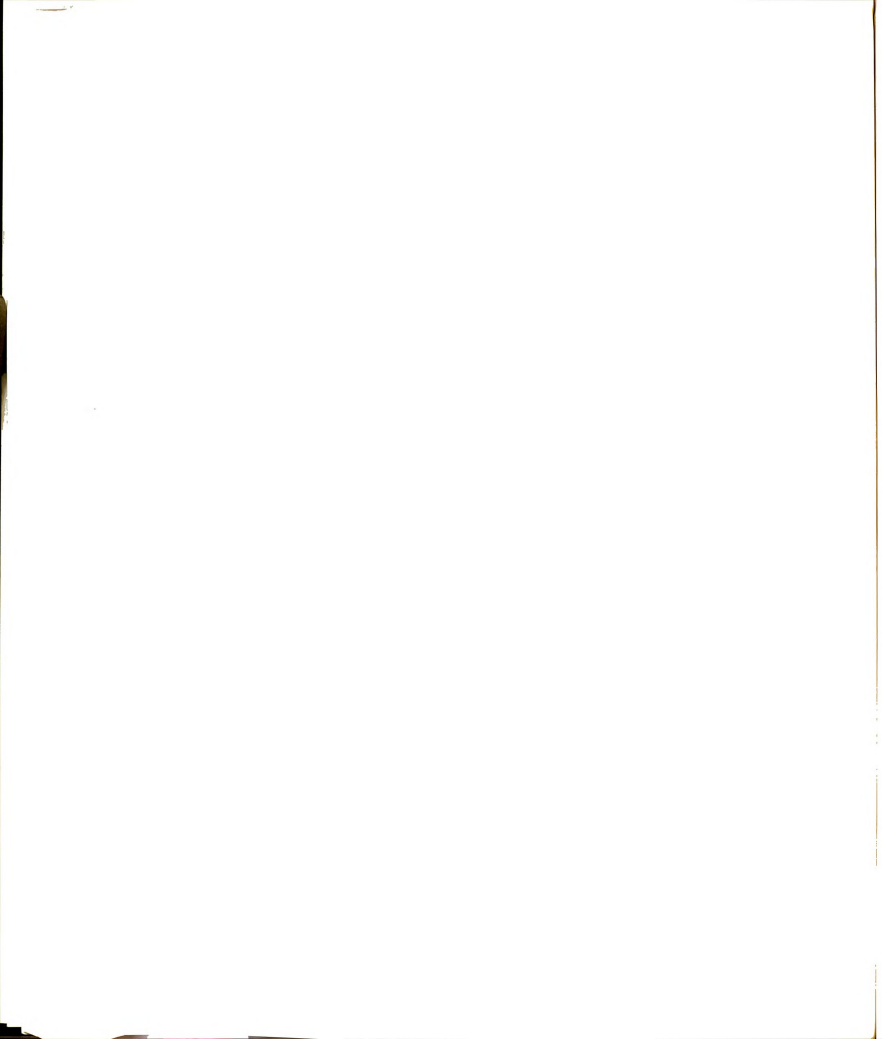
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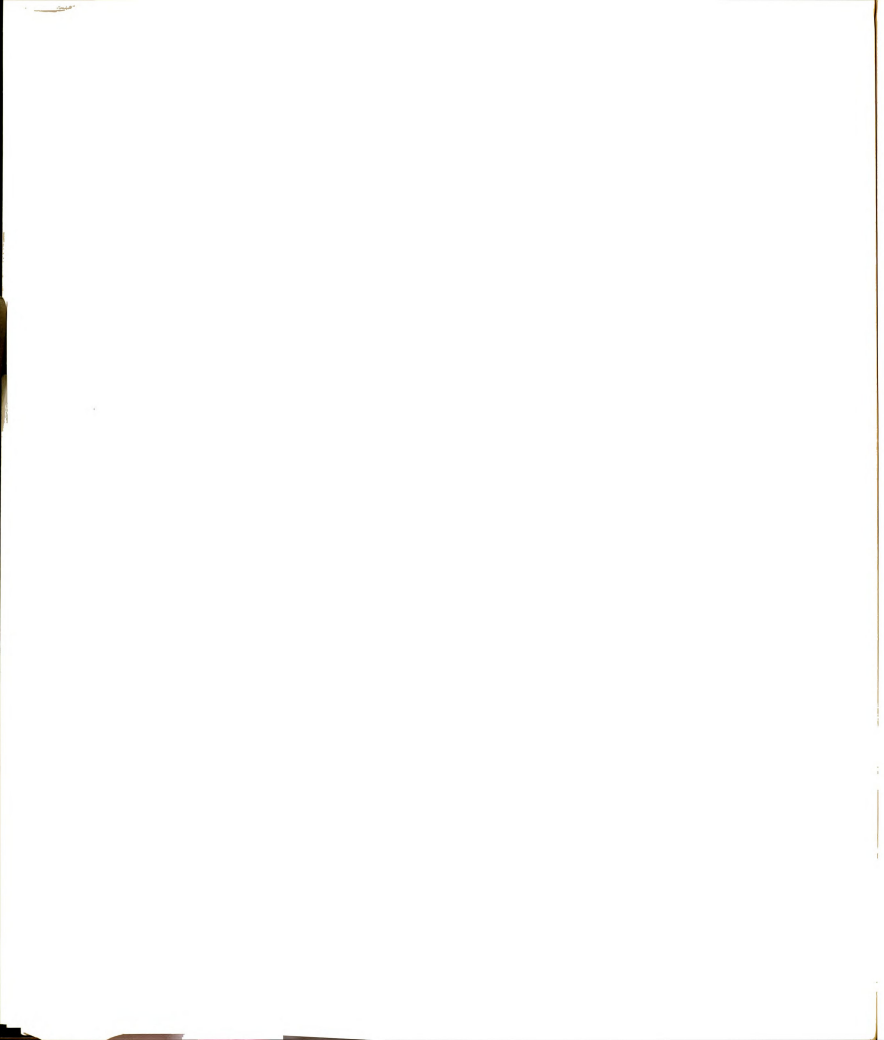
NZ=N/NX
JEN=NZ/2
LEE=JEN*2
IF(LEE.NE.NZ) JPH=JEN+1
JENM1=JEN-1
DO 30 K=1,JN
  DO 20 K=1,NX
    LEEQ(1)=NX*HK
    TPHA(1)=TPHA(1)+90.0
  20 CONTINUE
  30 CONTINUE
  CALL PLOT(0.0,0.5,-3)
  CALL SCALE(X,5.0,NX,1)
  DO (JEN,EG,1) GO TO 45
  DO 40 J=1,JENM1
    CALL MSCALE(TPHA(1,1),0,NX,1)
  40 CONTINUE
  45 CALL SCALE(TPHA(1,JEN),6.0,NX,1)
  IF(JEN.EQ.1) GO TO 55
  DO 50 J=1,JENM1
    TPHA(NX+1,J)=TPHA(NX+2,JEN)
    TPHA(NX+2,J)=TPHA(NX+2,JEN)
  50 CONTINUE
  55 CALL AXIS(0.0,0.0,11HX IN METERS,-1.5,0.0,X(NX+1),X(NX+2))
  CALL AXIS(0.0,0.0,37HREL. PHASE (O FOR INC. FIELD) IN DEG.,37,6,0,
  20.0,TPHAINX,JEN),TPHA(NX+2,JEN))
  DO 60 KK=1,NX
    CALL LINE(X,TPHA(1,KK),NX,1,0,0)
  60 CONTINUE
  CALL PLOT(12.0,0.0,999)
  RETURN
END

```



66

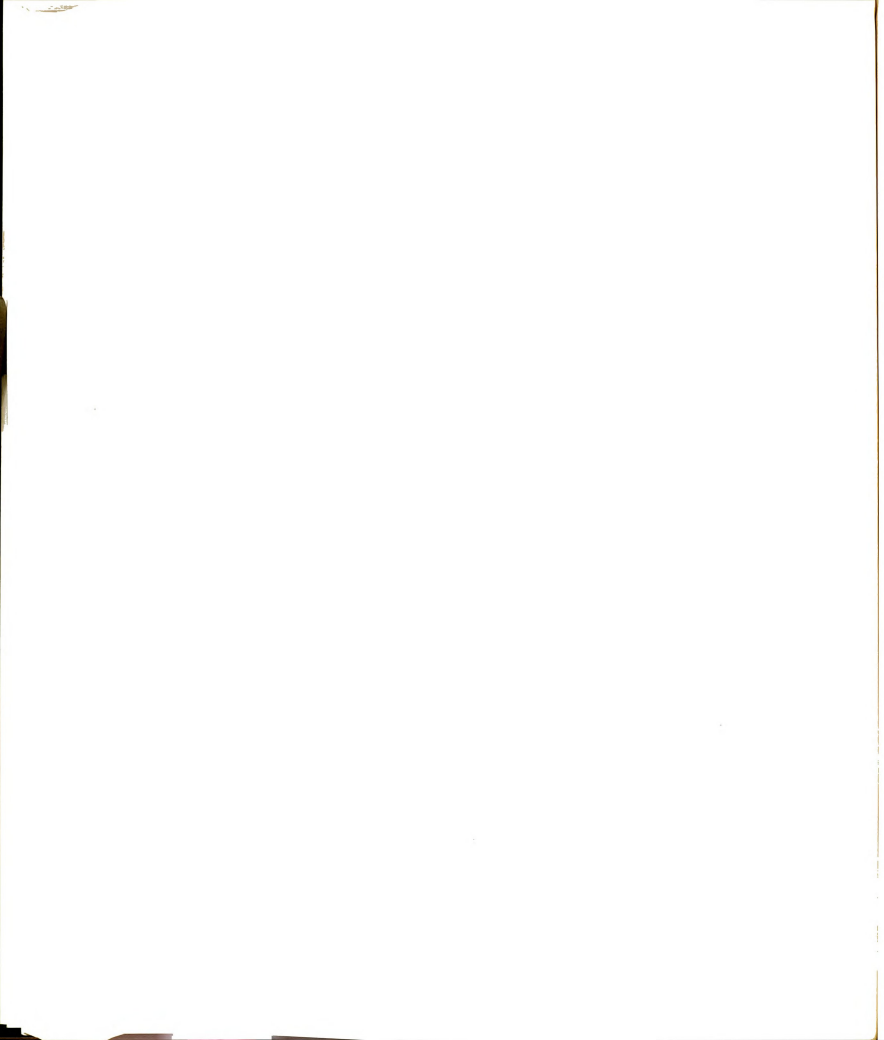
38

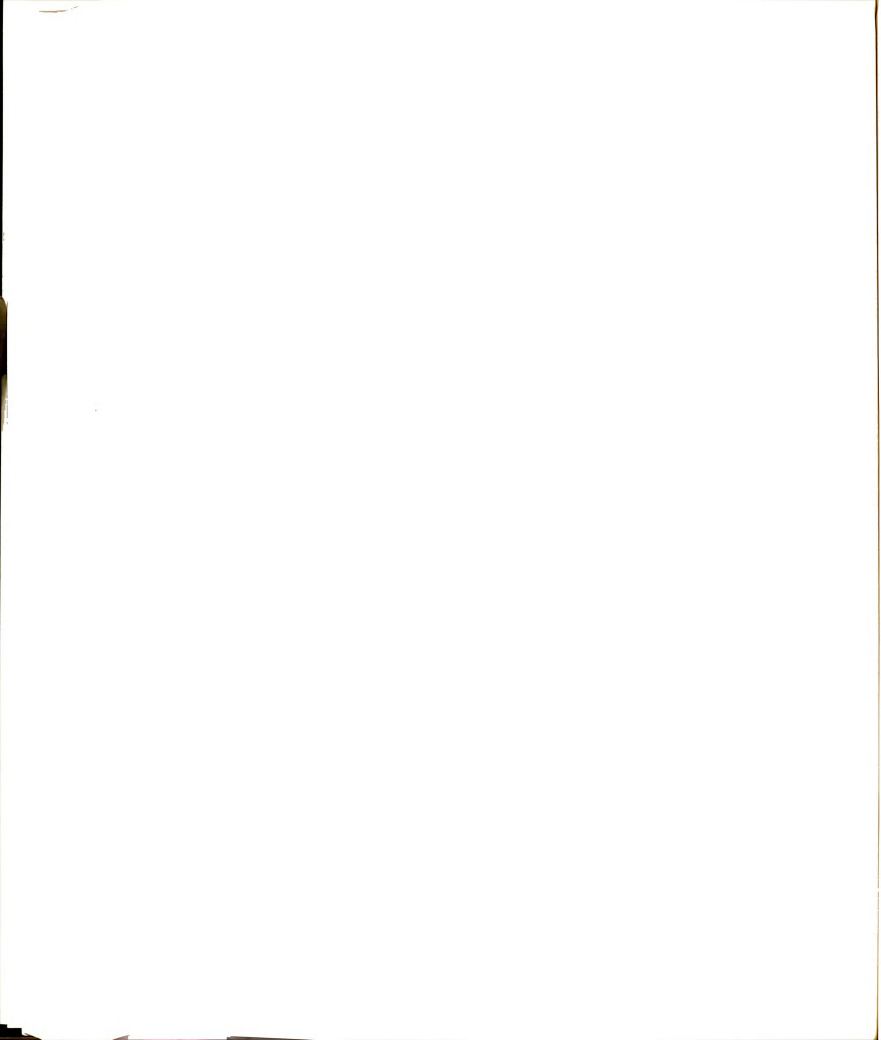


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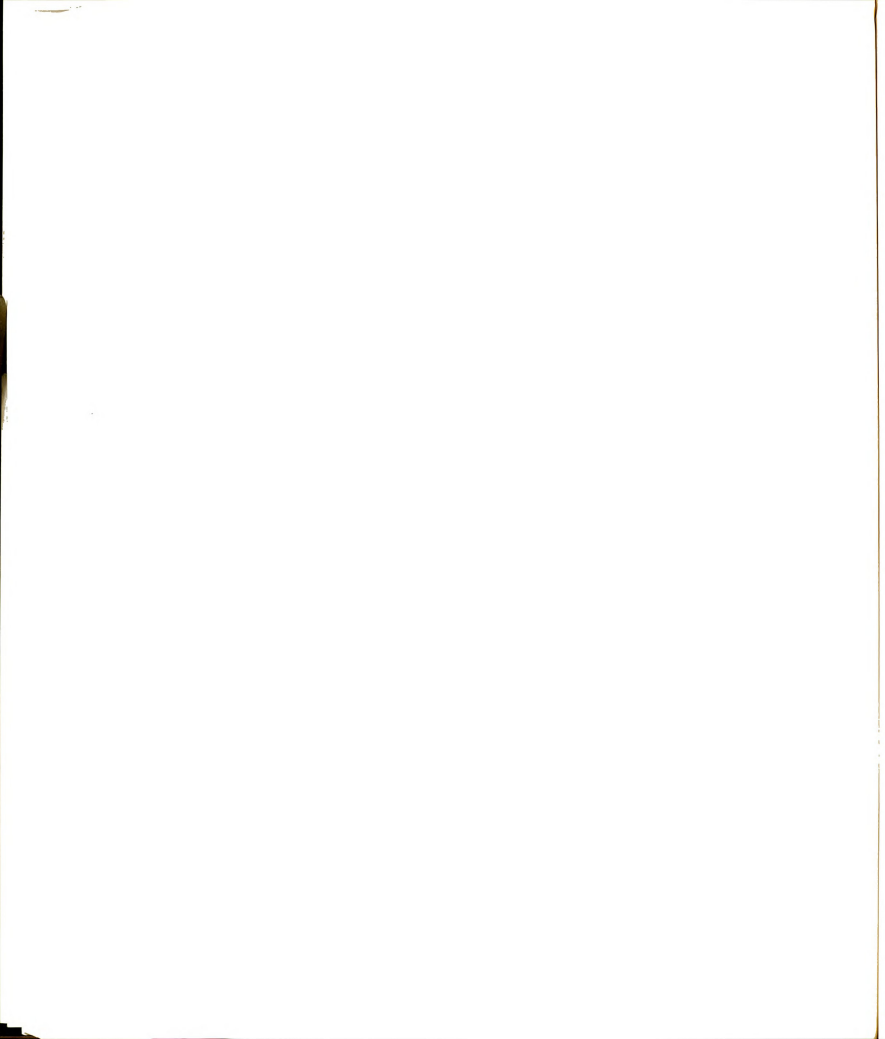
999          PRINT 999, MTIME
          GO TO 9999
          FORMAT(1H1, LUX, 2HILLEGAL IMPRESSED NUDE, 1X, A4, 1X,
+44HAS BEEN APPLIED; PROBLEM CAN NOT BE HANDLED)
          CONTINUE
          CALL MNSR(NCELL, X, Y, Z, DCELL, RLEP, SIG, FNEG, FREQ, OMEG,
+40, ARE, NCELL, L, MAT, MATP)
          CALL MATP1(X, Y, Z, ARE, NCELL, L, MAT, MATP)
          IF (MODE, EQ, EXPZ) AND, IQ, EQ, 1234) CALL, EXIM(L, MAT, MATP, Z, NCELL)
          IF (MODE, EQ, COSH) CALL, CHIM(L, MAT, MATP, Z, NCELL)
          IF (MODE, EQ, SINE) CALL, SEIM(L, MAT, MATP, Z, NCELL)
          IF (MODE, EQ, SINE) CALL, SEIM(L, MAT, MATP, Z, NCELL)
          CALL CRAT(L, MAT, MATP, SHC(L, MAT, MATP, Z, NCELL)
          CALL CRAT(L, MAT, MATP, SHC(L, MAT, MATP, Z, NCELL)
          CALL PRINT(INDU, CRNT, L, MAT, MATP)
          CALL T(L, MAT, MATP)
          CALL PRINT(INDU, FLD, L, MAT, MATP)
          IF (MODE, EQ, EXPZ) GO TO 9999
          IF (MODE, EQ, COSH) MTAPE=2
          IF (MODE, EQ, SINE) MTAPE=3
          IF (MODE, EQ, SINE) MTAPE=3
          IF (MODE, EQ, SINE) MTAPE=4
          BUFFER, OUT(MTAPE, 1) (L(1, MATP), L(1, MATP, MATP))
          DO 6, I=1, MAT
          CALL L, L, MATP
          PUNCH 7, L(1, MATP)
          CONTINUE
          6
          9999
          FORMAT(2E13, 6)
          CONTINUE
          END
          C***** BLOCK DATA*****
          BLOCK DATA
          REAL IMPR
          COMMON /IMPR/ IMPR, CRNT
          COMMON /IMPR/ IMPR, CRNT, SINE, SINH, EXPZ
          DATA IMPR/4H1IMPZ/, CRNT/4HCENT/
          DATA COSE/4HCOS/, COSH/4HCOSH/, SINE/4HSINE/, SINH/4HSINH/
          DATA EXPZ/4HEXPZ/
          END
          C***** SUBROUTINE MNSB*****
          SUBROUTINE MNSB(NCELL, X, Y, Z, DCELL, RLEP, SIG, FNEG, FREQ, OMEG,
+40, ARE, NCELL, L, MAT, MATP)
          REAL X, Y, Z, DCELL, RLEP, SIG, FNEG, FREQ, OMEG,
+40, ARE, NCELL, L, MAT, MATP
          COMPLEX FK, C, J, FHUN
          COMMON IQ
          590
          600
          610
          620
          630
          640

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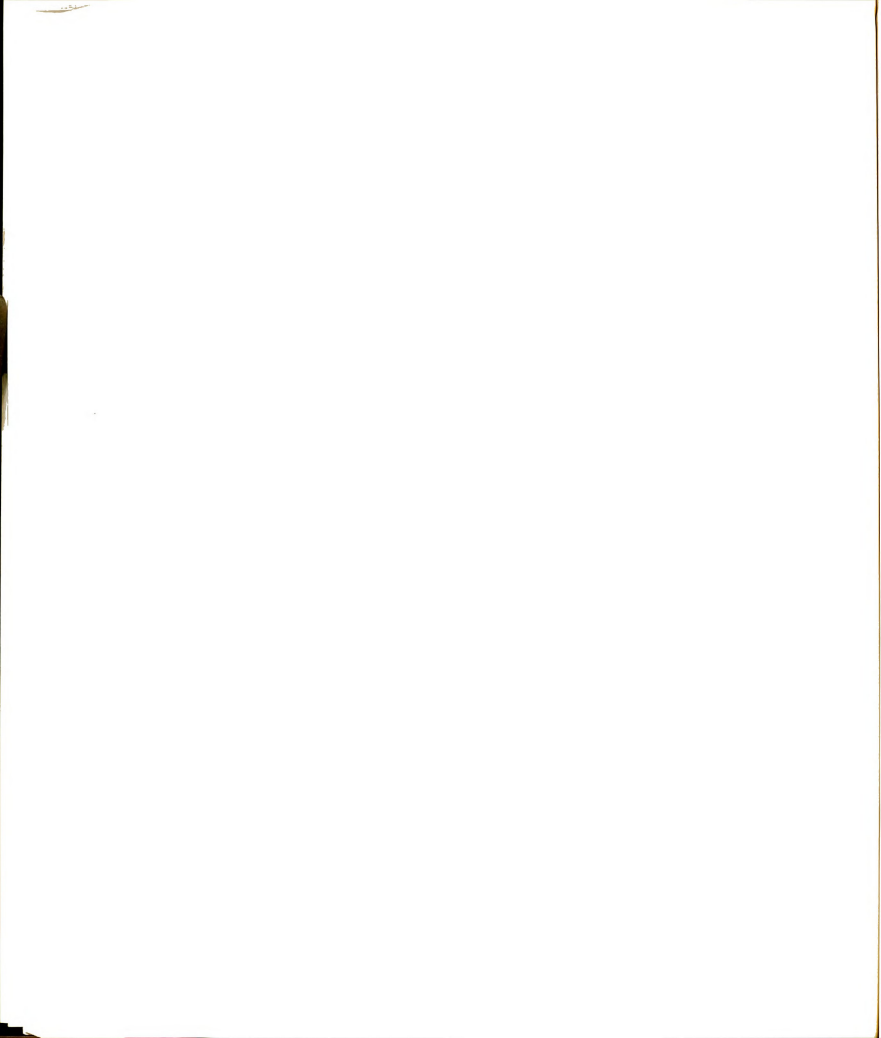




[illegible]



201

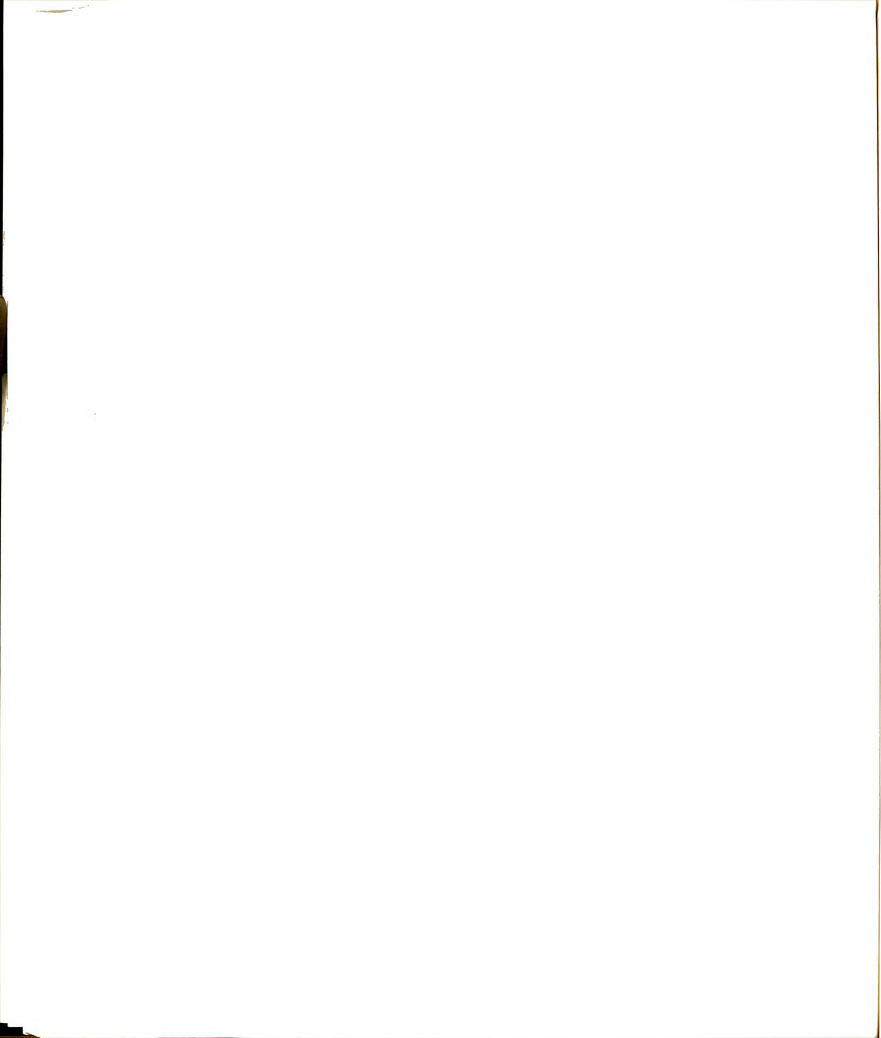


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GO TO 4
3 CONTINUE
  XN=X(MM)
  YN=Y(MM)
  ZM=Z(MM)
  XN=XN
  YN=YN
  ZN=ZN
  ARE=AREA(MM)
  AREA=AREA/PI
  P12=CJ*(P(2,A)+PI*P(1,A))
  PIM=PI
  FHUNM=FUHUN
  IF(MDIR.EG.4) PIM=-1.0*PI
  IF(MDIR.EG.4) FHUNM=-1.0*FHUN
  E(1)=4.0*PI*H
  E(2)=0.0
  E(3)=0.0
  E(4)=P12
  E(5)=0.0
  E(6)=E(1)
  E(7)=-1.0*P12
  E(8)=0.0
  E(9)=0.0
  E(10)=E(7)
  E(11)=FHUNM
  E(12)=0.0
  E(13)=P12
  E(14)=0.0
  E(15)=0.0
  E(16)=FHUNM
4 CONTINUE
  IF(IG.EG.1234) YN=-1.0*YN
  IF(IG.EG.1234) GO TO 102
  IF(IG.EG.1234) ZN=ZN
  CALL COEFF(MDIR,AREA,XM,YM,ZM,XN,YN,ZN,C)
  IF(MODE.EG.CDSH.OR.MODE.EG.SINE) GO TO 9
  DO 6 I=1,4
    E(I)=E(1)+C(I)
  6 CONTINUE
  DO 7 I=1,2
    E(I)=E(1)-C(I)
  7 CONTINUE
  DO 8 I=1,16
    E(I)=E(1)+C(I)
  8 CONTINUE
  GO TO 21

```

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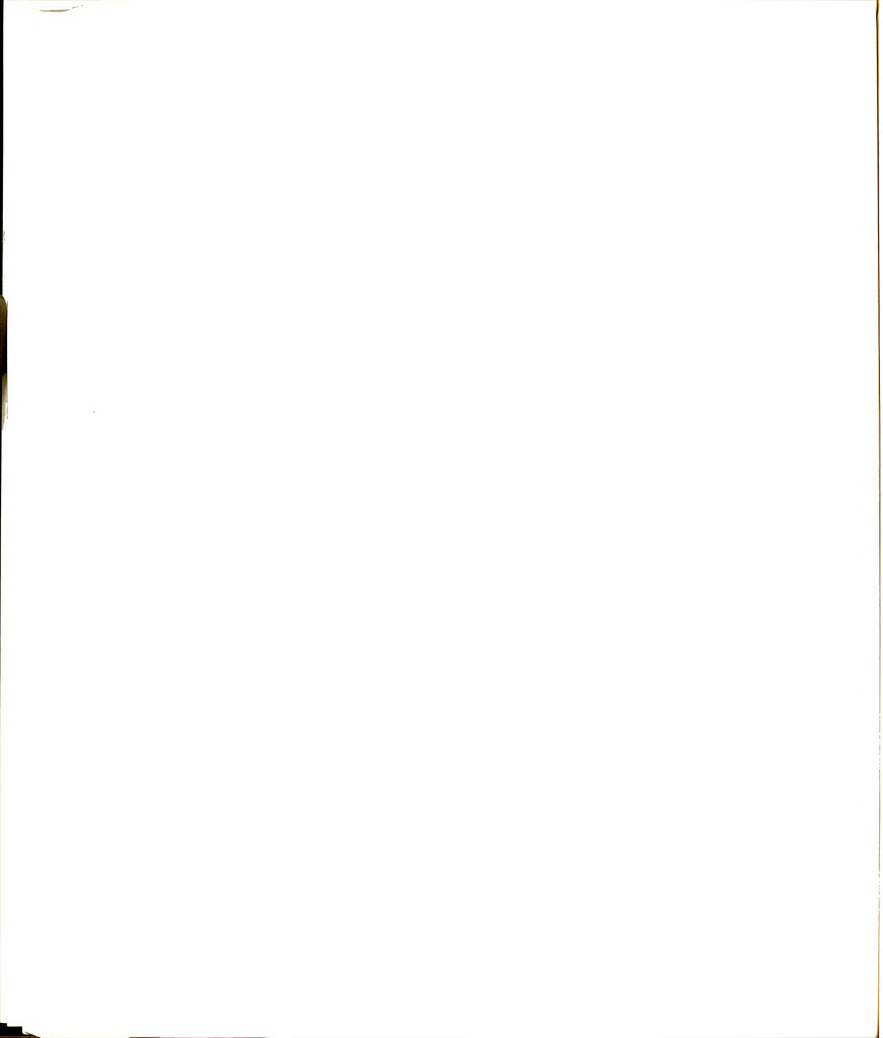


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9  CONTINUE
DO 10 I=1,4
E(I)=E(I)-C(I)
CONTINUE
10
DO 11 I=5,12
E(I)=E(I)+C(I)
CONTINUE
11
DO 12 I=13,16
E(I)=E(I)-C(I)
CONTINUE
12
CONTINUE
21
IF(IG.EQ.1256) ZN=-1.0*ZN
IF(IG.EQ.1256) YN=-1.0*YN
IF(IG.EQ.1256) GO TO 101
YN=-1.0*YN
CALL COEFX(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 24
DO 23 I=1,16
E(I)=E(I)-C(I)
CONTINUE
23
GO TO 32
24
DO 25 I=1,16
E(I)=E(I)+C(I)
CONTINUE
25
CONTINUE
30
CONTINUE
ZN=-1.0*ZN
102 CALL COEFX(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
DO 32 I=1,4
E(I)=E(I)-C(I)
CONTINUE
32
DO 33 I=5,12
E(I)=E(I)+C(I)
CONTINUE
33
DO 34 I=13,16
E(I)=E(I)-C(I)
CONTINUE
34
IF(IG.EQ.1458) GO TO 100
YN=-1.0*YN
101 CALL COEFX(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
DO 49 I=1,8
E(I)=E(I)-C(I)
CONTINUE
49
DO 50 I=9,16
E(I)=E(I)+C(I)
CONTINUE
50

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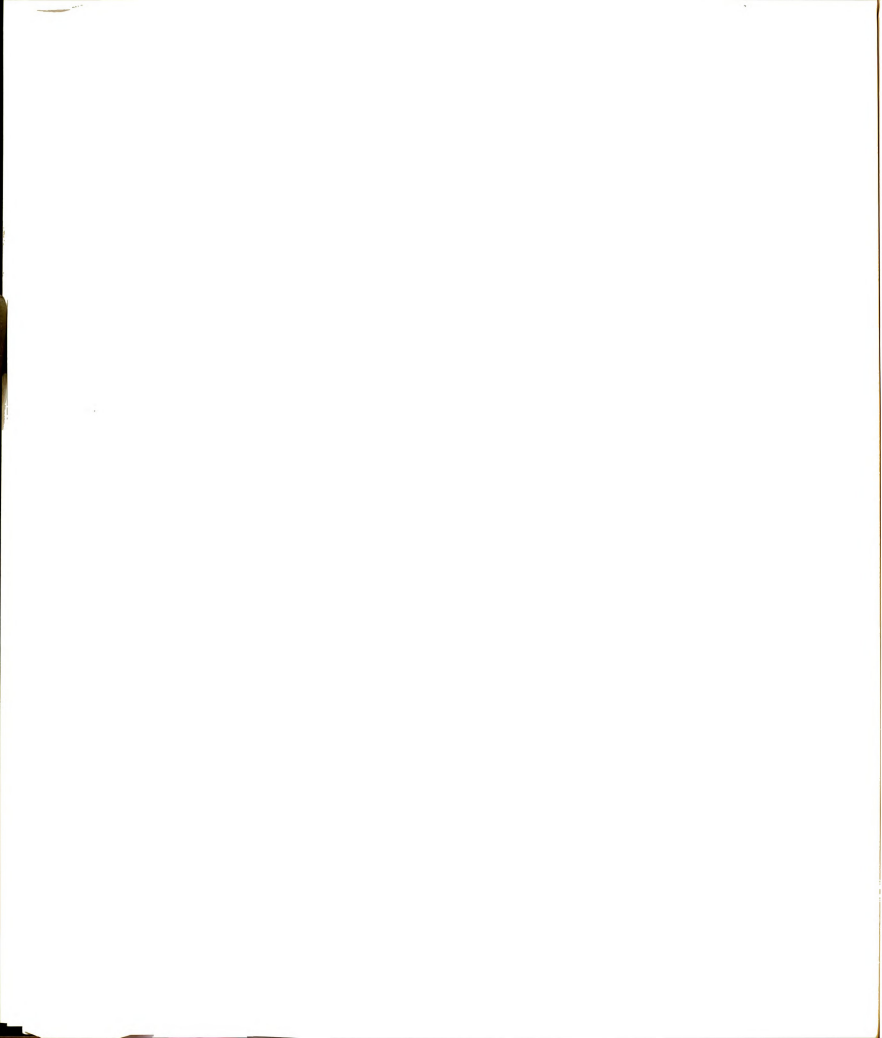



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IF(IQ.EQ.1234) YN=-1.0*YN
IF(IQ.EQ.1234) GO TO 103
ZN=-1.0*ZN
CALL COEFX(MDIR,AREA,XM,YM,ZM,XN,YN,ZN,C)
IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 66
DO 62 I=1,4
E(I)=E(I)-C(I)
62 CONTINUE
DO 63 I=5,8
E(I)=E(I)+C(I)
63 CONTINUE
DO 64 I=9,12
E(I)=E(I)-C(I)
64 CONTINUE
DO 65 I=13,16
E(I)=E(I)+C(I)
65 CONTINUE
GO TO 81
66 CONTINUE
DO 67 I=1,4
E(I)=E(I)+C(I)
67 CONTINUE
DO 68 I=5,8
E(I)=E(I)-C(I)
68 CONTINUE
DO 69 I=9,12
E(I)=E(I)+C(I)
69 CONTINUE
DO 70 I=13,16
E(I)=E(I)-C(I)
70 CONTINUE
81 IF(IQ.EQ.1256) GO TO 100
YN=-1.0*YN
CALL COEFX(MDIR,AREA,XM,YM,ZM,XN,YN,ZN,C)
IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 85
DO 83 I=1,8
E(I)=E(I)+C(I)
83 CONTINUE
DO 84 I=9,16
E(I)=E(I)-C(I)
84 CONTINUE
GO TO 94
85 CONTINUE
DO 86 I=1,8
E(I)=E(I)-C(I)
86 CONTINUE

```

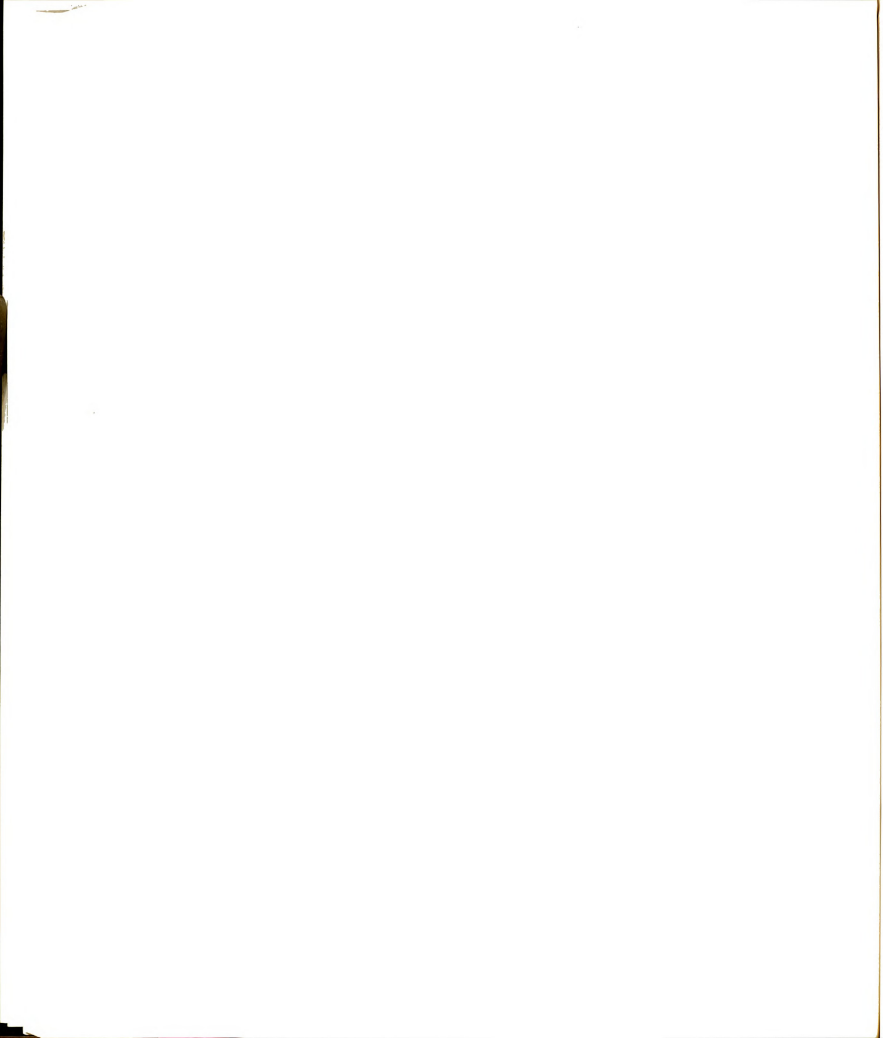
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```

DO 87 I=9,16,
E(I)=E(I)+C(I)
87 CONTINUE
94 CONTINUE
ZN=-1.0*ZN
103 CALL COEFX(MDIR,AREA,XM,YM,ZM,XN,YN,ZN,C)
DO 96 I=1,4
E(I)=E(I)+C(I)
96 CONTINUE
DO 97 I=5,8
E(I)=E(I)-C(I)
97 CONTINUE
DO 98 I=9,12
E(I)=E(I)+C(I)
98 CONTINUE
DO 99 I=13,16
E(I)=E(I)-C(I)
99 CONTINUE
100 CONTINUE
RETURN
END
C*****
SUBROUTINE ELEM(MDIR,NDIR,M,N,X,Y,Z,ARE,NCELL,E)
REAL X(NCELL),Y(NCELL),Z(NCELL),ARE(NCELL),MODE,MUO
COMPLEX E(16),C(16),PIZ,CJ,P,PHUN,FK,PHUNM
COMMON /GX7/NX,NY,NZ,NNX,NNY,NNZ
COMMON /CONS/P1,MUO,EPO
COMMON /CPLV/C1,PHUN
COMMON /PARA/FA,FK
COMMON /MODS/MODE
COMMON /CSEH/COSE,COSH,SINE,SINH,EXPZ
IF(MDIR.EQ.1) MM=M
IF(MDIR.EQ.2) MM=M+NX
IF(MDIR.EQ.3) MM=M+NX+NY
IF(MDIR.EQ.4) MM=M+NX+NY+NZ
IF(MDIR.EQ.5) MM=M+NX+NY+NZ+NNX
IF(MDIR.EQ.6) MM=M+NX+NY+NZ+NNX+NNY
IF(MDIR.EQ.2) NN=N+NX
IF(MDIR.EQ.5) NN=N+NX+NY+NZ+NNX
IF(MDIR.NE.NDIR) GO TO 20J
IF(P.EQ.N) GO TO 20E
20I XM=X(MM)
20J YM=Y(MM)
20E ZM=Z(MM)
20F XN=X(NN)
20G YN=Y(NN)
20H
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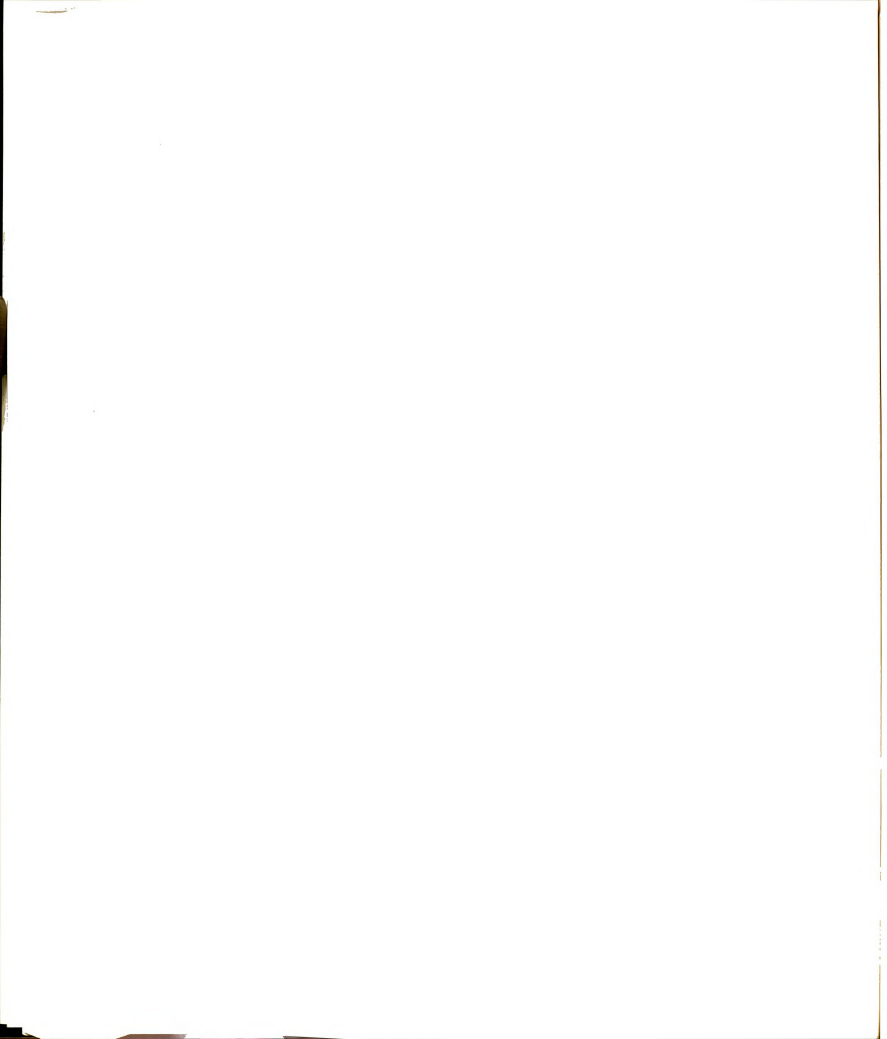
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```

3810 ZN=Z(NN)
3820 AREA=ARE(NN)
3830 CALL COEFY(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
3840 DO 2 I=1, 16
3850 E(I)=C(I)
3860 2 CONTINUE
3870 GO TO 203
3880 202 CONTINUE
3890 XM=X(NM)
3900 YM=Y(NM)
3910 ZM=Z(NM)
3920 XN=XM
3930 YN=YM
3940 ZN=ZM
3950 AREA=ARE(NM)
3960 AG=AG+AREA/PI)
3970 P1=CU*(P(2, A)+PI*P(1, A))
3980 P1M=P1
3990 FHUNM=EHUN
4000 IF(MDIR.EQ.5) P1M=-1.0*PI
4010 IF(MDIR.EQ.5) FHUNM=-1.0*FHUN
4020 E(1)=4.0*P1M
4030 E(2)=0.0
4040 E(3)=0.0
4050 E(4)=-1.0*P12
4060 E(5)=0.0
4070 E(6)=E(1)
4080 E(7)=P12
4090 E(8)=0.0
4100 E(9)=0.0
4110 E(10)=P12
4120 E(11)=FHUNM
4130 E(12)=0.0
4140 E(13)=E(4)
4150 E(14)=0.0
4160 E(15)=0.0
4170 E(16)=0.0
4180 E(17)=FHUNM
4190 203 CONTINUE
4200 IF(IQ.EQ.1458) ZN=-1.0*ZN
4210 IF(IQ.EQ.1458) GO TO 300
4220 XN=-1.0*XN
4230 CALL COEFY(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
4240 DO 4 I=1, 4
4250 E(I)=E(I)+C(I)
4260 4 CONTINUE
4270 DO 5 I=5, 12
4280 E(I)=E(I)-C(I)
5 CONTINUE

```

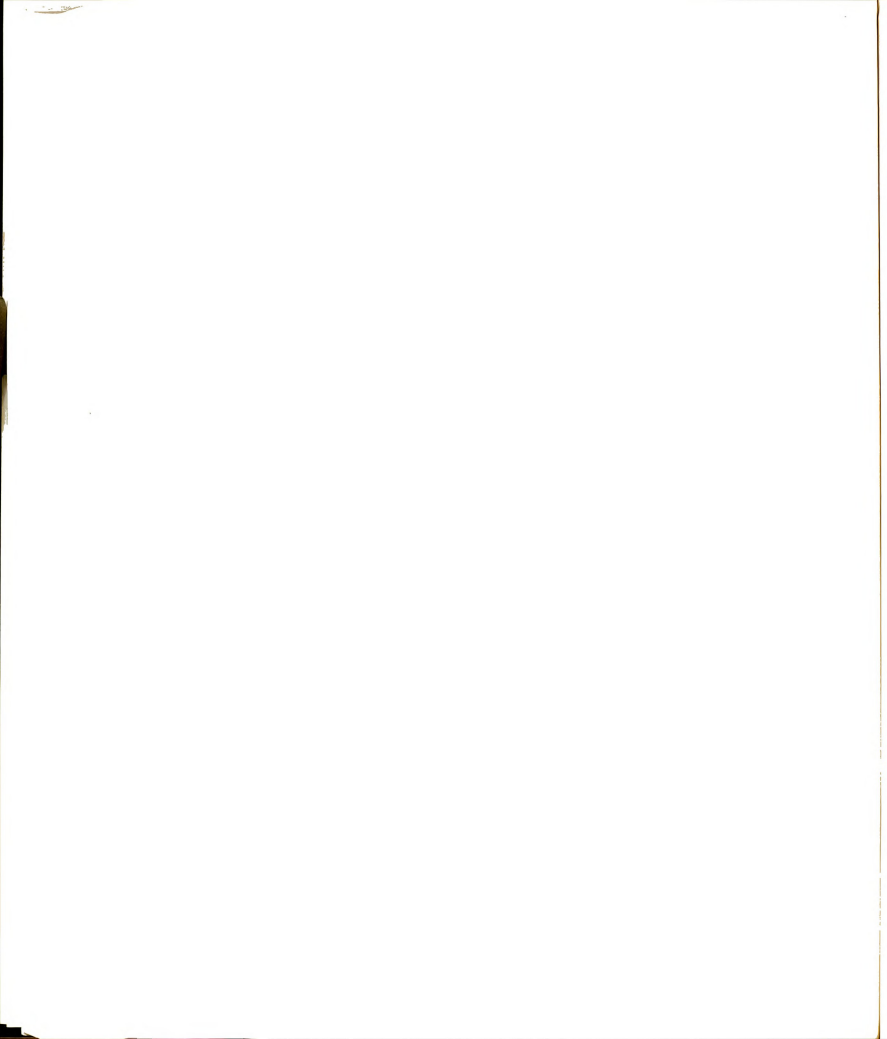


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DO 6 I=13,16,
E(I)=E(I)+C(I)
CONTINUE
6 IF(IQ.EQ.1234) XN=-1.0*YN
IF(IQ.EQ.1234) YN=-1.0*YN
IF(IQ.EQ.1234) GO TO 302
ZN=-1.0*ZN
CALL COEFF(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 22
DO 21 I=1,16,
E(I)=E(I)+C(I)
CONTINUE
21 GO TO 28
22 CONTINUE
DO 23 I=1,16,
E(I)=E(I)-C(I)
CONTINUE
23 CONTINUE
28 CONTINUE
300 XN=-1.0*XN
CALL COEFF(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 33
DO 30 I=1,4,
E(I)=E(I)+C(I)
CONTINUE
30 DO 31 I=5,12,
E(I)=E(I)-C(I)
CONTINUE
31 DO 32 I=13,16,
E(I)=E(I)+C(I)
CONTINUE
32 GO TO 45
33 CONTINUE
DO 34 I=1,4,
E(I)=E(I)-C(I)
CONTINUE
34 DO 35 I=5,12,
E(I)=E(I)+C(I)
CONTINUE
35 DO 36 I=13,16,
E(I)=E(I)-C(I)
CONTINUE
36 IF(IQ.EQ.1236) GO TO 113
45 ZN=-1.0*ZN
302 YN=-1.0*YN
CALL COEFF(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
DO 47 I=1,8

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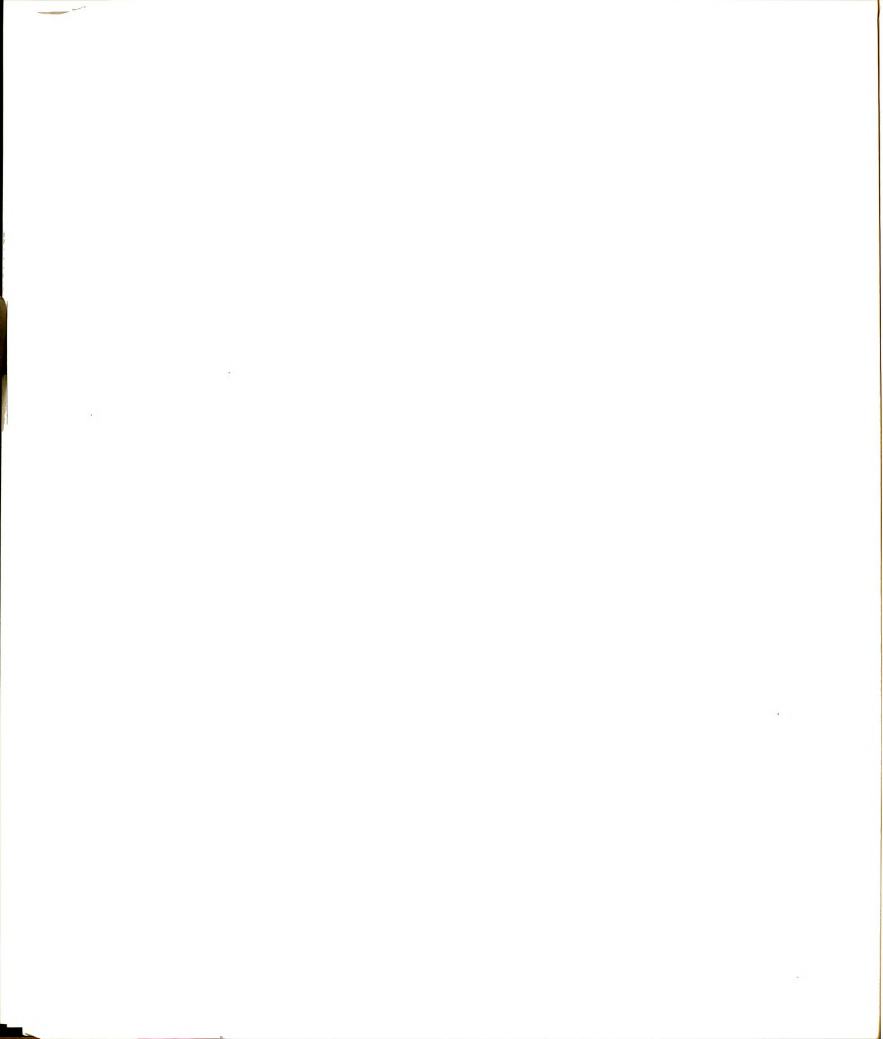



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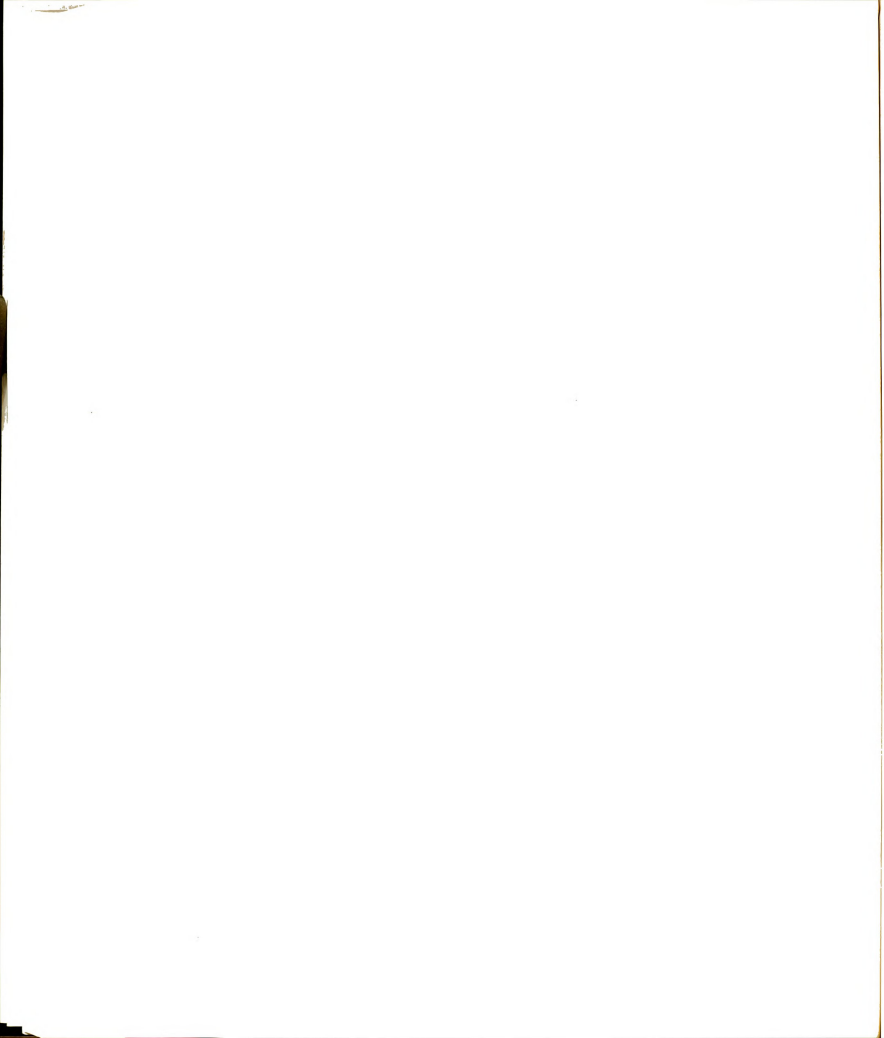
47  E(I)=E(I)+C(I)
    CONTINUE
    DO 48 I=9,16
    48  E(I)=E(I)-C(I)
    CONTINUE
    IF(IQ.EQ.1458) ZN=-1.0*ZN
    IF(IQ.EQ.1458) GO TO 301
    XN=-1.0*XN
    CALL COEFY(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
    DO 60 I=1,4
    60  E(I)=E(I)+C(I)
    DO 61 I=5,8
    61  E(I)=E(I)-C(I)
    CONTINUE
    DO 62 I=9,12
    62  E(I)=E(I)+C(I)
    DO 63 I=13,16
    63  E(I)=E(I)-C(I)
    CONTINUE
    IF(IQ.EQ.1234) GO TO 113
    ZN=-1.0*ZN
    CALL COEFY(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
    IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 83
    DO 81 I=1,8
    81  E(I)=E(I)+C(I)
    CONTINUE
    DO 82 I=9,16
    82  E(I)=E(I)-C(I)
    CONTINUE
    GO TO 92
    83  CONTINUE
    DO 84 I=1,8
    84  E(I)=E(I)-C(I)
    CONTINUE
    DO 85 I=9,16
    85  E(I)=E(I)+C(I)
    CONTINUE
    92  CONTINUE
    XN=-1.0*XN
    301 CALL COEFY(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
    IF(MODE.EQ.COSH.OR.MODE.EQ.SINE) GO TO 98
    DO 94 I=1,4
    94  E(I)=E(I)+C(I)
    CONTINUE
    DO 95 I=5,8
    95  E(I)=E(I)-C(I)

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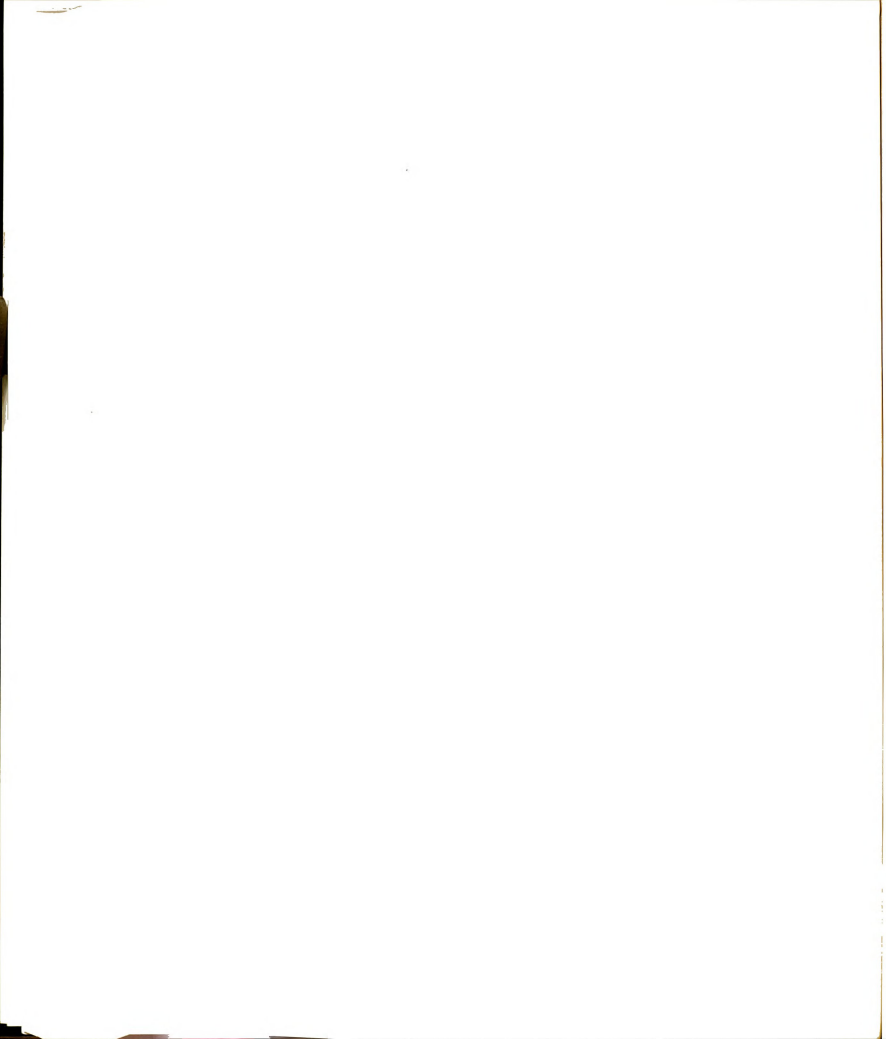
[illegible]



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5680 ZM=Z(MM)
5690 XN=X(NN)
5700 YN=Y(NN)
5710 ZN=Z(NN)
5720 AREA=ARE(NN)
5730 CALL COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
5740 DO 2 I=1, 16
5750 E(I)=C(I)
5760 GO TO 203
5770 CONTINUE
5780 XM=X(MM)
5790 YM=Y(MM)
5800 ZM=Z(MM)
5810 XN=X(MM)
5820 YN=Y(MM)
5830 ZN=Z(MM)
5840 AREA=ARE(MM)
5850 ARECO=(AREA/PI)
5860 PI=3.141592653589793
5870 PIM=PI
5880 FHUM=EHUM
5890 IF(MDIR.EQ.6) PIM=-1.0*PI
5900 IF(MDIR.EQ.6) FHUM=-1.0*FHUM
5910 E(1)=4.0*PIM
5920 E(2)=0.0
5930 E(3)=0.0
5940 E(4)=P12
5950 E(5)=0.0
5960 E(6)=E(1)
5970 E(7)=-1.0*P12
5980 E(8)=0.0
5990 E(9)=0.0
6000 E(10)=0.0
6010 E(11)=EHUM
6020 E(12)=0.0
6030 E(13)=P12
6040 E(14)=0.0
6050 E(15)=0.0
6060 E(16)=FHUM
6070 CONTINUE
6080
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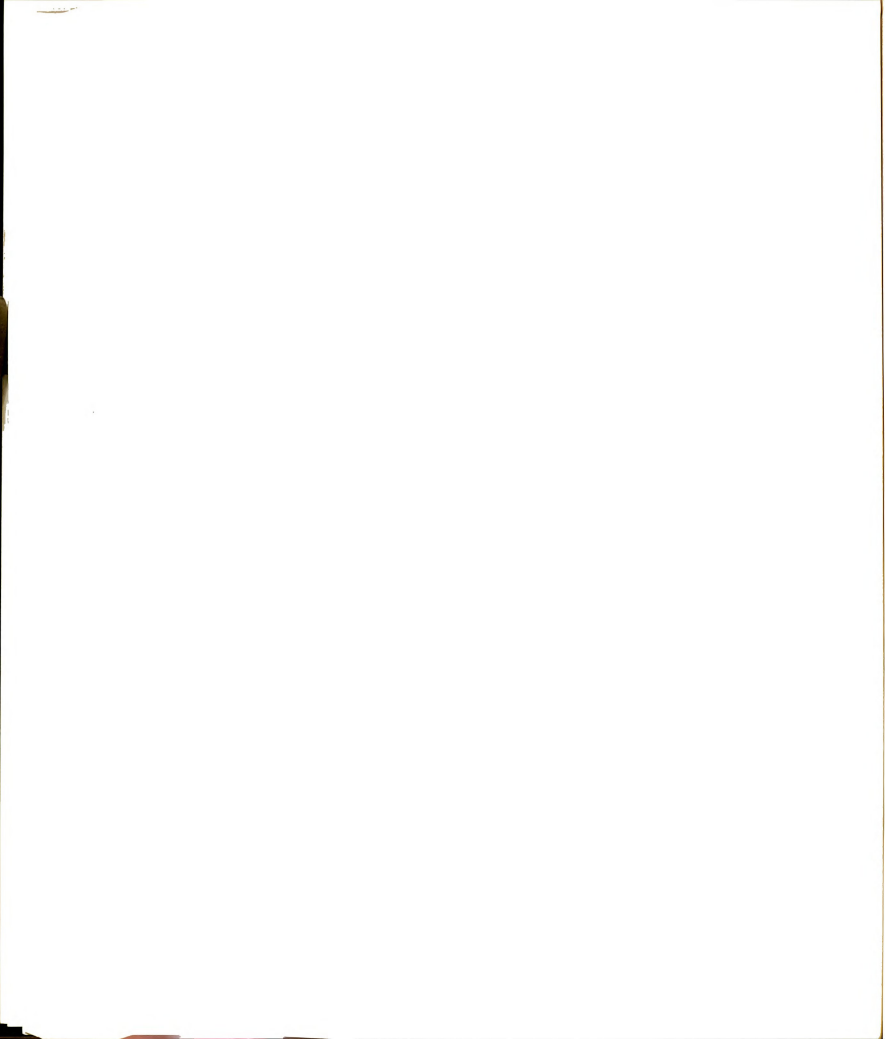


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DO 5 I=5,12
  E(I)=E(I)-C(I)
5  CONTINUE
DO 6 I=13,16
  E(I)=E(I)+C(I)
6  CONTINUE
  XN=-1.0*ZN
  IF(IG.EQ.1256) ZN=-1.0*ZN
  IF(IG.EQ.1256) GO TO 302
  YN=-1.0*YN
  CALL COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
  E(I)=E(I)+C(I)
DO 21 I=1,16
21 CONTINUE
  XN=1.0*ZN
  DO 30 I=1,16
  30  CALL COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
  E(I)=E(I)+C(I)
  CONTINUE
DO 31 I=5,12
  E(I)=E(I)-C(I)
31 CONTINUE
DO 32 I=13,16
  E(I)=E(I)+C(I)
32 CONTINUE
  IF(IG.EQ.1234) GO TO 113
  ZN=-1.0*ZN
  YN=-1.0*YN
302 CALL COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
  IF(MODE.EQ.COSH. OR. MODE.EQ.SINE) GO TO 49
DO 47 I=1,16
  47  E(I)=E(I)+C(I)
  CONTINUE
DO 48 I=9,16
  E(I)=E(I)-C(I)
48 CONTINUE
GO TO 58
49 CONTINUE
DO 50 I=1,16
  E(I)=E(I)-C(I)
50 CONTINUE
DO 51 I=9,16
  E(I)=E(I)+C(I)
51 CONTINUE
  YN=-1.0*YN
58 CONTINUE
  IF(IG.EQ.1458) GO TO 301
  XN=-1.0*ZN

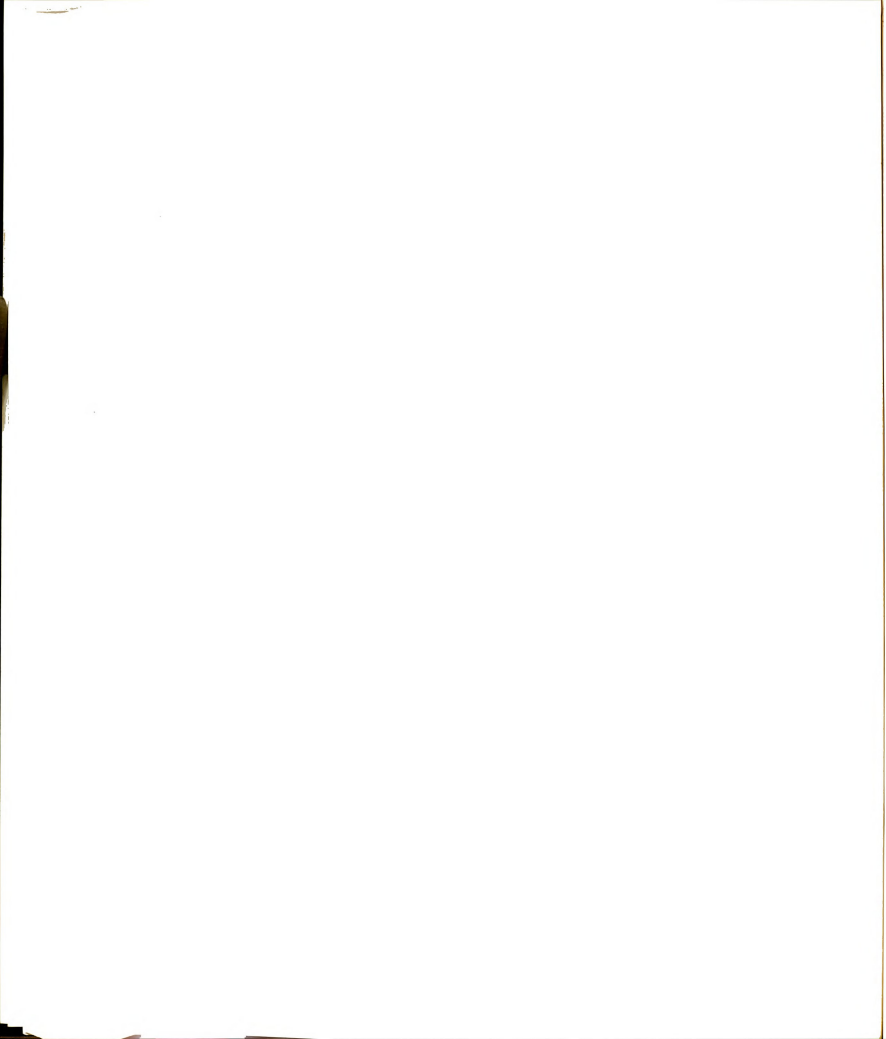
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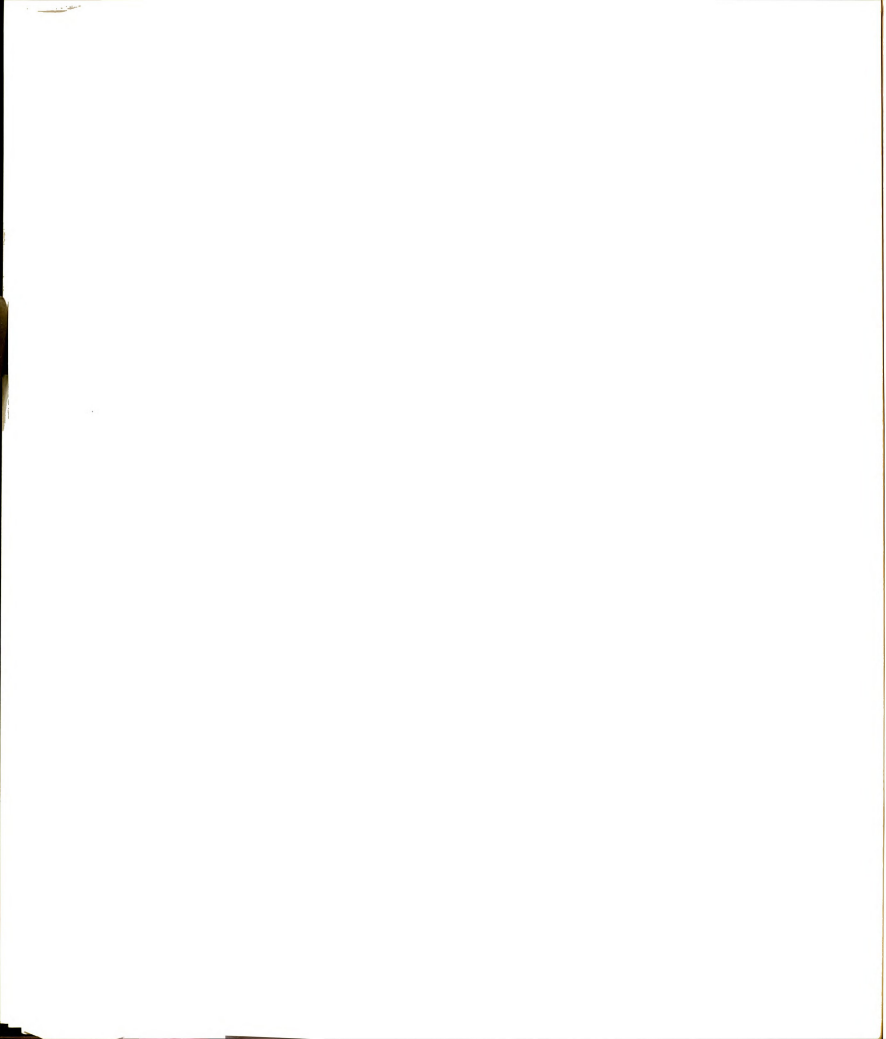
CALL COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
IF(MODE, EQ, COSH, OR, MODE, EQ, SINE) GO TO 84
DO 60 I=1,4
  E(I)=E(I)+C(I)
CONTINUE
60
DO 61 I=5,8
  E(I)=E(I)-C(I)
CONTINUE
61
DO 62 I=9,12
  E(I)=E(I)+C(I)
CONTINUE
62
DO 63 I=13,16
  E(I)=E(I)-C(I)
CONTINUE
63
DO 64 I=17,20
  E(I)=E(I)+C(I)
CONTINUE
64
DO 65 I=21,24
  E(I)=E(I)-C(I)
CONTINUE
65
DO 66 I=25,28
  E(I)=E(I)+C(I)
CONTINUE
66
DO 67 I=29,32
  E(I)=E(I)-C(I)
CONTINUE
67
DO 68 I=33,36
  E(I)=E(I)+C(I)
CONTINUE
68
DO 69 I=37,40
  E(I)=E(I)-C(I)
CONTINUE
69
IF(IQ, EQ, 1256) GO TO 113
YN=-1.0*YN
CALL COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
IF(MODE, EQ, COSH, OR, MODE, EQ, SINE) GO TO 83
DO 81 I=1,8
  E(I)=E(I)+C(I)
CONTINUE
81
DO 82 I=9,16
  E(I)=E(I)-C(I)
CONTINUE
82
DO 83 I=17,24
  E(I)=E(I)+C(I)
CONTINUE
83
DO 84 I=25,32
  E(I)=E(I)-C(I)
CONTINUE
84
DO 85 I=33,40
  E(I)=E(I)+C(I)
CONTINUE
85
DO 86 I=41,48
  E(I)=E(I)-C(I)
CONTINUE
86
DO 87 I=49,56
  E(I)=E(I)+C(I)
CONTINUE
87
DO 88 I=57,64
  E(I)=E(I)-C(I)
CONTINUE
88
DO 89 I=65,72
  E(I)=E(I)+C(I)
CONTINUE
89
DO 90 I=73,80
  E(I)=E(I)-C(I)
CONTINUE
90
DO 91 I=81,88
  E(I)=E(I)+C(I)
CONTINUE
91
DO 92 I=89,96
  E(I)=E(I)-C(I)
CONTINUE
92
DO 93 I=97,104
  E(I)=E(I)+C(I)
CONTINUE
93
DO 94 I=105,112
  E(I)=E(I)-C(I)
CONTINUE
94
DO 95 I=113,120
  E(I)=E(I)+C(I)
CONTINUE
95
DO 96 I=121,128
  E(I)=E(I)-C(I)
CONTINUE
96
DO 97 I=129,136
  E(I)=E(I)+C(I)
CONTINUE
97
DO 98 I=137,144
  E(I)=E(I)-C(I)
CONTINUE
98
DO 99 I=145,152
  E(I)=E(I)+C(I)
CONTINUE
99
DO 100 I=153,160
  E(I)=E(I)-C(I)
CONTINUE
100
DO 101 I=161,168
  E(I)=E(I)+C(I)
CONTINUE
101
DO 102 I=169,176
  E(I)=E(I)-C(I)
CONTINUE
102
DO 103 I=177,184
  E(I)=E(I)+C(I)
CONTINUE
103
DO 104 I=185,192
  E(I)=E(I)-C(I)
CONTINUE
104
DO 105 I=193,200
  E(I)=E(I)+C(I)
CONTINUE
105
DO 106 I=201,208
  E(I)=E(I)-C(I)
CONTINUE
106
DO 107 I=209,216
  E(I)=E(I)+C(I)
CONTINUE
107
DO 108 I=217,224
  E(I)=E(I)-C(I)
CONTINUE
108
DO 109 I=225,232
  E(I)=E(I)+C(I)
CONTINUE
109
DO 110 I=233,240
  E(I)=E(I)-C(I)
CONTINUE
110
DO 111 I=241,248
  E(I)=E(I)+C(I)
CONTINUE
111
DO 112 I=249,256
  E(I)=E(I)-C(I)
CONTINUE
112
DO 113 I=257,264
  E(I)=E(I)+C(I)
CONTINUE
113
DO 114 I=265,272
  E(I)=E(I)-C(I)
CONTINUE
114
DO 115 I=273,280
  E(I)=E(I)+C(I)
CONTINUE
115
DO 116 I=281,288
  E(I)=E(I)-C(I)
CONTINUE
116
DO 117 I=289,296
  E(I)=E(I)+C(I)
CONTINUE
117
DO 118 I=297,304
  E(I)=E(I)-C(I)
CONTINUE
118
DO 119 I=305,312
  E(I)=E(I)+C(I)
CONTINUE
119
DO 120 I=313,320
  E(I)=E(I)-C(I)
CONTINUE
120
DO 121 I=321,328
  E(I)=E(I)+C(I)
CONTINUE
121
DO 122 I=329,336
  E(I)=E(I)-C(I)
CONTINUE
122
DO 123 I=337,344
  E(I)=E(I)+C(I)
CONTINUE
123
DO 124 I=345,352
  E(I)=E(I)-C(I)
CONTINUE
124
DO 125 I=353,360
  E(I)=E(I)+C(I)
CONTINUE
125
DO 126 I=361,368
  E(I)=E(I)-C(I)
CONTINUE
126
DO 127 I=369,376
  E(I)=E(I)+C(I)
CONTINUE
127
DO 128 I=377,384
  E(I)=E(I)-C(I)
CONTINUE
128
DO 129 I=385,392
  E(I)=E(I)+C(I)
CONTINUE
129
DO 130 I=393,400
  E(I)=E(I)-C(I)
CONTINUE
130
DO 131 I=401,408
  E(I)=E(I)+C(I)
CONTINUE
131
DO 132 I=409,416
  E(I)=E(I)-C(I)
CONTINUE
132
DO 133 I=417,424
  E(I)=E(I)+C(I)
CONTINUE
133
DO 134 I=425,432
  E(I)=E(I)-C(I)
CONTINUE
134
DO 135 I=433,440
  E(I)=E(I)+C(I)
CONTINUE
135
DO 136 I=441,448
  E(I)=E(I)-C(I)
CONTINUE
136
DO 137 I=449,456
  E(I)=E(I)+C(I)
CONTINUE
137
DO 138 I=457,464
  E(I)=E(I)-C(I)
CONTINUE
138
DO 139 I=465,472
  E(I)=E(I)+C(I)
CONTINUE
139
DO 140 I=473,480
  E(I)=E(I)-C(I)
CONTINUE
140
DO 141 I=481,488
  E(I)=E(I)+C(I)
CONTINUE
141
DO 142 I=489,496
  E(I)=E(I)-C(I)
CONTINUE
142
DO 143 I=497,504
  E(I)=E(I)+C(I)
CONTINUE
143
DO 144 I=505,512
  E(I)=E(I)-C(I)
CONTINUE
144
DO 145 I=513,520
  E(I)=E(I)+C(I)
CONTINUE
145
DO 146 I=521,528
  E(I)=E(I)-C(I)
CONTINUE
146
DO 147 I=529,536
  E(I)=E(I)+C(I)
CONTINUE
147
DO 148 I=537,544
  E(I)=E(I)-C(I)
CONTINUE
148
DO 149 I=545,552
  E(I)=E(I)+C(I)
CONTINUE
149
DO 150 I=553,560
  E(I)=E(I)-C(I)
CONTINUE
150
DO 151 I=561,568
  E(I)=E(I)+C(I)
CONTINUE
151
DO 152 I=569,576
  E(I)=E(I)-C(I)
CONTINUE
152
DO 153 I=577,584
  E(I)=E(I)+C(I)
CONTINUE
153
DO 154 I=585,592
  E(I)=E(I)-C(I)
CONTINUE
154
DO 155 I=593,600
  E(I)=E(I)+C(I)
CONTINUE
155
DO 156 I=601,608
  E(I)=E(I)-C(I)
CONTINUE
156
DO 157 I=609,616
  E(I)=E(I)+C(I)
CONTINUE
157
DO 158 I=617,624
  E(I)=E(I)-C(I)
CONTINUE
158
DO 159 I=625,632
  E(I)=E(I)+C(I)
CONTINUE
159
DO 160 I=633,640
  E(I)=E(I)-C(I)
CONTINUE
160
DO 161 I=641,648
  E(I)=E(I)+C(I)
CONTINUE
161
DO 162 I=649,656
  E(I)=E(I)-C(I)
CONTINUE
162
DO 163 I=657,664
  E(I)=E(I)+C(I)
CONTINUE
163
DO 164 I=665,672
  E(I)=E(I)-C(I)
CONTINUE
164
DO 165 I=673,680
  E(I)=E(I)+C(I)
CONTINUE
165
DO 166 I=681,688
  E(I)=E(I)-C(I)
CONTINUE
166
DO 167 I=689,696
  E(I)=E(I)+C(I)
CONTINUE
167
DO 168 I=697,704
  E(I)=E(I)-C(I)
CONTINUE
168
DO 169 I=705,712
  E(I)=E(I)+C(I)
CONTINUE
169
DO 170 I=713,720
  E(I)=E(I)-C(I)
CONTINUE
170
DO 171 I=721,728
  E(I)=E(I)+C(I)
CONTINUE
171
DO 172 I=729,736
  E(I)=E(I)-C(I)
CONTINUE
172
DO 173 I=737,744
  E(I)=E(I)+C(I)
CONTINUE
173
DO 174 I=745,752
  E(I)=E(I)-C(I)
CONTINUE
174
DO 175 I=753,760
  E(I)=E(I)+C(I)
CONTINUE
175
DO 176 I=761,768
  E(I)=E(I)-C(I)
CONTINUE
176
DO 177 I=769,776
  E(I)=E(I)+C(I)
CONTINUE
177
DO 178 I=777,784
  E(I)=E(I)-C(I)
CONTINUE
178
DO 179 I=785,792
  E(I)=E(I)+C(I)
CONTINUE
179
DO 180 I=793,800
  E(I)=E(I)-C(I)
CONTINUE
180
DO 181 I=801,808
  E(I)=E(I)+C(I)
CONTINUE
181
DO 182 I=809,816
  E(I)=E(I)-C(I)
CONTINUE
182
DO 183 I=817,824
  E(I)=E(I)+C(I)
CONTINUE
183
DO 184 I=825,832
  E(I)=E(I)-C(I)
CONTINUE
184
DO 185 I=833,840
  E(I)=E(I)+C(I)
CONTINUE
185
DO 186 I=841,848
  E(I)=E(I)-C(I)
CONTINUE
186
DO 187 I=849,856
  E(I)=E(I)+C(I)
CONTINUE
187
DO 188 I=857,864
  E(I)=E(I)-C(I)
CONTINUE
188
DO 189 I=865,872
  E(I)=E(I)+C(I)
CONTINUE
189
DO 190 I=873,880
  E(I)=E(I)-C(I)
CONTINUE
190
DO 191 I=881,888
  E(I)=E(I)+C(I)
CONTINUE
191
DO 192 I=889,896
  E(I)=E(I)-C(I)
CONTINUE
192
DO 193 I=897,904
  E(I)=E(I)+C(I)
CONTINUE
193
DO 194 I=905,912
  E(I)=E(I)-C(I)
CONTINUE
194
DO 195 I=913,920
  E(I)=E(I)+C(I)
CONTINUE
195
DO 196 I=921,928
  E(I)=E(I)-C(I)
CONTINUE
196
DO 197 I=929,936
  E(I)=E(I)+C(I)
CONTINUE
197
DO 198 I=937,944
  E(I)=E(I)-C(I)
CONTINUE
198
DO 199 I=945,952
  E(I)=E(I)+C(I)
CONTINUE
199
DO 200 I=953,960
  E(I)=E(I)-C(I)
CONTINUE
200
DO 201 I=961,968
  E(I)=E(I)+C(I)
CONTINUE
201
DO 202 I=969,976
  E(I)=E(I)-C(I)
CONTINUE
202
DO 203 I=977,984
  E(I)=E(I)+C(I)
CONTINUE
203
DO 204 I=985,992
  E(I)=E(I)-C(I)
CONTINUE
204
DO 205 I=993,1000
  E(I)=E(I)+C(I)
CONTINUE
205
DO 206 I=1001,1008
  E(I)=E(I)-C(I)
CONTINUE
206
DO 207 I=1009,1016
  E(I)=E(I)+C(I)
CONTINUE
207
DO 208 I=1017,1024
  E(I)=E(I)-C(I)
CONTINUE
208
DO 209 I=1025,1032
  E(I)=E(I)+C(I)
CONTINUE
209
DO 210 I=1033,1040
  E(I)=E(I)-C(I)
CONTINUE
210
DO 211 I=1041,1048
  E(I)=E(I)+C(I)
CONTINUE
211
DO 212 I=1049,1056
  E(I)=E(I)-C(I)
CONTINUE
212
DO 213 I=1057,1064
  E(I)=E(I)+C(I)
CONTINUE
213
DO 214 I=1065,1072
  E(I)=E(I)-C(I)
CONTINUE
214
DO 215 I=1073,1080
  E(I)=E(I)+C(I)
CONTINUE
215
DO 216 I=1081,1088
  E(I)=E(I)-C(I)
CONTINUE
216
DO 217 I=1089,1096
  E(I)=E(I)+C(I)
CONTINUE
217
DO 218 I=1097,1104
  E(I)=E(I)-C(I)
CONTINUE
218
DO 219 I=1105,1112
  E(I)=E(I)+C(I)
CONTINUE
219
DO 220 I=1113,1120
  E(I)=E(I)-C(I)
CONTINUE
220
DO 221 I=1121,1128
  E(I)=E(I)+C(I)
CONTINUE
221
DO 222 I=1129,1136
  E(I)=E(I)-C(I)
CONTINUE
222
DO 223 I=1137,1144
  E(I)=E(I)+C(I)
CONTINUE
223
DO 224 I=1145,1152
  E(I)=E(I)-C(I)
CONTINUE
224
DO 225 I=1153,1160
  E(I)=E(I)+C(I)
CONTINUE
225
DO 226 I=1161,1168
  E(I)=E(I)-C(I)
CONTINUE
226
DO 227 I=1169,1176
  E(I)=E(I)+C(I)
CONTINUE
227
DO 228 I=1177,1184
  E(I)=E(I)-C(I)
CONTINUE
228
DO 229 I=1185,1192
  E(I)=E(I)+C(I)
CONTINUE
229
DO 230 I=1193,1200
  E(I)=E(I)-C(I)
CONTINUE
230
DO 231 I=1201,1208
  E(I)=E(I)+C(I)
CONTINUE
231
DO 232 I=1209,1216
  E(I)=E(I)-C(I)
CONTINUE
232
DO 233 I=1217,1224
  E(I)=E(I)+C(I)
CONTINUE
233
DO 234 I=1225,1232
  E(I)=E(I)-C(I)
CONTINUE
234
DO 235 I=1233,1240
  E(I)=E(I)+C(I)
CONTINUE
235
DO 236 I=1241,1248
  E(I)=E(I)-C(I)
CONTINUE
236
DO 237 I=1249,1256
  E(I)=E(I)+C(I)
CONTINUE
237
DO 238 I=1257,1264
  E(I)=E(I)-C(I)
CONTINUE
238
DO 239 I=1265,1272
  E(I)=E(I)+C(I)
CONTINUE
239
DO 240 I=1273,1280
  E(I)=E(I)-C(I)
CONTINUE
240
DO 241 I=1281,1288
  E(I)=E(I)+C(I)
CONTINUE
241
DO 242 I=1289,1296
  E(I)=E(I)-C(I)
CONTINUE
242
DO 243 I=1297,1304
  E(I)=E(I)+C(I)
CONTINUE
243
DO 244 I=1305,1312
  E(I)=E(I)-C(I)
CONTINUE
244
DO 245 I=1313,1320
  E(I)=E(I)+C(I)
CONTINUE
245
DO 246 I=1321,1328
  E(I)=E(I)-C(I)
CONTINUE
246
DO 247 I=1329,1336
  E(I)=E(I)+C(I)
CONTINUE
247
DO 248 I=1337,1344
  E(I)=E(I)-C(I)
CONTINUE
248
DO 249 I=1345,1352
  E(I)=E(I)+C(I)
CONTINUE
249
DO 250 I=1353,1360
  E(I)=E(I)-C(I)
CONTINUE
250
DO 251 I=1361,1368
  E(I)=E(I)+C(I)
CONTINUE
251
DO 252 I=1369,1376
  E(I)=E(I)-C(I)
CONTINUE
252
DO 253 I=1377,1384
  E(I)=E(I)+C(I)
CONTINUE
253
DO 254 I=1385,1392
  E(I)=E(I)-C(I)
CONTINUE
254
DO 255 I=1393,1400
  E(I)=E(I)+C(I)
CONTINUE
255
DO 256 I=1401,1408
  E(I)=E(I)-C(I)
CONTINUE
256
DO 257 I=1409,1416
  E(I)=E(I)+C(I)
CONTINUE
257
DO 258 I=1417,1424
  E(I)=E(I)-C(I)
CONTINUE
258
DO 259 I=1425,1432
  E(I)=E(I)+C(I)
CONTINUE
259
DO 260 I=1433,1440
  E(I)=E(I)-C(I)
CONTINUE
260
DO 261 I=1441,1448
  E(I)=E(I)+C(I)
CONTINUE
261
DO 262 I=1449,1456
  E(I)=E(I)-C(I)
CONTINUE
262
DO 263 I=1457,1464
  E(I)=E(I)+C(I)
CONTINUE
263
DO 264 I=1465,1472
  E(I)=E(I)-C(I)
CONTINUE
264
DO 265 I=1473,1480
  E(I)=E(I)+C(I)
CONTINUE
265
DO 266 I=1481,1488
  E(I)=E(I)-C(I)
CONTINUE
266
DO 267 I=1489,1496
  E(I)=E(I)+C(I)
CONTINUE
267
DO 268 I=1497,1504
  E(I)=E(I)-C(I)
CONTINUE
268
DO 269 I=1505,1512
  E(I)=E(I)+C(I)
CONTINUE
269
DO 270 I=1513,1520
  E(I)=E(I)-C(I)
CONTINUE
270
DO 271 I=1521,1528
  E(I)=E(I)+C(I)
CONTINUE
271
DO 272 I=1529,1536
  E(I)=E(I)-C(I)
CONTINUE
272
DO 273 I=1537,1544
  E(I)=E(I)+C(I)
CONTINUE
273
DO 274 I=1545,1552
  E(I)=E(I)-C(I)
CONTINUE
274
DO 275 I=1553,1560
  E(I)=E(I)+C(I)
CONTINUE
275
DO 276 I=1561,1568
  E(I)=E(I)-C(I)
CONTINUE
276
DO 277 I=1569,1576
  E(I)=E(I)+C(I)
CONTINUE
277
DO 278 I=1577,1584
  E(I)=E(I)-C(I)
CONTINUE
278
DO 279 I=1585,1592
  E(I)=E(I)+C(I)
CONTINUE
279
DO 280 I=1593,1600
  E(I)=E(I)-C(I)
CONTINUE
280
DO 281 I=1601,1608
  E(I)=E(I)+C(I)
CONTINUE
281
DO 282 I=1609,1616
  E(I)=E(I)-C(I)
CONTINUE
282
DO 283 I=1617,1624
  E(I)=E(I)+C(I)
CONTINUE
283
DO 284 I=1625,1632
  E(I)=E(I)-C(I)
CONTINUE
284
DO 285 I=1633,1640
  E(I)=E(I)+C(I)
CONTINUE
285
DO 286 I=1641,1648
  E(I)=E(I)-C(I)
CONTINUE
286
DO 287 I=1649,1656
  E(I)=E(I)+C(I)
CONTINUE
287
DO 288 I=1657,1664
  E(I)=E(I)-C(I)
CONTINUE
288
DO 289 I=1665,1672
  E(I)=E(I)+C(I)
CONTINUE
289
DO 290 I=1673,1680
  E(I)=E(I)-C(I)
CONTINUE
290
DO 291 I=1681,1688
  E(I)=E(I)+C(I)
CONTINUE
291
DO 292 I=1689,1696
  E(I)=E(I)-C(I)
CONTINUE
292
DO 293 I=1697,1704
  E(I)=E(I)+C(I)
CONTINUE
293
DO 294 I=1705,1712
  E(I)=E(I)-C(I)
CONTINUE
294
DO 295 I=1713,1720
  E(I)=E(I)+C(I)
CONTINUE
295
DO 296 I=1721,1728
  E(I)=E(I)-C(I)
CONTINUE
296
DO 297 I=1729,1736
  E(I)=E(I)+C(I)
CONTINUE
297
DO 298 I=1737,1744
  E(I)=E(I)-C(I)
CONTINUE
298
DO 299 I=1745,1752
  E(I)=E(I)+C(I)
CONTINUE
299
DO 300 I=1753,1760
  E(I)=E(I)-C(I)
CONTINUE
300
DO 301 I=1761,1768
  E(I)=E(I)+C(I)
CONTINUE
301
DO 302 I=1769,1776
  E(I)=E(I)-C(I)
CONTINUE
302
DO 303 I=1777,1784
  E(I)=E(I)+C(I)
CONTINUE
303
DO 304 I=1785,1792
  E(I)=E(I)-C(I)
CONTINUE
304
DO 305 I=1793,1800
  E(I)=E(I)+C(I)
CONTINUE
305
DO 306 I=1801,1808
  E(I)=E(I)-C(I)
CONTINUE
306
DO 307 I=1809,1816
  E(I)=E(I)+C(I)
CONTINUE
307
DO 308 I=1817,1824
  E(I)=E(I)-C(I)
CONTINUE
308
DO 309 I=1825,1832
  E(I)=E(I)+C(I)
CONTINUE
309
DO 310 I=1833,1840
  E(I)=E(I)-C(I)
CONTINUE
310
DO 311 I=1841,1848
  E(I)=E(I)+C(I)
CONTINUE
311
DO 312 I=1849,1856
  E(I)=E(I)-C(I)
CONTINUE
312
DO 313 I=1857,1864
  E(I)=E(I)+C(I)
CONTINUE
313
DO 314 I=1865,1872
  E(I)=E(I)-C(I)
CONTINUE
314
DO 315 I=1873,1880
  E(I)=E(I)+C(I)
CONTINUE
315
DO 316 I=1881,1888
  E(I)=E(I)-C(I)
CONTINUE
316
DO 317 I=1889,1896
  E(I)=E(I)+C(I)
CONTINUE
317
DO 318 I=1897,1904
  E(I)=E(I)-C(I)
CONTINUE
318
DO 319 I=1905,1912
  E(I)=E(I)+C(I)
CONTINUE
319
DO 320 I=1913,1920
  E(I)=E(I)-C(I)
CONTINUE
320
DO 321 I=1921,1928
  E(I)=E(I)+C(I)
CONTINUE
321
DO 322 I=1929,1936
  E(I)=E(I)-C(I)
CONTINUE
322
DO 323 I=1937,1944
  E(I)=E(I)+C(I)
CONTINUE
323
DO 324 I=1945,1952
  E(I)=E(I)-C(I)
CONTINUE
324
DO 325 I=1953,1960
  E(I)=E(I)+C(I)
CONTINUE
325
DO 326 I=1961,1968
  E(I)=E(I)-C(I)
CONTINUE
326
DO 327 I=1969,1976
  E(I)=E(I)+C(I)
CONTINUE
327
DO 328 I=1977,1984
  E(I)=E(I)-C(I)
CONTINUE
328
DO 329 I=1985,1992
  E(I)=E(I)+C(I)
CONTINUE
329
DO 330 I=1993,2000
  E(I)=E(I)-C(I)
CONTINUE
330
DO 331 I=2001,2008
  E(I)=E(I)+C(I)
CONTINUE
331
DO 332 I=2009,2016
  E(I)=E(I)-C(I)
CONTINUE
332
DO 333 I=2017,2024
  E(I)=E(I)+C(I)
CONTINUE
333
DO 334 I=2025,2032
  E(I)=E(I)-C(I)
CONTINUE
334
DO 335 I=2033,2040
  E(I)=E(I)+C(I)
CONTINUE
335
DO 336 I=2041,2048
  E(I)=E(I)-C(I)
CONTINUE
336
DO 337 I=2049,2056
  E(I)=E(I)+C(I)
CONTINUE
337
DO 338 I=2057,2064
  E(I)=E(I)-C(I)
CONTINUE
338
DO 339 I=2065,2072
  E(I)=E(I)+C(I)
CONTINUE
339
DO 340 I=2073,2080
  E(I)=E(I)-C(I)
CONTINUE
340
DO 341 I=2081,2088
  E(I)=E(I)+C(I)
CONTINUE
341
DO 342 I=2089,2096
  E(I)=E(I)-C(I)
CONTINUE
342
DO 343 I=2097,2104
  E(I)=E(I)+C(I)
CONTINUE
343
DO 344 I=2105,2112
  E(I)=E(I)-C(I)
CONTINUE
344
DO 345 I=2113,2120
  E(I)=E(I)+C(I)
CONTINUE
345
DO 346 I=2121,2128
  E(I)=E(I)-C(I)
CONTINUE
346
DO 347 I=2129,2136
  E(I)=E(I)+C(I)
CONTINUE
347
DO 348 I=2137,2144
  E(I)=E(I)-C(I)
CONTINUE
348
DO 349 I=2145,2152
  E(I)=E(I)+C(I)
CONTINUE
349
DO 350 I=2153,2160
  E(I)=E(I)-C(I)
CONTINUE
350
DO 351 I=2161,2168
  E(I)=E(I)+C(I)
CONTINUE
351
DO 352 I=2169,2176
  E(I)=E(I)-C(I)
CONTINUE
352
DO 353 I=2177,2184
  E(I)=E(I)+C(I)
CONTINUE
353
DO 354 I=2185,2192
  E(I)=E(I)-C(I)
CONTINUE
354
DO 355 I=2193,2200
  E(I)=E(I)+C(I)
CONTINUE
355
DO 356 I=2201,2208
  E(I)=E(I)-C(I)
CONTINUE
356
DO 357 I=2209,2216
  E(I)=E(I)+C(I)
CONTINUE
357
DO 358 I=2217,2224
  E(I)=E(I)-C(I)
CONTINUE
358
DO 359 I=2225,2232
  E(I)=E(I)+C(I)
CONTINUE
359
DO 360 I=2233,2240
  E(I)=E(I)-C(I)
CONTINUE
360
DO 361 I=2241,2248
  E(I)=E(I)+C(I)
CONTINUE
361
DO 362 I=2249,2256
  E(I)=E(I)-C(I)
CONTINUE
362
DO 363 I=2257,2264
  E(I)=E(I)+C(I)
CONTINUE
363
DO 364 I=2265,2272
  E(I)=E(I)-C(I)
CONTINUE
364
DO 365 I=2273,2280
  E(I)=E(I)+C(I)
CONTINUE
365
DO 366 I=2281,2288
  E(I)=E(I)-C(I)
CONTINUE
366
DO 367 I=2289,2296
  E(I)=E(I)+C(I)
CONTINUE
367
DO 368 I=2297,2304
  E(I)=E(I)-C(I)
CONTINUE
368
DO 369 I=2305,2312
  E(I)=E(I)+C(I)
CONTINUE
369
DO 370 I=2313,2320
  E(I)=E(I)-C(I)
CONTINUE
370
DO 371 I=2321,2328
  E(I)=E(I)+C(I)
CONTINUE
371
DO 372 I=2329,2336
  E(I)=E(I)-C(I)
CONTINUE
372
DO 373 I=2337,2344
  E(I)=E(I)+C(I)
CONTINUE
373
DO 374 I=2345,2352
  E(I)=E(I)-C(I)
CONTINUE
374
DO 375 I=2353,2360
  E(I)=E(I)+C(I)
CONTINUE
375
DO 376 I=2361,2368
  E(I)=E(I)-C(I)
CONTINUE
376
DO 377 I=2369,2376
  E(I)=E(I)+C(I)
CONTINUE
377
DO 378 I=2377,2384
  E(I)=E(I)-C(I)
CONTINUE
378
DO 379 I=2385,2392
  E(I)=E(I)+C(I)
CONTINUE
379
DO 380 I=2393,2400
  E(I)=E(I)-C(I)
CONTINUE
380
DO 381 I=2401,2408
  E(I)=E(I)+C(I)
CONTINUE
381
DO 382 I=2409,2416
  E(I)=E(I)-C(I)
CONTINUE
382
DO 383 I=2417,2424
  E(I)=E(I)+C(I)
CONTINUE
383
DO 384 I=2425,2432
  E(I)=E(I)-C(I)
CONTINUE
384
DO 385 I=2433,2440
  E(I)=E(I)+C(I)
CONTINUE
385
DO 386 I=2441,2448
  E(I)=E(I)-C(I)
CONTINUE
386
DO 387 I=2449,2456
  E(I)=E(I)+C(I)
CONTINUE
387
DO 388 I=2457,2464
  E(I)=E(I)-C(I)
CONTINUE
388
DO 389 I=2465,2472
  E(I)=E(I)+C(I)
CONTINUE
389
DO 390 I=2473,2480
  E(I)=E(I)-C(I)
CONTINUE
390
DO 391 I=2481,2488
  E(I)=E(I)+C(I)
CONTINUE
391
DO 392 I=2489,2496
  E(I)=E(I)-C(I)
CONTINUE
392
DO 393 I=2497,2504
  E(I)=E(I)+C(I)
CONTINUE
393
DO 394 I=2505,2512
  E(I)=E(I)-C(I)
CONTINUE
394
DO 395 I=2513,2520
  E(I)=E(I)+C(I)
CONTINUE
395
DO 396 I=2521,2528
  E(I)=E(I)-C(I)
CONTINUE
396
DO 397 I=2529,2536
  E(I)=E(I)+C(I)
CONTINUE
397
DO 398 I=2537,2544
  E(I)=E(I)-C(I)
CONTINUE
398
DO 399 I=2545,2552
  E(I)=E(I)+C(I)
CONTINUE
399
DO 400 I=2553,2560
  E(I)=E(I)-C(I)
CONTINUE
400
DO 401 I=2561,2568
  E(I)=E(I)+C(I)
CONTINUE
401
DO 402 I=2569,2576
  E(I)=E(I)-C(I)
CONTINUE
402
DO 403 I=2577,2584
  E(I)=E(I)+C(I)
CONTINUE
403
DO 404 I=2585,2592
  E(I)=E(I)-C(I)
CONTINUE
404
DO 405 I=2593,2600
  E(I)=E(I)+C(I)
CONTINUE
405
DO 406 I=2601,2608
  E(I)=E(I)-C(I)
CONTINUE
406
DO 407 I=2609,2616
  E(I)=E(I)+C(I)
CONTINUE
407
DO 408 I=2617,2624
  E(I)=E(I)-C(I)
CONTINUE
408
DO 409 I=2625,2632
  E(I)=E(I)+C(I)
CONTINUE
409
DO 410 I=2633,2640
  E(I)=E(I)-C(I)
CONTINUE
410
DO 411 I=2641,2648
  E(I)=E(I)+C(I)
CONTINUE
411
DO 412 I=2649,2656
  E(I)=E(I)-C(I)
CONTINUE
412
DO 413 I=2657,2664
  E(I)=E(I)+C(I)
CONTINUE
413
DO 414 I=2665,2672
  E(I)=E(I)-C(I)
CONTINUE
414
DO 415 I=2673,2680
  E(I)=E(I)+C(I)
CONTINUE
415
DO 416 I=2681,2688
  E(I)=E(I)-C(I)
CONTINUE
416
DO 417 I=2689,2696
  E(I)=E(I)+C(I)
CONTINUE
417
DO 418 I=2697,2704
  E(I)=E(I)-C(I)
CONTINUE
418
DO 419 I=2705,2712
  E(I)=E(I)+C(I)
CONTINUE
419
DO 420 I=2713,2720
  E(I)=E(I)-C(I)
CONTINUE
420
DO 421 I=2721,2728
  E(I)=E(I)+C(I)
CONTINUE
421
DO 422 I=2729,2736
  E(I)=E(I)-C(I)
CONTINUE
422
DO 423 I=2737,2744
  E(I)=E(I)+C(I)
CONTINUE
423
DO 424 I=2745,2752
  E(I)=E(I)-C(I)
CONTINUE
424
DO 425 I=2753,2760
  E(I)=E(I)+C(I)
CONTINUE
425
DO 426 I=2761,2768
  E(I)=E(I)-C(I)
CONTINUE
426
DO 427 I=2769,2776
  E(I)=E(I)+C(I)
CONTINUE
427
DO 428 I=2777,2784
  E(I)=E(I)-C(I)
CONTINUE
428
DO 429 I=2785,2792
  E(I)=E(I)+C(I)
CONTINUE
429
DO 430
```



```

XN=-1,O*YN
301 CALL -COEF(MDIR, AREA,XM,YM,ZM,XN,YN,ZN,C)
IF(MODE.EQ.COSH.OR.MODE.EG.SINE) GO TO 98
DO 94 I=1,4
E(I)=E(I)+C(I)
CONTINUE
94 DO 95 I=5,8
E(I)=E(I)-C(I)
CONTINUE
95 DO 96 I=9,12
E(I)=E(I)+C(I)
CONTINUE
96 DO 97 I=13,16
E(I)=E(I)-C(I)
CONTINUE
97 GO TO 113
98 CONTINUE
DO 99 I=1,4
E(I)=E(I)-C(I)
CONTINUE
99 DO 100 I=5,8
E(I)=E(I)+C(I)
CONTINUE
100 DO 101 I=9,12
E(I)=E(I)-C(I)
CONTINUE
101 DO 102 I=13,16
E(I)=E(I)+C(I)
CONTINUE
102 CONTINUE
113 RETURN
END
C*****SUBROUTINE EXIM(L,MAT,MATP,Z,NCELL)
SUBROUTINE EXIM(L,MAT,MATP,Z,NCELL)
REAL Z(NCELL),IMPR,MUO
COMPLEX L(MAT,MATP),CJ
COMMON /INCR/IMPR,CRNT
COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
DO 1 I=1,NX
IP=I+NX
L(IP,MATP)=CEXP(-1.O*CJ*Z(I))/377.O
CONTINUE
DO 11 I=1,NY
II=I+NX
L(IP,MATP)=CEXP(-1.O*CJ*Z(II))/377.O

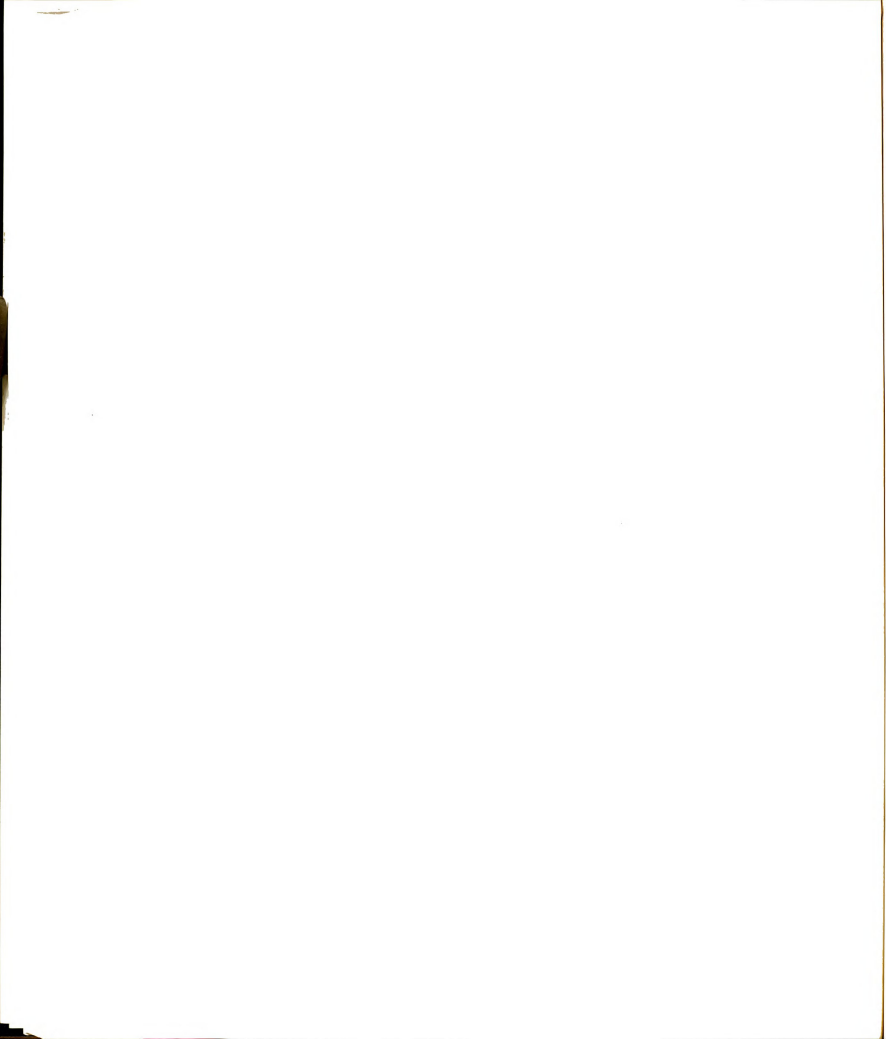
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```

11 CONTINUE
DO 2 I=1,NZ
  II=I+NX+NY
  IP=I+4*(NX+NY)
  IPP=IP+3*NZ
  L(IP,MATP)=-1.0*CEXP(-1.0*CJ*Z(II))/377.0
  L(IPP,MATP)=L(IP,MATP)
CONTINUE
IF (NNX.EQ.0) GO TO 6
DO 4 I=1,NNX
  II=I+NX+NY+NZ
  IP=I+4*(NX+NY+NZ)+NNX
  L(IP,MATP)=-1.0*CEXP(-1.0*CJ*Z(II))/377.0
CONTINUE
IF (NNY.EQ.0) GO TO 16
DO 14 I=1,NNY
  II=I+NX+NY+NZ+NNX
  IP=I+4*(NX+NY+NZ+NNX)+3*NNY
  L(IP,MATP)=-1.0*CEXP(-1.0*CJ*Z(II))/377.0
CONTINUE
IF (NNZ.EQ.0) GO TO 7
DO 5 I=1,NNZ
  II=I+NX+NY+NZ+NNX+NNY
  IP=I+4*(NX+NY+NZ+NNX+NNY)
  IPP=IP+3*NNZ
  L(IP,MATP)=CEXP(-1.0*CJ*Z(II))/377.0
  L(IPP,MATP)=L(IP,MATP)
CONTINUE
CONTINUE
DO 3 I=1,MAT
  DO 13 J=1,MAT
    CRNT=L(I,MATP)
    L(I,MATP)=4.0*PI*L(I,MATP)
  CONTINUE
  IF ((NNX.EQ.0).AND.(NNY.EQ.0).AND.(NNZ.EQ.0)) GO TO 9
  MATP=MATP+1
DO 8 I=MATP,MAT
  L(I,MATP)=-1.0*L(I,MATP)
CONTINUE
RETURN
END
*****
SUBROUTINE CEIM(L,MAT,MATP,Z,NCELL)
REAL Z(NCELL),IMPR,MUO
COMPLEX L(MAT,MATP),CJ
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7990
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8070
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8100
8110
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7460

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```

COMMON /IMCR/IMPR,CRNT
COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
COMMON /CONS/PI,MUO,EPO
COMMON /GPLX/CJ
DO 1 I=1,NY
  II=I+NX
  IP=I+4*NX+3*NY
  L(IP,MATP)=COS(Z(II))/377.0
1 CONTINUE
DO 2 I=1,NZ
  II=I+NX+NY
  IP=I+4*(NX+NY)+3*NZ
  L(IP,MATP)=(-1.0)*COS(Z(II))/377.0
2 CONTINUE
DO 3 I=1,NZ
  II=I+NX
  IP=I+4*(NX+NZ)+3*NNX
  L(IP,MATP)=(-1.0)*COS(Z(II))/377.0
3 CONTINUE
DO 4 I=1,NNZ
  II=I+NX+NY+NZ+NNX+NNY
  IP=I+4*(NX+NY+NZ+NNX+NNY)+3*NNZ
  L(IP,MATP)=COS(Z(II))/377.0
4 CONTINUE
DO 5 I=1,NNZ
  II=I+NX+NY+NZ+NNX+NNY
  IP=I+4*(NX+NY+NZ+NNX+NNY)+3*NNZ
  L(IP,MATP)=COS(Z(II))/377.0
5 CONTINUE
DO 6 I=1,MAT
  CALL PRNT(IMPR,CRNT,L,MAT,MATP)
  L(I,MATP)=4.0*PI*L(I,MATP)
6 CONTINUE
IF((NNX.EQ.0).AND.(NNY.EQ.0).AND.(NNZ.EQ.0)) GO TO 9
MATI=4*(NX+NY+NZ)
DO 7 I=1,MAT+1
  L(I,MATP)=MATI
7 CONTINUE
L(I,MATP)=-1.0*L(I,MATP)
8 CONTINUE
9 RETURN
END
C*****SUBROUTINE CHIM*****
SUBROUTINE CHIM(L,MAT,MATP,Z,NCELL)
REAL Z(NCELL),IMPR,MUO
COMPLEX L(MAT,MATP),CJ
COMMON /IMCR/IMPR,CRNT
COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
COMMON /CONS/PI,MUO,EPO
COMMON /GPLX/CJ

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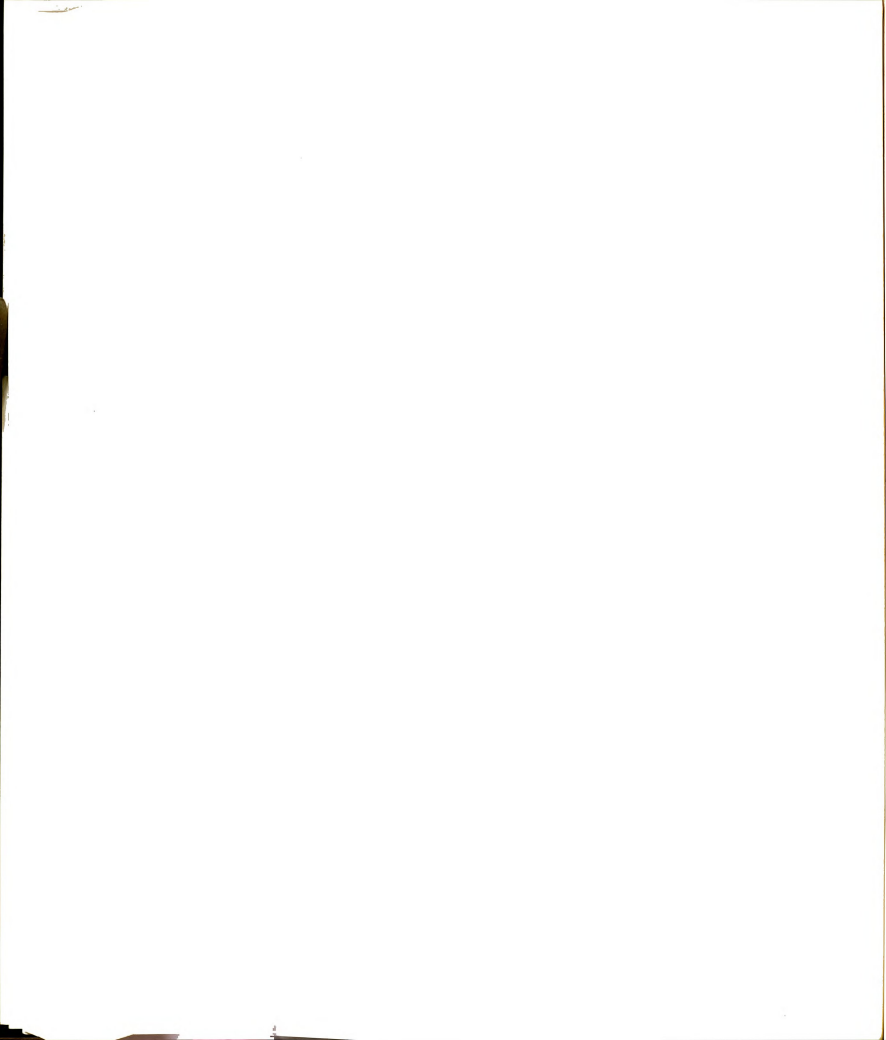
7470
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75507560
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75907600
76107620
76307640
76507660
76707680
76907700
77107720
77307740
77507760
77707780
7790

```

C*****SUBROUTINE CHIM*****
SUBROUTINE CHIM(L,MAT,MATP,Z,NCELL)
REAL Z(NCELL),IMPR,MUO
COMPLEX L(MAT,MATP),CJ
COMMON /IMCR/IMPR,CRNT
COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
COMMON /CONS/PI,MUO,EPO
COMMON /GPLX/CJ

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7800
7810
78207840
78507860
78707880
7890




```

DO 1 I=1,NX
IP=I+NX
L(IP,MATP)=COS(Z(I))/377.0
1 CONTINUE
DO 2 I=1,NZ
II=I+NX+NY
IP=I+4*(NX+NY)
L(IP,MATP)=(-1.0)*COS(Z(II))/377.0
2 CONTINUE
IF(NNX.EQ.0) GO TO 6
DO 4 I=1,NNX
II=I+NX+NY+NZ
IP=I+4*(NX+NY+NZ)+NNX
L(IP,MATP)=(-1.0)*COS(Z(II))/377.0
4 CONTINUE
6 CONTINUE
IF(NNZ.EQ.0) GO TO 7
DO 5 I=1,NNZ
II=I+NX+NY+NZ+NNX+NNY
IP=I+4*(NX+NY+NZ+NNX+NNY)
L(IP,MATP)=COS(Z(II))/377.0
5 CONTINUE
7 CONTINUE
CALL PRNT(IMPR,CRNT,L,MAT,MATP)
DO 3 I=1,MAT
L(I,MATP)=4.0*PI*L(I,MATP)
3 CONTINUE
IF(NNX.EQ.0).AND.(NNY.EQ.0).AND.(NNZ.EQ.0) GO TO 9
MATT=4*(NX+NY+NZ)
MATTP=MATT+1
DO 8 I=MATTP,MAT
L(I,MATP)=-1.0*L(I,MATP)
8 CONTINUE
9 RETURN
END
C*****SUBROUTINE SEIM*****
SUBROUTINE SEIM(L,MAT,MATP,Z,NCELL)
REAL Z(NCELL),IMPR,MUO
COMPLEX L(MAT,MATP),CJ
COMMON /IMCR/IMPR,CRNT
COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
COMMON /CONS/PI,MUO,EPO
DO 1 I=1,NY
II=I+NX
IP=I+4*NX+3*NY
L(IP,MATP)=(-1.0)*CJ*SIN(Z(II))/377.0
1 CONTINUE

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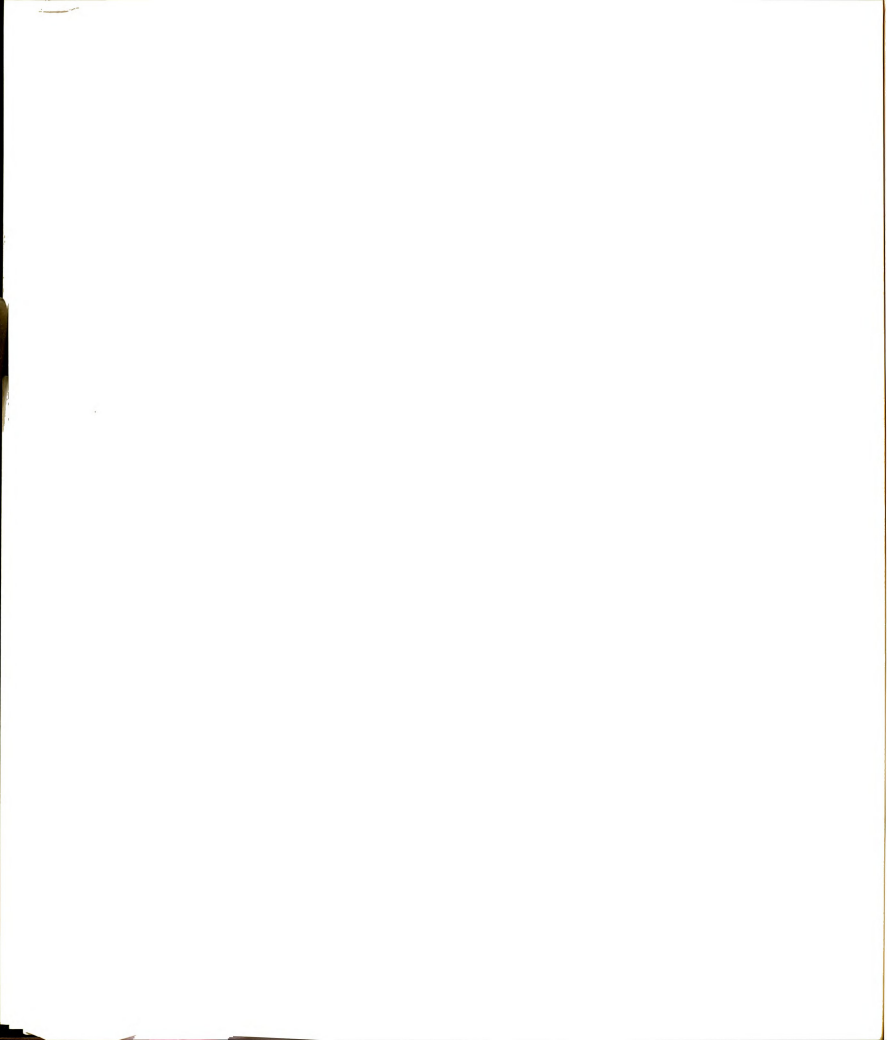
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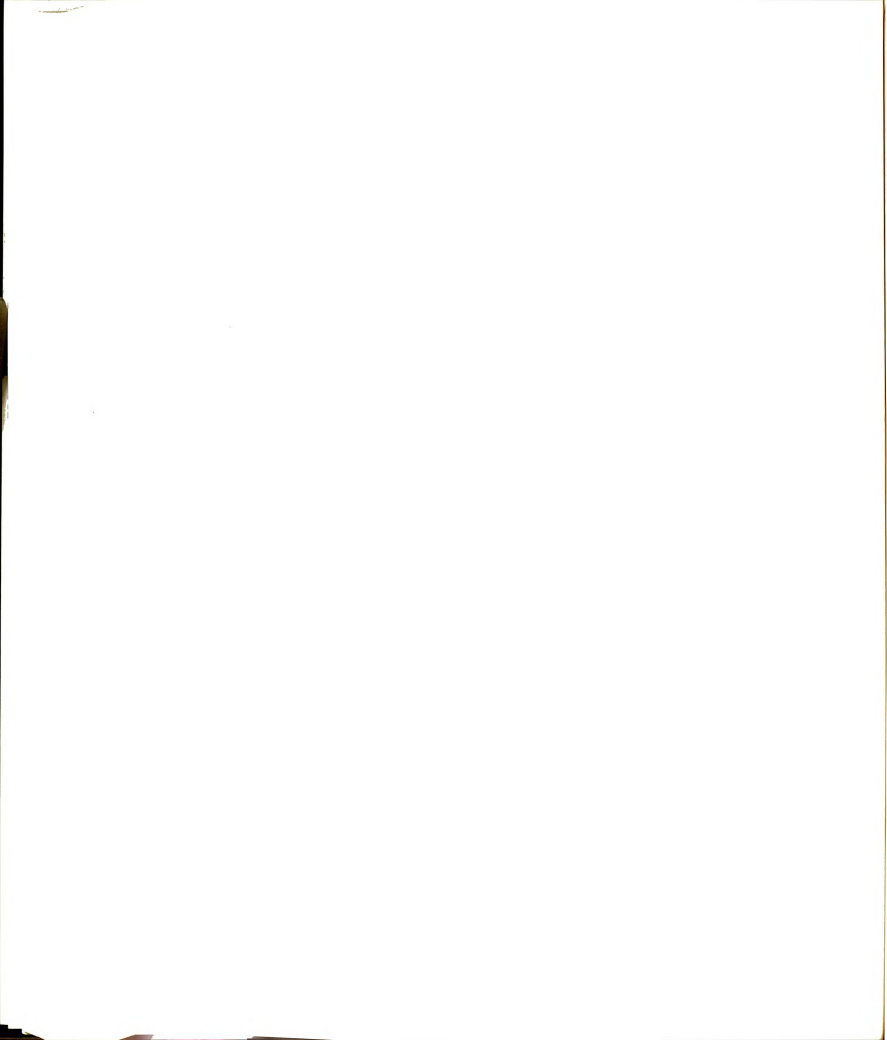


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DO 2 I=1,NZ
  I=I+NX+NY
  IP=I+4*(NX+NY)+3*NNZ
  L(IP,MATP)=CJ*SIN(Z(II))/377.0
2 CONTINUE
  IF(NNY.EQ.0) GO TO 6
  DO 4 I=1,NNY
    I=I+NX+NY+NZ+NNX
    IP=I+4*(NX+NY+NZ+NNX)+3*NNY
    L(IP,MATP)=CJ*SIN(Z(II))/377.0
4 CONTINUE
  IF(NNZ.EQ.0) GO TO 7
  DO 4 I=1,NNZ
    I=I+NX+NY+NZ+NNX+NNY
    IP=I+4*(NX+NY+NZ+NNX+NNY)+3*NNZ
    L(IP,MATP)=(-1.0)*CJ*SIN(Z(II))/377.0
7 CONTINUE
  CALL PRINT(IMPR,CRNT,L,MAT,MATP)
  DO 3 I=1,MAT
    L(I,MATP)=4.0*PI*L(I,MATP)
3 CONTINUE
  IF((NNX.EQ.0).AND.(NNY.EQ.0).AND.(NNZ.EQ.0)) GO TO 9
  MATP=4*(NX+NY+NZ)
  MATP=MATP+1
  DO 8 I=MATP,MAT
    L(I,MATP)=-1.0*L(I,MATP)
8 CONTINUE
9 RETURN
END
C*****
SUBROUTINE SHIM(L,MAT,MATP,Z,NCELL)
  REAL Z(NCELL),IMPR,MUO
  COMPLEX L(MAT,MATP),CJ
  COMMON /IMCR/IMPR,CRNT
  COMMON /NXZ/NX,NY,NZ,NNX,NNY,NNZ
  COMMON /CONS/PI,MUO,EPO
  COMMON /CPLX/CJ
  DO 1 I=1,NX
    IP=I+NX
    L(IP,MATP)=(-1.0)*CJ*SIN(Z(1))/377.0
1 CONTINUE
  DO 2 I=1,NY
    IP=I+NX+NY
    L(IP,MATP)=CJ*SIN(Z(II))/377.0
2 CONTINUE

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```

IF (NNX.EQ.0) GO TO 6
DO 4 I=1,NNX
  II=I+NX+NY+NZ
  IP=I+4*(NX+NY+NZ)+NNX
  L(IP,MATP)=CJ*SIN(Z(II))/377.0
4 CONTINUE
6 CONTINUE
IF (NNZ.EQ.0) GO TO 7
DO 5 I=1,NNZ
  II=I+NX+NY+NZ+NNX+NNY
  IP=I+4*(NX+NY+NZ+NNX+NNY)
  L(IP,MATP)=(-1.0)*CJ*SIN(Z(II))/377.0
5 CONTINUE
7 CONTINUE
CALL PRNT(IMPR,CRNT,L,MAT,MATP)
DO 3 I=1,MAT
  C(I,MATP)=4.0*PI*L(I,MATP)
3 CONTINUE
IF (NNX.EQ.0).AND.(NNY.EQ.0).AND.(NNZ.EQ.0) GO TO 9
L(I,MATP)=-1.0*L(I,MATP)
DO 8 I=1,MATP,MAT
  MATP=MATP+1
8 CONTINUE
9 RETURN
END
C*****SUBROUTINE I*****
SUBROUTINE I(L,MAT,MATP)
  COMPLEX L(MAT,MATP)
  COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
  DO 130 I=1,NX
    L(I,MATP)=-1.0*L(I,MATP)
  J=I+2*NX
  L(J,MATP)=L(J,MATP)*377.0
  K=J+NX
  L(K,MATP)=-377.0*L(K,MATP)
130 CONTINUE
DO 131 I=1,NY
  II=I+NY+4*NX
  L(II,MATP)=-1.0*L(II,MATP)
  JJ=II+NY
  L(JJ,MATP)=-377.0*L(JJ,MATP)
  KK=JJ+NY
  L(KK,MATP)=L(KK,MATP)*377.0
131 CONTINUE
DO 132 I=1,NZ
  II=I+4*(NX+NY)
  L(II,MATP)=-1.0*L(II,MATP)

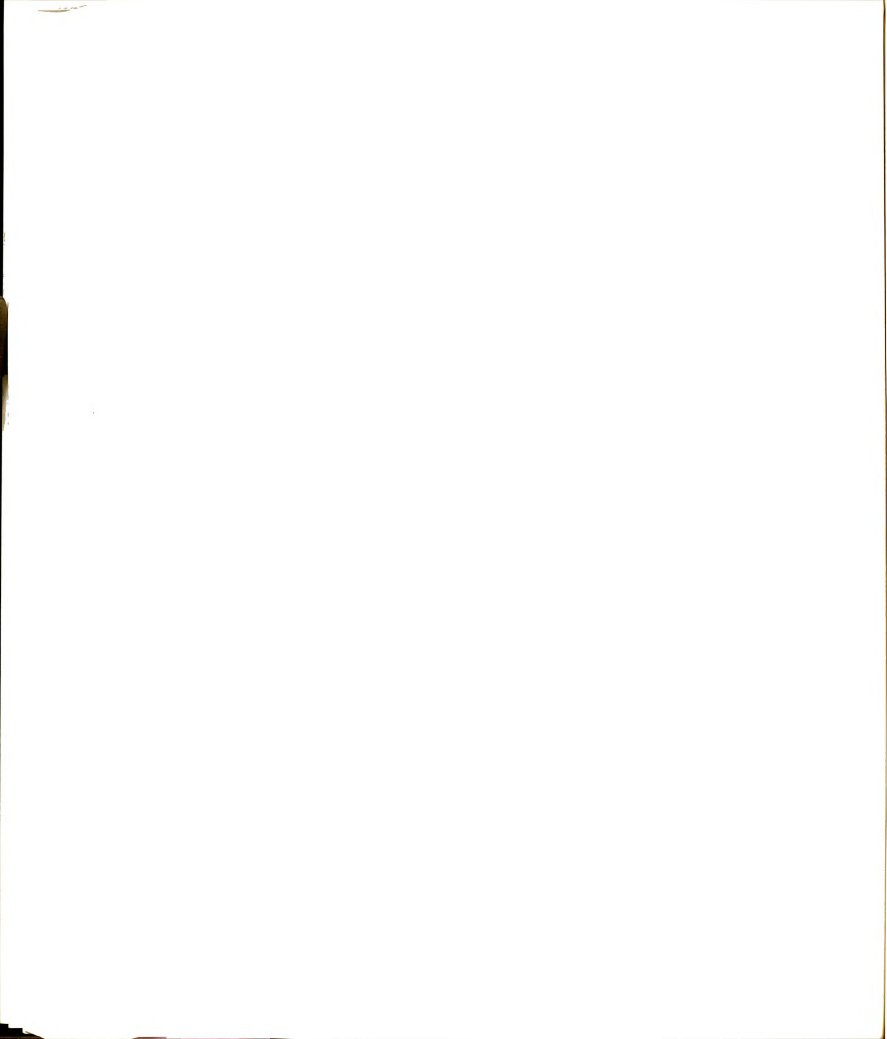
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8980
8990
9000

9010
9020
9030

C*****SUBROUTINE I*****
SUBROUTINE I(L,MAT,MATP)
COMPLEX L(MAT,MATP)
COMMON /XYZ/NX,NY,NZ,NNX,NNY,NNZ
DO 130 I=1,NX
L(I,MATP)=-1.0*L(I,MATP)
J=I+2*NX
L(J,MATP)=L(J,MATP)*377.0
K=J+NX
L(K,MATP)=-377.0*L(K,MATP)
130 CONTINUE
DO 131 I=1,NY
II=I+NY+4*NX
L(II,MATP)=-1.0*L(II,MATP)
JJ=II+NY
L(JJ,MATP)=-377.0*L(JJ,MATP)
KK=JJ+NY
L(KK,MATP)=L(KK,MATP)*377.0
131 CONTINUE
DO 132 I=1,NZ
II=I+4*(NX+NY)
L(II,MATP)=-1.0*L(II,MATP)

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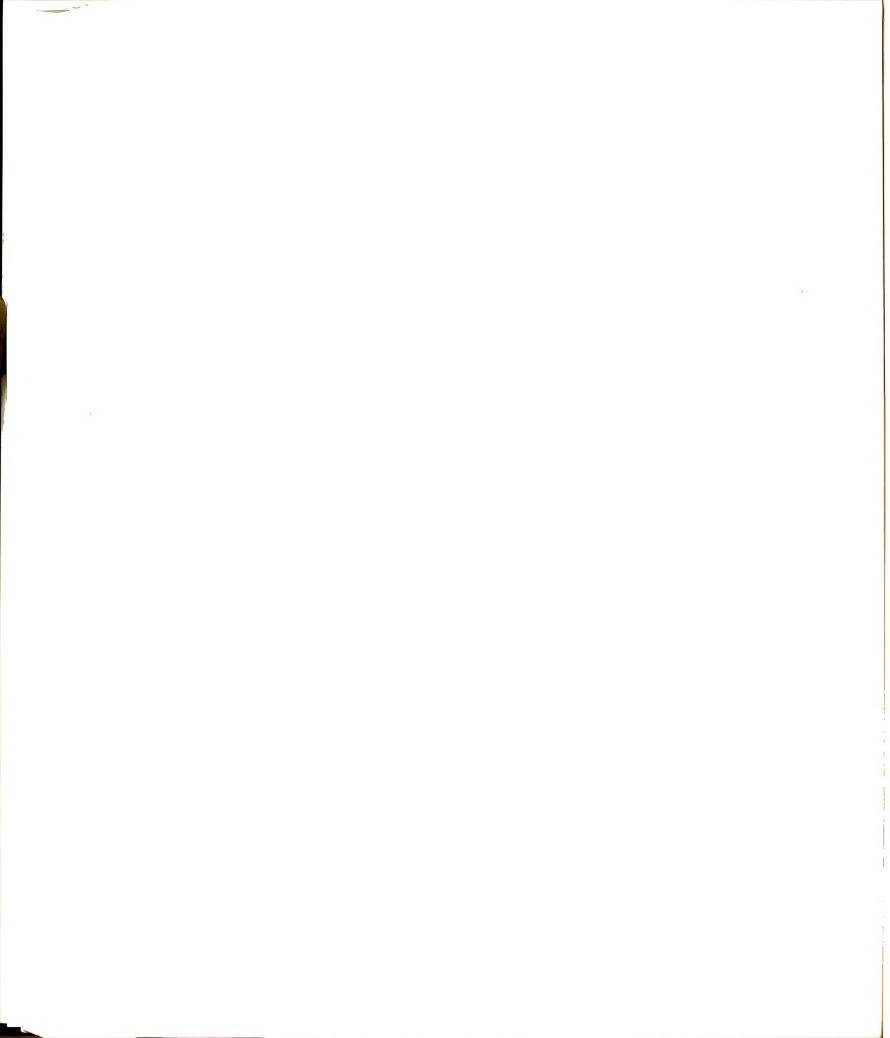
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J=I1+2*NZ
L(J,MATP)=L(J,MATP)*377.0
K=J+NZ
L(K,MATP)=-377.0*L(K,MATP)
CONTINUE
132 IF(NNX.EQ.0) GO TO 136
DO 133 I=1,NNX
  I1=I+4*(NX+NY+NZ)+NNX
  L(I1,MATP)=-1.0*L(I1,MATP)
  J=I1+NNX
  L(J,MATP)=-377.0*L(J,MATP)
  K=J+NNX
  L(K,MATP)=377.0*L(K,MATP)
CONTINUE
133 IF(NNY.EQ.0) GO TO 137
DO 134 I=1,NNY+NZ+NNX
  I1=I+4*(NX+NY+NZ)+NNX
  L(I1,MATP)=-1.0*L(I1,MATP)
  J=I1+2*NNY
  L(J,MATP)=377.0*L(J,MATP)
  K=J+NNY
  L(K,MATP)=-377.0*L(K,MATP)
CONTINUE
134 CONTINUE
137 CONTINUE
IF(NNZ.EQ.0) GO TO 138
DO 135 I=1,NNZ
  I1=I+4*(NX+NY+NZ+NNX+NNY)+NNZ
  L(I1,MATP)=-1.0*L(I1,MATP)
  J=I1+NNZ
  L(J,MATP)=-377.0*L(J,MATP)
  K=J+NNZ
  L(K,MATP)=377.0*L(K,MATP)
CONTINUE
135 CONTINUE
138 RETURN
END
C***** SUBROUTINE PRINT*****
SUBROUTINE PRINT(IMIN,CRFL,L,MAT,MATP)
  REAL IMIN,LEFT
  COMMON L(MAT,MATP)
  DATA TOP/4HTOP/,RIGHT/4HRIGH/,REAR/4HREAR/,FLDS/4HFLDS/
  DATA BTM/4HBTM/,LEFT/4HLEFT/,FRNT/4HFRNT/
  PRINT 111,IMIN,CRFL
  111 FORMAT(1H1,5X,4HTHE, A4,1X,A4,81H ARE LISTED AS FOLLOWS (AMP/METER
      2 FOR CRRNTS OR H FIELDS, VOL/METER FOR E FIELDS)/1H-.3H N ,6X,4HRE

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9630
9640
9650
9660
9670

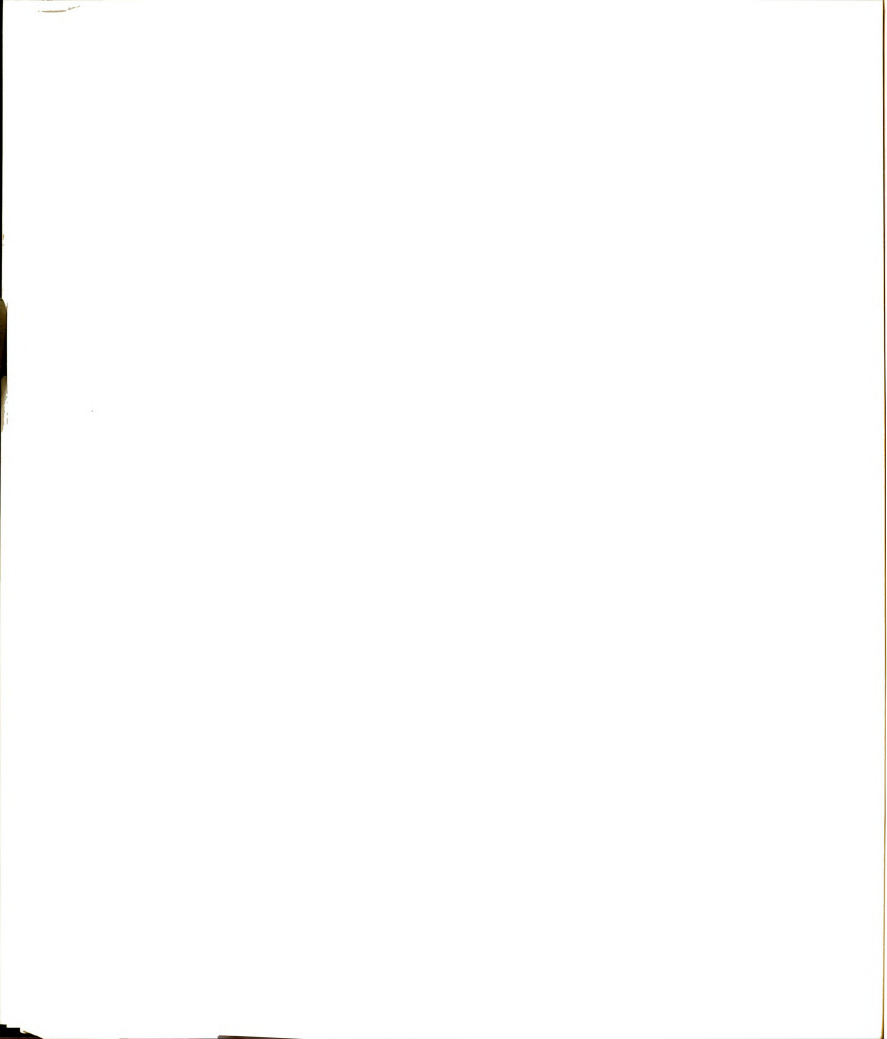
9700




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3AL, 11X, 4HIMAG, 11X, 4HREAL, 11X, 4HIMAG, 11X, 4HREAL, 11X, 4HRE
4AL, 11X, 4HIMAG)
PRINT 123, TOP
123 FORMAT(1H-, 15X, 7HON THE , A4, 28H SURFACE WE HAVE THE FOLLOWS) 9750
IF(CRFL.EQ.FLDS) GO TO 136
PRINT 112 9770
112 FORMAT(1H-, 14X, 10HJEY OR -HZ, 20X, 9HJEZ OR HY, 20X, 12HJMY OR EZ/WO, 1 9780
28X, 13HJMZ OR -EY/WO)
PRINT 136 9810
136 PRINT 137 9820
137 FORMAT(1H-, 14X, 10H , HZ , 20X, 9H HY , 20X, 12H EZ , 1
28X, 13HJMY OR -EY
138 PRINT 122
122 FORMAT(1H-, 5X, 4(28H*****2X)) 9850
DO 12 I=1, NX
I1=1
J1=1
IP=J+NX
IPP=IP+NX
IPPP=IPP+NX
12 PRINT 13, I1, L(J, MATP), L(IP, MATP), L(IPP, MATP), L(IPPP, MATP)
13 FORMAT(1H-, 13, 1X, 4(1X, 1H(, 2E13. 6, 1H), 1X))
IF(CRFL.EQ.FLDS) GO TO 139
PRINT 113 9960
113 FORMAT(1H-, 14X, 9HJEX OR HZ, 20X, 10HJEZ OR -HX, 20X, 13HJMX OR -EZ/WO, 9970
218X, 12HJMZ OR EX/WO)
GO TO 141 10020
139 PRINT 140
140 FORMAT(1H-, 14X, 9H , HZ , 20X, 10H HX , 20X, 13H EZ ,
218X, 12HJMX OR EX
141 PRINT 122 10050
I1=14, NX
I1=14, NX
IP=J+NX
IPP=IP+NX
IPPP=IPP+NX
14 PRINT 13, I1, L(J, MATP), L(IP, MATP), L(IPP, MATP), L(IPPP, MATP) 10060
PRINT 122, REAR 10070
IF(CRFL.EQ.FLDS) GO TO 142 10080
PRINT 114 10090
114 FORMAT(1H-, 14X, 10HJEX OR -HY, 20X, 9HJEY OR HX, 20X, 12HJMX OR EY/WO, 1 10100
28X, 13HJMY OR -EX/WO) 10110
GO TO 144 10120
142 PRINT 143 10130
10140
10150
10160
10190
10200

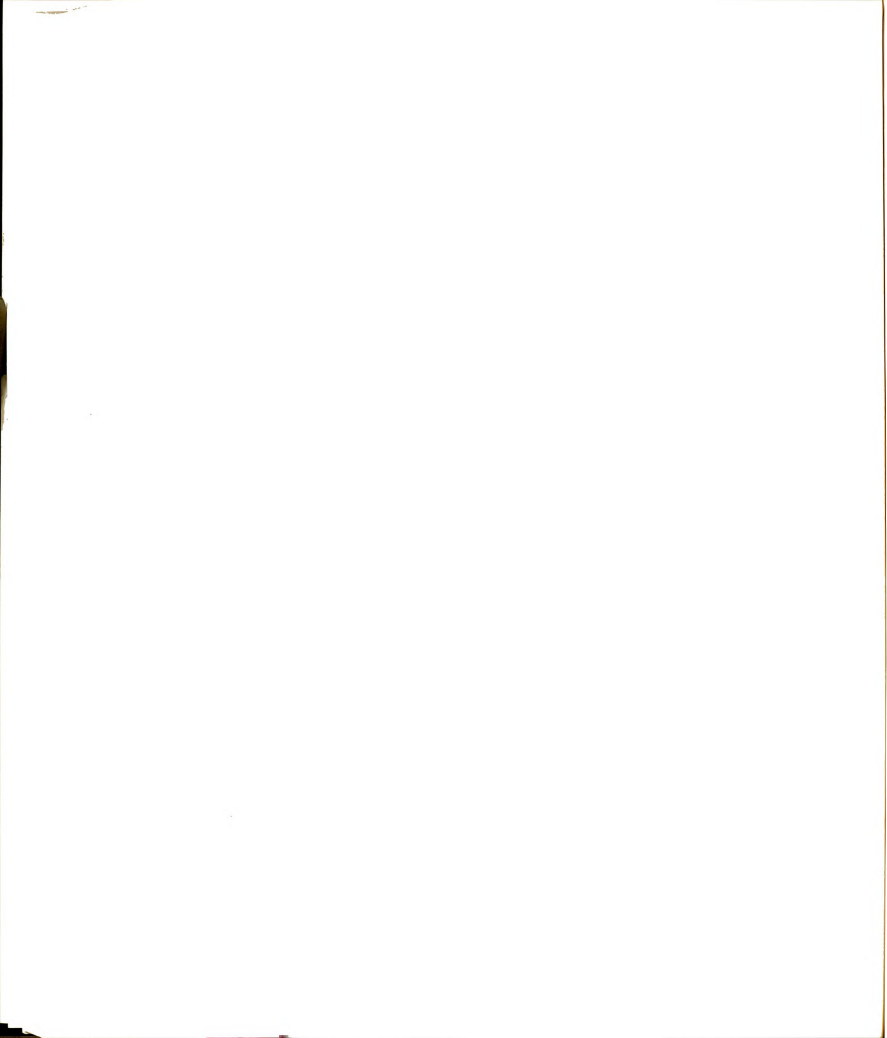
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143 FORMAT(1H-,14X,10H , HY , 20X,9H HX , 20X,12H EY , 1
144 28X,13H EX
DO 15 I=1,NZ
I=1+NX+NY
J=I+4*(NX+NY)
IP=J+NZ
IPP=IP+NZ
IPPP=IPP+NNX
15 PRINT 13,II,L(J,MATP),L(IP,MATP),L(IPP,MATP),L(IPPP,MATP)
PRINT 122
IF(NNX.EQ.O) GO TO 19
PRINT 123,BTTM
IF(CRFL.EQ.FLDS) GO TO 145
PRINT 146
146 FORMAT(1H-,14X,9HJEY OR HZ,20X,10HJEZ OR -HY,20X,
213HJMY OR -EZ/WO,18X,12HJNZ OR EY/WO)
GO TO 147
145 PRINT 137
DO 16 I=1,NNX
I=1+NX+NY+NZ
J=I+4*(NX+NY+NZ)
IP=J+NNX
IPP=IP+NNX
IPPP=IPP+NNX
16 PRINT 13,II,L(J,MATP),L(IP,MATP),L(IPP,MATP),L(IPPP,MATP)
19 CONTINUE
IF(NNY.EQ.O) GO TO 20
PRINT 122,LEFT
IF(CRFL.EQ.FLDS) GO TO 149
PRINT 150
150 FORMAT(1H-,14X,10HJEX OR -HZ,20X,9HJEZ OR HX,20X,
212HJMY OR EZ/WO,18X,13HJNZ OR -EX/WO)
GO TO 151
149 PRINT 140
DO 17 I=1,NNY
I=1+NX+NY+NZ+NNX
J=I+4*(NX+NY+NZ+NNX)
IP=J+NNY
IPP=IP+NNY
IPPP=IPP+NNY
17 PRINT 13,II,L(J,MATP),L(IP,MATP),L(IPP,MATP),L(IPPP,MATP)
PRINT 122
20 CONTINUE
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10670

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```

IF (NNZ.EQ.0) GO TO 21
PRINT 123,FRNT
IF (CRFL.EQ.FLDS) GO TO 153
PRINT 154
154 FORMAT(1H-,14X,9HJEX OR HY,20X,10HJEY OR -HX,20X,
213HJMX OR -EY/WO,18X,12HJMY OR EX/WO)
GO TO 155
153 PRINT 143
155 PRINT 122
DO 18 I=1,NNZ
I1=I+NX+NY+NZ+NNX+NNY
J1=I+*(NX+NY+NZ+NNX+NNY)
IP=J+NNZ
IPP=IP+NNZ
18 PRINT 13,11,L(J,MATP),L(IP,MATP),L(IPP,MATP),L(IPP,MATP)
21 CONTINUE
21 RETURN
END
C*****SUBROUTINE COEFFX*****
SUBROUTINE COEFFX(MDIR,AREA,XM,YM,ZM,XN,YN,ZN,C)
REAL D(3),R(5),MUO
COMPLEX C(16),G(4),G(4),F(10),CJ,FK,FHUN
COMPLEX H(5),HT(5),FT(10)
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
COMMON /PARA/FK,FHUN
COMMON /NUMDV/NDIV
NDIV2=NDIV*NDIV
DL=SQRT(AREA)
DL=DL/NDIV
MDIRM=MDIR
IF (MDIR.GE.4) MDIRM=MDIR-3
IF (MDIRM-2) 100,200,300
DO 1 I=1,5
HT(I)=0.0
CONTINUE
1 HT(I)=DL*(DL+DL)/2
DO 4 J=1,NDIV
YT=YT-DI,NDIV
DO 3 K=1,NDIV
ZT=ZT-DI,NDIV
CALL DRG(XM,YM,ZM,XN,YT,ZT,D,R,G)
CALL GG(R,G)
CALL FCNXX(AREA,D,G,H)
DO 2 I=1,5
HT(I)=HT(I)+H(I)
2 CONTINUE

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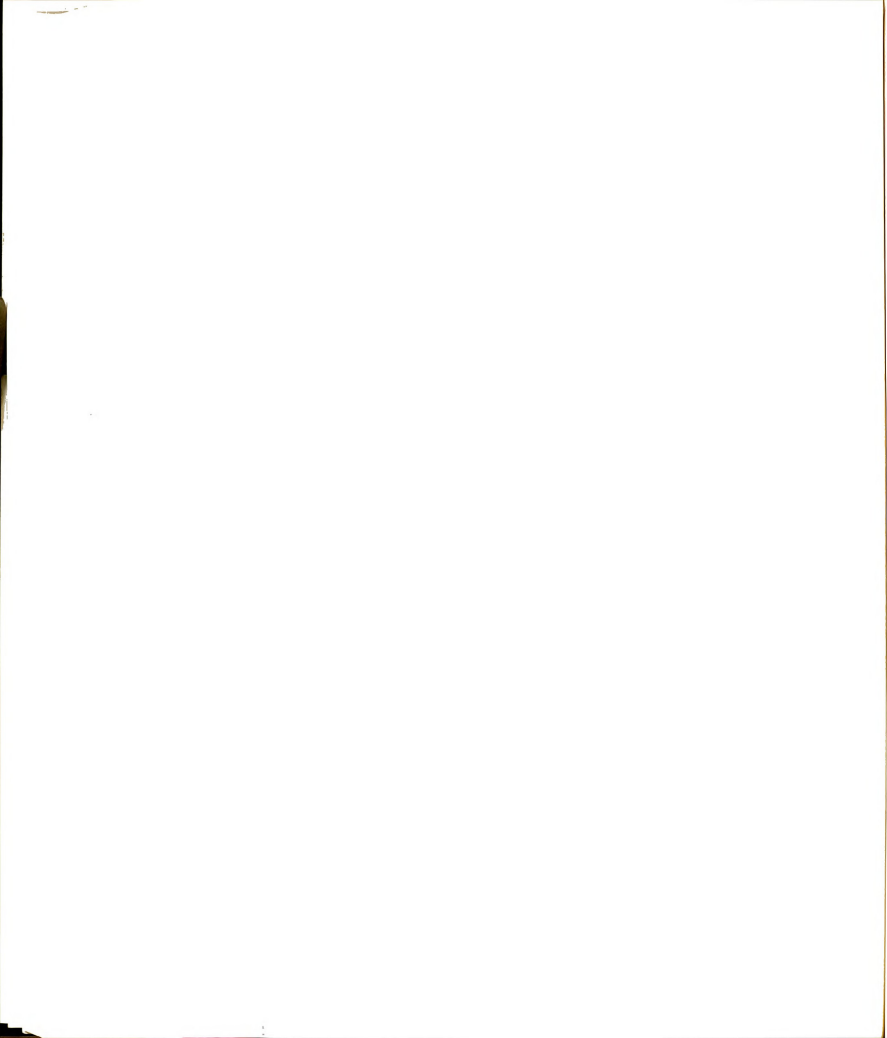
10680
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11990
12000

2



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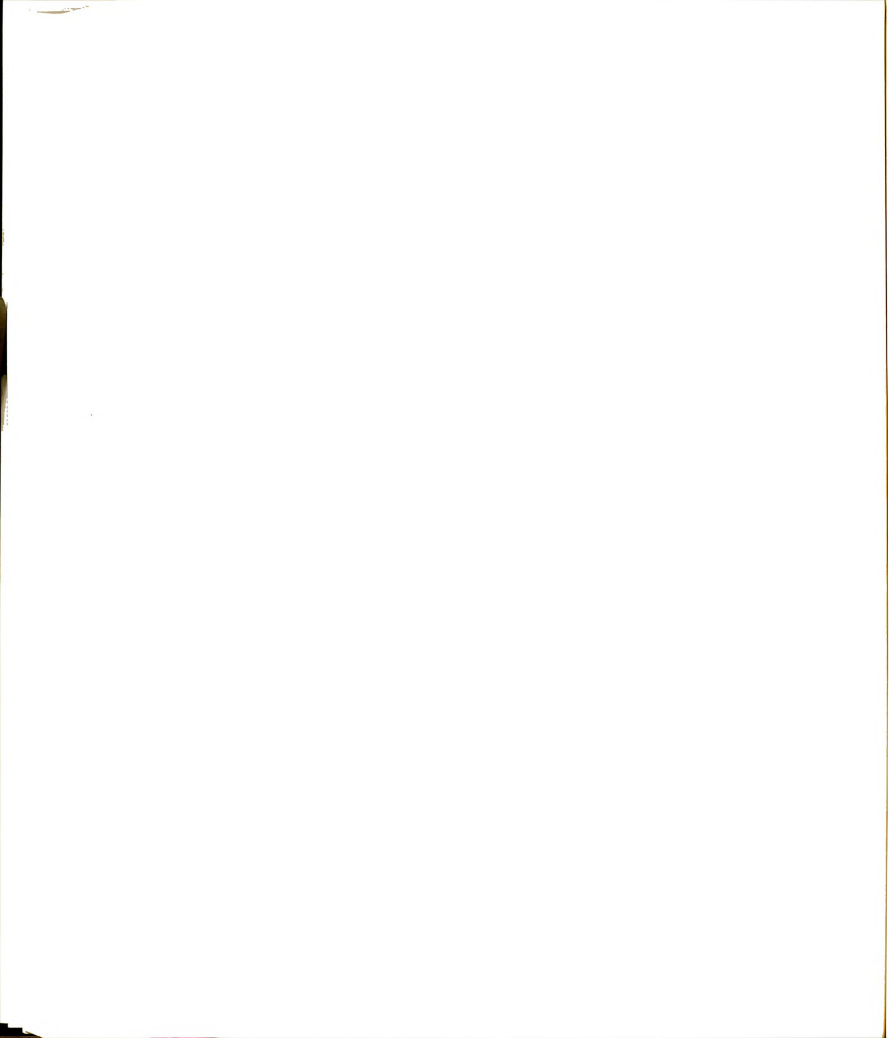
3    CONTINUE
4    ZT=ZT+DLEN
5    CONTINUE
   DO 5 I=1, 5
   HT(I)=HT(I)/NDIV2
   CONTINUE
   CALL CEFXX(HT, C)
   RETURN
200  DO 101 I=1, 10
   FT(I)=0
   CONTINUE
   ZT=ZT+(DLEN+DL)/2
   DO 104 J=1, NDIV
   VT=VT-DL
   DO 103 K=1, NDIV
   ZT=ZT-DL
   CALL DRG(XM, YM, ZM, XN, YT, ZT, D, R, G)
   CALL GG(R, G, G)
   CALL FCNXX(AREA, D, G, F)
   DO 102 I=1, 10
   FT(I)=FT(I)+F(I)
   CONTINUE
102  CONTINUE
103  ZT=ZT+DLEN
104  CONTINUE
   DO 105 I=1, 10
   FT(I)=FT(I)/NDIV2
   CONTINUE
105  CALL CEFXX(FT, C)
   RETURN
300  DO 201 I=1, 10
   FT(I)=0
   CONTINUE
201  YZ=YZ+(DLEN+DL)/2
   ZT=ZN+(DLEN+DL)/2
   DO 204 J=1, NDIV
   VT=VT-DL
   DO 203 K=1, NDIV
   ZT=ZT-DL
   CALL DRG(XM, YM, ZM, XN, YT, ZT, D, R, G)
   CALL GG(R, G, G)
   CALL FCNXX(AREA, D, G, F)
   DO 202 I=1, 10
   FT(I)=FT(I)+F(I)
   CONTINUE
202  CONTINUE
203

```

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12290
12300
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12470

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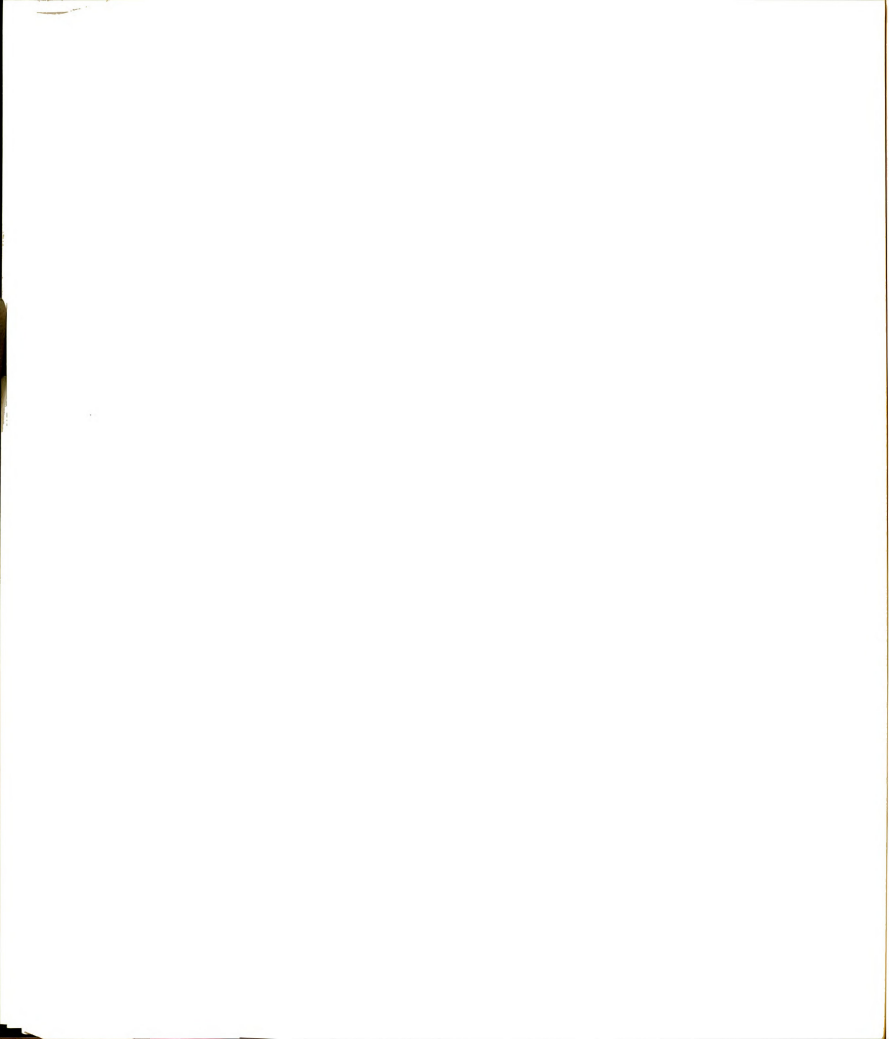
204 ZT=ZT+DLEN
CONTINUE
DO 205 I=1,10
FT(I)=FT(I)/NDIV2
205 CONTINUE
CALL CEFXY(FT,C)
RETURN
END
C***SUBROUTINE COEFY*****
SUBROUTINE COEFY(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
REAL D(3), R(5), MUO
COMPLEX C(16), G(4), F(10), CJ, FK, FHUN
COMPLX H(5), HT(5), FT(10)
COMMON /CONS/PI, MUO, EPO
COMMON /CPLX/CJ
COMMON /PARA/FK, FHUN
COMMON /NUMDV/NDIV
NDIV2=NDIV*NDIV
DLEN=SGRT(AREA)
DL=DLEN/NDIV
MDIRM=MDIR
IF(MDIR.GE.4) MDIRM=MDIR-3
IF(MDIRM-2) 100,200,300
DO 1 I=1,10
FT(I)=0.0
CONTINUE
ZT=ZN+(DLEN+DL)/2
XT=XN+(DLEN+DL)/2
DO 4 J=1,NDIV
ZT=ZT-DL
DO 3 K=1,NDIV
XT=XT-DL
CALL DRG(XM, YM, ZM, XT, YN, ZT, D, R, G)
CALL GG(R, G, G)
CALL FCNXY(AREA, D, G, F)
DO 2 I=1,10
FT(I)=FT(I)+F(I)
CONTINUE
XT=XT+DLEN
CONTINUE
DO 5 I=1,10
FT(I)=FT(I)/NDIV2
CONTINUE
CALL CEFXY(FT,C)
RETURN
DO 101 I=1,5
HT(I)=0.0

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12480
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12500
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12650
12660
12670
12680
12690
12700
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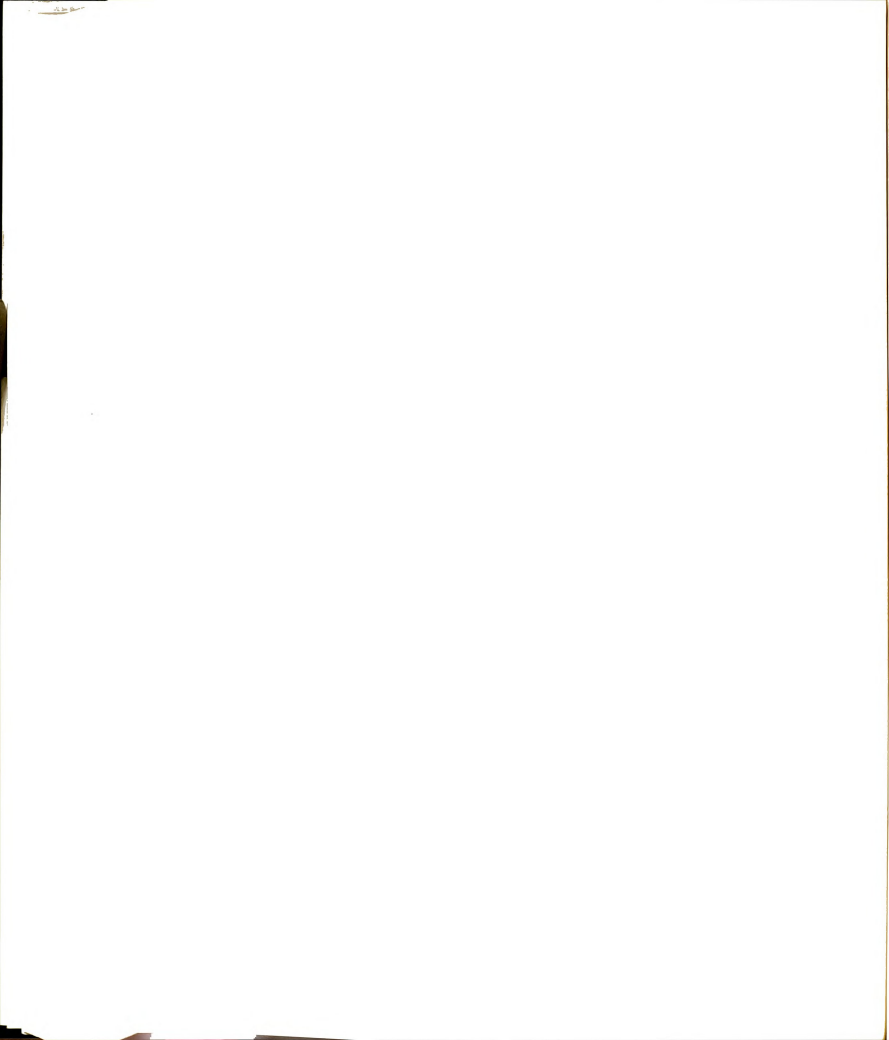
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101 CONTINUE
   ZT=ZN+(DLEN+DL)/2
   XT=XN+(DLEN+DL)/2
   DO 104 J=1,NDIV
     ZT=ZT-DL
     DO 103 K=1,NDIV
       XT=XT-DL
       CALL DRG(XM, YM, ZM, XT, YN, ZT, D, R, G)
       CALL GG(R, G, G)
       CALL FCNXY(AREA, D, G, H)
     DO 102 I=1,5
       HT(I)=HT(I)+H(I)
     CONTINUE
   CONTINUE

102 XT=XT+DLEN
103 CONTINUE

104 CONTINUE
   DO 105 I=1,5
     HT(I)=HT(I)/NDIV2
   CONTINUE
   CALL CEFYY(HT, C)
   RETURN
300 FT(I)=0,0
201 CONTINUE
   ZT=ZN+(DLEN+DL)/2
   XT=XN+(DLEN+DL)/2
   DO 204 J=1,NDIV
     ZT=ZT-DL
     DO 203 K=1,NDIV
       XT=XT-DL
       CALL DRG(XM, YM, ZM, XT, YN, ZT, D, R, G)
       CALL GG(R, G, G)
       CALL FCNXY(AREA, D, G, F)
     DO 202 I=1,10
       FT(I)=FT(I)+F(I)
     CONTINUE
   CONTINUE
   YN=YN+DLEN
202 CONTINUE
203 CONTINUE
204 DO 205 I=1,10
   FT(I)=FT(I)/NDIV2
205 CONTINUE
   CALL CEFZY(FT, C)
   RETURN
   END
12940
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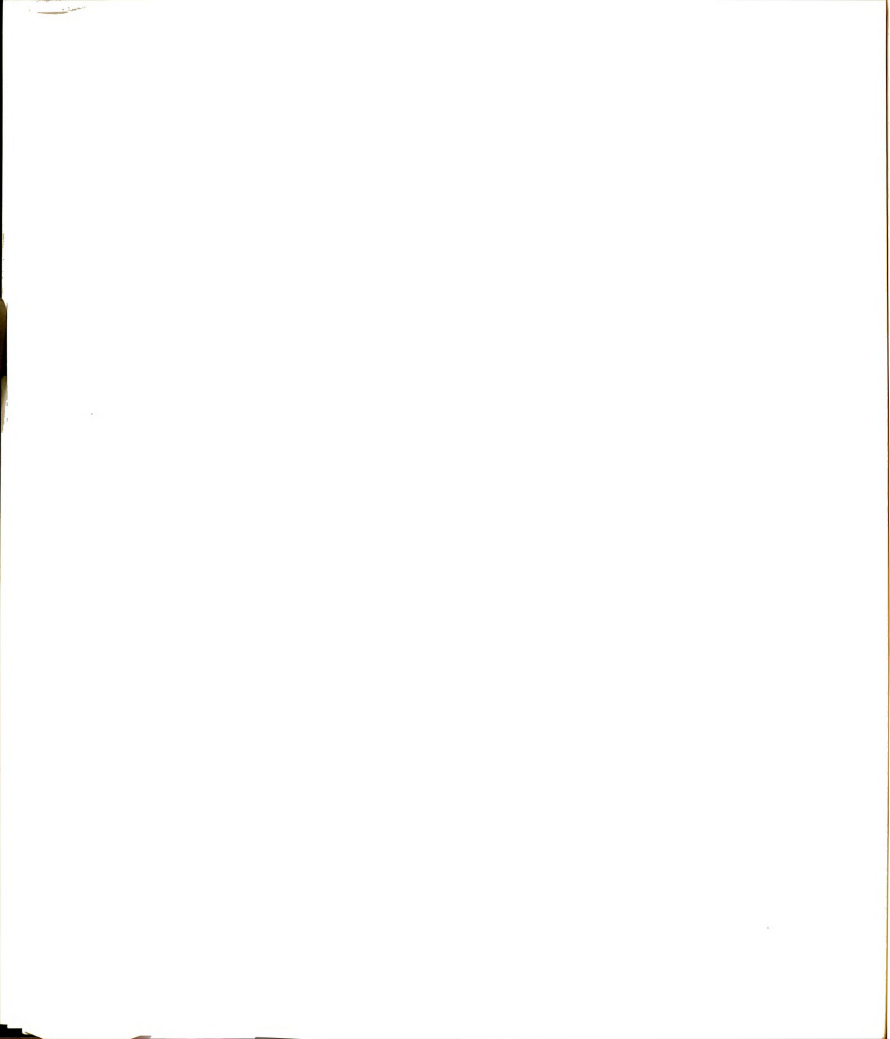
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*****
C*****SUBROUTINE COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
SUBROUTINE COEFZ(MDIR, AREA, XM, YM, ZM, XN, YN, ZN, C)
REAL D(3), R(5), MUO
COMPLEX C(16), G(4), F(10), C(4), F(10), C(4), F(10)
COMPLEX H(5), HI(5), FI(10)
COMMON /CONS/PI, MUO, EPO
COMMON /CPLX/CJ
COMMON /PARA/FK, FHUN
COMMON /NUMDV/NDIV
NDIV2=NDIV*NDIV
DLEN=SGRT(AREA)
DL=DLEN/NDIV
MDIRM=MDIR
IF(MDIR.GE.4) MDIRM=MDIR-3
IF(MDIRM-2) 100, 200, 300
DO 1 I=1, 10
  FT(I)=0.0
  CONTINUE
  100
  1
  YI=YN+(DLEN+DL)/2
  XT=XN+(DLEN+DL)/2
  DO 4 J=1, NDIV
    DO 3 K=1, NDIV
      YI=XT-DL
      XT=XI-DL
      CALL DRG(XM, YM, ZM, XT, YT, ZN, D, R, G)
      CALL GG(R, G, Q)
      CALL FCNZ(AREA, D, G, F)
      DO 2 I=1, 10
        FT(I)=FT(I)+F(I)
      CONTINUE
      CONTINUE
      XT=XI+DLEN
      CONTINUE
      2
      3
      4
      DO 5 I=1, 10
        FT(I)=FT(I)/NDIV2
      CONTINUE
      CALL CEFXZ(FT, C)
      RETURN
      5
      DO 101 I=1, 10
        FT(I)=0.0
      CONTINUE
      101
      YI=YN+(DLEN+DL)/2
      XT=XN+(DLEN+DL)/2
      DO 104 J=1, NDIV
        YI=XT-DL
        XT=XI-DL
        DO 103 K=1, NDIV
          YI=XT-DL
          XT=XI-DL
          103
          104
          101
          100
          1
          100
          1
          2
          3
          4
          5
          200
          101
          13390
          13400
          13410
          13420
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          13690
          13700
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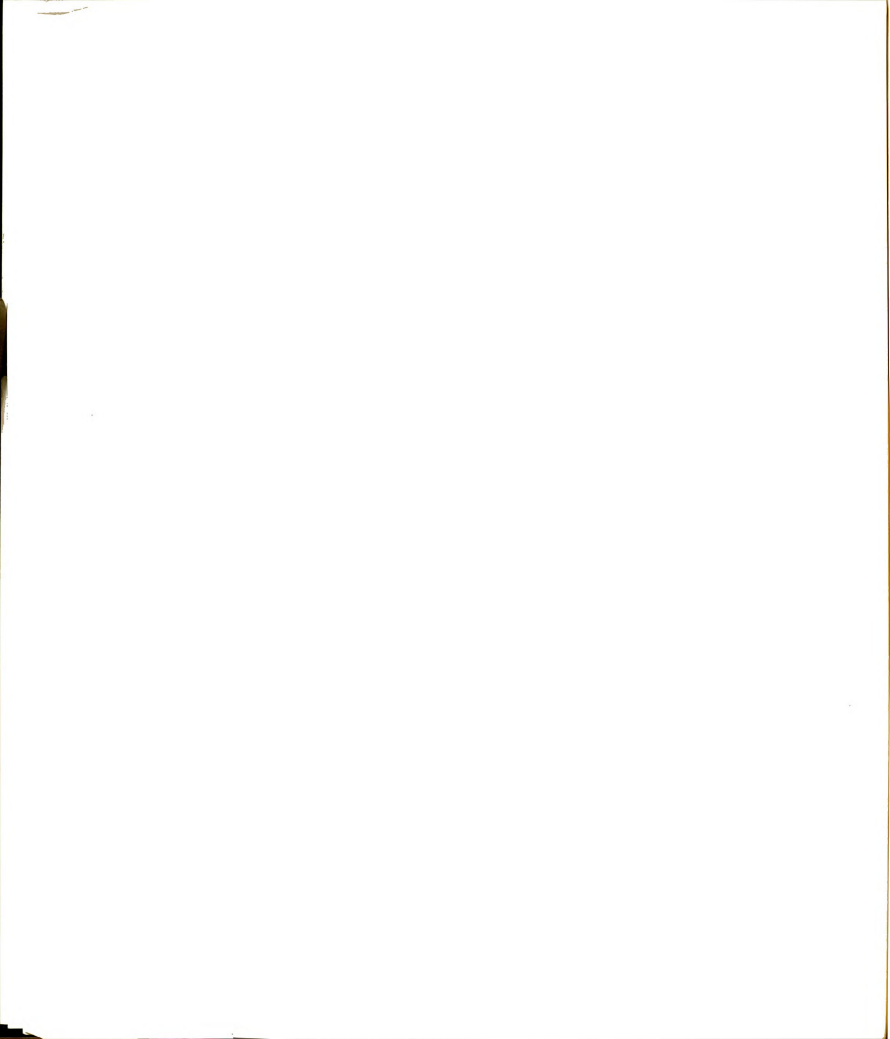
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102 CALL DRG(XM, YM, ZM, X1, YT, ZN, D, R, G)
103 CALL GG(R, G, Q)
104 CALL FCNZZ(AREA, D, G, F)
105 DO 102 I=1, 10
    FT(I)=FT(I)+F(I)
    CONTINUE
106 CONTINUE
107 XT=XT+DLEN
108 CONTINUE
109 DO 105 I=1, 10
    FT(I)=FT(I)/NDIV2
110 CONTINUE
111 CALL CEFYZ(FT, C)
112 RETURN
113 DO 201 I=1, 5
    HT(I)=0
114 CONTINUE
115 YI=YN+(DLEN+DL)/2
116 XT=XN+(DLEN+DL)/2
117 DO 204 J=1, NDIV
    YT=YT-DL
118 DO 203 K=1, NDIV
    XT=XT-DL
119 CALL DRG(XM, YM, ZM, XT, YT, ZN, D, R, G)
120 CALL GG(R, G, Q)
121 CALL FCNZZ(AREA, D, G, H)
122 DO 202 I=1, 5
    HT(I)=HT(I)+H(I)
123 CONTINUE
124 CONTINUE
125 XT=XT+DLEN
126 DO 205 I=1, 5
    HT(I)=HT(I)/NDIV2
127 CONTINUE
128 CALL CEFZZ(HT, C)
129 RETURN
130 *****
131 SUBROUTINE GG(R, G, Q)
132 SUBROUTINE MUO
133 REAL R(5), MUO
134 COMPLEX G(4), Q(4), CJ
135 COMMON /CONS/PI, MUO, EPO
136 COMMON /CPLX/CJ
137 G(1)=CJ*G(2)/R(2)+G(1)/R(3)
138 G(2)=G(3)/R(3)-3.0*CJ*G(2)/R(4)-3.0*G(1)/R(5)
139 G(3)=CJ*G(4)/R(2)+G(3)/R(3)
140 G(4)=-1.0*G(3)/R(1)
141 RETURN
142 END
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C*****COMPLEX FUNCTION P*****
COMPLEX FUNCTION P(I,A)
REAL MUO
COMPLEX CJ,FK,FHUN
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
COMMON /PARA/FA,FK,FHUN
IF(I.EQ.2) GO TO 1
P=CJ*CEXP(-1.0*CJ*FA)-CJ*FK*CEXP(-1.0*CJ*FK*FA)+(CEXP(-1.0*CJ*FA)-CEX
2P(-1.0*CJ*FK*FA))/A
RETURN
1 P=2.0*PI*CJ*((FK*CEXP(-1.0*CJ*FK*FA)-CEXP(-1.0*CJ*FA))-(FK-1.0))
RETURN
END
C*****SUBROUTINE DRG*****
SUBROUTINE DRG(XM,YM,ZM,XN,YN,ZN,D,R,G)
REAL D(3),R(5),MUO
COMPLEX G(4),EX(2),CJ,FK,FHUN
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
COMMON /PARA/FA,FK,FHUN
D(1)=XN-XM
D(2)=YN-YM
D(3)=ZN-ZM
DD=D(1)*D(1)+D(2)*D(2)+D(3)*D(3)
R(1)=SQRT(DD)
R(2)=R(1)*R(1)
R(3)=R(2)*R(1)
R(4)=R(3)*R(1)
R(5)=R(4)*R(1)
EX(1)=CEXP(-1.0*CJ*R(1))
EX(2)=CEXP(-1.0*CJ*FK*R(1))
G(1)=EX(1)-EX(2)
G(2)=EX(1)-FK*EX(2)
G(3)=EX(1)-FK*FK*EX(2)
G(4)=EX(1)-FK*FK*FK*EX(2)
RETURN
END
C*****SUBROUTINE LEE*****
SUBROUTINE LEE(E,MT,MP,MPP,NPP,NP,NT,NP,NPP,NPP,L,MAT,MATP)
COMPLEX E(16),L(MAT,MATP)
L(MT,NT)=E(1)
L(MP,NT)=E(2)
L(MPP,NT)=E(3)
L(MPP,NP)=E(4)
L(MT,NP)=E(5)
L(MP,NP)=E(6)

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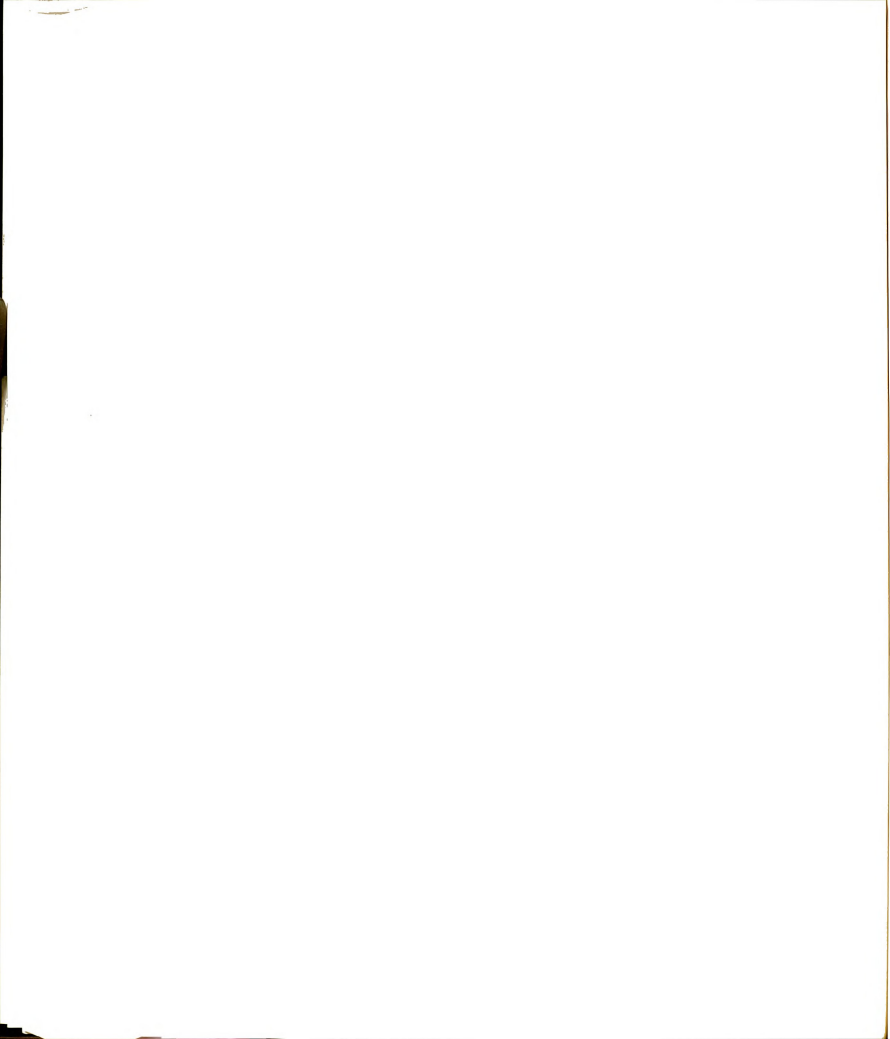
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C*****
L(MPP,NP)=E(7)
L(MPPP,NP)=E(8)
L(MT,NPP)=E(9)
L(MP,NPP)=E(10)
L(MPP,NPP)=E(11)
L(MPPP,NPP)=E(12)
L(MT,NPPP)=E(13)
L(MP,NPPP)=E(14)
L(MPP,NPPP)=E(15)
L(MPPP,NPPP)=E(16)
RETURN
END
SUBROUTINE FCNXX(AREA,D,G,F)
REAL D(3),MUO
COMPLEX G(4),F(5),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(1)*G(1)
F(2)=CJ*AREA*D(2)*D(3)*G(2)
F(3)=CJ*AREA*(G(1)+D(3)*D(3)*G(2)+G(4))
F(4)=CJ*AREA*(G(1)+D(2)*D(2)*G(4))
F(5)=AREA*D(1)*G(3)
RETURN
C*****

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C*****
SUBROUTINE FCNXY(AREA,D,G,F)
REAL D(3),MUO
COMPLEX G(4),F(10),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(2)*G(1)
F(2)=CJ*AREA*(1)*D(3)*G(2)
F(3)=CJ*AREA*(G(1)+D(3)*D(3)*G(2)+G(4))
F(4)=AREA*D(3)*G(1)
F(5)=AREA*D(1)*G(1)
F(6)=CJ*AREA*D(2)*D(1)*G(2)
F(7)=CJ*AREA*D(2)*D(3)*G(2)
F(8)=AREA*D(2)*G(3)
F(9)=AREA*D(3)*G(3)
F(10)=AREA*D(1)*G(3)
RETURN
END
SUBROUTINE FCNXX(AREA,D,G,F)
REAL D(3),MUO
COMPLEX G(4),F(10),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ

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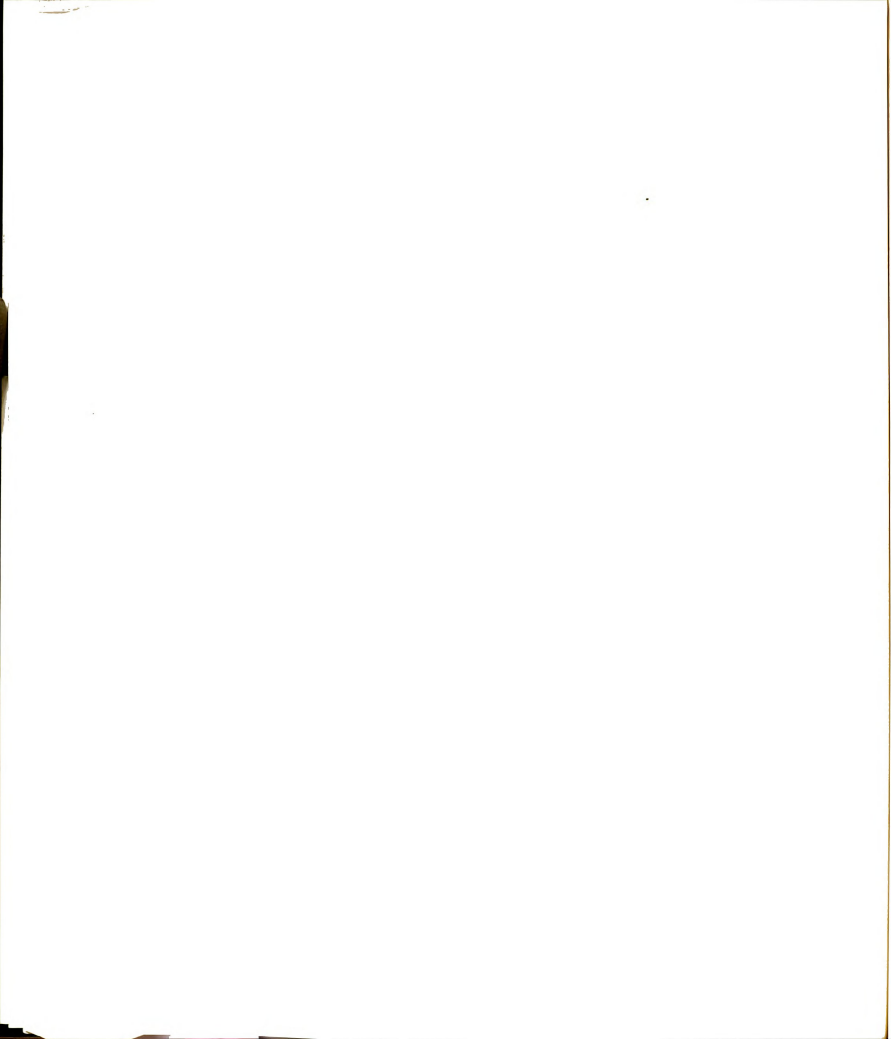
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C*****
SUBROUTINE FCNXX(AREA,D,G,F)
REAL D(3),MUO
COMPLEX G(4),F(10),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ

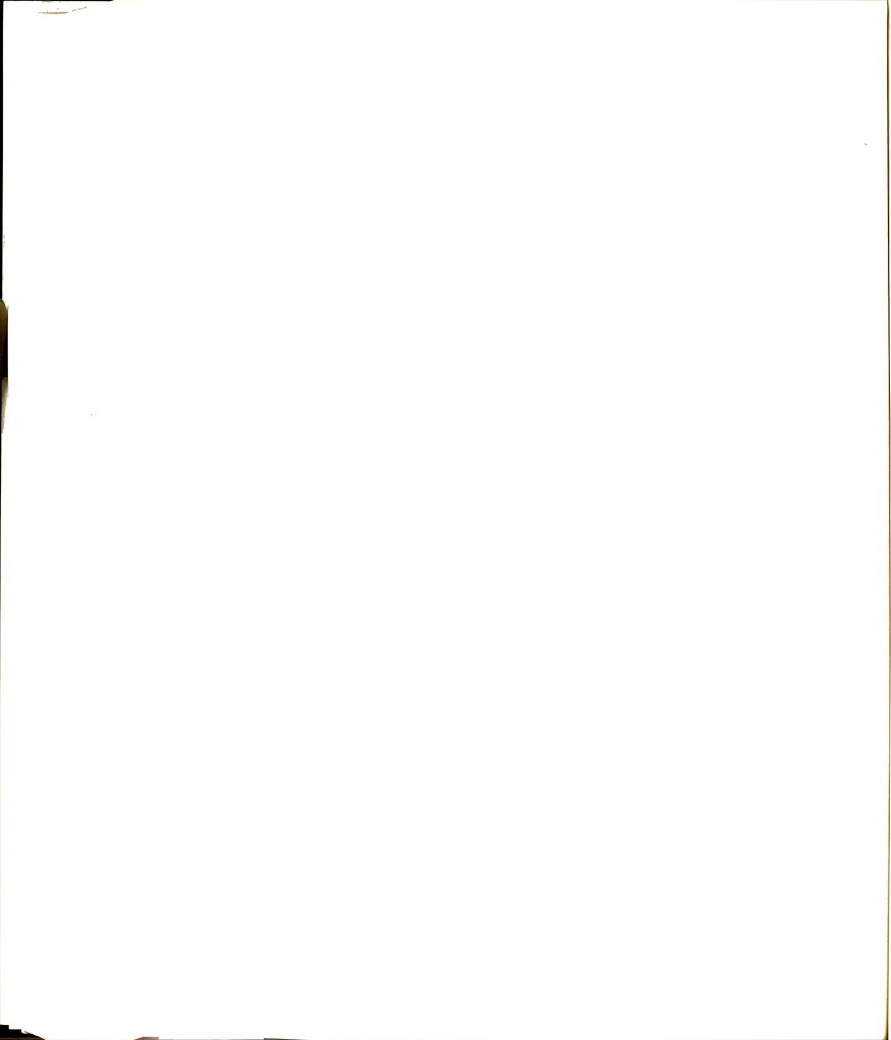
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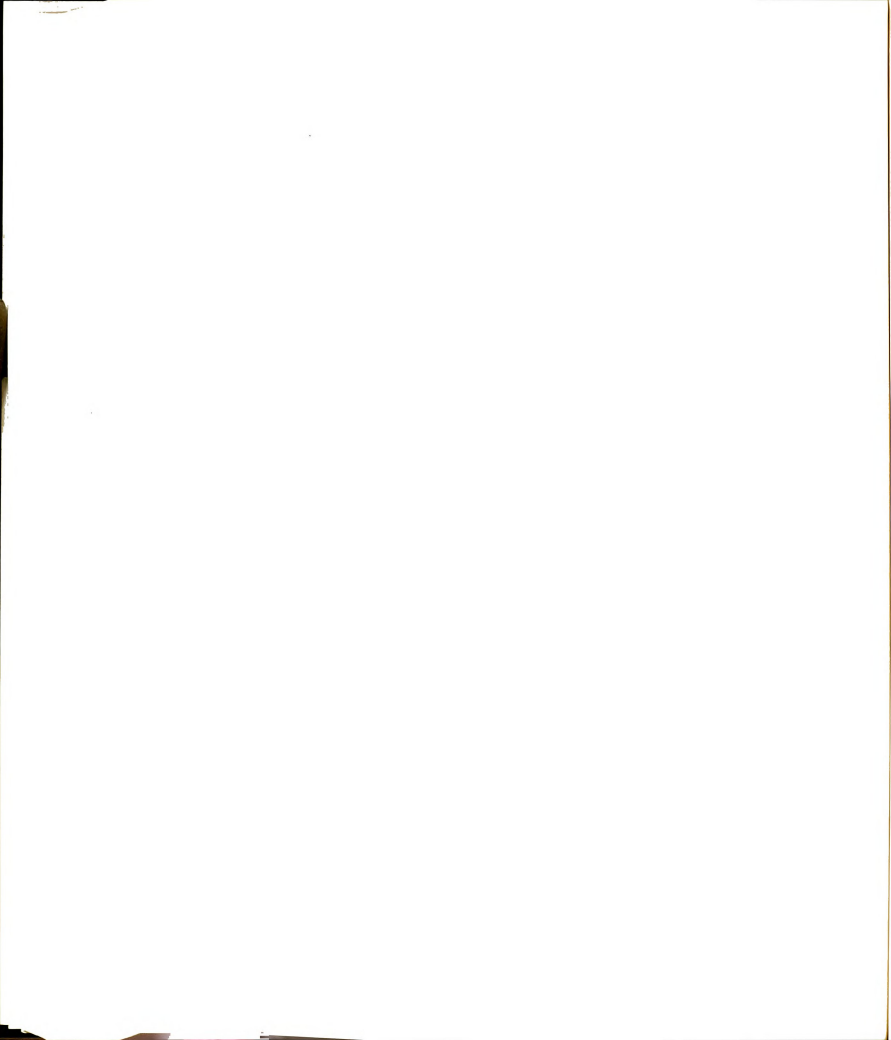
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F(1)=AREA*D(2)*Q(1)
F(2)=AREA*D(1)*Q(1)
F(3)=CJ*AREA*D(1)*D(3)*Q(2)
F(4)=CJ*AREA*D(2)*D(3)*Q(2)
F(5)=AREA*D(3)*Q(1)
F(6)=CJ*AREA*D(2)*D(1)*Q(2)
F(7)=CJ*AREA*(Q(1)+D(2)*D(2)*Q(2)+Q(4))
F(8)=AREA*D(2)*Q(3)
F(9)=AREA*D(1)*Q(3)
F(10)=AREA*D(3)*Q(3)
RETURN
END
C*****SUBROUTINE FCNYX*****
SUBROUTINE FCNYX(AREA,D,Q,F)
REAL D(3),MUO
COMPLEX Q(4),F(10),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(1)*Q(1)
F(2)=CJ*AREA*D(2)*D(3)*Q(2)
F(3)=CJ*AREA*(Q(1)+D(3)*D(3)*Q(2)+Q(4))
F(4)=AREA*D(3)*Q(1)
F(5)=AREA*D(2)*Q(1)
F(6)=CJ*AREA*D(1)*D(2)*Q(2)
F(7)=CJ*AREA*D(1)*D(3)*Q(2)
F(8)=AREA*D(1)*Q(3)
F(9)=AREA*D(3)*Q(3)
F(10)=AREA*D(2)*Q(3)
RETURN
END
C*****SUBROUTINE FCNYY*****
SUBROUTINE FCNYY(AREA,D,Q,F)
REAL D(3),MUO
COMPLEX Q(4),F(5),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(2)*Q(1)
F(2)=CJ*AREA*D(1)*D(3)*Q(2)
F(3)=CJ*AREA*(Q(1)+D(3)*D(3)*Q(2)+Q(4))
F(4)=CJ*AREA*(Q(1)+D(1)*D(1)*Q(2)+Q(4))
F(5)=AREA*D(2)*Q(3)
RETURN
END
C*****SUBROUTINE FCNYZ*****
SUBROUTINE FCNYZ(AREA,D,Q,F)
REAL D(3),MUO
COMPLEX Q(4),F(10),CJ
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ

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```
C*****
SUBROUTINE FCNZX(AREA,D,G,F)
REAL D(3),MUO
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(2)*Q(1)
F(2)=AREA*D(1)*Q(1)
F(3)=CJ*AREA*D(2)*D(3)*Q(2)
F(4)=CJ*AREA*D(2)*D(3)*Q(2)
F(5)=AREA*D(3)*Q(1)+D(1)*D(1)*Q(2)+Q(4))
F(6)=CJ*AREA*(Q(1)+D(1))*D(1)*Q(2)
F(7)=CJ*AREA*D(1)*D(2)*Q(2)
F(8)=AREA*D(2)*Q(3)
F(9)=AREA*D(1)*Q(3)
F(10)=AREA*D(3)*Q(3)
RETURN
END
C*****
SUBROUTINE FCNZY(AREA,D,G,F)
REAL D(3),MUO
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(1)*Q(1)
F(2)=CJ*AREA*(Q(1)+D(2)*D(2)*Q(2)+Q(4))
F(3)=CJ*AREA*D(2)*D(3)*Q(2)
F(4)=AREA*D(3)*Q(1)
F(5)=AREA*D(3)*Q(1)
F(6)=CJ*AREA*D(1)*D(2)*Q(2)
F(7)=CJ*AREA*D(1)*D(3)*Q(2)
F(8)=AREA*D(1)*Q(3)
F(9)=AREA*D(2)*Q(3)
F(10)=AREA*D(3)*Q(3)
RETURN
END
C*****
SUBROUTINE FCNZZ(AREA,D,G,F)
REAL D(3),MUO
COMMON /CONS/PI,MUO,EPO
COMMON /CPLX/CJ
F(1)=AREA*D(3)*Q(1)
F(2)=AREA*D(1)*Q(1)
F(3)=CJ*AREA*D(1)*D(2)*Q(2)
F(4)=CJ*AREA*D(2)*D(3)*Q(2)
F(5)=AREA*D(2)*Q(1)
F(6)=CJ*AREA*(Q(1)+D(1))*D(1)*Q(2)+Q(4))
F(7)=CJ*AREA*D(1)*D(3)*Q(2)
F(8)=AREA*D(3)*Q(3)
F(9)=AREA*D(1)*Q(3)
F(10)=AREA*D(2)*Q(3)
RETURN
END
C*****
```




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C**** SUBROUTINE FCNZZ*****
SUBROUTINE FCNZZ (AREA, D, Q, F)
REAL D(3), MUO
COMPLEX G(4), F(3), CJ
COMMON /CONS/PI, MUO, EPO
COMMON /CPLX/CJ
F(1)=AREA*D(3)*G(1)
F(2)=CJ*AREA*D(1)*D(2)*G(2)
F(3)=CJ*AREA*(G(1)+D(2)*D(2)*G(2)+G(4))
F(4)=CJ*AREA*(G(1)+D(1)*D(1)*G(2)+G(4))
RETURN
END
C**** SUBROUTINE CEFXX*****
SUBROUTINE CEFXX(F, C)
COMPLEX F(5), C(16)
C(1)=F(1)
C(2)=0.0
C(3)=-1.0*F(2)
C(4)=F(4)
C(5)=0.0
C(6)=F(1)
C(7)=-1.0*F(3)
C(8)=F(2)
C(9)=F(2)
C(10)=-1.0*F(4)
C(11)=F(5)
C(12)=0.0
C(13)=F(3)
C(14)=1.0*F(2)
C(15)=F(5)
C(16)=F(5)
RETURN
END
C**** SUBROUTINE CEFXY*****
SUBROUTINE CEFXY(F, C)
COMPLEX F(10), C(16)
C(1)=1.0*F(1)
C(2)=-1.0*F(4)
C(3)=-1.0*F(2)
C(4)=F(6)
C(5)=0.0
C(6)=F(5)
C(7)=-1.0*F(3)
C(8)=F(7)
C(9)=F(2)
C(10)=-1.0*F(6)
C(11)=-1.0*F(8)

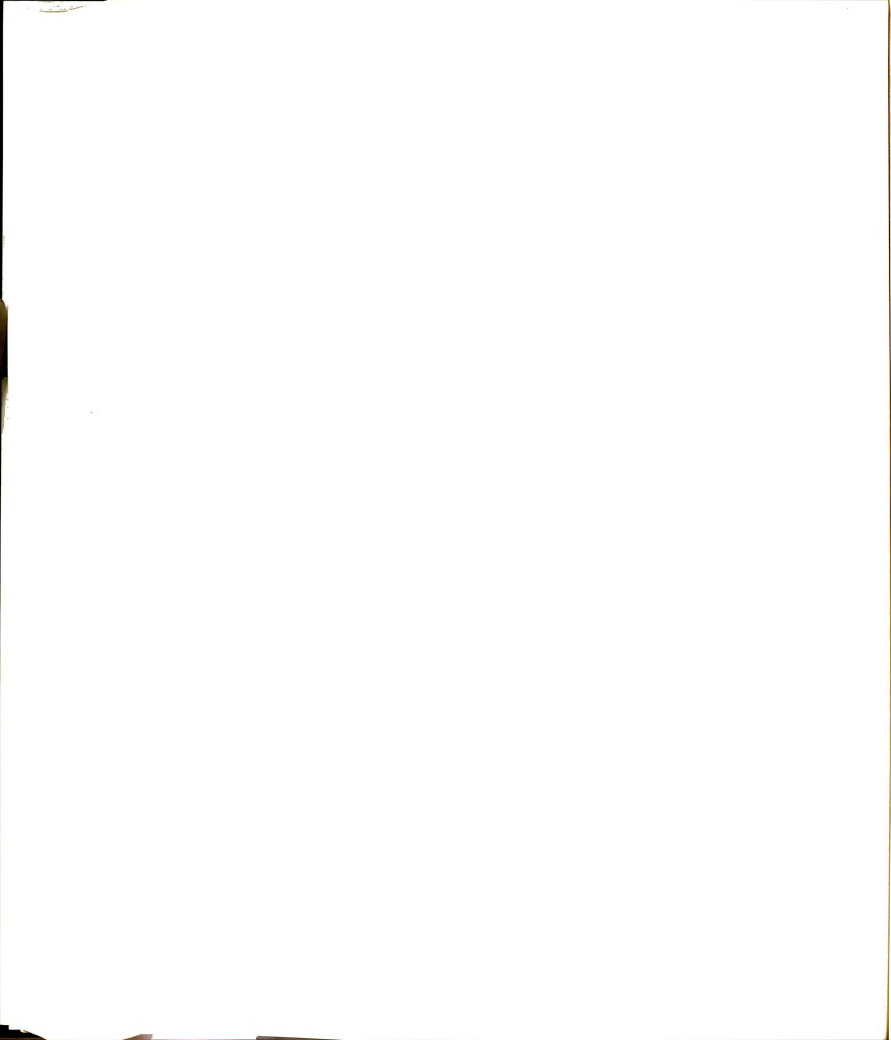
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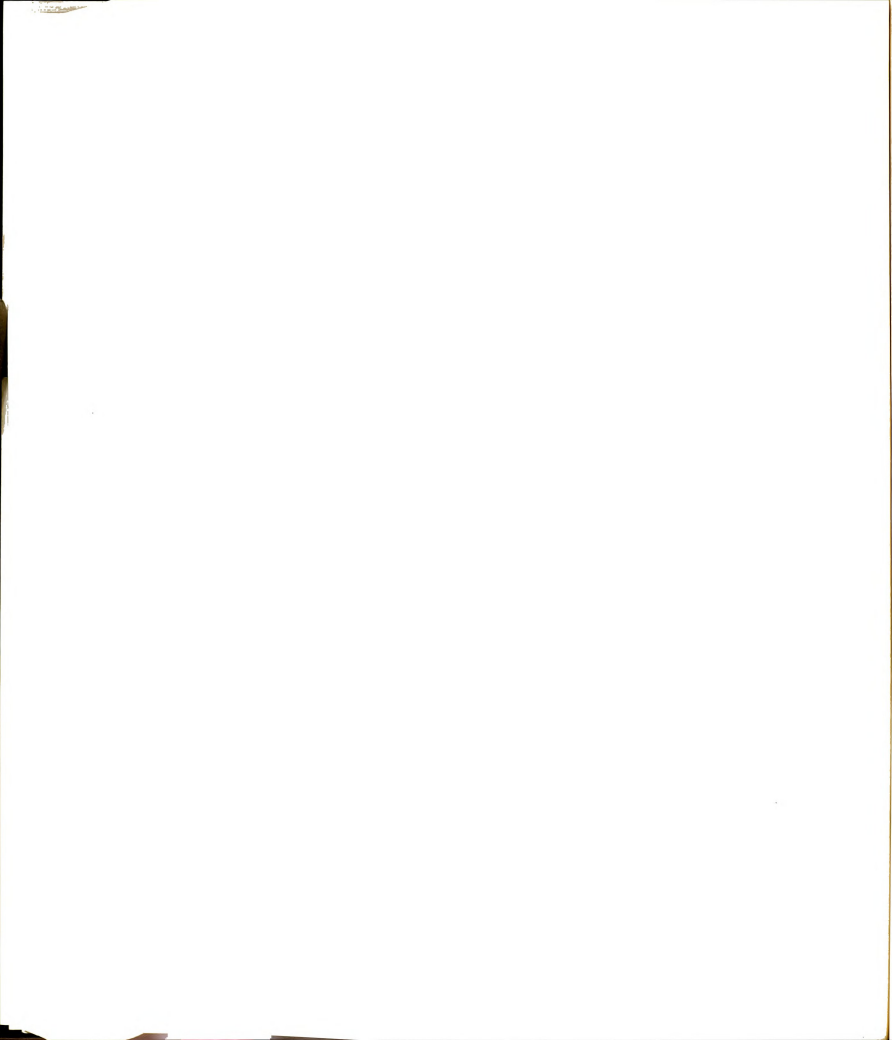
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[illegible]

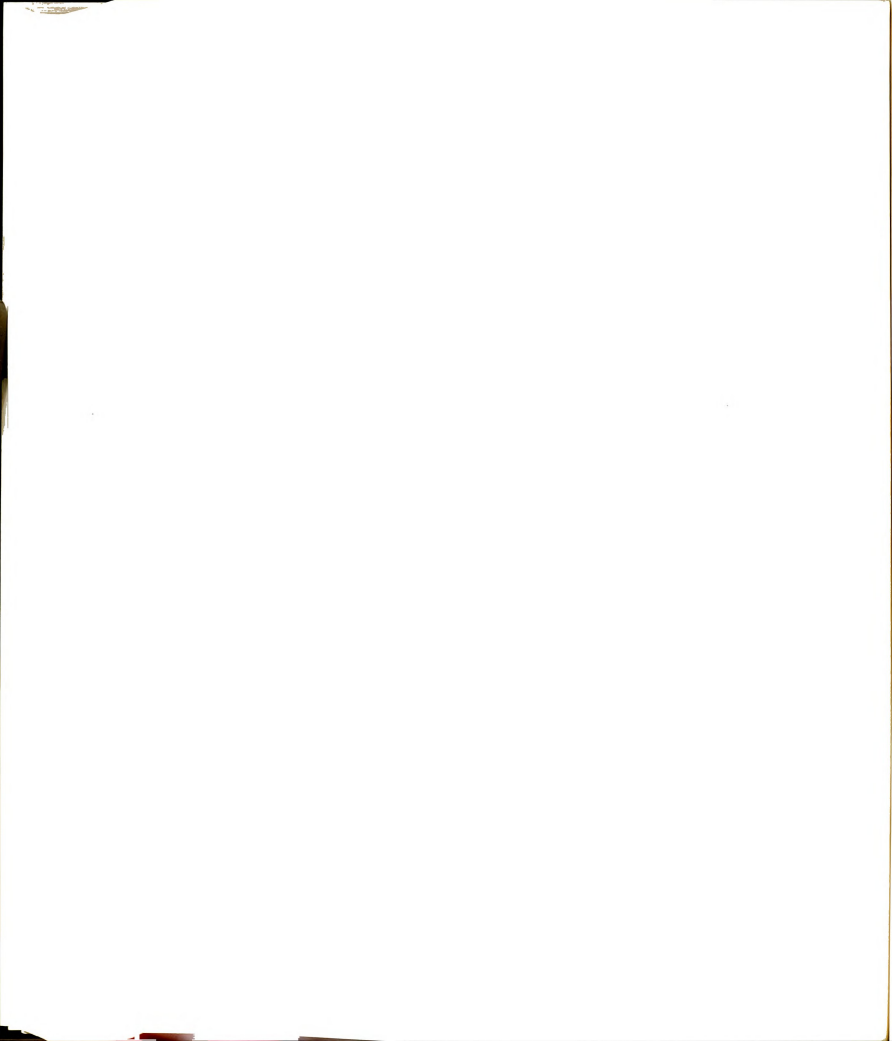


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COMPLEX F(5), C(16)
C(1)=F(1)
C(2)=0.0
C(3)=F(2)
C(4)=-1.0*F(4)
C(5)=0.0
C(6)=F(1)
C(7)=F(3)
C(8)=-1.0*F(2)
C(9)=-1.0*F(2)
C(10)=F(4)
C(11)=F(5)
C(12)=0.0
C(13)=-1.0*F(3)
C(14)=F(2)
C(15)=0.0
C(16)=F(5)
RETURN
END
C*****SUBROUTINE CEFYZ(F,C)
SUBROUTINE CEFYZ(F,C)
COMPLEX F(10), C(16)
C(1)=F(1)
C(2)=0.0
C(3)=F(3)
C(4)=-1.0*F(6)
C(5)=1.0*F(2)
C(6)=-1.0*F(5)
C(7)=F(4)
C(8)=-1.0*F(7)
C(9)=-1.0*F(3)
C(10)=F(6)
C(11)=F(8)
C(12)=0.0
C(13)=-1.0*F(4)
C(14)=F(7)
C(15)=-1.0*F(9)
C(16)=-1.0*F(10)
RETURN
END
C*****SUBROUTINE CEFZX(F,C)
SUBROUTINE CEFZX(F,C)
COMPLEX F(10), C(16)
C(1)=0.0
C(2)=F(4)
C(3)=-1.0*F(2)
C(4)=F(6)
C(5)=-1.0*F(1)

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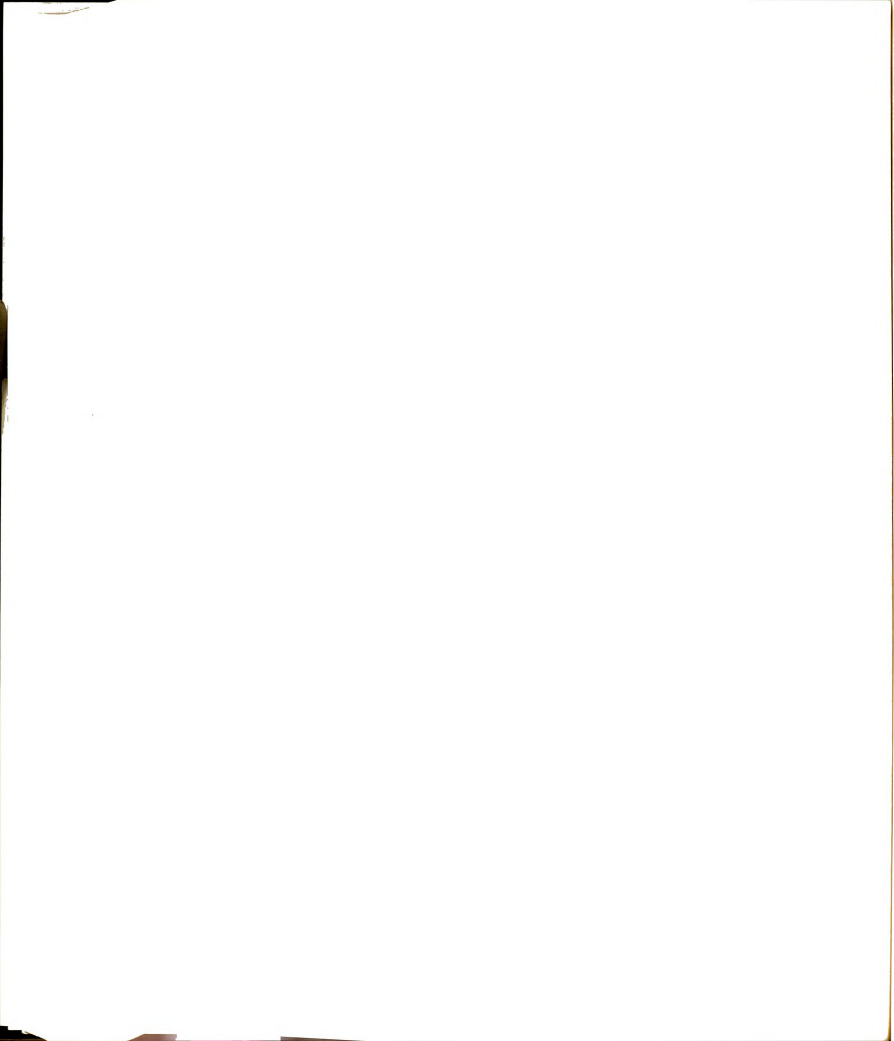
C(6)=-1.0*F(5)
C(7)=-1.0*F(3)
C(8)=F(7)
C(9)=F(2)
C(10)=-1.0*F(6)
C(11)=0.0
C(12)=F(9)
C(13)=F(3)
C(14)=-1.0*F(7)
C(15)=-1.0*F(8)
C(16)=-1.0*F(10)
RETURN
END
C*****
SUBROUTINE CEFZY*****
SUBROUTINE CEFZY(F,C)
COMPLEX F(10),C(16)
C(1)=F(1)
C(2)=0.0
C(3)=-1.0*F(3)
C(4)=F(6)
C(5)=-1.0*F(2)
C(6)=-1.0*F(5)
C(7)=-1.0*F(4)
C(8)=F(7)
C(9)=F(3)
C(10)=-1.0*F(6)
C(11)=F(8)
C(12)=0.0
C(13)=F(4)
C(14)=-1.0*F(7)
C(15)=-1.0*F(9)
C(16)=-1.0*F(10)
RETURN
END
C*****
SUBROUTINE CEFZZ*****
SUBROUTINE CEFZZ(F,C)
COMPLEX F(5),C(16)
C(1)=F(1)
C(2)=0.0
C(3)=-1.0*F(2)
C(4)=F(4)
C(5)=0.0
C(6)=F(1)
C(7)=-1.0*F(3)
C(8)=F(2)
C(9)=F(2)
C(10)=-1.0*F(4)
C(11)=F(5)

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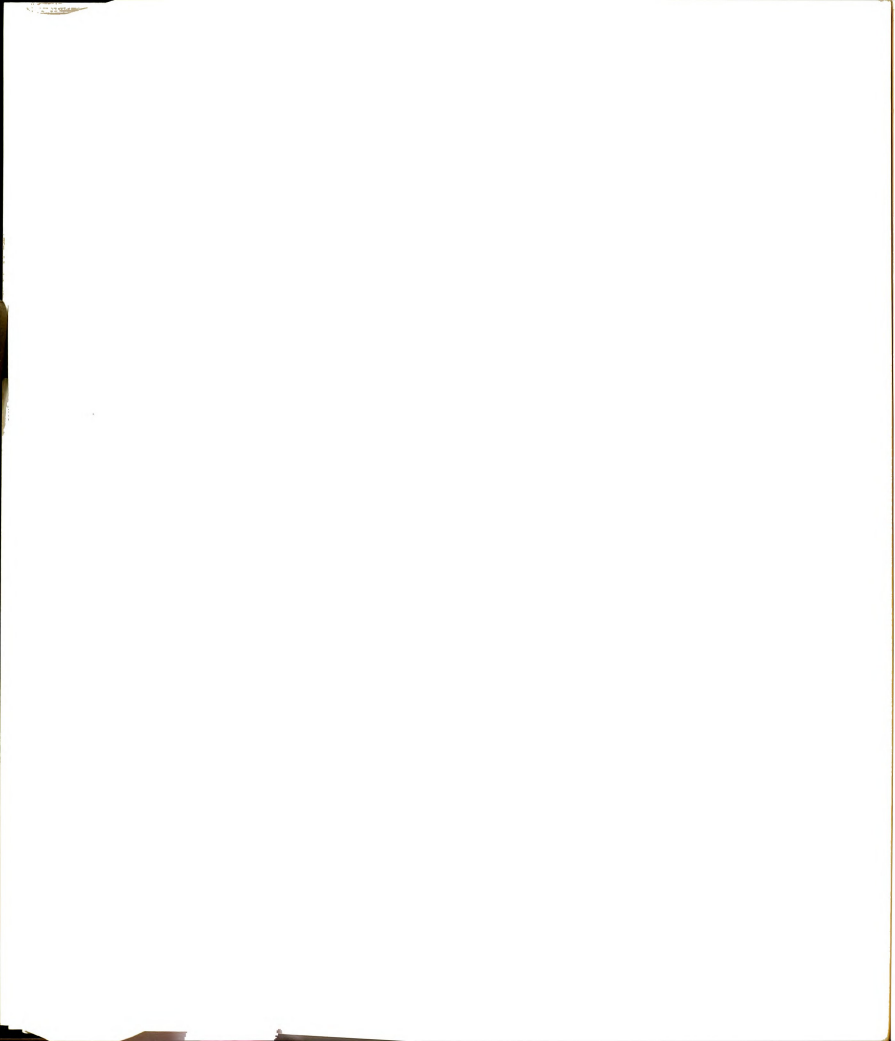
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15170 C(12)=0.0
15180 C(13)=F(3)
15190 C(14)=-1.0*F(2)
15200 C(15)=0.0
15210 C(16)=F(5)
15220 RETURN
15230 END
*****
C*****SUBROUTINE CMATP(A,N,NP,N,DET,EP)
15240 SUBROUTINE CMATP(A,N,NP,N,DET,EP)
15250 COMPLEX A,B,DET,CONST,S
15260 DIMENSION A(N,NP)
15270 DET=CMPLX(1.0,0.0)
15280 NP1=N+1
15290 NPM=N+M
15300 NM1=N-1
15310 DO 4 J=1,NM1
15320 C=CABS(A(J,J))
15330 JP1=J+1
15340 DO 5 I=JP1,N
15350 D=CABS(A(I,J))
15360 IF (C-D) 6,5,5
15370 6 DET=-DET
15380 DO 7 K=J,NPM
15390 B=A(I,K)
15400 A(I,K)=A(J,K)
15410 A(J,K)=B
15420 7 CONTINUE
15430 5 CONTINUE
15440 15 DO 4 J=JP1,N/A(J,J)
15450 CONST=A(I,J)/A(J,J)
15460 DO 4 K=JP1,NPM
15470 4 A(I,K)=A(I,K)-CONST*A(J,K)
15480 IF (CABS(A(N,N))-EP) 14,18,18
15490 14 DET=CMPLX(0.0,0.0)
15500 PRINT 30
15510 RETURN
15520 18 CONTINUE
15530 DO 12 I=1,N
15540 K=N-I+1
15550 KP1=K+1
15560 DO 12 L=NP1,NPM
15570 S=CMPLX(0.0,0.0)
15580 IF (N-KP1) 12,19,19
15590 12 DO 13 J=KP1,N
15600 13 S=S+A(K,J)*A(J,L)
15610 12 A(K,L)=(A(K,L)-S)/A(K,K)
15620 30 FORMAT(1X,42HTHE DETERMINANT OF THE SYSTEM EQUALS ZERO. /
15630 11X,36HTHE PROGRAM CANNOT HANDLE THIS CASE. //)
15640 RETURN
15650 END

```



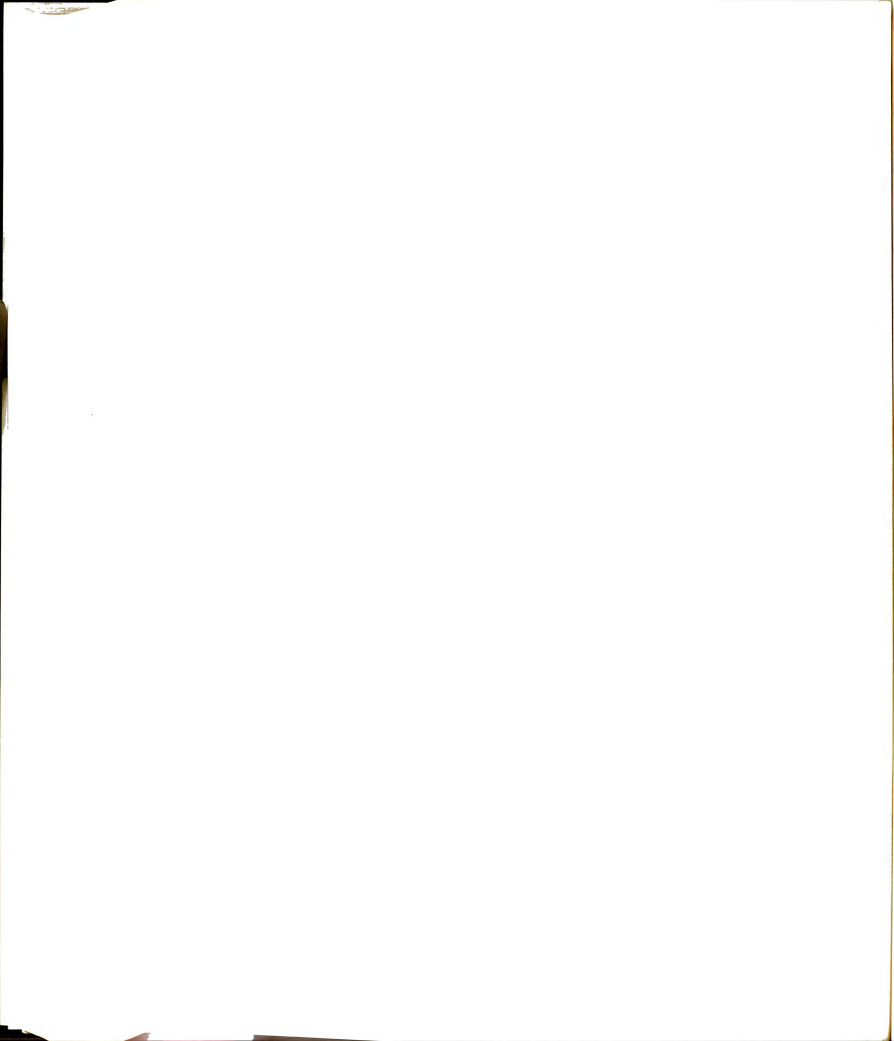
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C*****PROGRAM COMBINE IS TO COMBINE RESULTS FROM COSH, SINE, SINH MODES*****
C*****AND TO OBTAIN THE TOTAL FIELDS ON THE SURFACES OF FIRST AND FIFTH QUADRANT*****
C*****
PROGRAM COMBINE INPUT, OUTPUT, TAPE1, TAPE2, TAPE3, TAPE4)
COMPLEX SUM1(208), SUM5(208), ECOS(208), HCSIN(208)
COMMON /XYZ/NX, NY, NZ, NNX, NNY, NNZ
DIMENSION X(52), Y(52), Z(52), DCELL(52)
DIMENSION ABS1(208), ABS5(208), ANG1(208), ANG5(208)
DIMENSION NV(6)
DATA FIRST, 'FIRS', CMPL, 'CMPL', COSE, 'COSE', COSH, 'COSH', SINE, 'SIN', SIN150
+E, 'SINH', 'SINH', EXPZ, 'EXPZ', FIFT, 'FIFT', ABSL, 'ABSL' /
CALL NOBLANK
READ 500, NDIV
FORMAT(12)
500
5
FORMAT(14)
PRINT 10, 10
PRINT 11
PRINT 12
PRINT 13
READ 1, NX, NY, NZ, NNX, NNY, NNZ
DCELL=NX+NY+NZ+NNX+NNY+NNZ
DO 14 I=1, NCELL
READ 2, X(I), Y(I), Z(I), DCELL(I)
2
CONTINUE
READ 4, RLEP, SIG, FMEG
DO 14 I=1, NCELL
PRINT 15, I, X(I), Y(I), Z(I), DCELL(I)
14
PRINT 16, NX, NY, NZ, NNX, NNY, NNZ
PRINT 101, NDIV, NDIV
FORMAT(1H-5X, 33#EACH SURFACE CELL IS DIVIDED INTO,
21X, I2, 1X, 2#BY, 1X, I2, 1X, 17#SUB-SURFACE CELLS)
PRINT 17, RLEP, SIG, FMEG
MAT=4*NCELL
DO 18 J=1, MAT
SUM1(J)=CMPLX(0, 0)
SUM5(J)=CMPLX(0, 0)
18
CONTINUE
BUFFER IN (1, 1) (ECOS(1), ECOS(MAT))
CALL UNIT(1)
BUFFER IN (2, 1) (HCSIN(1), HCSIN(MAT))
CALL UNIT(2)
BUFFER IN (3, 1) (ESIN(1), ESIN(MAT))
CALL UNIT(3)
BUFFER IN (4, 1) (HSIN(1), HSIN(MAT))
CALL UNIT(4)

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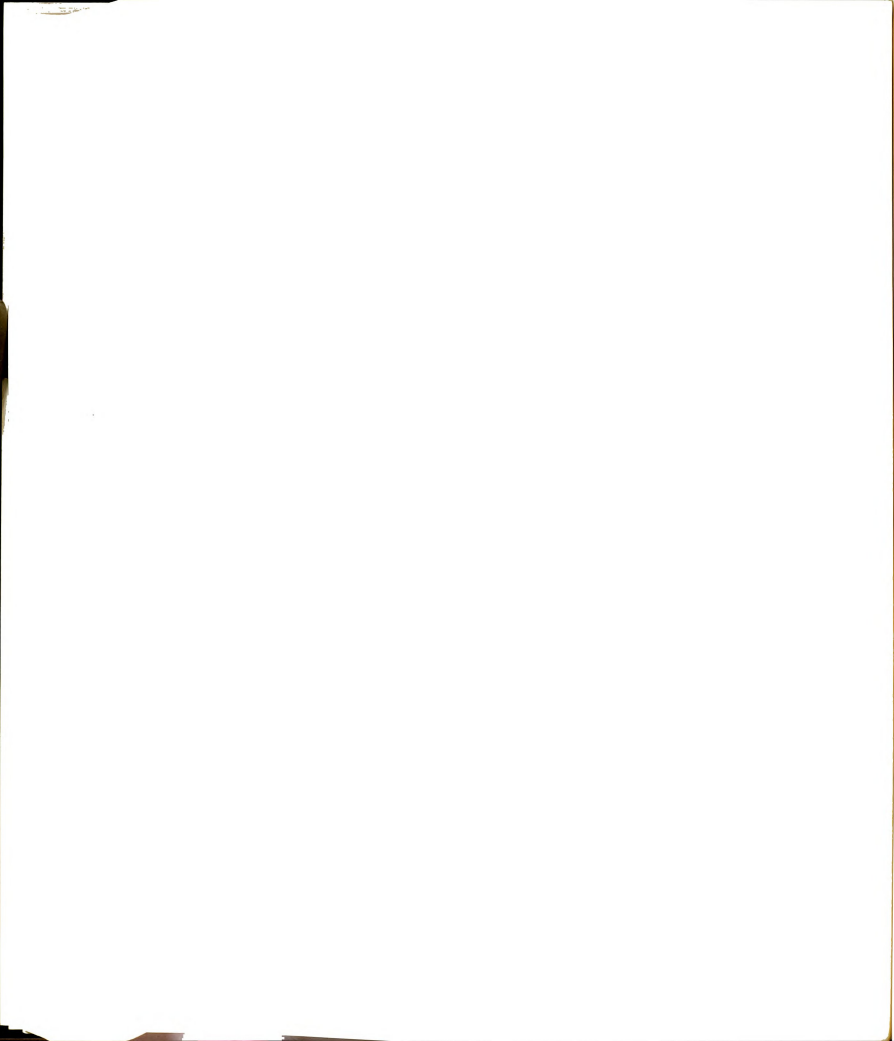


```

CALL PREL(EIRS, CMPL, COSE, ECOS, MAT)
CALL PREL(EIRS, CMPL, COSH, HCOS, MAT)
CALL PREL(EIRS, CMPL, SINE, ESIN, MAT)
CALL PREL(EIRS, CMPL, SINH, HSIN, MAT)
DO 19 I=1, MAT
  SUM1(I)=ECOS(I)+HCOS(I)+ESIN(I)+HSIN(I)
  CONTINUE
19 IF (IG.EQ.1234) GO TO 34
  NV(1)=NV
  NV(2)=NV
  NV(3)=NZ
  NV(4)=NNX
  NV(5)=NNY
  NV(6)=NNZ
  L=1
  M=2
  N=3
  J=1
  K=NV(L)
35 CONTINUE
  DO 20 I=J, K
    SUM5(I)=ECOS(I)-HCOS(I)-ESIN(I)+HSIN(I)
    CONTINUE
    J=J+NV(L)
    K=K+NV(L)
  DO 21 I=J, K
    SUM5(I)=-ECOS(I)+HCOS(I)+ESIN(I)-HSIN(I)
    CONTINUE
    J=J+NV(L)
    K=K+NV(L)
  DO 22 I=J, K
    SUM5(I)=-ECOS(I)+HCOS(I)+ESIN(I)-HSIN(I)
    CONTINUE
    J=J+NV(L)
    K=K+NV(L)
  DO 23 I=J, K
    SUM5(I)=ECOS(I)-HCOS(I)-ESIN(I)+HSIN(I)
    CONTINUE
    J=J+NV(L)
    K=K+NV(L)
  IF (NV(M).EQ.0) GO TO 37
  DO 24 I=J, K
    SUM5(I)=ECOS(I)-HCOS(I)-ESIN(I)+HSIN(I)
    CONTINUE
    J=J+NV(M)
    K=K+NV(M)
  24

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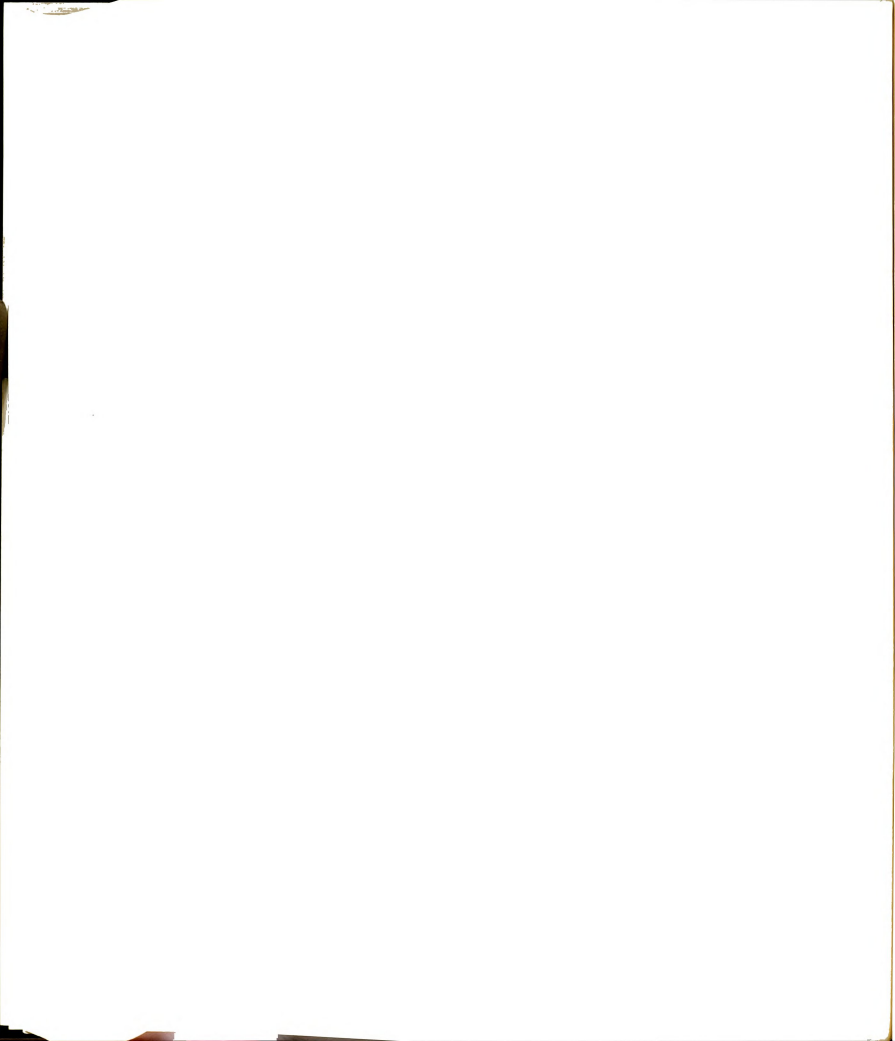


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DO 25 I=J,K
SUM5(I)=-ECOS(I)+HCOS(I)+ESIN(I)-HSIN(I)
25 CONTINUE
J=J+NV(M)
K=K+NV(M)
DO 26 I=J,K
SUM5(I)=-ECOS(I)+HCOS(I)+ESIN(I)-HSIN(I)
26 CONTINUE
J=J+NV(M)
K=K+NV(M)
DO 27 I=J,K
SUM5(I)=-ECOS(I)-HCOS(I)-ESIN(I)+HSIN(I)
27 CONTINUE
37 CONTINUE
J=J+NV(M)
K=K+NV(N)
IF(NV(N).EQ.0) GO TO 38
DO 28 I=J,K
SUM5(I)=-ECOS(I)+HCOS(I)+ESIN(I)-HSIN(I)
28 CONTINUE
J=J+NV(N)
K=K+NV(N)
DO 29 I=J,K
SUM5(I)=-ECOS(I)+HCOS(I)+ESIN(I)-HSIN(I)
29 CONTINUE
J=J+NV(N)
K=K+NV(N)
DO 30 I=J,K
SUM5(I)=-ECOS(I)-HCOS(I)-ESIN(I)+HSIN(I)
30 CONTINUE
J=J+NV(N)
K=K+NV(N)
DO 31 I=J,K
SUM5(I)=-ECOS(I)-HCOS(I)-ESIN(I)+HSIN(I)
31 CONTINUE
38 IF(L.EQ.4) GO TO 34
L=4
M=5
N=6
J=J+NV(3)
K=K+NV(L)
IF(NV(L).EQ.0) GO TO 36
GO TO 35
34 CALL PREL(FIRS,CMP,L,EXPZ,SUM1,MAT)
PRINT 32,FIRS
CALL PREL(FIFT,CMP,L,EXPZ,SUM5,MAT)
PRINT 32,FIFT

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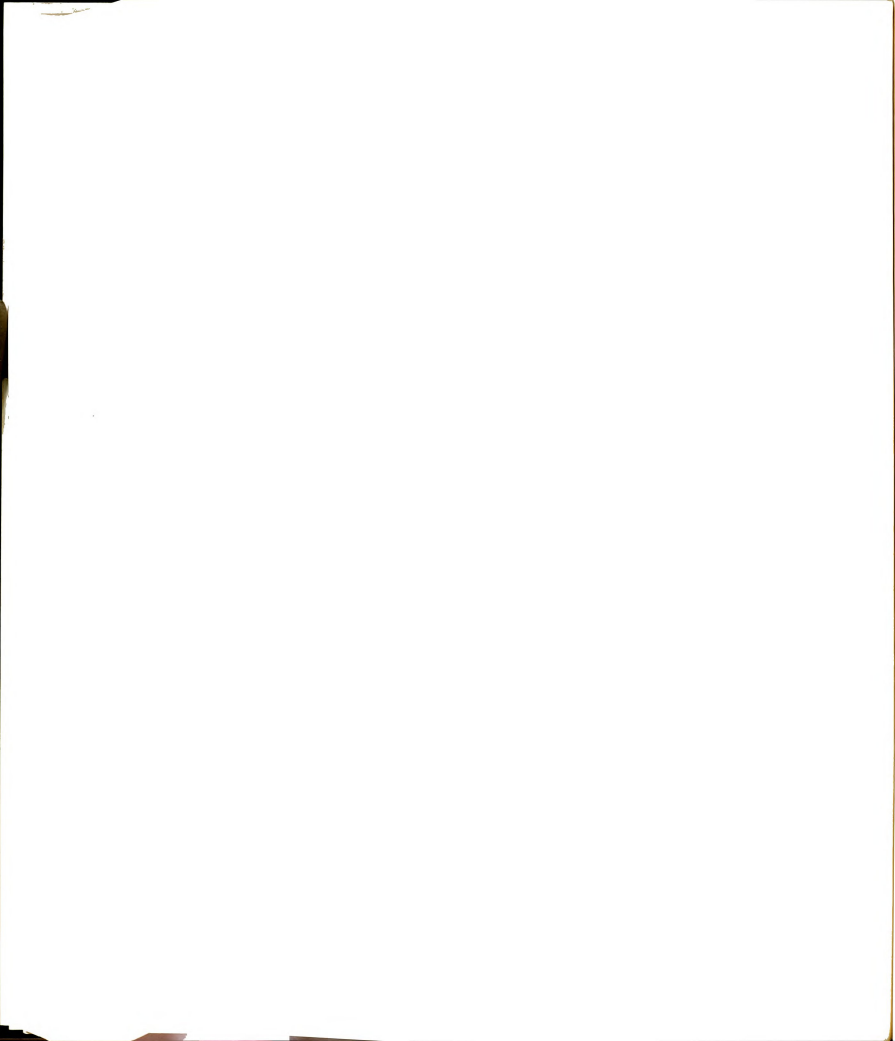
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DO 33 I=1,MAT
  ABS1(I)=CABS(SUM1(I))
  ABS5(I)=CABS(SUM5(I))
  CALL ANGLE(SUM1(I),ANG1(I))
  CALL ANGLE(SUM5(I),ANG5(I))
  SUM1(I)=CMPLX(ABS1(I),ANG1(I))
  SUM5(I)=CMPLX(ABS5(I),ANG5(I))
33 CONTINUE
CALL PREF(FIRS,ABSL,EXPZ,SUM1,MAT)
PRINT 32,FIRS
CALL PREF(FIFT,ABSL,EXPZ,SUM5,MAT)
PRINT 32,FIFT
10 FORMAT("1, 10X,"THIS IS THE TOTAL INDUCED FIELDS ON THE SURFACES
+ OF AN",1X,14,1X,"G'S SYMMETRICAL BODY")
11 FORMAT("-","THE RESULTS ARE OBTAINED BY COMBINING RESULTS DUE
+ TO FOUR DIFFERENT TYPES OF IMPRESSED FIELDS WHICH ARE COSE, COSH,
+SINE, AND SINH")
12 FORMAT("-","10X,"THE LOCATION AND DIMENSION OF EACH SURFACE-CELL
+ ARE LISTED BELOW")
13 FORMAT("-","10X," N "7X,"X (IN CM)",10X,"Y (IN CM)",10X,
+"Z (IN CM)",10X,"D (IN CM)")
1 FORMAT(6I3)
3 FORMAT(4F10.5)
4 FORMAT(3E13.6)
15 FORMAT("-","5X,13,7X,F10.5,10X,F10.5,10X,F10.5,10X,F10.5)
16 FORMAT("-","5X,"THERE ARE NX=",13,2X,"NV=",13,2X,"NZ=",13,2X,
+"NNZ=",13,2X,"NNV=",13,2X,"AND NNZ=",13,2X,
+"CELLS ON SURFACES OF THE FIRST QUADRANT")
17 FORMAT("-","5X,"DIELECTRIC CONST.=",E13.6/"-",5X,"CONDUCTIVITY=",
+E13.6,"MHG/METER"/"-",5X,"FREQUENCY=",E13.6,"NEG HZ"/"-",5X,"THE
+ TOTAL INCIDENT FIELD IS A PLANE WAVE.")
32 FORMAT("-","10X,"ABOVE ARE TOTAL FIELDS ON THE ",A4," QUADRANT")
END
***** SUBROUTINE PREF *****
SUBROUTINE PREF(GUAD,COAB,MODE,INFL,MAT)
REAL MODE,INFL,MAT
COMMON AX,YZ,NX,NY,NZ,NNX,NNY,NNZ
DATA TOP,"TOP",RIGHT,"RIGHT",REAR,"REAR",CMPL/"CMPL"/,ABSL/"ABSL"/
+,"FIRS","FIRS",FIFT/"FIFT"/,FRNT/"FRNT"/
DATA BTM,"BTM",LEFT/"LEFT"/
PRINT 111,MODE,QUAD
IF (COAB.EQ.ABSL) PRINT 112
IF (COAB.EQ.CMPL) PRINT 113
PRINT 123,TOP
PRINT 137,
PRINT 122

```

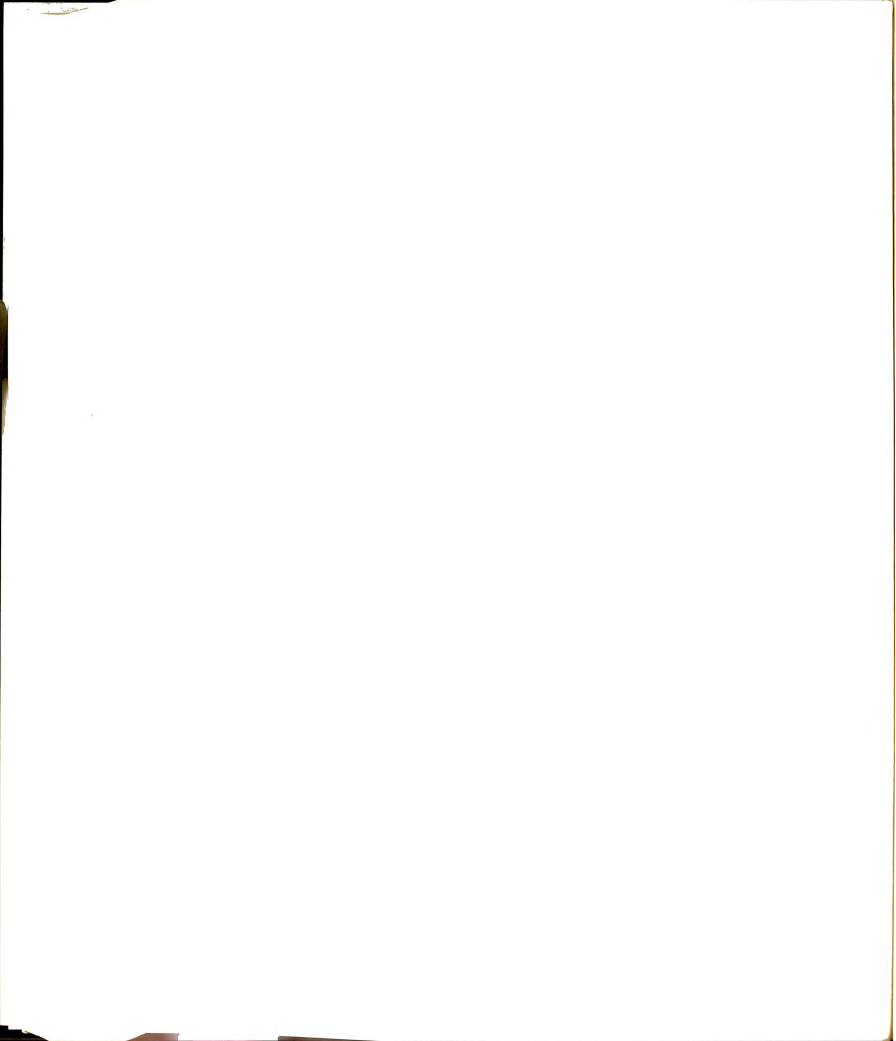


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DO 12 I=1, NX
  I=1
  Y=1
  J=J+NX
  IPP=IPP+NX
12 PRINT 13, I, INFL(J), INFL(IP), INFL(IPP), INFL(IPPP)
  PRINT 123, RIGHT
  PRINT 123
  PRINT 122
  DO 14 I=1, NY
    I=1+NX
    J=J+4*NX
    IPP=IPP+NY
    IP=IP+NY
    IPPP=IP+NY
14 PRINT 13, I, INFL(J), INFL(IP), INFL(IPP), INFL(IPPP)
  PRINT 122
  IF(QUAD.EQ.FIRST) PRINT 123, REAR
  IF(QUAD.EQ.FIFT) PRINT 123, FRNT
  PRINT 143
  PRINT 122
  DO 15 I=1, NZ
    I=1+NX+NY
    J=J+4*(NX+NY)
    IPP=J+NZ
    IPPP=IP+NZ
15 PRINT 13, I, INFL(J), INFL(IP), INFL(IPP), INFL(IPPP)
  PRINT 122
  IF(NNX.EQ.0) GO TO 19
  PRINT 123, BTM
  PRINT 137
  PRINT 122
  DO 16 I=1, NNZ
    I=1+NX+NY+NZ
    J=J+4*(NX+NY+NZ)
    IPP=J+NNX
    IPPP=IP+NNX
16 PRINT 13, I, INFL(J), INFL(IP), INFL(IPP), INFL(IPPP)
  PRINT 122
19 CONTINUE
  IF(NNY.EQ.0) GO TO 20
  PRINT 123, LEFT
  PRINT 140
  PRINT 122

```

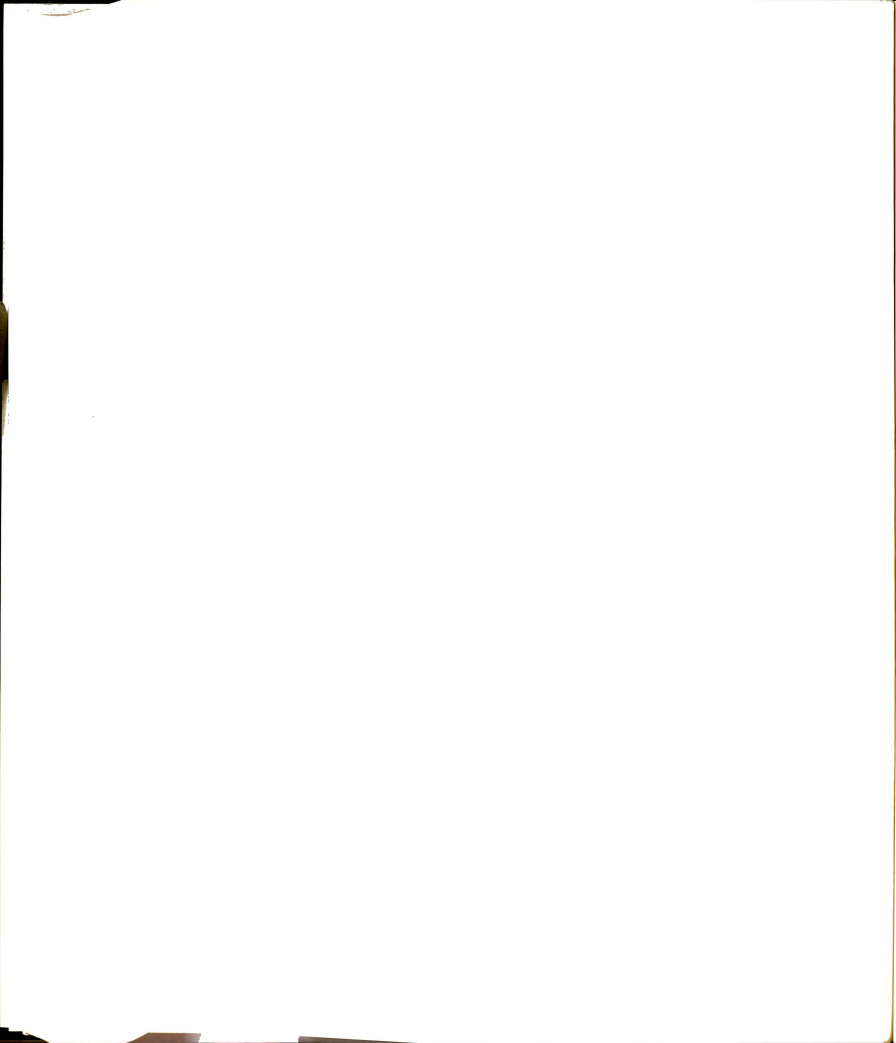
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DO 17 I=1,NNY
  J1=1+NX+NY+NZ+NNX
  J=1+4*(NX+NY+NZ+NNX)
  IP=J+NNY
  IPP=IP+NNY
  IPPP=IPP+NNY
  17 PRINT 13,11,INFL(J),INFL(IP),INFL(IPP),INFL(IPPP)
  20 CONTINUE
  IF (NNZ.EQ.0) GO TO 21
  IF (QUAD.EG.FIRST) PRINT 123,FRNT
  IF (QUAD.EG.FIFT) PRINT 123,REAR
  PRINT 143
  PRINT 122
  DO 18 I=1,NNZ
    J1=1+NX+NY+NZ+NNX+NNY
    J=1+4*(NX+NY+NZ+NNX+NNY)
    IP=J+NNZ
    IPP=IP+NNZ
    IPPP=IPP+NNZ
    18 PRINT 13,11,INFL(J),INFL(IP),INFL(IPP),INFL(IPPP)
  CONTINUE
  21 FORMAT(1,"THE FIELDS INDUCED BY "A4," TYPE OF IMPRESSED FIELD,"
    +1X,"IN "A4," QUADRANT ARE LISTED BELOW (AMP/METER FOR H FIELDS),"
    +1X,"VOL/METER FOR E FIELDS)")
  112 FORMAT(1,"N "6X,"ABSL",11X,"PHAS",11X,"ABSL",11X,"PHAS",
    +1X,"ABSL",11X,"PHAS",11X,"ABSL",11X,"PHAS")
  113 FORMAT(1,"N "6X,"REAL",11X,"IMAG",11X,"REAL",11X,"IMAG",
    +1X,"REAL",11X,"IMAG",11X,"REAL",11X,"IMAG")
  123 FORMAT(1,"N "15X,"ON THE "A4," SURFACE WE HAVE THE FOLLOWS")
  157 FORMAT(1,"N "14X," HZ "20X," HY "20X," HX "20X,"
    +18X," EY ")
  122 FORMAT(1,"N "5X,4(28H"*****2X"))
  13 FORMAT(1,"N "13,1X,4(1X,"(",2E13.6,"),1X))
  140 FORMAT(1,"N "14X," HZ "20X," HX "20X," EY "20X,"
    +18X," EX ")
  143 FORMAT(1,"N "14X," HZ "20X," HX "20X," EY "20X,"
    +18X," EX ")
  RETURN
END
C*****SUBROUTINE ANGLE*****
SUBROUTINE ANGLE(G,EF)
  COMPLEX G
  F=57.2957795131
  IF (AIMAG(G)) 10,20,30
  10 IF (REAL(G)) 40,50,40

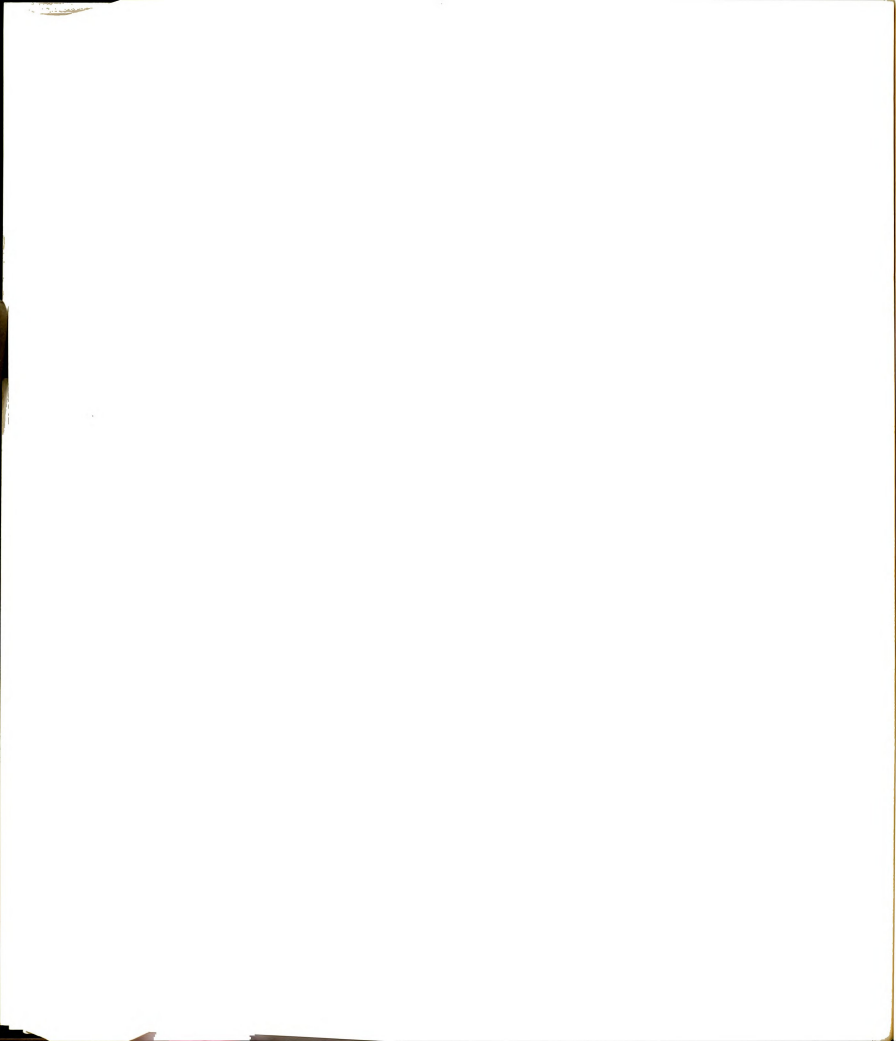
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40 EF=F*TAN2(AIMAG(G),REAL(G))
50 GO TO 100
50 EF=-90.0
50 GO TO 100
20 IF (REAL(G)) 60,65,65
60 EF=-180.0
60 GO TO 100
65 EF=0.0
65 GO TO 100
30 IF (REAL(G)) 40,75,40
75 EF=90.0
100 CONTINUE
    RETURN
    END
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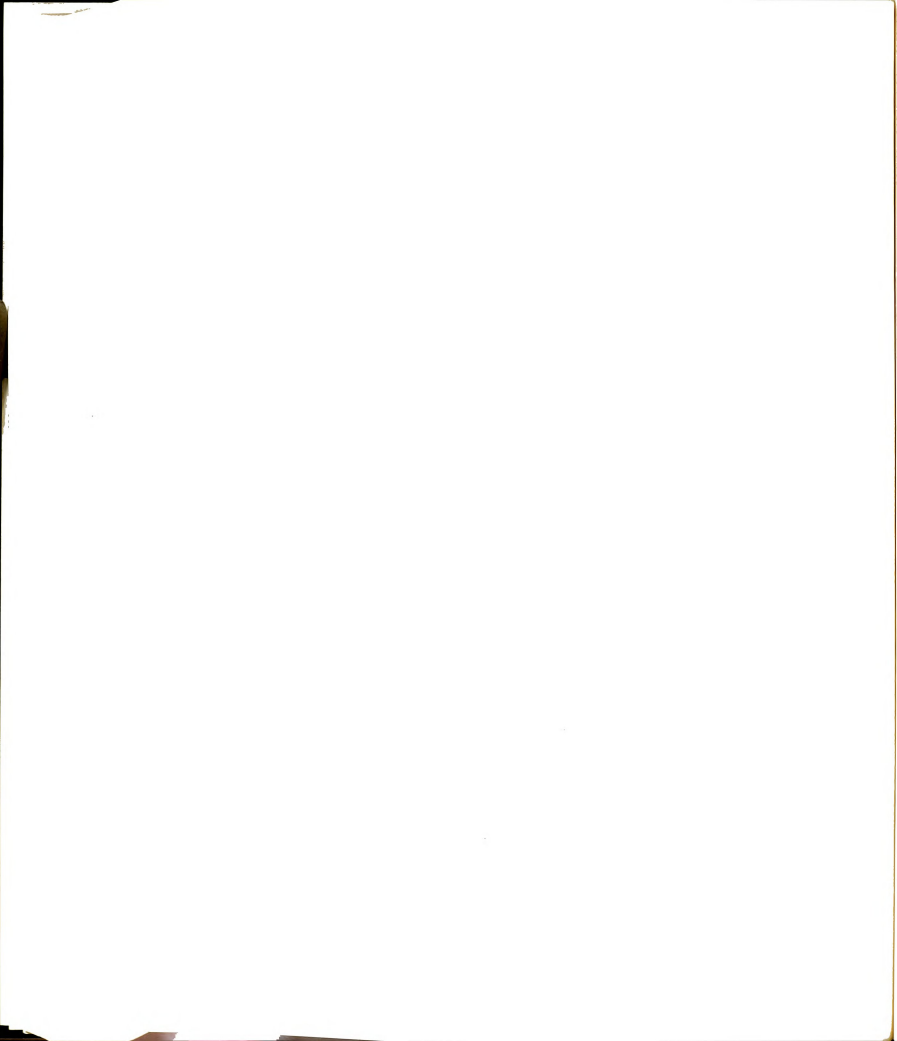
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CHAPTER V

SUMMARY

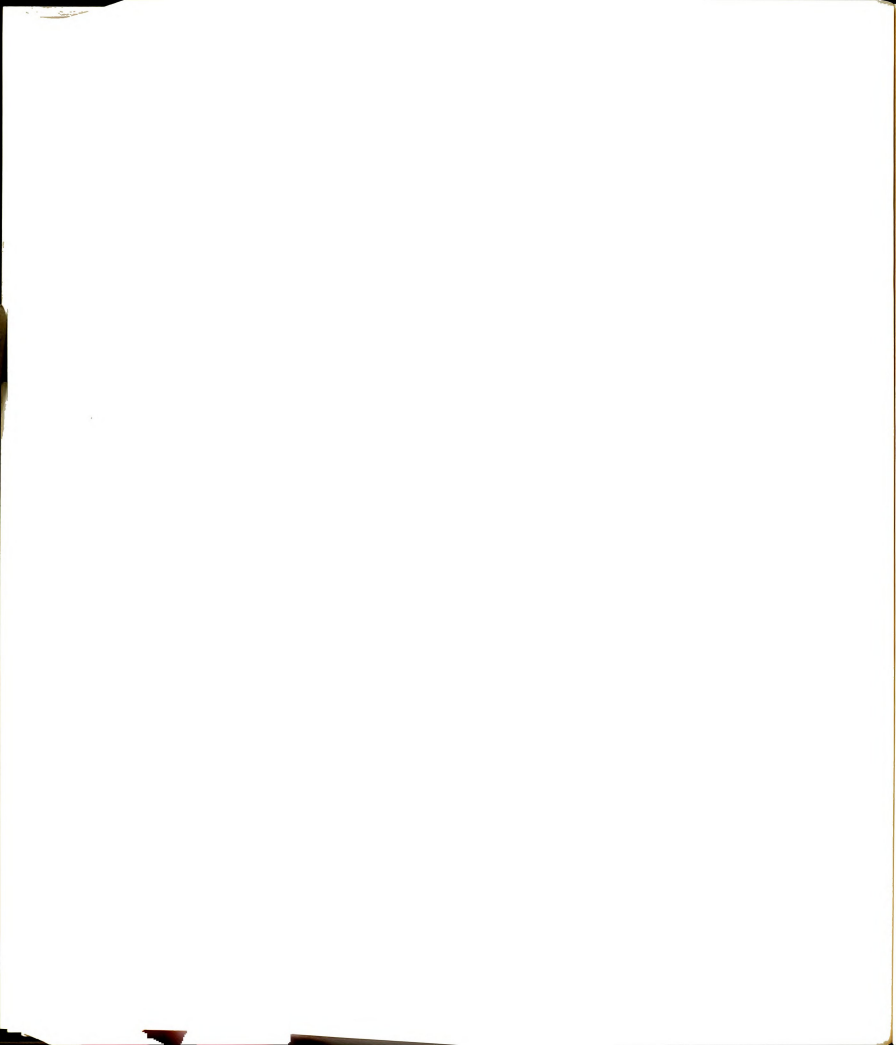
A new theoretical method for determining the eddy current induced by an oscillating magnetic field in a finite conducting body with rotational symmetry was presented in Chapter II. This study was motivated mainly by two reasons: (1) More biological research and medical applications utilize the irradiation of RF magnetic fields. (2) When a thick biological body is irradiated by a HF-VHF EM wave, the induced current can be divided into the electric and magnetic modes. The magnetic mode of induced current, also called the eddy current, is predominant and is difficult to determine with a numerical method. The new theoretical method was derived by utilizing the theory of vector potential, and the numerical results were obtained by using the point-matching and matrix inversion techniques. Two different forms of magnetic fields were considered in Chapter II, namely, a uniformly impressed magnetic field and a beam of uniform magnetic field. For both cases, the results obtained by the new method were found to deviate significantly from the conventionally used, quasi-static solution. An experiment for measuring the magnitudes of the induced



electric fields inside rotationally symmetrical bodies was conducted, and some experimental results were presented. Also, a closed form solution for the induced electric field inside a finite conducting sphere was given in that chapter. The accuracy of the new method was verified by comparing the present solution with the experimental measurements as well as the closed form solution for a finite conducting sphere.

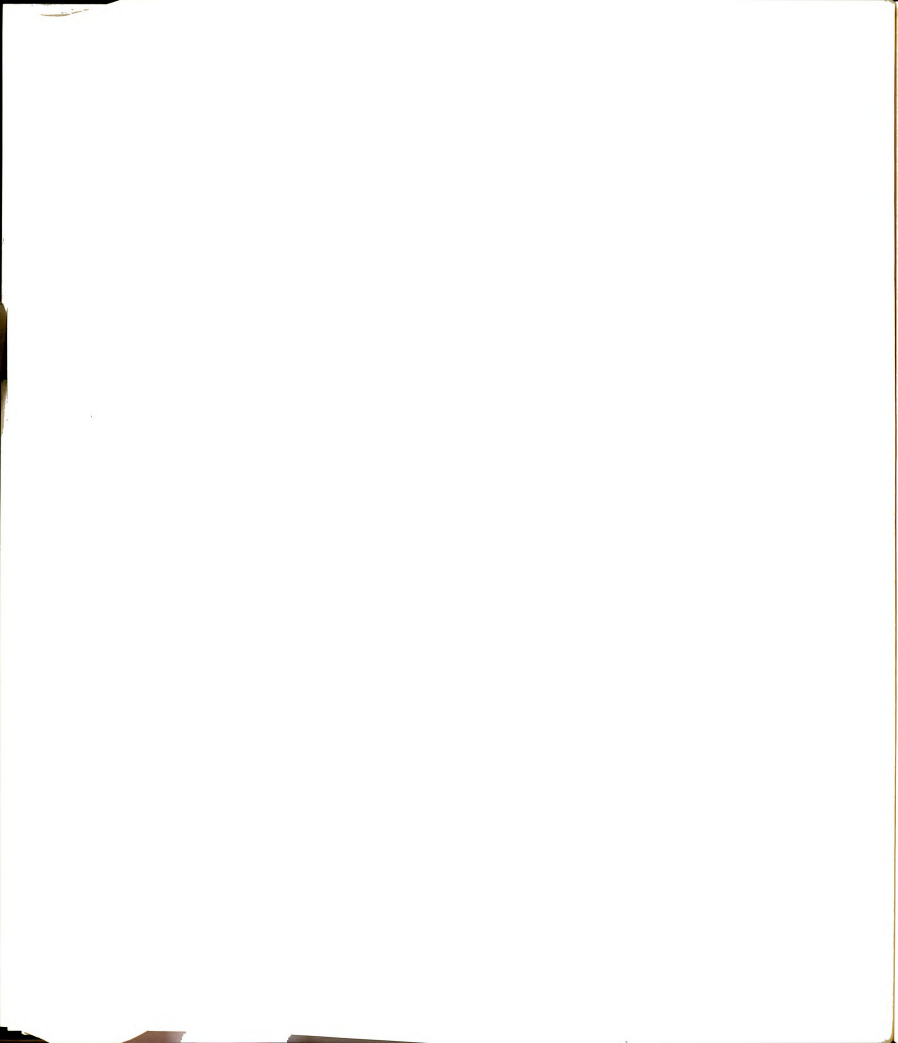
In the quantification of the magnetic mode of the induced current in an irradiated biological body, one often encounters the difficulty of numerical convergence. The new method presented in Chapter II shows an excellent numerical convergence and accuracy, and should help solve this problem. It is noted that the applicability of this new method is restricted to the cases of rotationally symmetrical bodies which are constructed with nonmagnetic materials. For the cases of magnetic materials without rotational symmetry, a different technique should be developed.

In Chapter III, the microwave interactions with finite conducting, biological bodies were studied. As the first step, we derived two coupled, surface integral equations which relate the induced fields to the incident fields. The validity of these equations was then verified by considering the special case of a two-dimensional, infinite interface between free space and a finite conducting medium.

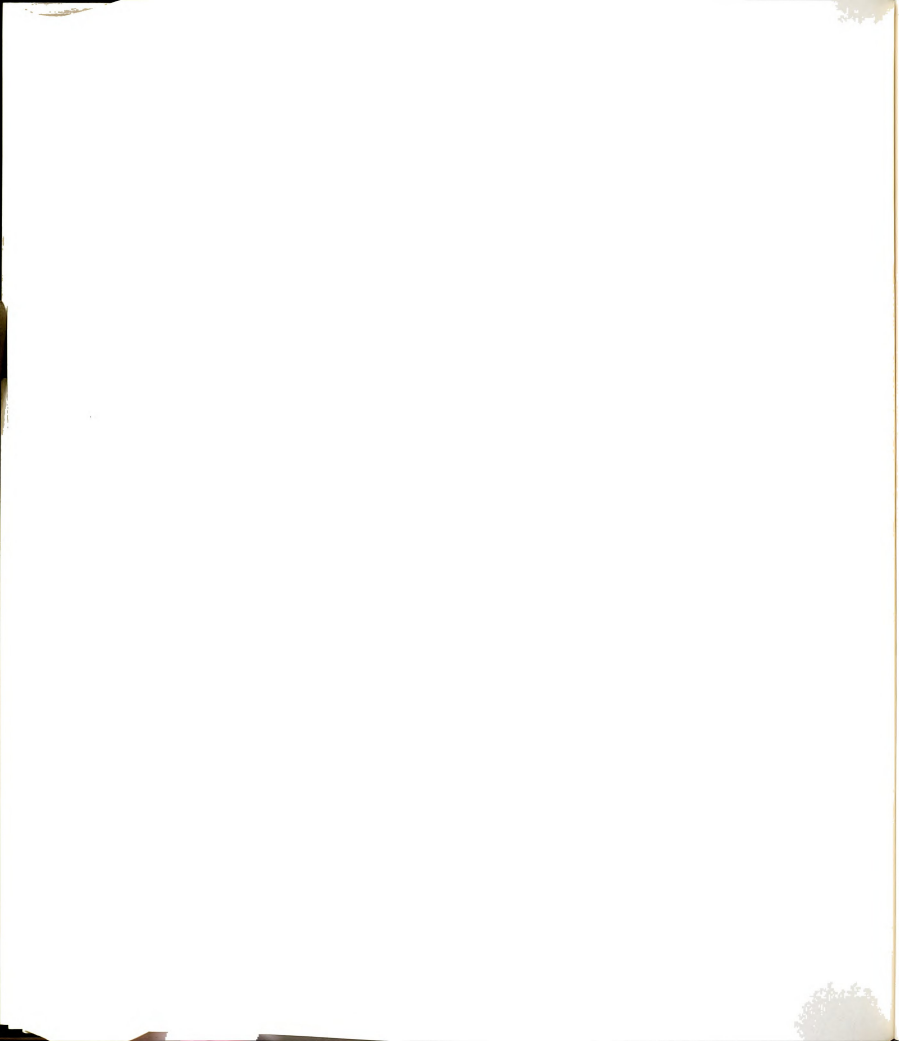


These two surface integral equations were solved numerically by use of moment-method. The results so obtained agree quite well with the solutions obtained by a volume integral equation method--Tensor Integral Equation Method. For the purpose of reducing the number of unknowns, only the cases involved with electrically small bodies were demonstrated in Chapter III. It is believed that this surface integral equation method is potentially more efficient than the volume integral equation method since when an electrically large body is irradiated by EM waves, the fields induced inside the body mainly concentrate in the region near the body surface. Under this condition, solving the internal fields through the volume integral equation becomes impractical.

Moreover, with a finer subdivision of the body for the method of moment, the number of unknowns encountered in solving the surface integral equations is usually less than that encountered for the volume integral equation, if the number of unknowns exceeds a certain number as in an electrically large body. It is pointed out that the validity of the surface integral equation method is limited to the cases of homogeneous bodies due to the nature of deriving the equations.

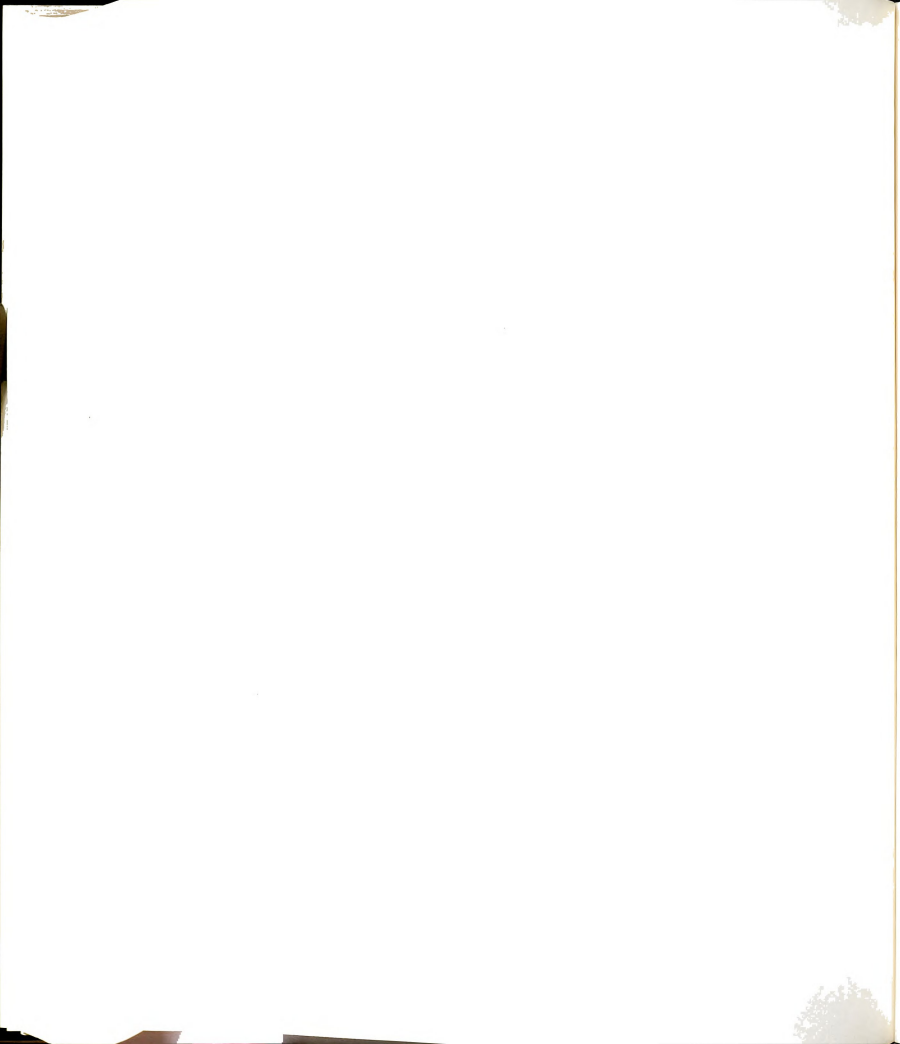


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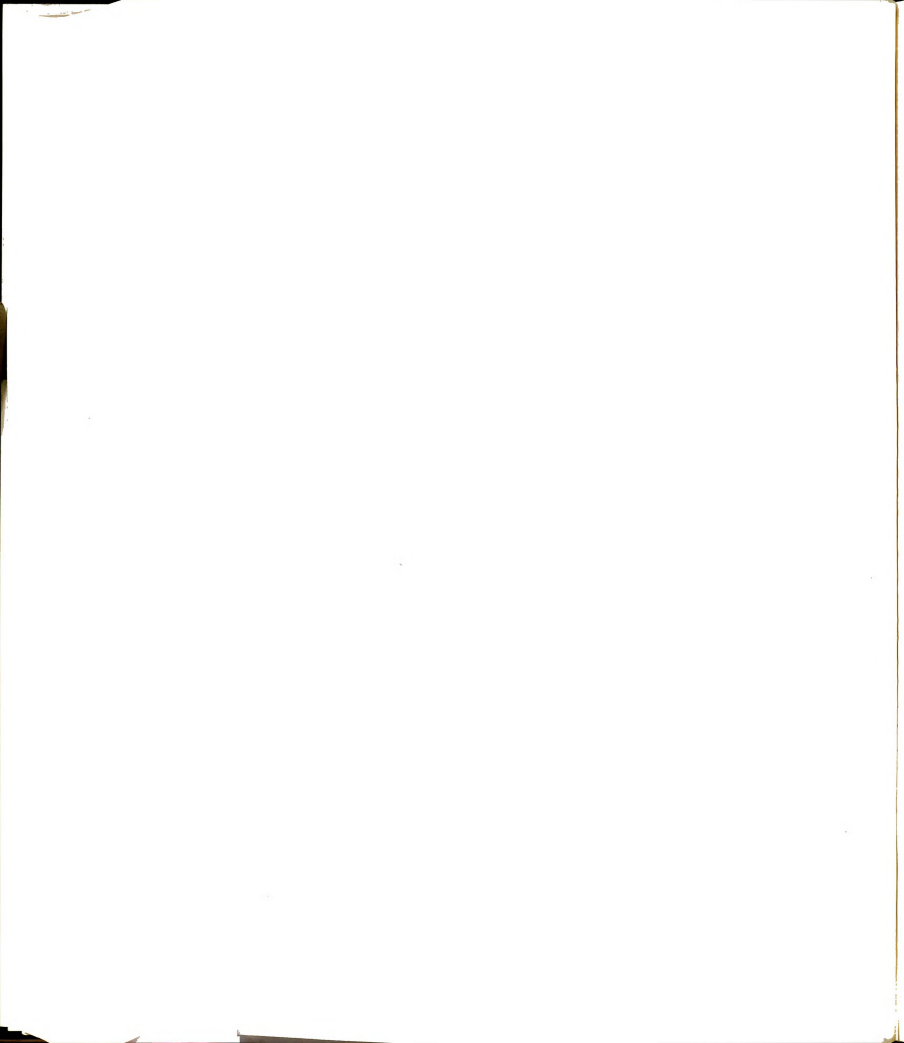


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