

THERMAL PROPERTIES OF THE GAS IN EARLY-TYPE GALAXIES AND GALAXY
CLUSTERS

By

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ABSTRACT

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Most of the baryons, or “normal” matter, found in galaxies and galaxy clusters are found in the hot, X-ray emitting gas known as the circumgalactic medium (CGM) or intracluster medium (ICM). The hot gas traces the gravitational potential well and is affected by both thermal and gravitational processes, so we use observations of the hot gas to explore changes across the galaxy or cluster’s radius. Heating and cooling in the central regions of galaxies and clusters is primarily driven by feedback processes, including Active Galactic Nuclei (AGNs) and Type Ia supernovae. We can use X-ray observations of the hot gas to understand its thermal history and how the various feedback mechanisms affect the gas at small and large radii. Furthermore, we use X-ray gas properties (temperature, density, entropy, concentration, centroid shift, and power ratios) to characterize galaxies and clusters, understand their evolution, and classify them in meaningful ways. The combination of observations along with theoretical models and simulations explored in this thesis provides key insight into understanding how feedback processes affect the hot gas.

I begin by presenting gas property results for a uniformly reduced sample of 348 galaxy clusters and show how those results can be used to characterize the sample and for further galaxy cluster science. I will then turn my focus to early-type galaxies for the remainder of this work. I examine a sample of 12 nearby early-type galaxies with powerful radio sources and find that IC 4296 exhibits unusually low central entropy as previously observed in NGC 4261. We also find some evidence that the minimum of the ratio between the cooling time and free-fall time, if it occurs at the galaxy center, may indicate the presence of a powerful radio source. Finally, I examine the galactic atmospheres of a sample of 49 early-type galaxies. I will show that the equilibrium pressure and density radial profiles for single- and multiphase galaxies agree with the Voit et al. (2020) theoretical

model. I also find evidence for a correlation between the central velocity dispersion and entropy profile slope of the galaxies in the sample that agrees with the theoretical model.

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For Dustin and Torrey

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TABLE OF CONTENTS

LIST OF TABLES	xi
LIST OF FIGURES	xii
CHAPTER 1 INTRODUCTION	1
1.1 Galaxy Clusters and Early-Type Galaxies	1
1.2 Cosmology Primer	5
1.3 X-ray Observations	8
1.3.1 A Brief History of X-ray Astronomy	8
1.3.1.1 <i>Chandra</i> X-ray Observatory	10
1.3.2 X-ray Observables	10
1.3.3 Deprojection as a Tool for X-ray Spectroscopy	11
1.3.4 Using Entropy to Understand X-ray Gas	13
1.3.5 Morphology	13
1.4 The Gas in Clusters and Early-Type Galaxies	14
1.4.1 Multiphase Gas	15
1.4.2 SNIa Feedback in Clusters vs. Galaxies	16
1.5 Structure of this Dissertation	16
CHAPTER 2 ACCEPT 2.0 AND XMM HERITAGE	18
2.1 Introduction	18
2.2 Entropy Profile Fitting	20
2.2.1 Comparison to ACCEPT	23
2.2.2 Morphology Calculations and K_0	24
2.2.2.1 Calculating Centroid Shift, Concentration, and Power Ratios	25
2.3 Science with ACCEPT 2.0	29
2.3.1 Morphological Properties and K_0 for Sample Comparisons	30
2.3.2 ACCEPT 2.0 Comparison with REFLEX and XMM Heritage	32
2.4 Summary	36
CHAPTER 3 PROPERTIES OF THE CGM IN EARLY-TYPE GALAXIES WITH POWERFUL RADIO SOURCES	38
3.1 Abstract	38
3.2 Introduction	38
3.3 Sample Selection and Data Analysis	42
3.3.1 Sample Selection and Distances	42
3.3.2 <i>Chandra</i> data reduction	44
3.3.3 Spectral Analysis	45
3.3.4 Thermodynamic Properties	47
3.3.4.1 Electron Density Profiles	47
3.3.4.2 Entropy and $t_{\text{cool}}/t_{\text{ff}}$ Profiles	47

3.4	Discussion	48
3.4.1	$t_{\text{cool}}/t_{\text{ff}}$ Profiles and Multiphase Gas	48
3.4.1.1	Comparison with Previous X-ray Analysis	50
3.4.2	Radio Luminosity and $t_{\text{cool}}/t_{\text{ff}}$	50
3.4.3	Comparison to Simulations	51
3.4.4	Metallicity Analysis	55
3.5	Conclusions	57
CHAPTER 4 RELATIONSHIPS BETWEEN CENTRAL VELOCITY DISPERSIONS AND ATMOSPHERES OF EARLY-TYPE GALAXIES		58
4.1	Abstract	58
4.2	Introduction	58
4.3	Methods	61
4.3.1	Sample Description	61
4.3.2	Theoretical Model	64
4.3.3	Entropy Profiles	65
4.3.3.1	Distribution of Central Entropy	66
4.3.4	Sub-Sample Selection	67
4.4	Discussion	68
4.4.1	Low Central Entropy, Restricted σ_v , and the Analytical Prediction	68
4.4.1.1	The Black-hole Feedback Valve Prediction	68
4.4.1.2	Comparison to the Analytic Prediction	68
4.4.2	Comparison to the Analytic Model and Numerical Integration Results	71
4.4.3	Best fit Entropy Profile Slope, Multiphase gas extent, and $\min(t_{\text{cool}}/t_{\text{ff}})$	72
4.4.4	Comments on Individual Galaxies	76
4.4.4.1	M87	76
4.4.4.2	NGC 4636	76
4.4.4.3	NGC 1521	77
4.4.4.4	NGC 4125	77
4.4.4.5	NGC 1404	78
4.4.4.6	NGC 533	78
4.4.5	Predictions for Equilibrium Pressure and Density Profiles	78
4.5	Conclusions	81
CHAPTER 5 SUMMARY		83
5.1	Summary	83
5.2	Future Work	84
APPENDICES		85
APPENDIX A	ACCEPT 2.0 PIPELINE DESCRIPTION	86
APPENDIX B	ACCEPT 2.0 CENTRAL ENTROPY FITTING RESULTS	91
APPENDIX C	ACCEPT 2.0 RADIAL ENTROPY PROFILES	104
APPENDIX D	ACCEPT 2.0 MORPHOLOGICAL PROPERTIES	154
APPENDIX E	RADIAL PROFILES OF EARLY-TYPE GALAXIES WITH POW- ERFUL RADIO SOURCES	171

APPENDIX F	RADIAL PROFILES FOR THE GALAXIES IN LAKHCHAURA	
	(2018)	177
BIBLIOGRAPHY		199

LIST OF TABLES

Table 2.1: ACCEPT vs. ACCEPT 2.0	19
Table 2.2: Density Errata in ACCEPT	24
Table 3.1: <i>Chandra</i> Observations of Early-type Galaxies	43
Table 3.2: Sample of Radial Profile Properties	46
Table 4.1: Galaxy Parameters	63
Table 4.2: Entropy Profile Slope and Velocity Dispersion Relationship Results	68
Table B.1: Fit parameters, redshift, and profile center for ACCEPT 2.0 clusters	91
Table D.1: Morphological properties for ACCEPT 2.0 clusters	155
Table E.1: Radial profiles for early-type galaxies with powerful radio sources	171
Table F.1: Radial profiles for the HQ sample	177

LIST OF FIGURES

Figure 1.1: <i>Hubble</i> Deep Field	2
Figure 1.2: Hubble Tuning Fork	3
Figure 1.3: SDSS DR1 Cosmic Voids	4
Figure 1.4: Virgo Cluster SDSS Optical	5
Figure 1.5: Hubble H_0 diagram	6
Figure 1.6: Deprojection	12
Figure 1.7: Composite Image of Hercules A	15
Figure 2.1: MCMC fit results	22
Figure 2.2: Central Entropy for ACCEPT vs. ACCEPT 2.0	25
Figure 2.3: Distribution of Central Entropy for clusters present in both ACCEPT and ACCEPT 2.0	26
Figure 2.4: Distribution of Central Entropy for clusters in ACCEPT with all clusters with a measured Central Entropy in ACCEPT 2.0	27
Figure 2.5: Morphological Property Comparison for ACCEPT 2.0	31
Figure 2.6: Central Entropy vs. concentration for ACCEPT 2.0	32
Figure 2.7: Morphological Property Comparison for the REFLEX sub-sample of AC- CEPT 2.0	33
Figure 2.8: Central Entropy vs. concentration for XMM Heritage and REFLEX sub- samples of ACCEPT 2.0	34
Figure 2.9: Morphological Property comparison for the XMM Heritage sub-sample of ACCEPT 2.0	35
Figure 3.1: Entropy Profiles for early-type galaxies with powerful radio sources	48
Figure 3.2: Radial profiles of $t_{\text{cool}}/t_{\text{ff}}$	49

Figure 3.3: Radio luminosity, $\min(t_{\text{cool}}/t_{\text{ff}})$, and radius at $\min(t_{\text{cool}}/t_{\text{ff}})$	52
Figure 3.4: Entropy profiles of NGC 4261 and IC 4296 compared to simulations	53
Figure 3.5: Comparison of the inferred entropy profiles for NGC 4261 for different assumed values of abundance	55
Figure 4.1: Stellar velocity dispersion vs. X-ray luminosity	61
Figure 4.2: Distribution of K_0 values for the HQ sample	66
Figure 4.3: Stellar velocity dispersion vs. entropy profile slope for the HQ sub-sample of galaxies	70
Figure 4.4: Stellar velocity dispersion vs. entropy profile slope for the low K_0 sub-sample of galaxies	71
Figure 4.5: α_K vs. $\min(t_{\text{cool}}/t_{\text{ff}})$ for the main sample and α_K by gas extent for the HQ sample	73
Figure 4.6: α_K vs. $\min(t_{\text{cool}}/t_{\text{ff}})$ and α_K by gas extent for the low K_0 sample	74
Figure 4.7: Equilibrium pressure vs. radius for single phase and multiphase galaxies in the HQ sample	79
Figure 4.8: Equilibrium electron density vs. radius for single phase and multiphase galaxies in the HQ sample	80
Figure C.1: ACCEPT 2.0 Entropy Profiles and Fit Information	105

CHAPTER 1

INTRODUCTION

1.1 Galaxy Clusters and Early-Type Galaxies

The Hubble Deep Field image (Figure 1.1, Williams et al. (1996)) revealed that there are 3,000 galaxies visible in just one twenty-four-millionth of the sky, indicating that the universe is full of galaxies. The masses, stellar populations, and shapes of galaxies vary widely, and astronomers still often use the early Hubble classifications of galaxies to group them by their defining characteristics. Broadly, there are elliptical, spiral, and irregular galaxies. The Hubble “tuning fork” (see Figure 1.2) was developed because Hubble believed that elliptical galaxies would eventually evolve into spiral galaxies. The belief that ellipticals evolved into spirals turned out to be incorrect, but the naming convention of referring to generally elliptical shaped galaxies as “early-type” and spiral shaped galaxies as “late-type” has prevailed and will appear in this dissertation. While the tuning fork does not sort galaxies by evolutionary stage, it does sort them by their angular momentum with spiral galaxies generally rotating faster than elliptical galaxies. Elliptical galaxies are generally red in color, contain mostly older, low mass stars, have little active star formation, and are the most massive galaxies. Spiral galaxies, like our own Milky Way, are generally blue in color, have all types of stars, and are actively forming stars. Irregular galaxies have no specific shape, contain all types of stars, are usually actively forming stars, and are typically the least massive of the three types. Particularly low mass galaxies are known as dwarf galaxies, but they are beyond the scope of this dissertation.

More recently, sky surveys of galaxies (the 2dF Galaxy survey (Colless et al., 2001) and the Sloan Digital Sky Survey (Stoughton et al., 2002)) revealed that galaxies are clumped together along large scale filaments around large voids with diameters of ~ 150 million light years (see Figure 1.3). The large scale filaments trace the distribution of dark matter in the universe. Dark matter and dark energy are so named because we cannot directly observe them; we can only



Figure 1.1: ***Hubble Deep Field*** The *Hubble Deep Field* image shows the 3,000 galaxies found in one twenty-four-millionth of the sky. The image was composed of data from the *Hubble Space Telescope* taken over ten days in 1995 and published in 1996 (Williams et al., 1996).

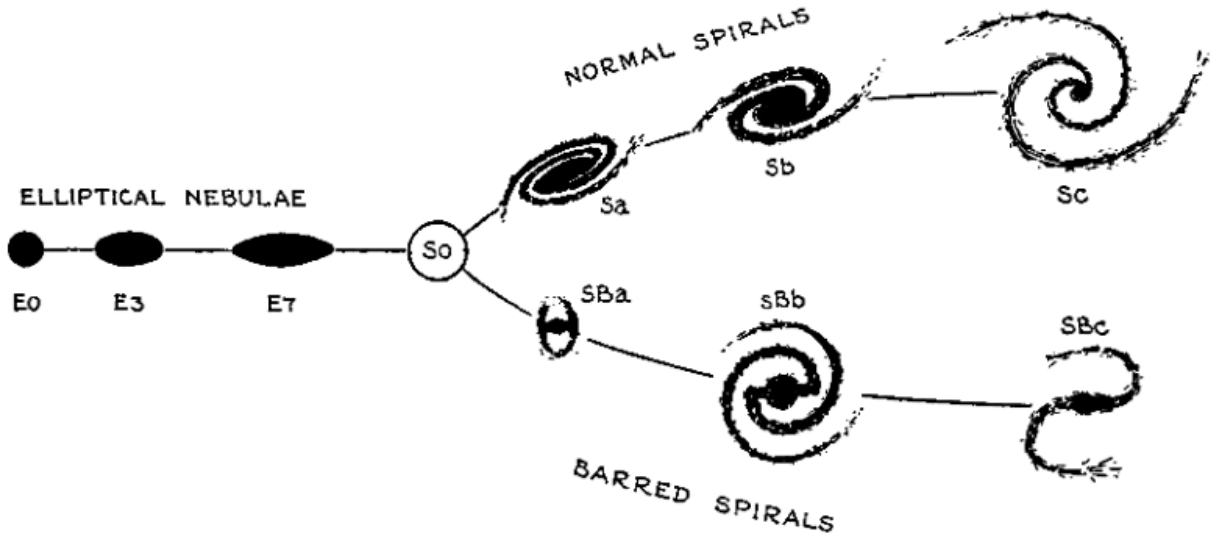


Figure 1.2: **Hubble Tuning Fork** The original tuning fork diagram from Hubble’s 1936 book, *The Realm of the Nebulae* (Hubble, 1936). Hubble developed the tuning fork to classify galaxies by their “family traits” and thought that galaxies evolved along the tuning fork from ellipticals to spirals. While we now know that the classifications are not evolutionary, but rather by angular momentum, astronomers still classify galaxies in this way.

observe their gravitational influence on the universe. Dark matter and dark energy represent most of the content of the universe, with dark energy comprising $\sim 70\%$ and dark matter comprising $\sim 25\%$ (Planck Collaboration et al., 2016). However, the focus of this dissertation is on the normal, or “baryonic,” matter in the universe. Baryons make up the remaining $\sim 5\%$ of the matter in the universe. However, there are some baryons that can be difficult to detect, leading to the “missing baryon problem” in Cosmology.

The “missing baryon problem” refers to the disparity between the baryonic mass density inferred from primordial nucleosynthesis via Cosmic Microwave Background (CMB) measurements and the baryonic mass density of galaxies, where the baryonic mass from galaxies falls far short of the baryonic mass from the CMB. The “Warm-Hot-Intergalactic-Medium” (WHIM) model was proposed to account for the missing baryons (see Cen & Ostriker 1999; Bregman 2007). The WHIM is characterized by a low-density, hot ($10^5 - 10^6$ K, Davé et al. (1999)) plasma which produces a weak signal and would be challenging to detect. Bregman et al. (2018) showed that

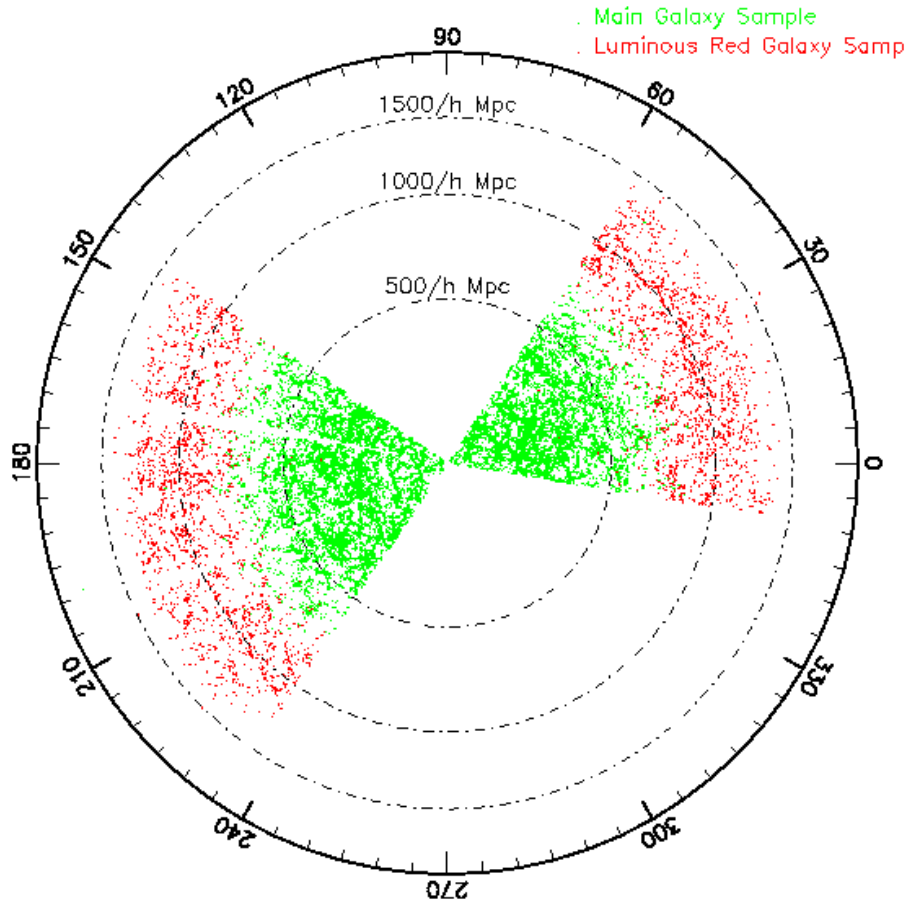


Figure 1.3: **SDSS DR1 Cosmic Voids** The distribution of galaxies from the first Sloan Digital Sky Survey (SDSS) data release (Stoughton et al., 2002). The green galaxies are from the main SDSS galaxy sample, and the red galaxies are from the luminous red galaxy sample (LRG).

baryon density estimates could be made, even with a weak signal, by stacking X-ray observations of many early-type galaxies, scaled by their radii (R_{200}). The stacked observations revealed that most, if not all, of the “missing baryons” are hot and located beyond R_{200} . While the source of missing baryons for early-type galaxies can be accounted for with the stacked X-ray observations, Bregman et al. (2018) also showed that the observed signal from early-type galaxies would be too high for spiral galaxies. Therefore, the hot halos of spiral galaxies may be different, and further constraints on the hot gas content of early-type galaxies are needed from next generation X-ray observatories.

The large scale structure of the universe likely formed in a “bottom-up” fashion, meaning that

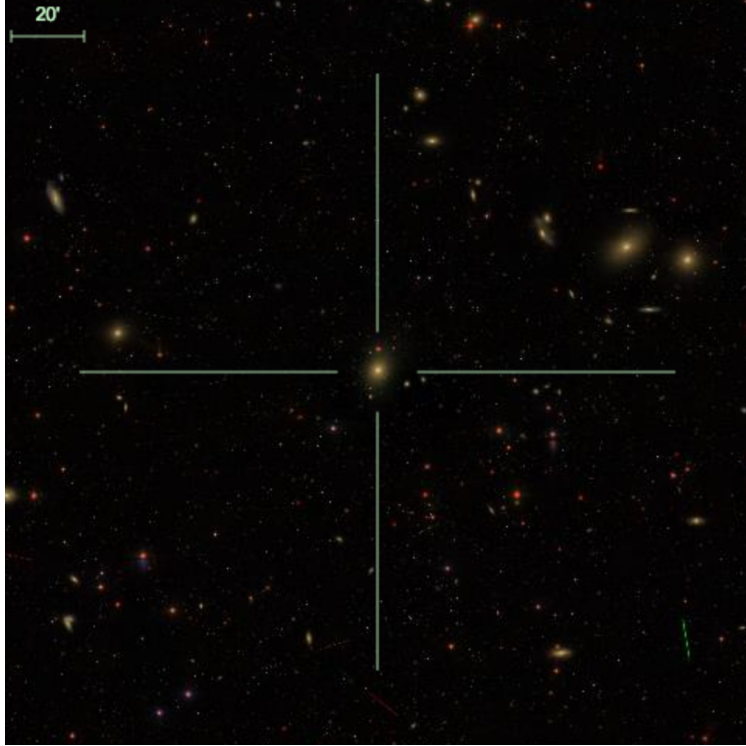


Figure 1.4: **Virgo Cluster SDSS Optical** Image of the Virgo Cluster assembled from SDSS DR15 optical data (Aguado et al., 2019). The Brightest Cluster Galaxy (BCG) is M87 (NGC 4486) and is marked by the cross-hairs, and the 20' scale corresponds to 100 kpc.

smaller scale structures merge and join together to create larger structures, along the distribution of dark matter. The largest gravitationally bound structures in the universe are galaxy clusters, composed of 100s-1000s of galaxies, all within a large clump of dark matter known as a dark matter potential well (see Figure 1.4). Galaxy clusters can contain all types of galaxies, but the majority of galaxy clusters are dominated by one galaxy called the Brightest Cluster Galaxy (BCG). BCGs are usually the brightest, most massive galaxies in a cluster, are centrally located in the cluster, and are usually elliptical galaxies. The BCG and the cluster evolve together, and understanding how BCGs (and massive elliptical galaxies) work is crucial for understanding how galaxy clusters evolve.

1.2 Cosmology Primer

Hubble (1929) observed a sample of galaxies and found that all galaxies are moving away from our own, and their recessional velocities are proportional to their distance from our galaxies

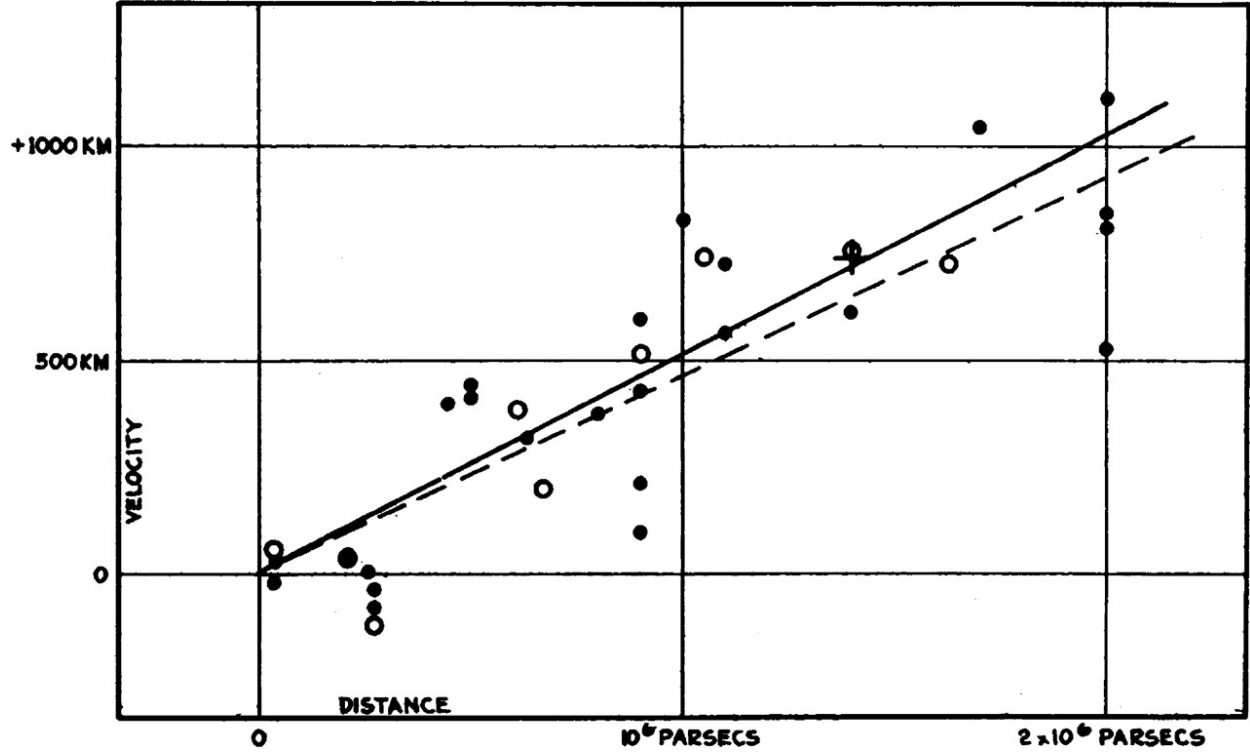


Figure 1.5: **Hubble H_0 diagram** Figure 1 from Hubble (1929) shows the distances and velocities to the sample of galaxies Hubble used to show that the universe is expanding. The y-axis are the radial velocities, corrected for solar motion, and the x-axis are the distances estimated from stars and mean luminosities in the galaxies. The slope of this distance-velocity relationship is the Hubble parameter, H_0 .

(see Figure 1.5). The velocity-distance relation, $v = H_0 r$ (*Hubble's Law*), where H_0 is referred to as *Hubble's constant*, shows that the universe is expanding uniformly. *Hubble's constant* characterizes the expansion of the universe and remains an active area of research. Type Ia supernovae measurements give $H_0 = 73.8 \pm 2.4 \text{ km s}^{-1} \text{ Mpc}^{-1}$ (Riess et al., 2011), while Cosmic Microwave Background measurements give $H_0 = 67.4 \pm 0.5 \text{ km s}^{-1} \text{ Mpc}^{-1}$ (Planck Collaboration et al., 2018). In this dissertation, we will use the widely accepted value for a cold-dark matter cosmology of $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$. The effect of choosing the Type Ia supernovae value or Cosmic Microwave background value is much smaller than the statistical uncertainty for our relevant measurements, so our choice to use single digit precision for H_0 is negligible for this work. Furthermore, for the work in Chapters 3 and 4, the galaxies are nearby enough to have distance measurements independent of H_0 . In modern astronomy, when it became possible to obtain spectra

of galaxies, the distances to galaxies could be described more precisely with a quantity known as redshift, z . Redshift is measured by the shift of the galaxy's spectrum due to its motion away from us in the the expanding universe. The shifting of a receding galaxy's optical spectrum is similar to how an ambulance's siren appears to drop in pitch as it drives away. In the case of the ambulance, the frequency of the receding siren is lowered and can be measured by the Doppler effect for a receding source. For a receding galaxy, the shift is observed in the optical (rather than sound) spectrum, but the effect is similar. The spectrum shifts to longer (redder) wavelengths the farther the galaxy is from us because we know from Hubble's Law that the farther a galaxy is from us, the faster it is moving. The redshift, z of the galaxy's spectrum is therefore

$$z = \frac{\lambda_{\text{obs}} - \lambda_{\text{rest}}}{\lambda_{\text{rest}}} \quad (1.1)$$

where λ_{obs} is the observed wavelength, and λ_{rest} is the emitted wavelength. A small redshift indicates that the galaxy is in the local universe, and a larger redshift indicates that the galaxy is far away. Because higher redshifts correspond to larger distances, which in turn correspond to longer light travel times, high redshift galaxies also provide us with an idea of how the universe looked at earlier times. For context, the most distant confirmed lensed quasars¹ are around $z \sim 6.5$ (Fan et al., 2019). The largest spectroscopic redshift obtained from a galaxy (GN-z11) is at $z = 11.1$ (Oesch et al., 2016). However, the nearby galaxies in this dissertation range from $z = 0.001 - 0.02$ (corresponding to a luminosity distance, D_L , of $\sim 4 - 90$ Mpc), and the galaxy clusters extend to $z \sim 1.5$ ($D_L \sim 10^4$ Mpc).

The expansion of the universe provides insight into what sort of universe we live in, and key cosmological parameters describe fundamental characteristics of that universe. As the quality of observations has increased along with our theoretical understanding, our ability to describe the universe with cosmological parameters has improved. Cosmic Microwave Background (CMB) measurements showed that, in general, the universe is homogeneous and isotropic, and thus, H_0

¹“Quasi-stellar objects,” are extremely luminous Active Galactic Nuclei (AGN) found in the centers of early galaxies. AGNs are compact regions at the centers of galaxies with higher than normal luminosity and have spectra that indicate the excess luminosity is not from stars.

should be constant in all directions. However, because the expansion of the universe is accelerating due to the presence of dark energy, H_0 is dependent on redshift by:

$$H(z) = H_0 E(z), \quad (1.2)$$

where $E(z) = \sqrt{\Omega_M(1+z)^3 + \Omega_\Lambda}$, and the cosmological parameters $\Omega_M \simeq 0.3$ and $\Omega_\Lambda \simeq 0.7$ refer to the matter and dark energy content of the universe, respectively. For this dissertation, we will assume the stated values for Ω_M and Ω_Λ when determining distances and spatial properties of galaxy clusters and early-type galaxies, except when redshift-independent distance measures are available in the case of nearby galaxies (see Chapter 3). The effect of using lower precision values for Ω_M and Ω_Λ affect $E(z)$ by $\sim 2\%$, which is much lower than our statistical uncertainty, so we are safe to make this assumption for this dissertation.

1.3 X-ray Observations

1.3.1 A Brief History of X-ray Astronomy

Much of the work in this dissertation builds on a long history of discoveries in the X-ray universe. We can trace the historical context and motivation for my study of early-type galaxies and galaxy clusters from the early days of X-ray astronomy to now. In contrast to many other fields of observational astronomy that date back hundreds to even thousands of years, the ~ 70 year history of X-ray astronomy is comparatively short. Because X-rays are almost entirely blocked by the Earth's atmosphere, the only way to observe the X-ray universe is from space. As a result, observations of the X-ray universe were out of reach until detectors could be launched sufficiently high in the atmosphere. Giacconi et al. (1962) launched a rocket that detected Sco X-1, the first X-ray source besides the Sun (first observed in the late 1940s (Burnight, 1949)), and detected an isotropic X-ray background. The discovery of the first X-ray source was monumental and inspired further exploration of the X-ray universe.

The first observations of extra-galactic X-ray sources were made with the *Uhuru* X-ray satellite (Giacconi et al., 1971), launched in 1970. In a series of 4 letters (Giacconi et al. 1971; Tananbaum et al. 1971; Gursky et al. 1971; Kellogg et al. 1971), presented observations from *Uhuru* that

confirmed the existence of extra-galactic X-ray sources. Furthermore, they provided observations of the spectral features and the variability of X-ray sources and the structure of X-ray emitting regions at a resolution of $30'$. The observations also led to the first observations of extended emission from galaxy clusters from the Perseus cluster (Forman et al., 1972). The results from the *Uhuru* satellite planted the first seeds for a great X-ray observatory that would eventually lead to the launch of *Chandra* in 1999.

One major question X-ray astronomers sought to address when developing the successors to *Uhuru* was whether the extended emission from extra-galactic X-ray sources was due to the integrated contributions of several discrete X-ray point sources or from diffuse processes, particularly in galaxies. To address the question of the nature of the extended emission, the *Einstein* X-ray telescope (Giacconi et al., 1979b) was developed and launched in 1978 with significantly increased spatial resolution and sensitivity ($\sim 10^6$ more sensitive than the early X-ray detectors). *Einstein* showed that many clusters of galaxies were “young” in dynamical age and involved in mergers rather than “old” and dynamically relaxed (Jones et al. 1979; Jones & Forman 1984). Forman et al. (1985) showed that early-type galaxies have hot gaseous coronae, a discovery that still influences our exploration of the hot gas in early-type galaxies today. *Einstein* confirmed the presence of extended X-ray emission in galaxy clusters, but nearby, early-type galaxies remained an area of debate (see Sarazin (1986) for a review). While observations of the emission from early-type galaxies from *Einstein* did contribute to understanding the extended emission from nearby early-type galaxies (Giacconi et al., 1979a), the debate over its origin continued and would continue until *Chandra* began collecting data and showed clear evidence for diffuse, extended, X-ray gas from the Intracluster Medium (ICM) and Circumgalactic Medium (CGM).

In between the *Einstein* and the launch of finer spatial resolution X-ray telescopes, the *Roentgen Satellite* (*ROSAT*; Pfeffermann et al. (1987)) provided the first spatially-resolved X-ray all-sky survey. *ROSAT* was launched in 1990 and completed its all-sky survey in the first 6 months and continued to take pointed observations for the next nine years. The satellite was sensitive to the “soft” X-rays between 0.1–2 keV, and contributed to mapping the diffuse galactic X-ray background

(Snowden et al., 1995) among many other discoveries. *ROSAT* remains the best X-ray all-sky survey to this day, though it will soon be succeeded by *eROSITA* (Merloni et al., 2012) which will complete an X-ray all sky survey by 2023 at similar angular resolution and ten times greater sensitivity than *ROSAT*.

1.3.1.1 *Chandra* X-ray Observatory

NASA’s *Chandra* X-ray Observatory (Weisskopf et al., 2002) is a telescope specially designed to detect X-ray emission from very hot regions of the universe such as exploded stars, clusters of galaxies, and matter around black holes. The telescope orbits above the Earth’s atmosphere, up to an altitude of 139,000 km, to capture the X-rays normally blocked by the atmosphere. The telescope was launched on July 23, 1999 and has been providing unparalleled X-ray observations for almost 20 years.

Chandra carries four precisely-constructed mirrors nested inside each other. Using data acquired with X-ray CCD detectors, detailed spectroscopic images of the cosmic source can be made and analyzed. *Chandra* has unparalleled spatial resolution and has provided countless insights for the X-ray universe. One of the richest contributions of *Chandra* is the publicly available *Chandra* Data Archive of observations, and the analysis software (CIAO) developed to reduce the data (Fruscione et al., 2006). The *Chandra* X-ray observations used in this dissertation are entirely from the archive, and this work would not be possible without the available data and the software tools to reduce and analyze it.

1.3.2 X-ray Observables

In this dissertation, we are primarily concerned with the radial properties of extended X-ray sources: galaxies and galaxy clusters. X-ray telescopes, like *Chandra* and *XMM*, obtain temporal, spatial, and energy information about X-ray photons emitted from the hot gas in galaxy clusters and early-type galaxies. The X-ray emission from the hot gas is primarily in the form of thermal bremsstrahlung radiation, due to interactions where electrons pass close to ions and lose energy.

While the temperature of the hot gas in galaxy clusters is hotter than groups or individual galaxies, they all exhibit extended X-ray emission, meaning that the X-ray emission is coming from diffuse, hot gas, rather than point sources (see Sarazin (1986) for a review). In galaxies, the extended X-ray gas is called the circumgalactic medium (CGM), and in galaxy clusters, it is called the intracluster medium (ICM). The hot gas traces the gravitational potential well and is affected by both thermal and gravitational processes, so we use observations of the hot gas to explore changes across the galaxy or cluster’s radius.

The temperature of the hot gas can be directly measured from X-ray observations, and the density can be measured by modeling the hot gas properties (e.g. cooling, emissivity, metallicity) and accounting for the hydrogen column density (N_H). Because clusters and nearby early-type galaxies are large enough to be spatially resolved by *Chandra* and *XMM*, we can break up the images of extended sources into concentric annuli, usually centered on the peak or centroid of the X-ray emission, to obtain independent X-ray spectra. Then we use the spectra to derive radial profiles of temperature and density. The width of the annuli is usually set by a signal-to-noise threshold based on the science goals for a particular analysis but must be larger than the observatory’s point-spread-function (PSF) to avoid significant corrections for light scattered from outside the annuli or from other annuli. Profiles constructed from two-dimensional annuli are known as “projected” radial profiles because the emission is assumed to be two-dimensional, and the annuli are independent of each other. Projected profiles can be useful, particularly because of their low computational cost, but they are limited in their ability to capture the three-dimensional nature of the hot gas. Therefore, we rely on deprojection techniques to extract radial properties.

1.3.3 Deprojection as a Tool for X-ray Spectroscopy

If we can assume that the emissivity of the gas is constant and optically thin within a spherical shell, we can use deprojection to obtain three-dimensional source properties from a two-dimensional image. Figure 1.6 shows a geometric view of how deprojection works. The ICM/CGM are approximately spherical, the gas is optically thin, and variations in the gas properties are small

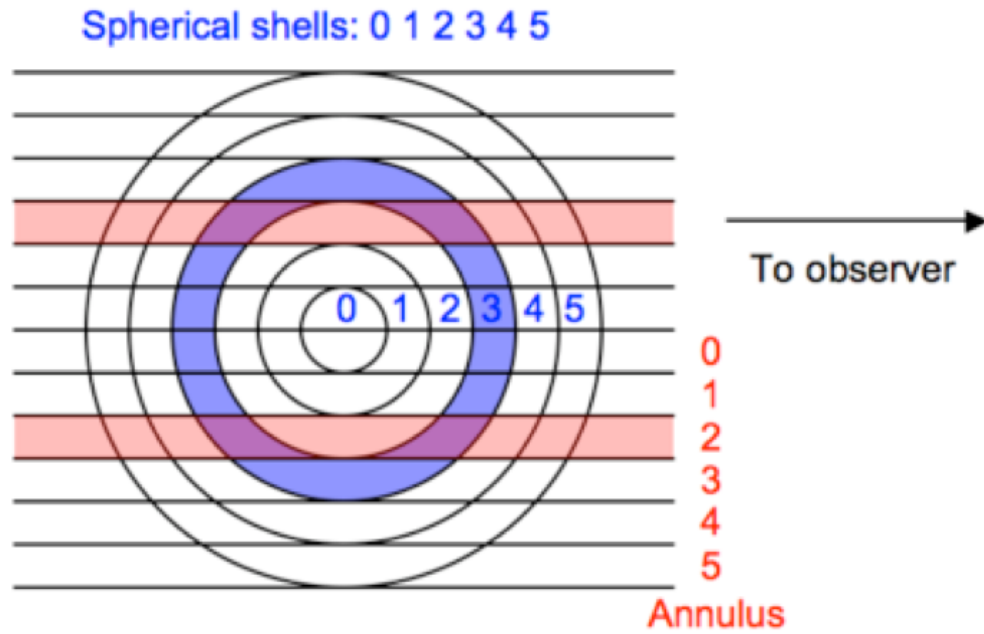


Figure 1.6: **Deprojection** Geometric view of how deprojection works. The two-dimensional annuli are converted to spherical shells to obtain two-dimensional source properties of a three-dimensional object, in this case, a galaxy or galaxy cluster. Image courtesy of deproject documentation (<https://deproject-test.readthedocs.io>).

across annuli, so deprojection is well suited for obtaining radial profiles. Like projected profiles, the image is broken up into annuli, centered approximately on the peak of the X-ray emission, based on a signal-to-noise threshold for the annuli. However, rather than each annulus being independent, the spectra of the annuli are fit from the outermost annulus inwards, with each bin accounting for the properties of the previous bins. The technique is often referred to as “onion-peeling” because it treats the source as an onion where each annulus is representative of a spherical shell of emission, so the emission in an annulus includes contributions from all external annuli, like a core sample from an onion would contain contributions from all the layers.

Deprojection requires more computational time because of the iterative spectral fitting process and does not work as well for emission that is not centrally peaked or radii where the emission is close to the contributions from the background. However, it does provide more accurate measurements for temperature and density in the centers of galaxy clusters and early-type galaxies where we are most interested in the radial gas properties. This dissertation contains multiple uses of deprojected

profiles to better understand the ICM and CGM at smaller radii.

1.3.4 Using Entropy to Understand X-ray Gas

The conventional definition of entropy, $dS = dQ/T$, is not easily determined from X-ray observations, so we must resort to a simpler surrogate for entropy, an adiabat. For an ideal monoatomic gas, the adiabatic equation of state is $P = K \rho^{5/3}$, where K is the adiabatic constant and ρ is the number density. Recasting that equation in terms of the X-ray observables temperature and electron density, we get the quantity we call entropy:

$$K = kT_{\text{keV}} n_e^{-2/3}, \quad (1.3)$$

where kT_{keV} is the temperature of the gas in keV, and n_e is the electron density. It is the preferred physical quantity for capturing feedback in galaxies because the density and temperature of the gas can change independently. Feedback processes such as thermal cooling, supernovae, or AGN outbursts do not necessarily heat the gas, but they can change the rate at which it radiates energy away which can change the time it takes the gas to cool. Entropy tracks gains and losses of energy in the X-ray gas. The galaxy potential well serves as an entropy sorting device, with higher entropy gas at outer radii and lower entropy gas at small radii. The lowest entropy gas is the densest and brightest in the X-ray.

1.3.5 Morphology

In addition to determining the gas properties from X-ray observations, we can also quantify what the galaxies and clusters “look” like, or their morphology. For a discussion of how to calculate morphological properties, see Section 2.2.2.1, but here I will introduce the ways in which we quantify the X-ray emission distribution qualitatively. We can describe how “peaked” the emission is with a parameter referred to as concentration. Concentration is a ratio between the luminosity interior to some inner radius to the total luminosity inside a larger radius. The value of the ratio is between 0 and 1, where 1 indicates that all of the emission is within the inner radius and 0 indicates

that all of the emission would be outside the inner radius. X-ray luminosity is an observationally easy measurement to make for most clusters, so determining how it relates to other properties of the X-ray gas could allow us to make estimates for measurements that ordinarily require much longer observations.

We can also determine how spherically symmetric, or “relaxed” the hot gas is by calculating the centroid shift and power ratios. Full details of these calculations are provided in Section 2.2.2. Centroid shift is measured by calculating the distance between the peak and centroid of the X-ray emission and how that distance changes as the aperture for the measurement changes in size. Small centroid shift indicates that the gas is more compact, and large centroid shift indicates that it is more diffuse. For measurements of power ratios, any number of moments can be calculated for the hot gas, but the most relevant are the 0th and 3rd moments because their ratio provides an indication of asymmetries in the gas. The 0th moment is simply the total flux, the 1st moment should be near zero if the aperture is centered on the centroid, the 2nd moment estimates ellipticity, and the 3rd moment is sensitive to asymmetries in the surface brightness. If the 3rd moment is large in comparison to the 0th moment, there is more substructure in the gas, meaning it is less symmetric and less relaxed.

1.4 The Gas in Clusters and Early-Type Galaxies

In the center of (almost) every galaxy, there is a supermassive black hole that is tightly coupled to the evolution of the galaxy (e.g. Kormendy & Richstone 1995; Haehnelt et al. 1998; Magorrian et al. 1998; Ferrarese & Merritt 2000; Gebhardt et al. 2000; Kormendy & Ho 2013; Reines & Volonteri 2015; Saglia et al. 2016; Ricarte et al. 2019). Many galaxies also have an Active Galactic Nucleus (AGN) powered by accretion onto the black hole (Brandt & Hasinger, 2005). AGNs are very small spatially in comparison to their host galaxy yet they are able to affect the hot gas on much larger scales (see Figure 1.7). How exactly the accretion fueling of the black hole is coupled to the surrounding medium is still an unanswered question, but it may be through precipitation driven feedback (e.g. Pizzolato & Soker 2005; McCourt et al. 2012; Voit et al. 2015b; Sharma et al. 2012; Voit et al. 2017). The precipitation model posits that feedback from the central black hole

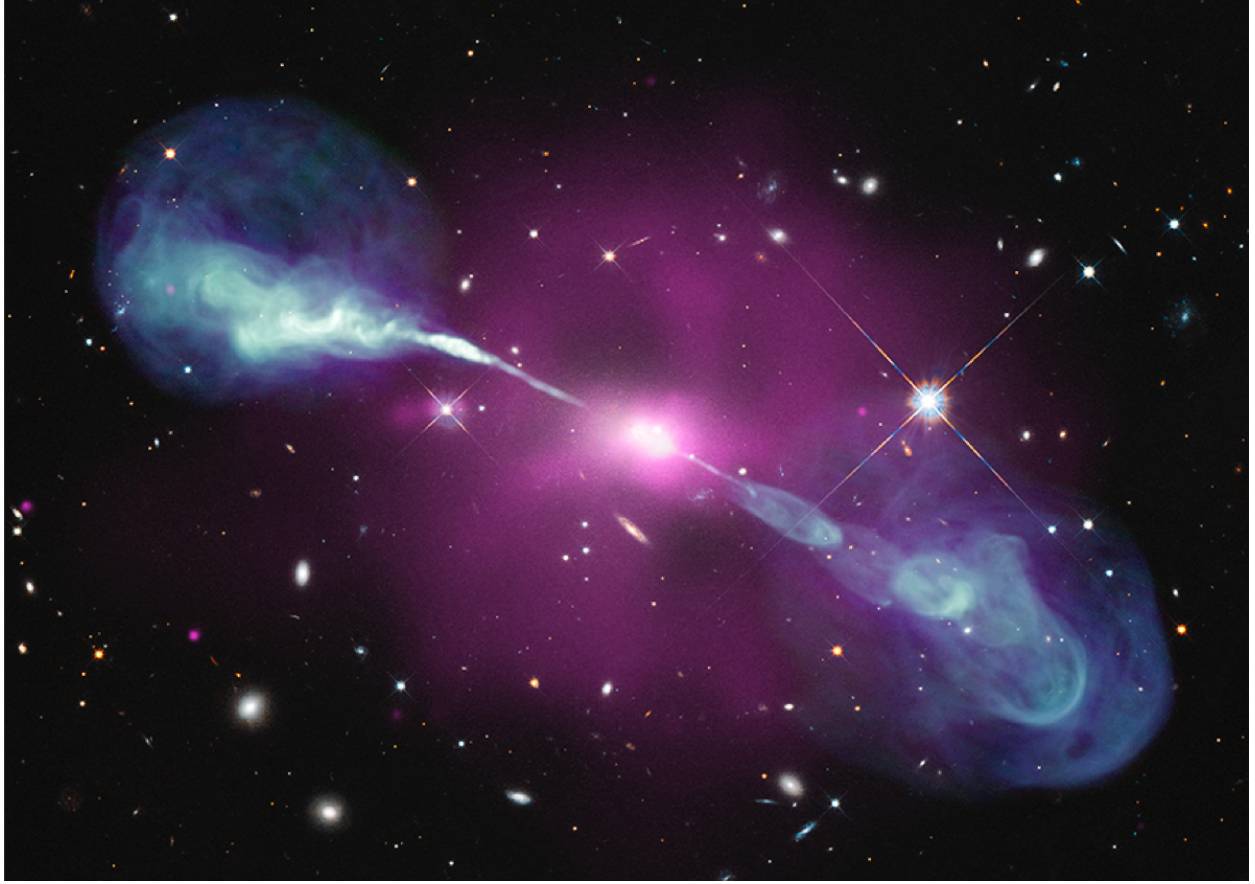


Figure 1.7: **Composite Image of Hercules A** A composite image (X-ray (pink), optical, and radio emission (blue)) for the nearby early-type galaxy Hercules A. The composite image illustrates the importance of a multi-wavelength approach to studying the phenomena in galaxies and galaxy clusters. *NASA/CXC/SAO*

holds the CGM in a state marginally unstable to condensation, so with the right conditions gas can precipitate out. As gas cools and falls into the black hole, the AGN will turn on, heating the gas, thus lengthening its cooling time and diminishing precipitation. In this dissertation, I will explore how AGNs couple to the galactic atmosphere and affect the gas entropy (see Chapters 3 and 4).

1.4.1 Multiphase Gas

Up to this point, I have focused on the hot X-ray emitting gas in galaxies and clusters, but I also explored some of the other gas found in the ICM/CGM. At smaller radii in the ICM/CGM, we find multiphase gas: gas at different temperatures and ionization states found in close proximity to each other. The hot X-ray emitting gas is the volume-filling “ambient” phase of the atmosphere,

while the molecular gas is generally found towards the cluster/galaxy center. The presence of $H\alpha$ or CO emission usually indicates abundant molecular gas (Edge, 2001) and the potential for active star formation. Some galaxies, particularly early-type or elliptical galaxies, have little to no active star formation and thus have no extended multiphase gas, so we call these “single phase” galaxies. Chapters 3 and 4 examine how the multiphase gas extent relates to other galaxy properties. We also expect clusters to have multiphase gas if their gas entropy is low in the center.

1.4.2 SNIa Feedback in Clusters vs. Galaxies

In addition to feedback from AGN, this dissertation is also concerned with feedback from Type Ia supernovae (SNIa), particularly with respect to its effect on the hot atmospheres of early-type galaxies (Voit et al., 2015b, 2020). Despite the generally older, low mass stellar populations of elliptical galaxies, SNIa are found in early-type galaxies because SNIa progenitors are long-lived white dwarf stars in binary systems. Unlike core-collapse supernovae, SNIa result from systems where a two white dwarfs merge or a white dwarf accretes matter from a low-mass companion star merge until oxygen fusion begins and the white dwarf explodes, rather than collapsing. Because SNIa are found in older stellar populations, they can also be used as standard candles to measure the expansion of the universe (see Section 1.2 and e.g. Riess et al. (2011)).

In galaxy clusters, the contributions to the sweeping of gas from inner radii via SNIa can usually be neglected because the effect is small in comparison to the influence of gravitational accretion and AGN feedback processes. However, contributions from SNIa via stellar winds can be sufficient to remove gas from the inner radii of an early-type galaxy. Chapter 4 discusses how stellar winds from SNIa contribute to the balance of feedback and cooling in early-type galaxies.

1.5 Structure of this Dissertation

The structure of this dissertation is as follows. In Chapter 2, I will discuss the ACCEPT 2.0 entropy profile and morphology measurements and present an early science application of the ACCEPT 2.0 database. In Chapter 3, I will present the results of Frisbie et al. (2020) exploring the thermal properties of the gas in a small sample of early-type galaxies with powerful radio

sources and compare them to simulations. In Chapter 4, I will present an observational test of the black-hole feedback valve model for galactic atmospheres using early-type galaxies. Chapter 5 contains a summary of this dissertation as well as potential future projects.

CHAPTER 2

ACCEPT 2.0 AND XMM HERITAGE

2.1 Introduction

Most of the baryonic mass in a galaxy clusters is actually not in the galaxies but in the hot ($10^7 - 10^8$ K) Intracluster Medium (ICM). The X-ray emission we observe comes from the radiation of the ICM gas, and the baryons track the dark matter halo. In a galaxy cluster, the cooling of gas, winds, and the heating of gas due to feedback from Active Galactic Nuclei (AGN) drive the cluster cores away from hydrostatic equilibrium. To study the non-gravitational processes of the X-ray gas, we primarily use the gas entropy (K) (see Section 1.3.4). Convection in the hot ICM, bound by the gravitational potential of a cluster, causes high entropy gas to rise and low entropy gas to sink, creating a positive entropy gradient across the cluster radius ($dK/dr > 0$). If the physics of the hot ICM is dominated only by gravitational processes and gas accretion, the entropy should be a single power law, so departures from power law entropy allow us to measure the effect of feedback and radiative cooling in the X-ray gas.

Cavagnolo et al. (2009) presented the ACCEPT (Archive of Chandra Cluster Entropy Profile Tables) project and showed that, generally, every galaxy cluster has an entropy excess near its center. To quantify the entropy excess, they fit a mathematical model of a power law with a core excess:

$$K(r) = K_0 + K_{100} \left(\frac{r}{100 \text{ kpc}} \right)^\alpha, \quad (2.1)$$

where K_0 is a characteristic central entropy, K_{100} is the best fit entropy at a radius of 100 kpc, and α is the best fit power law slope. They measured K_0 for a sample of 239 clusters and found that, while the model to fit the entropy profiles was a purely mathematical model, the distribution of K_0 values in the sample was bimodal with peaks at 15 keV cm^2 and 150 keV cm^2 and a threshold entropy of 30 keV cm^2 . Furthermore, they found that clusters with “low” central entropy ($K_0 < 30 \text{ keV cm}^2$) generally also had multiphase gas present while clusters with “high” central entropy ($K_0 > 30 \text{ keV cm}^2$) never have multiphase gas. Therefore, central entropy serves as a

convenient way to sort galaxy clusters into those with BCGs that might have multi-phase gas and those with BCGs that never have it.

The sorting of clusters by central entropy also fairly neatly separates clusters into cool core (CC) and non-cool core (NCC). The division between cool-core and non-cool core arises from the radial temperature profiles of clusters. Cool-core clusters have a drop in temperature towards the center that aligns with a sharper peak in surface brightness and an increase in density. In non-cool core clusters, however, the surface brightness peaks less dramatically, and their temperature profiles are flatter.

While the *Chandra* telescope remains the same, computational power and the number of clusters observed have greatly increased since 2009. Therefore, the ACCEPT 2.0 project was established in 2015. The X-ray data reduction was completed by Alessandro Baldi and full details of the reduction pipeline can be found Section A.1. The morphological measurements were completed with a pipeline written and run by Megan Donahue in 2018 (see Section 2.2.2.1). The main improvements from ACCEPT to ACCEPT 2.0 are summarized in Table 2.1, but I will highlight the most significant for my thesis work here.

Table 2.1: **ACCEPT vs. ACCEPT 2.0** Comparison between the data products of ACCEPT and ACCEPT 2.0

	ACCEPT	ACCEPT 2.0
# of Clusters	239	606
# of Profiles	239	348
Deprojected Density	Yes	Yes
Deprojected Temperature	No	Yes
Global T, L, Z	No	Yes
Morphology	No	Yes

ACCEPT 2.0 contains entropy (and the associated temperature and density) profiles for 348 clusters and global measurements for up to 606 clusters. Deprojection is computationally intensive because of the iterative fitting process required, and in 2009, sufficient computational resources were not available, so Cavagnolo et al. (2009) used projected temperature profiles instead. ACCEPT 2.0

uses deprojected profiles which provide a more accurate measure of the temperature and density than projected temperatures. Rather than determining the temperature from two-dimensional annulus, deprojected profiles use three-dimensional spherical annuli (see Section 1.3.3 for a full description), accounting for the three-dimensional nature of galaxy clusters and providing more accurate profile measurements. In addition to more robust profile measurements, ACCEPT 2.0 includes morphology (here referring to the shape and distribution of the X-ray gas) measurements while ACCEPT did not. Finally, ACCEPT 2.0 contains global property measurements (temperature, luminosity, and metallicity) for the sample while ACCEPT did not.

2.2 Entropy Profile Fitting

The X-ray gas in clusters tends to be relatively spherically symmetric and centrally peaked (usually on the BCG), so deprojection generally works well. However, deprojection does have its limitations. The primary limitation that is relevant in this work is that when the difference between the background and source emission is small, specifically in the outer edges of the cluster, the contents of the deprojected bins can be over- or under- subtracted, resulting in a jagged profile. The jagged profile results because one bin compensates for the estimated contents of an outer bin. For that reason, direct deprojection is far more stable when the inner bins are far brighter than the outer bins as in extended but centrally-peaked source distributions. Deprojection generally assumes spherical or ellipsoidal symmetry, so clusters that have mergers, shocks, or other significant asymmetries in the X-ray gas will have more uncertain profiles. Finally, there is some covariance in the radial profile that would not be present in a projected profile because projected annuli do not attempt to consider contributions from other annuli as deprojection does, but it generally does not need to be considered for my analysis.

Because entropy traces gains and losses of energy in the gas, we can use an entropy profile to gain an understanding of the thermal history of the cluster. The goal of fitting entropy profiles is to characterize the core excess of entropy because it provides a simple way of examining sub-populations of galaxy clusters with similar characteristics. Specifically, Cavagnolo et al. (2009) showed that there is some correlation between the presence of multiphase gas and cool-core clusters

and the absence of multiphase gas in non cool-core clusters. Furthermore, galaxy clusters with low central entropy may have more peaked surface brightness profiles and may be more relaxed. Therefore, central entropy measurements can provide additional ways to examine the morphological properties of galaxy clusters, particularly with a sample as large as ACCEPT 2.0.

The entropy profiles are fit with the functional form in Equation 2.1. While the functional form describes the shape of the entropy profile well, it is not a physically motivated model. Because the data are not always smooth and not uniform across clusters, we elected to fit the entropy profiles Markov-Chain Monte Carlo fitting, rather than simpler, less computationally expensive methods. The best fit parameters were determined using the `emcee` Python package. Because the best fit parameters for K_0 , K_{100} , and α are not expected to have a particular a priori distribution, we chose to initially limit parameter space broadly in log space for K_0 and K_{100} ($-10^2 < K_0 < 10^3$, $0 < K_{100} < 10^3$) and linear space for α ($0 < \alpha < 2$). To manage the computational time required, we utilized 1000 parallel random walkers, each taking 1500 steps, to obtain a statistical distribution for the best fit parameters. See Figure 2.1 shows the results for one galaxy cluster (Abell 209). Errors were determined from gaussian 1σ contours in two dimensions (16-84%). The complete results of these fits are available in a table in Appendix B and radial entropy profiles are in Appendix C.

For the purposes of classifying the clusters of ACCEPT 2.0 by central entropy, our priority was to characterize the shape of the entropy profiles in the innermost regions of the galaxy clusters. In 80 of the clusters, to obtain a reasonable fit in the central region, we restricted the radial range of the fit to the central radial bins (see Appendix C for radial ranges), rather than requiring a good fit to points at radii greater than ~ 100 kpc. If restricting the radial range of the fit did not improve the statistical significance of the fit, or if there were insufficient radial bins (less than 4) to fit inside the region of interest, we removed the profile from the sample. Of the 606 galaxy clusters in ACCEPT 2.0, 348 had sufficient counts to fit an entropy profile. Of those 348 profiles, 39 were removed from the sample because of insufficient data resolution in the region of interest or the lack of a statistically significant fit (reduced $\chi^2 > 1.1$) in the region of interest.

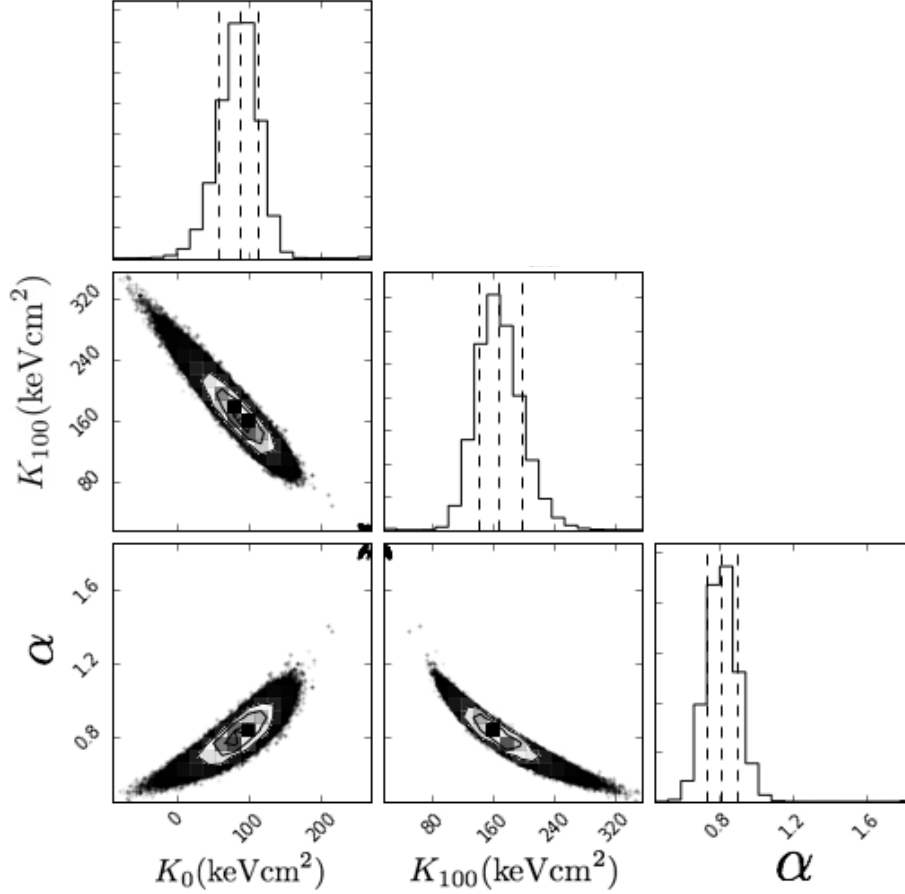


Figure 2.1: **MCMC fit results** Results of MCMC fitting for the entropy profiles using `emcee` for Abell 209. Abell 209 is used to represent the fitting process because it has average data quality for our sample and is present in both ACCEPT and ACCEPT 2.0 for validation purposes. Two dimensional, marginal distributions of K_0 , K_{100} , and α are given with respect to each other as well as the distribution of each parameter independently with two dimensional gaussian 1σ errors, represented by vertical dashed lines in one dimension and contours in two dimensions (lighter gray in the center of the distributions).

2.2.1 Comparison to ACCEPT

As a first validation of the ACCEPT 2.0 data reduction pipeline and central entropy fitting and a test of the conclusions from Cavagnolo et al. (2009), we compared the K_0 results for clusters present in both ACCEPT and ACCEPT 2.0 and the K_0 results for all clusters in ACCEPT 2.0. While in general the best fit central entropies are slightly lower in ACCEPT 2.0 than in ACCEPT, there is strong agreement between the clusters present in both samples, providing convincing validation of both the entropy profile fitting procedure and the data reduction pipeline (see Figures 2.2 and 2.3). In Figure 2.4, we compare the best fit K_0 for all ACCEPT clusters with the best fit K_0 for all ACCEPT 2.0 clusters. Cavagnolo et al. (2009) showed that the distribution of K_0 was bimodal with peaks at 15 keV cm^2 and 150 keV cm^2 . They also showed that the characteristic central entropy, $K_0 = 30 \text{ keV cm}^2$, divides clusters into cool core and non-cool core clusters, and we recover that same threshold entropy in the ACCEPT 2.0 cluster sample. Furthermore, ACCEPT showed that $K_0 = 30 \text{ keV cm}^2$ divides clusters into those with central radio sources and evidence for multiphase gas ($K_0 < 30 \text{ keV cm}^2$) and those without ($K_0 > 30 \text{ keV cm}^2$), and we would expect the same behavior from ACCEPT 2.0 clusters.

While it does not strongly affect the agreement between ACCEPT and ACCEPT 2.0, there is a small offset between the results of ACCEPT and ACCEPT 2.0. Deprojected inner temperatures (as we used in ACCEPT 2.0) are generally lower than projected temperatures (as we used in ACCEPT) because deprojected temperatures are obtained by subtracting off the contribution of a spherical shell rather than a two-dimensional annulus (as discussed in Section 2.2). Therefore, the slightly lower characteristic central entropies in ACCEPT 2.0 are reasonable, and in fact expected. In my comparison of ACCEPT to ACCEPT 2.0 profiles, I found that 21 (see Table 2.2 ACCEPT profiles) were systematically offset from ACCEPT 2.0 as a result of an extraneous correction factor of ~ 1.2 applied to the electron density computation in the ACCEPT profiles. However, the systematic offset in density was small relative to the uncertainty in temperature and the statistical uncertainty in determining K_0 , so clusters with the offset in ACCEPT remained in the comparison.

Table 2.2: **Density Errata in ACCEPT** Clusters listed in this table (by ACCEPT 2.0 name) are clusters where the density provided by ACCEPT is systematically lower than the ACCEPT 2.0 density by a factor of ~ 1.2 .

ACCEPT 2.0 Name
MFGC_06756
Abell_223
ABELL_0402
ABELL_0611
ABELL_0963
ABELL_2069
ABELL_2813
ABELL_3088
ABELL_3444
ABELL_S0592
MACS_J2214-1359
MCXC_J0220.9-3829
MCXC_J0439.0+0520
MCXC_J0454.1-0300
MCXC_J0547.0-3904
MCXC_J1000.5+4409
MCXC_J1010.5-1239
MCXC_J1022.0+3830
MCXC_J1130.0+3637
ZwCl_0857.9+2107
ZwCl_0949.6+5207

2.2.2 Morphology Calculations and K_0

Morphological measurements of galaxy clusters give us a way to characterize the shape and distribution of the X-ray gas. The different morphology measurements provide insight into the history of galaxy clusters, including significant mergers, star formation, and feedback processes. We look at morphology measurements in relation to K_0 for two primary reasons. First, because K_0 is a convenient way to divide a sample of galaxy clusters into groups with respect to the presence of multiphase gas, we want to understand what correlations, if any, exist with respect to the morphological properties (power ratio, P_3/P_0 , centroid shift, w , and concentration, c). Second, K_0 is an observationally expensive measurement to make because it requires sufficient counts to

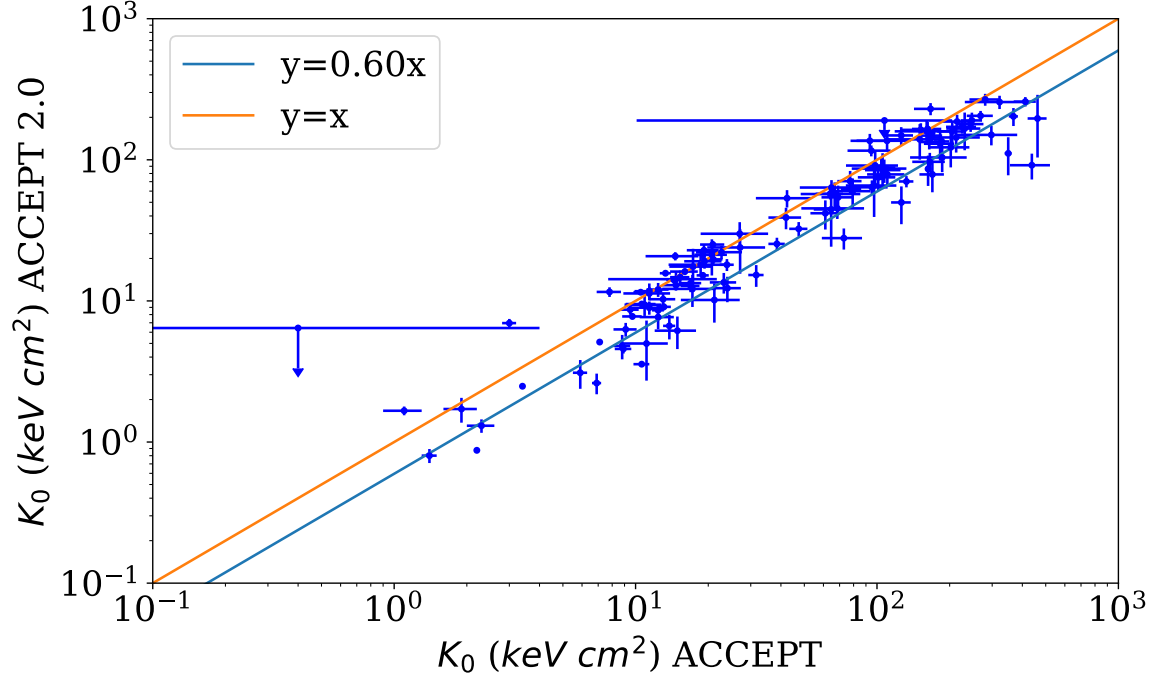


Figure 2.2: **Central Entropy for ACCEPT vs. ACCEPT 2.0** K_0 for ACCEPT vs. K_0 for ACCEPT 2.0 to visualize consistency between the best fit K_0 values that overlap between the two samples. Errors are 1-sigma. The blue line is a linear fit to the data where the slope is the average of the ratio between K_0 for ACCEPT and K_0 for ACCEPT 2.0, weighted by their statistical errors, and the red line plots $y = x$ (the results if the K_0 values were perfectly consistent), for comparison.

construct a robust deprojected entropy profile, and concentration requires far fewer counts, so we want to determine if concentration could be used to sort cluster populations instead of K_0 .

2.2.2.1 Calculating Centroid Shift, Concentration, and Power Ratios

The morphological calculations were completed by Alessandro Baldi in 2015 (concentration) and Megan Donahue in 2018 (centroid shift and power ratio) but the details are included here because of their relevance to my thesis work. The full table of morphological properties can be found in Appendix D.

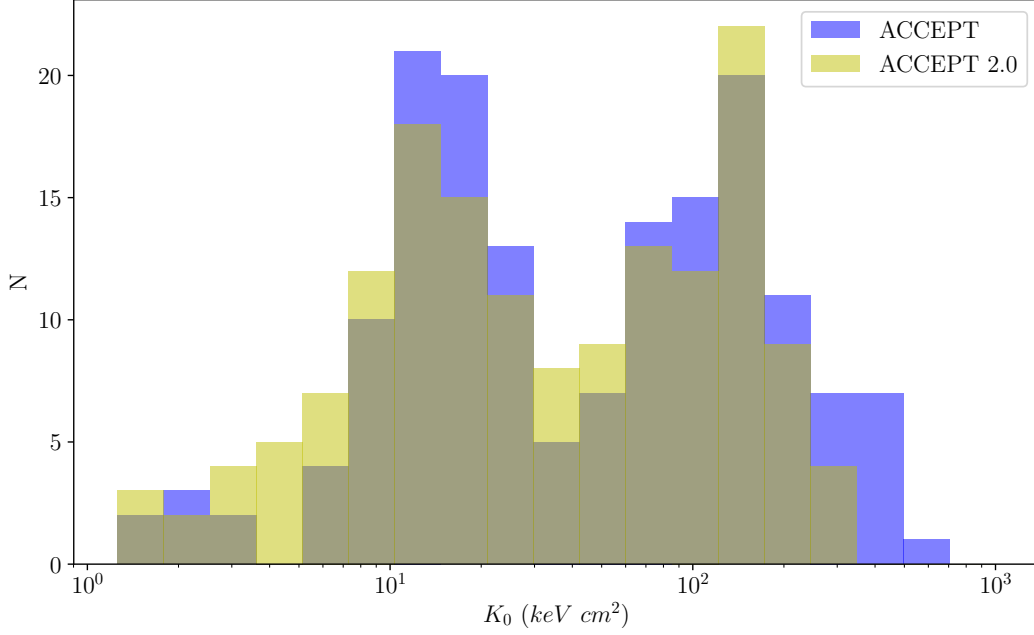


Figure 2.3: **Distribution of Central Entropy for clusters present in both ACCEPT and ACCEPT 2.0** Distribution of K_0 for 164 clusters present in both ACCEPT and ACCEPT 2.0, with statistically significant fits in ACCEPT 2.0. The ACCEPT clusters are in blue, and the ACCEPT 2.0 clusters are in yellow.

2.2.2.1.1 Centroid Shift

The centroid provides a measure of the center of gravity of the cluster and is calculated by:

$$c_x = \frac{\sum x_i f(x_i, y_j)}{\sum f(x_i, y_j)} c_y = \frac{\sum y_j f(x_i, y_j)}{\sum f(x_i, y_j)}, \quad (2.2)$$

where (x_i, y_j) are the coordinates of the pixels, and $f(x_i, y_j)$ are the pixel values at those coordinates.

The dimensionless centroid shift, w , used in this work is based on the definition from Cassano et al. (2010):

$$w = \left[\frac{1}{N-1} \sum (\Delta_i - \langle \Delta \rangle)^2 \right]^{1/2} \frac{1}{R_{max}}, \quad (2.3)$$

where the index i is for each sub-aperture (i runs from 1 to N , and in this case, $N = 20$), Δ_i is the distance between the X-ray peak within R_{max} and the centroid of the i -th aperture, and $\langle \Delta \rangle$ is the average of this separation for all the apertures. For this morphology analysis, $R_{max} = R_{2500}$ from the ACCEPT 2.0 core-excised global temperatures.

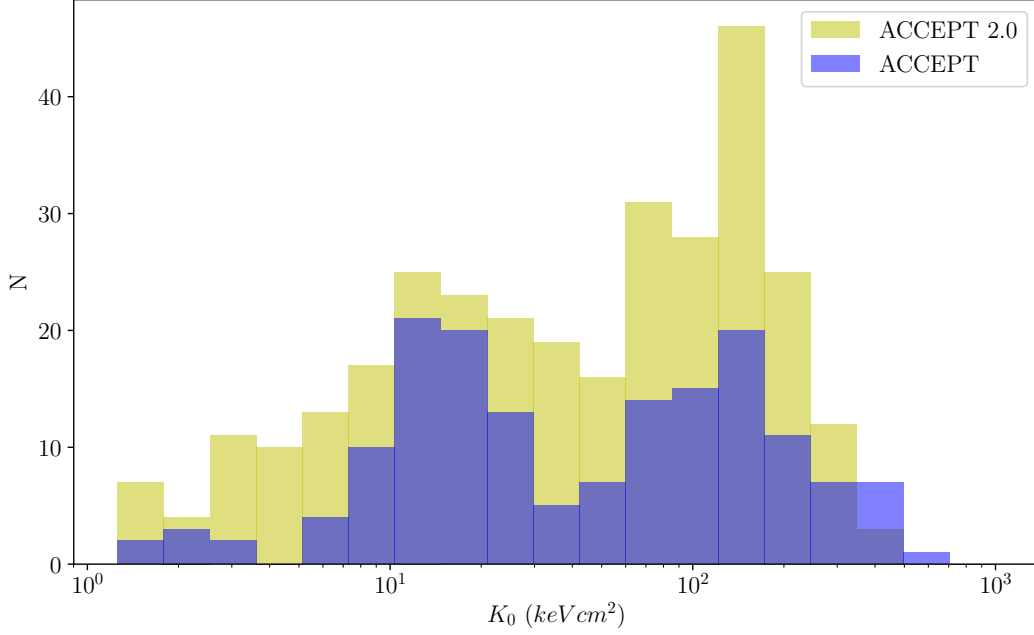


Figure 2.4: **Distribution of Central Entropy for clusters in ACCEPT with all clusters with a measured Central Entropy in ACCEPT 2.0** Distribution of K_0 for 164 clusters in ACCEPT with 348 clusters with a measured K_0 in ACCEPT 2.0. The ACCEPT clusters are in blue, and the ACCEPT 2.0 clusters are in yellow.

The centroid derived from the largest aperture was used for the power ratio estimations. The first moment, P_1 , is not particularly interesting but if it is close to zero, it verifies that the centroid is reasonable. The second moment, P_2 , gives the ellipticity and position angle, and the third moment, P_3 , indicates asymmetries in the surface brightness. Following the treatment in Buote & Tsai (1995), and noting that the surface brightness maps are in units of counts per image pixel (where an image pixel for these maps was a double-binned physical pixel), $S_x(x, y)$, is given as

$$S_{int} = \sum S_x(x^2 + y^2 < R_{max}^2), \quad (2.4)$$

$$P_{1x} = \frac{\sum S_x(x, y)x}{S_{int}}, \quad (2.5)$$

$$P_{1y} = \frac{\sum S_x(x, y)y}{S_{int}}, \quad (2.6)$$

where x, y are defined to be the horizontal and vertical offset from the nominal centroid position.

If the centroid is correct, $P_{1x} \sim P_{1y} \sim 0$. We used this as an internal verification for the computation of centroids. The second moments lead to the computation of the ellipticity and position angle inside 500 kpc or R_{max} as allowed by the field of view, similar to the procedure in Donahue et al. (2014). We computed 3 terms,

$$A_{xx,yy,xy} = \frac{\sum S_x(x, y) \times (xx, yy, xy)}{S_{int}}, \quad (2.7)$$

and then diagonalized the matrix,

$$\begin{bmatrix} A_{xx} & A_{xy} \\ A_{xy} & A_{yy} \end{bmatrix} \quad (2.8)$$

to obtain the second moment. The morphological properties for ACCEPT 2.0 can be found in Appendix D.

2.2.2.1.2 Concentration

Two surface brightness concentration parameters, $c_{500 \text{ kpc}}$ and c_{R500} , are computed to measure the concentration of the X-ray-emission in the ACCEPT 2.0 clusters. Concentration has been defined in various ways in the literature. In this work, we adopt a common convention, that concentration is a ratio, ranging from 0.0 to 1.0, between the total X-ray luminosity interior to an inner radius, r_{inner} , to the total luminosity inside a larger radius, r_{outer} . If all the detected flux is inside both radii, then the concentration is close to 1.0, and if the source were somehow shaped like a donut, with an empty center, the concentration would be zero. We have defined two interpretations of concentration: $c_{500 \text{ kpc}}$, where $r_{inner} = 100 \text{ kpc}$ and $r_{outer} = 500 \text{ kpc}$ (eg, Cassano et al. (2010)), and c_{R500} , where $r_{inner} = 0.1r_{500}$ and $r_{outer} = 0.5r_{500}$ (as in Rasia et al. (2012)).

2.2.2.1.3 Power Ratios

Power ratios (Buote & Tsai, 1995), mimic a multiple decomposition of the 2-D projected mass distribution inside a certain aperture, R_{ap} (R_{2500} for this analysis), but it is much simpler to apply this decomposition to the X-ray surface brightness images S , instead of the mass. The m -th order

power ratio ($m > 0$) is defined as P_m/P_0 with

$$P_m = \frac{1}{2m^2 R_{ap}^{2m}} (a_m^2 + b_m^2); P_0 = a_0 \ln(R_{ap}), \quad (2.9)$$

where a_0 is the total intensity within the aperture radius, and a_m and b_m are expressed in polar coordinates (R and ϕ) and given by

$$a_m(r) = \int_{R' \leq R_{ap}} S(x')(R')^m \cos(m\phi') d^2 x', \quad (2.10)$$

and

$$b_m(r) = \int_{R' \leq R_{ap}} S(x')(R')^m \sin(m\phi') d^2 x'. \quad (2.11)$$

The power ratio P_2/P_0 gives information about the cluster ellipticity and P_3/P_0 is an indicator of bimodal distribution in the surface brightness and therefore is the most sensitive to detecting asymmetries or substructures. P_4/P_0 is similar to P_2/P_0 but more sensitive to smaller scales. The ACCEPT 2.0 pipeline computes all power ratios P_m/P_0 with $1 \leq m \leq 6$, but we will focus on P_3/P_0 in this work because of its sensitivity to substructure and because it is less noisy than higher order moments.

2.3 Science with ACCEPT 2.0

The primary goal of ACCEPT 2.0 was to provide a uniformly-reduced database of galaxy clusters with as many X-ray observable properties as possible, given data reduction constraints. While there are multiphase gas measurements for the clusters of ACCEPT, the sample in ACCEPT 2.0 can be used to further examine the presence or absence of multiphase gas and its correlation with central entropy. The selection function of ACCEPT 2.0 is quite complicated because it is a purely archival sample and therefore not only holds potential bias in mass, luminosity, and other observables, but also in the selection of targets themselves. Future work could attempt to characterize the selection function to answer scientific questions such as how common certain types of clusters are and cosmological questions about the cluster mass function using the largest possible sample. However, because ACCEPT 2.0 is a large, uniformly reduced sample, it is currently well suited for drawing well-defined sub-samples. As an example of this type of work, I have used

ACCEPT 2.0 data for an initial exploration of the sample of galaxy clusters observed through the XMM Heritage project. The science uses for ACCEPT 2.0 discussed in this thesis represent just a few of the countless projects that could use ACCEPT 2.0 to address key questions in cosmology and cluster and galaxy evolution.

2.3.1 Morphological Properties and K_0 for Sample Comparisons

Broadly, morphological properties describe what the galaxy clusters “look” like, including how the X-ray gas is distributed (see Section 2.2.2.1 for details). Cassano et al. (2010) showed that, for a small (32 clusters) sample, the power ratio P_3/P_0 is correlated with the centroid shift, w , and the concentration, c ; and concentration is correlated with centroid shift. Here, we will examine the morphological properties for the ACCEPT 2.0 sample and two sub-samples; ROSAT-ESO Flux-Limited X-ray (REFLEX) Galaxy Cluster Survey (Böhringer et al., 2004) and XMM Heritage.

In Figure 2.5, we show the results for correlation between P_3/P_0 , w , and c_{500} for the entire ACCEPT 2.0 sample. As in Cassano et al. (2010), we see correlation between P_3/P_0 and w , but the correlation between P_3/P_0 and c_{500} is less clear. Overall, the clusters exhibit a correlation between centroid shift and power ratio, and there are fewer clusters in the less relaxed, less symmetric parameter space.

In Figure 2.6, we plot the central entropy fit values and the concentration parameter. The hope is that concentration could serve as a low-signal proxy for central entropy because we expect low entropy clusters to be highly concentrated and high entropy clusters to not be as concentrated. That is, if we know the concentration for a galaxy cluster with data quality insufficient to get a deprojected entropy profile, we could make a prediction about the central entropy. While there is weak correlation between c_{500} and K_0 , it is not strong enough to provide strong predictions for a K_0 measurement based on measured concentration. One weakness of c_{500} is that it is based on r_{500} , which may be too large a radius to be captured in the *Chandra* field of view.

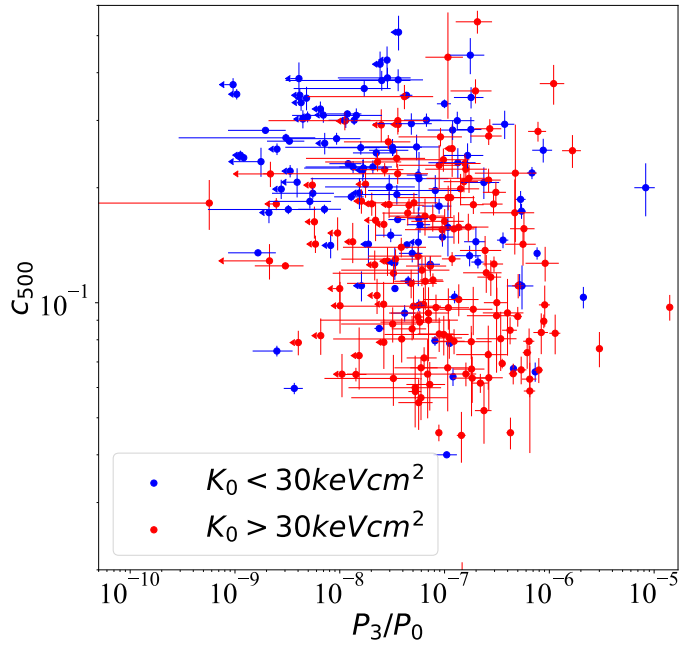
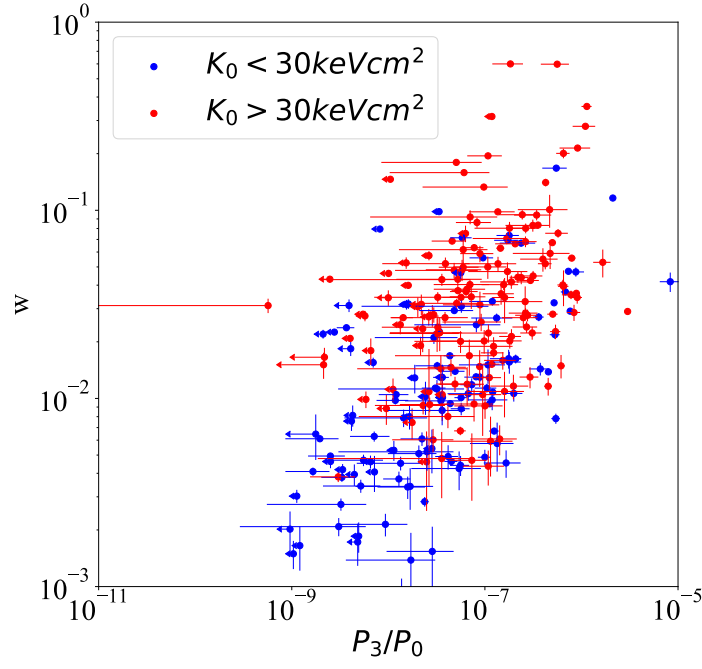


Figure 2.5: **Morphological Property Comparison for ACCEPT 2.0** Power ratios compared to centroid shift, w , and concentration, c_{500} , colored by central entropy. Red points are clusters with $K_0 > 30 \text{ keV cm}^2$ and blue points are clusters with $K_0 < 30 \text{ keV cm}^2$ for all of ACCEPT 2.0.

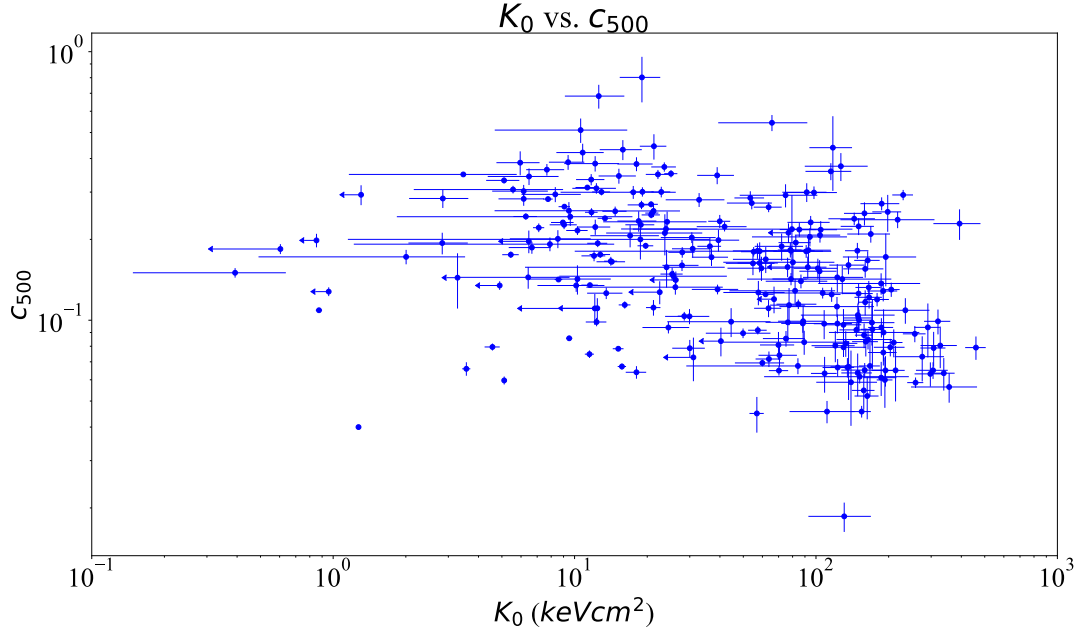


Figure 2.6: **Central Entropy vs. concentration for ACCEPT 2.0** Central entropy, K_0 , is plotted with c_{500} for clusters in ACCEPT 2.0 to illustrate the potential for concentration to serve as a noisy estimate of K_0 .

2.3.2 ACCEPT 2.0 Comparison with REFLEX and XMM Heritage

The XMM Heritage project seeks to create a signal to noise limited sample of deep observations from XMM of 118 clusters from the Planck PSZ2 cosmological catalog. The targets include a local sample at $z < 0.2$ with a mass range of $10^{14} M_{\odot} < M_{500} < 9 \times 10^9 M_{\odot}$ and the most massive ($M_{500} > 10^{14} M_{\odot}$) clusters at $z < 0.6$. The ACCEPT 2.0 is an archival sample of galaxy clusters observed by Chandra and is an expansion and analysis improvement of the ACCEPT project (Cavagnolo et al., 2009). Here we combine our data products for clusters present in both samples to gain insight into the properties of the X-ray gas in the XMM Heritage sample.

Bringing together the morphological properties and central entropies with the common clusters between the two samples, we see that they exhibit a correlation between centroid shift and power ratio, and the lower entropy clusters are not present at the less relaxed, less symmetric parameter space (Figure 2.9). Shown in a different way, we also find that as in Bauer et al. (2005), cool core clusters are more likely to be compact and symmetric.

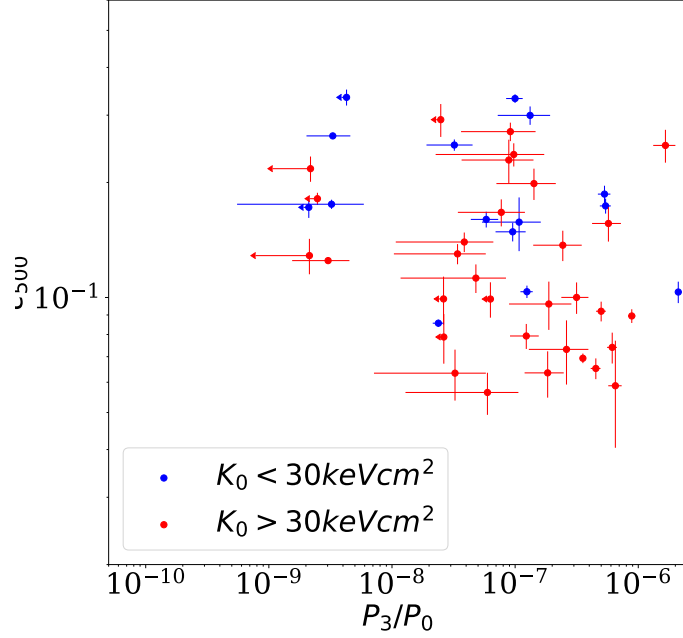
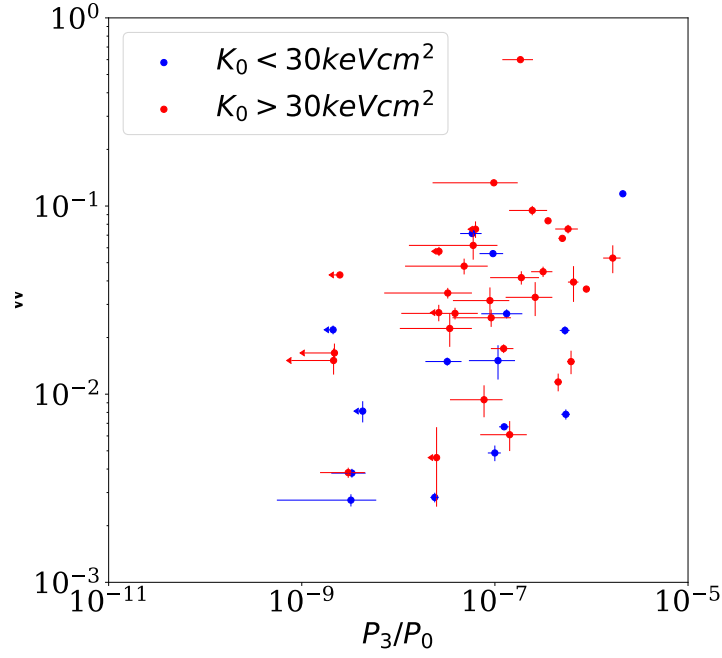


Figure 2.7: **Morphological Property Comparison for the REFLEX sub-sample of ACCEPT 2.0** Power ratios compared to centroid shift, w , and concentration, c_{500} , colored by central entropy. Red points are clusters with $K_0 > 30 \text{ keV cm}^2$ and blue points are clusters with $K_0 < 30 \text{ keV cm}^2$ for the REFLEX sub-sample of ACCEPT 2.0.

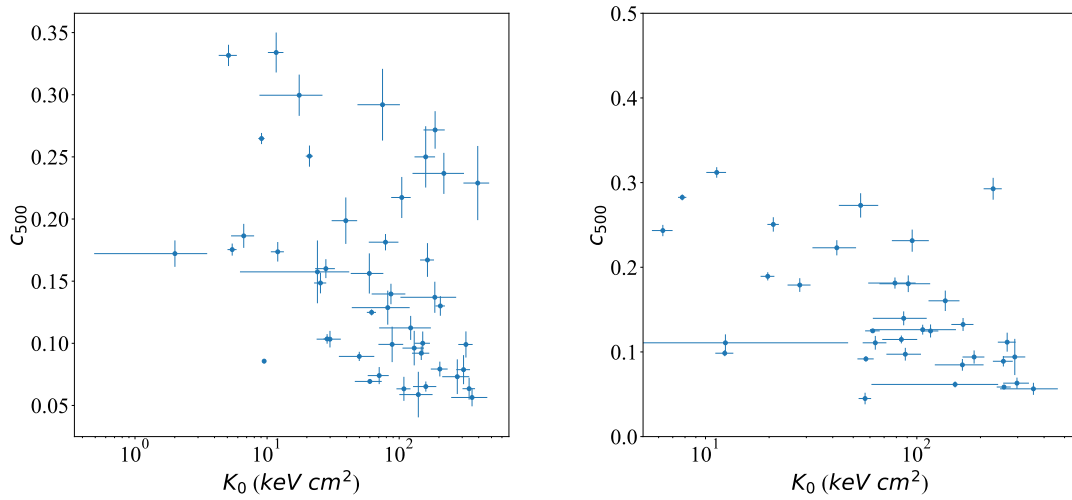


Figure 2.8: **Central Entropy vs. concentration for XMM Heritage and REFLEX sub-samples of ACCEPT 2.0** Central entropy, K_0 , is plotted with c_{500} for clusters in ACCEPT 2.0 and REFLEX (left) and ACCEPT 2.0 and XMM Heritage (right).

Figure 2.7 shows the same plots as in Figure 2.9 but made for the common clusters between ACCEPT 2.0 and the REFLEX sample (Böhringer et al., 2004). Between the two sub-samples of ACCEPT 2.0, we see small differences in the plots for selection between X-ray flux selected (REFLEX) and high X-ray pressure selected (XMM Heritage/SZ) clusters. The pressure selected sample that overlaps with ACCEPT 2.0 contains fewer cool core clusters and fewer symmetric, relaxed clusters than the flux selected sample overlapping with ACCEPT 2.0.

This work represents one of the early tests of the broad applications of the ACCEPT 2.0 project. As more data for the XMM Heritage project are taken, we will be able to use the insights learned from ACCEPT 2.0 to learn more about the thermal properties of the X-ray gas in these clusters. We are also able to identify and confirm Brightest Cluster Galaxy (BCG) coordinates for the overlapping targets. Through our comparison of the XMM Heritage-ACCEPT 2.0 and REFLEX-ACCEPT 2.0 samples, we can begin to see the X-ray selection effects for samples of galaxy clusters as well as the value of archival data. The scientific applications for both ACCEPT 2.0 and the XMM Heritage projects are widespread and extensive.

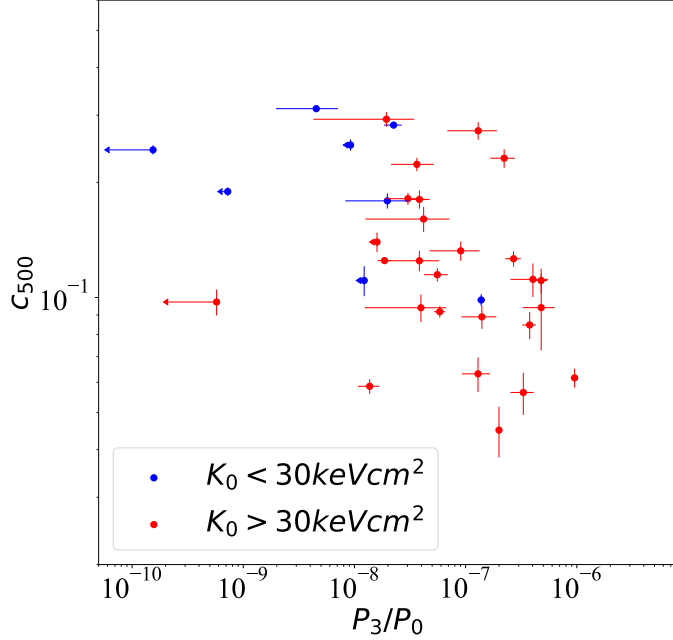
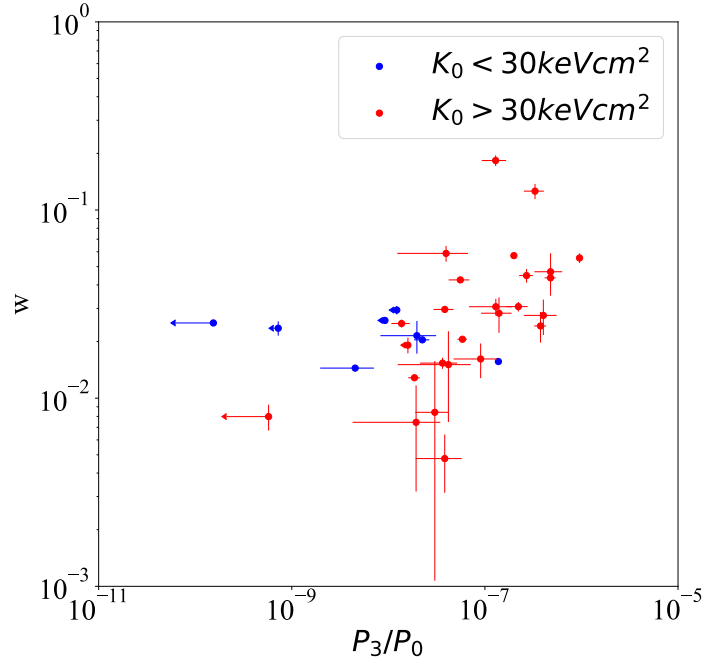


Figure 2.9: **Morphological Property comparison for the XMM Heritage sub-sample of ACCEPT 2.0** Power ratios compared to centroid shift, w , and concentration, c_{500} , colored by central entropy. Red points are clusters with $K_0 > 30 \text{ keV cm}^2$ and blue points are clusters with $K_0 < 30 \text{ keV cm}^2$ for the XMM Heritage sub-sample of ACCEPT 2.0.

2.4 Summary

ACCEPT 2.0 expands and improves upon the work done by Cavagnolo et al. (2009). In this chapter, I have presented the initial pipeline verification via comparison to ACCEPT and a few early science applications of the ACCEPT 2.0 data. Initial verification was accomplished by comparing the deprojected temperature, density, and entropy profiles for the common clusters between ACCEPT and ACCEPT 2.0 as well as comparing the overall sample characteristics of the profiles and central entropy between ACCEPT and all ACCEPT 2.0 clusters. As in Cavagnolo et al. (2009), I found that the distribution of central entropy was bimodal, with a break at 30 keV cm^2 and peaks at $\sim 15 \text{ keV cm}^2$ and $\sim 150 \text{ keV cm}^2$. The distribution of central entropy suggests that almost all galaxy clusters have a core excess of entropy, and that central entropy can approximately divide galaxy clusters into cool core and non-cool core clusters, a classification scheme that proves to be useful when examining cluster properties for large samples.

We applied the classification via central entropy to the XMM Heritage and comparison samples to understand some of the selection characteristics, and we found that XMM Heritage has more non-cool core clusters in the less relaxed, less symmetric parameters space compared to cool core clusters. The distribution of clusters in the XMM Heritage sample by central entropy and morphology may be the result of the sample being selected by X-ray pressure rather than X-ray luminosity. Pressure selected samples are selected based on measurements of the Sunyaev-Zeldovich effect which is more significant for less-relaxed clusters, like those that have undergone recent mergers. We do find that in the ACCEPT 2.0 sample, and more strongly in the XMM Heritage sample, the power ratio and centroid shift do exhibit some correlation, supporting the work of Cassano et al. (2010) on a smaller sample.

The ACCEPT 2.0 pipeline provides a large, uniformly reduced sample of galaxy cluster properties, and the measurements largely support claims about the distribution of galaxy cluster parameters from previous works. Furthermore, ACCEPT 2.0 is an invaluable resource for comparisons to other X-ray samples, particularly for understanding sample selection characteristics. We expect that AC-

CEPT 2.0 will provide a springboard for galaxy cluster studies for years to come.

CHAPTER 3

PROPERTIES OF THE CGM IN EARLY-TYPE GALAXIES WITH POWERFUL RADIO SOURCES

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3.1 Abstract

We present an archival analysis of *Chandra* X-ray observations for 12 nearby early-type galaxies hosting radio sources with radio power $> 10^{23} \text{ W Hz}^{-1}$ at 1.4 GHz, similar to the radio power of the radio source in NGC 4261. Previously, in a similar analysis of eight nearby X-ray and optically bright elliptical galaxies, Werner et al. (2012), found that NGC 4261 exhibited unusually low central gas entropy compared to the full sample. In the central 0.3 kpc of NGC 4261, the ratio of cooling time to freefall time ($t_{\text{cool}}/t_{\text{ff}}$) is less than 10, indicating that cold clouds may be precipitating out of the hot ambient medium and providing fuel for accretion in the central region. NGC 4261 also hosts the most powerful radio source in the original sample. Because NGC 4261 may represent an important phase during which powerful feedback from a central active galactic nucleus (AGN) is fueled by multiphase condensation in the central kiloparsec, we searched the *Chandra* archive for analogs to NGC 4261. We present entropy profiles of those galaxies as well as profiles of $t_{\text{cool}}/t_{\text{ff}}$. We find that one of them, IC 4296, exhibits properties similar to NGC 4261, including the presence of only single-phase gas outside of $r \sim 2$ kpc and a similar central velocity dispersion. We compare the properties of NGC 4261 and IC 4296 to hydrodynamic simulations of AGN feedback fueled by precipitation. Over the course of those simulations, the single-phase galaxy has an entropy gradient that remains similar to the entropy profiles inferred from our observations.

3.2 Introduction

Over the past two decades, *Chandra* has been used to observe the ambient medium of early-type galaxies because of its high sensitivity in the soft X-ray band (0.5-2.0 keV) and its spatial resolution,

resulting in 2D spectroscopy of unprecedented quality (e.g. Kim et al. 2018; Diehl & Statler 2007, 2008a,b; Lakhchaura et al. 2018; Sun 2009). The hot atmospheres of those early-type galaxies have provided key clues about the energetic processes known as “feedback” (McNamara & Nulsen, 2012; Soker, 2016; Fabian, 2012). X-ray signatures of feedback processes observed in the hot atmospheres of nearby, early-type galaxies are also commonly and prominently observed in the hot atmospheres of Brightest Cluster Galaxies (BCG), the brightest and most massive galaxies in galaxy clusters. The supermassive black holes at the center of BCGs in clusters interact with the surrounding medium, inflating bubbles of relativistic plasma (e.g. Boehringer et al. 1993; Churazov et al. 2000; Fabian et al. 2003, 2006; Bîrzan et al. 2004; Dunn & Fabian 2006, 2008; Dunn et al. 2005; Forman et al. 2005, 2007; Rafferty et al. 2006; McNamara & Nulsen 2007). One insight from studying feedback processes in galaxy clusters is that the activity state of the central Active Galactic Nucleus (AGN) in a BCG is closely coupled to the thermodynamic state of the Intracluster Medium (ICM) (e.g. Cavagnolo et al., 2008; Rafferty et al., 2008; Voit & Donahue, 2015; Voit et al., 2015a). However, in individual early-type galaxies in groups, like those we discuss in this work, the relationship may be a little more complex (Sun, 2009; Connor et al., 2014). Because the gravitational potential depths are shallower for galaxies in groups than galaxies in clusters, supernova explosions and galactic winds are energetically more important for galaxies than for galaxy clusters. Furthermore, while nearly any reasonable amount of kinetic AGN output can be contained in a cluster atmosphere, the question of whether or not a powerful AGN jet thermalizes its energy output near or far from the AGN depends on the external gas pressure. In turn, the external gas pressure may depend on the large-scale structure the galaxy inhabits.

McNamara & Nulsen (2007, 2012) have summarized the evidence suggesting that black holes suppress the star formation in massive galaxies, but how the accretion onto the black hole is affected by the surrounding hot gas is less clear. Precipitation-regulated feedback models hypothesize that feedback suspends the ambient medium in a state that is marginally stable to multiphase condensation. Feedback input affects the thermodynamic state and susceptibility of the ambient gas to condensation. Feedback output depends sensitively on the rate at which cold clouds precipitate

out of the hot medium (Pizzolato & Soker, 2005; Sharma et al., 2012; Gaspari et al., 2012, 2013, 2015, 2017; Voit et al., 2015b, 2017; Wang et al., 2019). Such a system is self-regulated, and finds a balance at the marginally stable point.

Spatially resolved X-ray spectroscopy of the hot ambient medium provides insight into its thermal evolution. The normalization and shape of an X-ray spectrum yields gas electron density (n_e), temperature (T_X), and metallicity (Z). Broadly, for early-type galaxies, the temperature of the hot gas is ~ 1 keV with a nearly isothermal radial profile, and the radial profile of the electron density approximately follows a power law. The temperature and density of the X-ray gas, considered independently, do not reveal the thermal history because heating and cooling of gravitationally confined gas can cause it to expand or contract without much change in temperature. However, combining these two X-ray observables to make the quantity $K = kT_X n_e^{-2/3}$ provides us with more direct information about thermal history, because changes in $kT_X n_e^{-2/3}$ correspond directly to changes in the specific entropy of the gas. Only gains and losses of heat energy in the gas can change the entropy, so we can trace the thermal history of the ambient gas of a galaxy cluster by observing the profile $K(r)$, which we will call an entropy profile.

In addition to what we have learned from X-ray observations, numerical simulations show that cool clouds can precipitate out of a galaxy's hot gas atmosphere via thermal instability even if the galaxy is in a state of global thermal balance, with heating approximately equal to cooling (Gaspari et al., 2012; McCourt et al., 2012; Sharma et al., 2012). The critical criterion for precipitation is the ratio between the cooling and freefall times of the gas. Here, the cooling time (t_{cool}) is defined to be the time needed for a gas at temperature T to radiate an energy $3kT/2$ per particle, and the free fall time from a galactocentric radius r at the local gravitational acceleration g is defined to be $t_{\text{ff}} = (2r/g)^{1/2}$. We note that these models do not presume to claim that the gas must be freely falling. The parameter t_{ff} merely specifies a useful dynamical timescale that characterizes gravitationally driven motions. The freefall time does not assume anything about the turbulence, viscosity, or other fluid properties and is based on galaxy properties that can be inferred from observations of the stellar light.

In both observations and in simulations (McCourt et al., 2012; Sharma et al., 2012; Gaspari et al., 2012; Li & Bryan, 2014a), cooling appears to be fast enough for a fraction of the hot gas to condense into cold clouds and precipitate out of the hot medium if $t_{\text{cool}}/t_{\text{ff}} \sim 10$. Precipitation may therefore play an essential role in maintaining the required state of global thermal balance if gas cooled from the hot phase boosts the fuel supply for accretion (Pizzolato & Soker, 2010; Gaspari et al., 2013, 2015; Li & Bryan, 2014a,b). In numerical simulations, accretion of precipitating clouds can produce a black hole fueling rate two orders of magnitude greater than the Bondi accretion rate of ambient gas. Such strong accretion then produces a feedback response that heats the gas, bringing the system into approximate balance near $t_{\text{cool}}/t_{\text{ff}} \approx 10$. Voit et al. (2015b) showed that early-type galaxies do indeed have $\min(t_{\text{cool}}/t_{\text{ff}}) \approx 10$.

The hot atmosphere of an early-type galaxy can be broadly categorized as single-phase gas or multiphase gas, depending on the extent of the $\text{H}\alpha$ and $[\text{N II}]$ emission. Observationally, galaxies with multiphase atmospheres have extended $\text{H}\alpha$ and $[\text{N II}]$ emission present outside their centers (central ~ 1 kpc), whereas galaxies with single-phase atmospheres have no evidence for extended $\text{H}\alpha$ emission outside of ~ 2 kpc. X-ray observations of giant ellipticals from Werner et al. (2012, 2014) showed that single- and multiphase galaxies are distinctly bimodal from 1–10 kpc. The entropy profiles of single-phase galaxies scale as $K \propto r$, while in multiphase galaxies the entropy scales as $K \propto r^{2/3}$. However, both types exhibit excess entropy in the innermost kiloparsec equivalent to $\sim 2 \text{ keV cm}^2$.

While Werner et al. (2012, 2014) showed that both single- and multiphase galaxies tend to have entropy excesses relative to a power law in the central kiloparsec, one galaxy differed from the rest. X-ray observations of NGC 4261 from Werner et al. (2012) revealed that the entropy profile of NGC 4261 follows a single power law ($K \propto r$), but instead of exhibiting an excess within the central kpc, the power law continues into the central ~ 0.5 kpc ($\sim 4''$). The unusually low entropy in the center ($K \approx 0.8 \text{ keV cm}^2$) results in $t_{\text{cool}}/t_{\text{ff}} < 10$, putting it slightly below the limit at which precipitation appears inevitable. NGC 4261’s radio luminosity is 2 orders of magnitude greater than the rest of the Werner et al. (2012) sample, and the central jet power is $10^{44} \text{ erg s}^{-1}$. Adopting

a central black hole mass of $M_{\text{BH}} = 5 \times 10^8 M_{\odot}$ (Gaspari et al., 2013) would require an implausible 30% mass energy to jet power conversion efficiency for the radio source to be powered by Bondi accretion alone (Voit et al., 2015b). Simulations from Gaspari et al. (2013, 2015) showed that a transition to chaotic cold accretion could boost the jet power by up to 100 times over what Bondi accretion of hot ambient gas could achieve and occurs when $t_{\text{cool}}/t_{\text{ff}} \approx 10$.

Because this transitional regime has not been extensively investigated, we decided to explore it by looking for other galaxies like NGC 4261. To that end, we analyzed an archival sample of *Chandra* observations of 12 additional early-type galaxies with powerful radio sources. In this paper, we present a summary of our findings for this archival study, which yielded at least one additional nearby analog, IC 4296, that similarly has both a steep entropy profile with $t_{\text{cool}}/t_{\text{ff}} < 10$ at small radii and a powerful radio source.

The structure of our paper is as follows. Section 3.3 describes our sample selection, data analysis, and our measurements of the thermodynamic properties. Section 3.4 presents a comparison of $t_{\text{cool}}/t_{\text{ff}}$ profiles to previous works, a comparison with simulations, and an analysis of the effects of metallicity assumptions on our measurements. Section 3.5 concludes by discussing how our sample adds to the paradigm of precipitation-regulated feedback in massive galaxies. We assume a Λ CDM cosmology with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ and $\Omega_M = 0.3$ ($\Omega_{\Lambda} = 0.7$) throughout.

3.3 Sample Selection and Data Analysis

3.3.1 Sample Selection and Distances

NGC 4261 exhibits an unusually low central entropy, as well as $t_{\text{cool}}/t_{\text{ff}} < 10$ at $r < 0.3 \text{ kpc}$ (Voit et al., 2015b). It also has a powerful radio source emitting $2.3 \times 10^{24} \text{ W Hz}^{-1}$ in the 1.4 GHz band, which may be powered by chaotic cold accretion onto the central supermassive black hole (Gaspari et al., 2012, 2013, 2015) fed by precipitation of cold clouds out of the hot atmosphere (Voit et al., 2015b).

In search of other systems similar to NGC 4261, we compiled a *Chandra* archival sample of nearby ($z < 0.02$) massive early-type galaxies hosting radio sources with a power output

Table 3.1: **Chandra Observations of Early-type Galaxies** Column (1): galaxy name; Column (2): redshift obtained from NED^a; Column (3): distance calculated from z_{spec} with the exception of NGC 4374 and NGC 7626, for which we use redshift-independent distances from Tonry et al. (2001) and Cantiello et al. (2007), respectively. Column (4): galactic neutral hydrogen column densities from Kalberla et al. (2005) and HI4PI Collaboration et al. (2016); Column (5): radio fluxes from VLA or NVSS (Condon et al., 1998) except for NGC 4261 (PKS, Brown et al. 2011); Column (9) whether there were sufficient counts to make deprojected temperature and density profiles for a galaxy. Column (10): power-law entropy slope α determined by fitting the relation $K \propto r^\alpha$ in the 1–10 kpc interval; Column (11): central velocity dispersion from Makarov et al. (2014); Column (12): $H\alpha$ + $[\text{N } 2]$ morphology reported by Lakhchaura et al. (2018) from Connor (in preparation) and Sun (in preparation), classified as follows: N: no cool gas emission; NE: $H\alpha$ + $[\text{N } 2]$ extent < 2 kpc; E: $H\alpha$ + $[\text{N } 2]$ extent $\geq$ 2 kpc; U: galaxies for which the presence/absence of $H\alpha$ + $[\text{N } 2]$ could not be confirmed with current observations.

Galaxy	z_{spec}	D (Mpc)	$N_{\text{H,HI}}$ (10^{20} cm^{-2})	1.4 GHz ($10^{24} \text{ W Hz}^{-1}$)	ObsID	Exp (ks)	Net Counts per Bin	Profile (Y/N)	α (1-10 kpc)	σ_v (km s^{-1})	$H\alpha$ + $[\text{N } 2]$ Morph ^d
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)
NGC 193	0.015	63.08	2.46	0.468	11389	93.13	300	N	--	197.6	E
NGC 193	0.015	63.08	2.46	0.468	11389	93.13	300	N	-	197.6 $\pm$ 4.8	E
NGC 315	0.016	67.20	5.88	0.973	4156	53.84	930	N	-	293.6 $\pm$ 10.1	U
NGC 741	0.018	75.42	4.24	0.327	17198	91.02	1500	Y	-	287.4 $\pm$ 9.3	N
NGC 1316	0.006	25.51	1.99	9.75	2022	29.86	450	Y	0.80 $\pm$ 0.08	223.1 $\pm$ 3.3	E
IC 1459	0.006	25.51	0.94	0.1	2196	58.00	300	Y	-	296.1 $\pm$ 6.4	E
NGC 3801	0.011	46.48	1.99	0.296	6843	59.20	330	N	-	191.8 $\pm$ 16.6	-
NGC 3894	0.011	46.48	1.83	0.125	10389	38.54	300	N	-	252.8 $\pm$ 11.3	-
NGC 4261	0.007	31.32	1.61	2.58	9569	100.34	1600	Y	1.09 $\pm$ 0.07	296.7 $\pm$ 4.3	NE
IC 4296	0.012	50.64	3.95	5.52	3394	24.84	800	Y	1.12 $\pm$ 0.12	327.4 $\pm$ 5.4	NE ^b
NGC 4374	0.003	18.37	2.90	0.125	803	28.46	650	Y	0.75 $\pm$ 0.05	277.6 $\pm$ 2.4	NE
NGC 4782	0.015	63.08	3.10	3.33	3220	49.33	320	Y	-	310.0 $\pm$ 11.3	NE
NGC 5419	0.014	58.94	5.40	0.146	5000	14.81	320	Y	-	344.3 $\pm$ 5.4	N
NGC 7626	0.011	58.34	4.59	0.222	2074	26.54	370	Y	-	266.6 $\pm$ 3.7	E

^aThe NASA/IPAC Extragalactic Database (NED) is funded by the National Aeronautics and Space Administration and operated by the California Institute of Technology.

^bIC 4296 was identified as E in Lakhchaura et al. (2018). Since its multiphase gas is at < 2 kpc, we classify it here as NE.

^cBased on observations obtained at the Southern Astrophysical Research (SOAR) telescope, which is a joint project of the Ministério da Ciência, Tecnologia, Inovações e Comunicações (MCTIC) do Brasil, the U.S. National Optical Astronomy Observatory (NOAO), the University of North Carolina at Chapel Hill (UNC), and Michigan State University (MSU)

^dBased on observations obtained with the Apache Point Observatory 3.5-meter telescope, which is owned and operated by the Astrophysical Research Consortium.

$> 10^{23} \text{ W Hz}^{-1}$ at 1.4 GHz (Condon et al., 1998, see Table 3.1). Other recent studies of *Chandra* observations of early-type galaxies (e.g. Lakhchaura et al., 2018; Grossová et al., 2019; Juráňová et al., 2019) include some of the same galaxies, but our sample emphasizes powerful radio sources in order to identify galaxies similar to NGC 4261.

In our analysis, we used distances derived from the redshifts of the galaxies when calculating the electron density, because the effect of small uncertainties in distance on the inferred density, entropy, and cooling time is small. However, for NGC 4374 and NGC 7626 we used the redshift-independent measurements because the differences between the best redshift-independent distance measurements and the redshift-dependent distances are large (20 – 30%; see Table 3.1).

In this work, we pay particular attention to NGC 4374 (M84), NGC 1316 (Fornax A), and IC 4296 because *Chandra* observations of the central 10 kpc of those galaxies have the best signal-to-noise ratios among those in our sample, and we are most interested in atmospheric properties closest to the center. Each galaxy represents a different manifestation of a powerful radio source. M84 hosts an FR I radio jet¹ (Harris et al., 2002). Fornax A has a weak core in the radio (250 mJy), but its radio lobes are some of the brightest radio sources in the sky (125,000 mJy, Ekers et al., 1983). IC 4296 is the Brightest Group Galaxy (BGG) in a nearby galaxy group (Abell 3565), and *Hubble Space Telescope* (*HST*) spectroscopy indicates a central black hole mass of $\sim 10^9 M_{\odot}$ (Dalla Bontà et al., 2009). Recent *Very Large Array* (*VLA*) D-configuration observations show improved mapping of the 160 kpc diameter radio lobes, first discovered by Killeen et al. (1986), located over 230 kpc from the AGN host galaxy, as well as *X-ray Multi-Mirror Mission* (*XMM*) observations that reveal a corresponding X-ray cavity (Grossová et al., 2019).

3.3.2 *Chandra* data reduction

All of the data used in this work are archival *Chandra* data taken between 2000 May and 2015 December. All observations were taken with the ACIS-S detector, except for NGC 5419 and NGC 7626, which were obtained with the ACIS-I detector. We reprocessed the archival *Chandra*

¹Defined as a radio source in which the low-brightness regions of the jet are farther from the galaxy than the high-brightness regions (Fanaroff & Riley, 1974)

data listed in Table 3.1 using CIAO 4.9 and CALDB version 4.7.4. For simplicity, in the case of targets with multiple observations, we chose to analyze the one with the longest net exposure time.

The time intervals containing data with anomalously high background were identified and removed using the `deflare` script in CIAO. Bright point sources were identified and removed using the `wavdetect` script (Freeman et al., 2002). We opted to account for the effect of central point sources in our spatially resolved spectral analysis. Background images and spectra were derived from the blank-sky fields available from the *Chandra* X-ray Center. The background files contain both particle and photon backgrounds and were filtered and reprojected to match the target observations. We rescaled the reprojected background rates to match the particle count rates, gauged from the event rate between 10.0 and 12.0 keV (Hickox & Markevitch, 2006). Because our analyses are based on regions of the galaxy where the signal is much higher than the background, our results are insensitive to the details of the background scaling.

3.3.3 Spectral Analysis

We derived deprojected radial profiles of the X-ray gas properties: temperature, density, and gas entropy. To prepare the spectra, we defined radial annuli each containing at least 300 counts after background subtraction (at temperatures around 0.7 – 1 keV, a minimum of ~ 300 counts between 0.5 and 7 keV are required for a robust X-ray temperature estimate). We used the definitions of these radial bins to extract radially binned X-ray event spectra for each galaxy and background spectrum from the scaled and reprojected deep background data.

For each galaxy, we fit all radial bins simultaneously with XSPEC v.12.9 (Arnaud, 1996) using the `projct` model together with the X-ray thermal emission model `apec` and Galactic absorption column model `phabs`. Because the spectral band above 2 keV is more likely to be dominated by emission from X-ray background and unresolved point sources in typical X-ray spectra of early-type galaxies, we restricted the energy range for the spectral fits to 0.6 – 2.0 keV.

For each galaxy, the Galactic column density and redshift were fixed to the values in Table 3.1, and the gas metallicity was fixed at a solar abundance. We will discuss the impact of this abundance

Table 3.2: **Sample of Radial Profile Properties** A portion of this table is printed here for form and content, Additional profiles can be found in Appendix E. Errors given for radius represent bin widths; all other errors are 1σ . Column (1): galaxy name. Column (2): radial bin center. Column (3): half-width of the radial bin. Column (4): grouping of temperature bins. Columns (5)-(6): best-fit temperatures and their errors. Column (7): electron density bin number. Columns (8)-(9): best-fit densities and their errors. in units of 10^{-2} cm^{-3} for compactness. Columns (10)-(11): calculated entropies and their errors.

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
IC 4296	0.48	0.24	1	0.75	0.02	1	16.50	0.57	2.48	0.09
IC 4296	0.72	0.12	1	0.75	0.02	2	10.10	0.53	3.44	0.16
IC 4296	0.97	0.12	2	0.78	0.03	3	5.62	0.39	5.29	0.33
IC 4296	1.45	0.24	2	0.78	0.03	4	3.52	0.20	7.23	0.41
IC 4296	1.93	0.24	3	0.84	0.03	5	2.50	0.18	9.82	0.63
IC 4296	2.66	0.36	3	0.84	0.03	6	1.19	0.06	16.08	0.86
IC 4296	3.87	0.60	4	0.89	0.05	7	0.49	0.03	30.82	1.98
IC 4296	6.28	1.21	4	0.89	0.05	8	0.38	0.03	36.60	2.73
IC 4296	9.42	1.57	5	2.10	1.07	9	0.24	0.03	116.60	59.92
IC 4296	12.56	1.57	5	2.10	1.07	10	0.16	0.03	152.23	79.11
IC 4296	15.46	1.45	6	1.29	0.21	11	0.17	0.03	91.31	18.74
IC 4296	17.88	1.21	6	1.29	0.21	12	0.17	0.05	91.62	22.83
			...		...		...		...	

assumption in Section 3.4.4. Because the X-ray temperature gradient across the radial range we are interested in is small, we can produce better statistical fits with deprojection by fitting a single temperature across multiple (two to five) adjacent annuli while allowing the spectral normalization to be free in each annulus. The full tabulated results of these fits including uncertainties are provided in Table 3.2.

NGC 193, NGC 3801, and NGC 3894 were removed from our sample because there were not enough counts to obtain a deprojected temperature profile with three or more radial bins. NGC 4782 had sufficient counts to extract a profile but had a bright central point source resulting in large uncertainties in the central bins. For NGC 1316 and IC 4296, we do not attempt to fit the central point sources because our primary goal is to assess the shape of the entropy profile and the data quality for future work. Therefore, the central $2''$ from IC 4296 and NGC 1316, 0.25 and 0.12 kpc, respectively, were excluded from our deprojection analyses of these two galaxies.

3.3.4 Thermodynamic Properties

3.3.4.1 Electron Density Profiles

To estimate the electron density within a given concentric shell i , we use the best-fit spectral normalization from the deprojection model in XSPEC,

$$\eta_i = \frac{10^{-14}}{4\pi D^2(1+z)^2} \int n_{e,i} n_{p,i} dV_i. \quad (3.1)$$

The `project` model performs the projection from 3D to 2D and the total emission measure within the extraction volume as shown in Equation 3.1, in which D is the angular diameter distance to the galaxy in centimeters (Table 1), n_e and n_p are the electron and proton number densities, respectively, in cm^{-3} , and V_i is the volume of the concentric shell in cm^3 . With this definition of normalization, the expression

$$n_e(\text{shell}) = \sqrt{\frac{4\pi\eta(\text{shell})D^2(1+z)^2}{10^{-14}(n_e/n_p)V(\text{shell})}} \quad (3.2)$$

gives us the deprojected radial electron density profile for each galaxy.

3.3.4.2 Entropy and $t_{\text{cool}}/t_{\text{ff}}$ Profiles

We plot the entropy profiles of the galaxies in our sample in Figure 3.1. Radial profiles of the $t_{\text{cool}}/t_{\text{ff}}$ ratio are shown in Figure 3.2, with t_{cool} defined by

$$t_{\text{cool}} = \frac{3}{2} \frac{nkT}{n_e n_H \Lambda(T, Z)} \quad (3.3)$$

where n is the total number density of particles, n_e is the electron density, n_p is the hydrogen density (where we assume $n_p = n_e/1.2$), and $\Lambda(T, Z)$ is the temperature-dependent cooling function for plasma of metallicity Z . Our fiducial cooling function, from Schure et al. (2009), assumes a solar-metallicity ($Z_\odot$) plasma. The freefall time is calculated assuming a singular isothermal sphere with velocity dispersions found in Table 3.1. We calculated $t_{\text{cool}}/t_{\text{ff}}$ for NGC 4261 and three additional galaxies with the best data quality (NGC 1316, NGC 4374, and IC 4296).

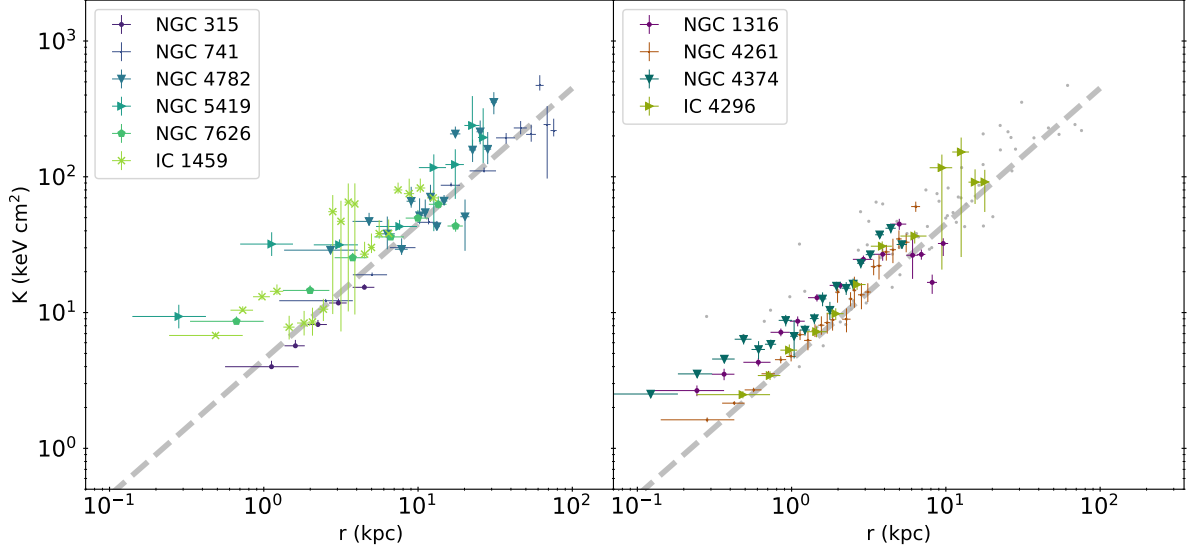


Figure 3.1: **Entropy Profiles for early-type galaxies with powerful radio sources** Left panel: entropy profiles for the galaxies in our sample with sufficient data counts to extract a deprojected radial profile but insufficient data to isolate the central ~ 0.5 kpc. Right panel: deprojected entropy profiles of the four galaxies with the best data quality (NGC 4261, IC 4296, NGC 1316, NGC 4374). For comparison, the gray dots are the data points from the galaxies in the left panel. Gray dashed lines on both plots show power-law profiles with $K \propto r$ to illustrate that NGC 4261 and IC 4296 differ from the other galaxies with comparable data quality (NGC 1316 and NGC 4374) by approximately following a similar power law into the central kiloparsec, rather than exhibiting a small excess like the other single-phase galaxies.

3.4 Discussion

3.4.1 $t_{\text{cool}}/t_{\text{ff}}$ Profiles and Multiphase Gas

Figure 3.2 shows the $t_{\text{cool}}/t_{\text{ff}}$ profiles for the four galaxies with entropy profiles that come closest to probing the inner ~ 0.5 kpc of the galaxy. The profiles of IC 4296 and NGC 4374 are of particular interest. While the data are not of the resolution of NGC 4261, they still allow us to see the shape of the $t_{\text{cool}}/t_{\text{ff}}$ and entropy profiles near the central ~ 0.5 kpc of the galaxy. We also note that while the X-ray structure of NGC 315 is not resolved inside ~ 1 kpc, its gas entropy profile appears to follow a single power law like IC 4296 and NGC 4261. Furthermore, from the spectra reported in Ho et al. (1997, 1993), its multiphase gas appears to be confined to the nucleus, making it another promising candidate for a system in this powerful but possibly short-lived state.

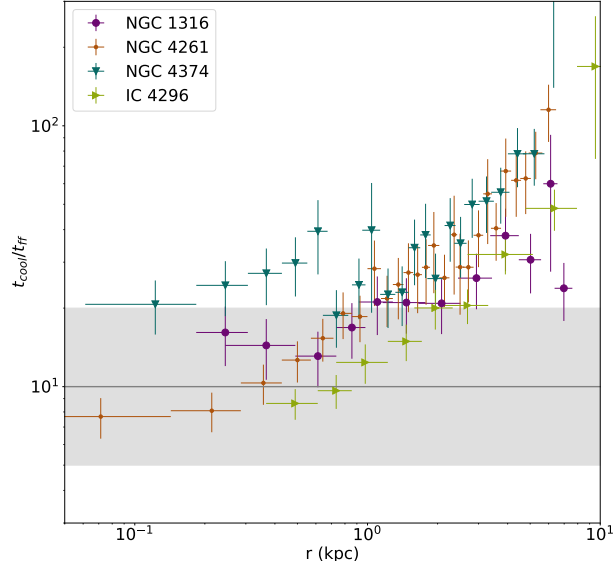


Figure 3.2: **Radial profiles of $t_{\text{cool}}/t_{\text{ff}}$** Radial profiles of $t_{\text{cool}}/t_{\text{ff}}$ for the four galaxies with the best S/N. The shaded region ($t_{\text{cool}}/t_{\text{ff}} = 5 - 20$) represents the precipitation zone where multiphase gas is found for $r = 1 - 10$ kpc. We find that, like NGC 4261, IC 4296 reaches $t_{\text{cool}}/t_{\text{ff}} < 10$ in the central ~ 1 kpc while the other galaxies do not.

Voit et al. (2015b) showed that $t_{\text{cool}}/t_{\text{ff}}$ in the central ~ 1 kpc of both single- and multiphase galaxies usually remains above the apparent precipitation limit at $t_{\text{cool}}/t_{\text{ff}} \sim 10$. Farther out from the center (1–10 kpc), galaxies with multiphase gas have $t_{\text{cool}}/t_{\text{ff}}$ profiles that approximately track this precipitation limit, whereas galaxies with single-phase gas generally lie above the precipitation zone at $t_{\text{cool}}/t_{\text{ff}} \sim 5 - 20$ (blue shaded region in Figure 3.2). Voit et al. (2015b) found that, in a sample of morphologically relaxed, X-ray-bright galaxies (Werner et al., 2012), only the radial profile for NGC 4261 dipped below $t_{\text{cool}}/t_{\text{ff}} \sim 10$ in the center.

In our sample, the $t_{\text{cool}}/t_{\text{ff}}$ profile for NGC 4374 remains above the precipitation zone, and NGC 1316 is consistent with the multiphase galaxy pattern from Voit et al. (2015b). However, IC 4296 goes down to $t_{\text{cool}}/t_{\text{ff}} \sim 10$ near the center, as in NGC 4261, suggesting that the AGN feedback occurring in IC 4296 has interesting similarities to that of NGC 4261. The data were sufficient to probe the inner ~ 0.5 kpc of NGC 4261, but in general the profiles more closely follow a single power law than a power law with an excess inside the central kiloparsec.

The $\text{H}\alpha$ emission in NGC 4261 is nuclear rather than extended (Ferrarese et al., 1996;

Lakhchaura et al., 2018), consistent with the picture of giant galaxies with single-phase gas having entropy profiles that scale as $K(r) \propto r$. Of our studied galaxies, IC 4296 most closely resembles NGC 4261, and Grossová et al. (2019) reported that in narrowband images from the *Hubble* and Southern Astrophysical Research (SOAR) telescopes, IC 4296 also has no H α emission beyond $r \sim 2$ kpc.

3.4.1.1 Comparison with Previous X-ray Analysis

In an independent analysis, Lakhchaura et al. (2018) report entropy profiles for a sample of 49 elliptical galaxies, including eight of the galaxies analyzed in this paper: NGC 315, NGC 741, NGC 1316, NGC 4261, NGC 4374, NGC 4782, IC 4296, and NGC 5419. While there are small variations among bin sizes and radial ranges, we verified that our results are nevertheless mutually consistent within the measurement uncertainties. However, the work of Lakhchaura et al. (2018) treated gas metallicity differently, which we address in Section 3.4.4. In Section 3.4.2, we include some of the results of Lakhchaura et al. (2018) in our discussion. Additionally, in Table 3.1, we report the multiphase gas classifications from Lakhchaura et al. (2018) as well as additional results from Connor (in preparation) and Sun (in preparation) that use observations carried out using the SOAR optical Imager (SOI) and Goodman High Throughput Spectrograph of the 4.1m SOAR telescope and the Apache Point Observatory (APO) Astrophysics Research Consortium (ARC) 3.5m telescope.

3.4.2 Radio Luminosity and $t_{\text{cool}}/t_{\text{ff}}$

Figure 3.3 shows the minimum values of $t_{\text{cool}}/t_{\text{ff}}$ for NGC 4261, IC 4296, NGC 1316, and NGC 4374, along with the giant ellipticals from Lakhchaura et al. (2018), plotted as a function of the radius at which $t_{\text{cool}}/t_{\text{ff}}$ reaches its minimum value. We have adjusted the $\min(t_{\text{cool}}/t_{\text{ff}})$ values reported by Lakhchaura et al. (2018) for uniform comparison with our work, using the correction factor estimated in Section 3.4.4. The typical amplitude and direction of that correction are plotted in Figure 3.3 in the form of a purple arrow. This adjustment typically decreased the t_c/t_{ff} estimates

from Lakhchaura et al. (2018) by a factor of 1.6. Points vary in size according to radio power in the 1.4 GHz band.

Notice that NGC 4261, IC 4296, and NGC 1316 have a lower $\min(t_{\text{cool}}/t_{\text{ff}})$ at a smaller radius than most of the other giant elliptical galaxies in the Lakhchaura et al. (2018) sample. Furthermore, the $t_{\text{cool}}/t_{\text{ff}}$ profiles in NGC 4261 and IC 4296, reach their minimum values in the central radial bin, raising the possibility that $\min(t_{\text{cool}}/t_{\text{ff}})$ is overestimated in these galaxies because of limited spatial resolution. However, it is also possible for those $\min(t_{\text{cool}}/t_{\text{ff}})$ values to be slight underestimates. In well-resolved galaxies that reach $\min(t_{\text{cool}}/t_{\text{ff}})$ outside the central radial bin, statistical fluctuations tend to cause the measurement of $\min(t_{\text{cool}}/t_{\text{ff}})$ to be biased low. Figure 3.2 shows why the magnitude of that bias in the galaxies we are focusing on is likely to be small. In all four galaxies, the second-lowest value of $t_{\text{cool}}/t_{\text{ff}}$ is nearly identical to the minimum value, well within the 1-sigma statistical uncertainties. Also, the $t_{\text{cool}}/t_{\text{ff}}$ profiles of those four galaxies are not constant in the 1–10 kpc range but only within the central ~ 1 kpc, where there are only a few radial bins, reducing the likelihood of an unrepresentative statistical fluctuation. Consequently, the fact that NGC 4261 and IC 4296 have unusually low $\min(t_{\text{cool}}/t_{\text{ff}})$ and greater radio power than most other galaxies in the sample suggests that there may be a correlation between high radio power and $t_{\text{cool}}/t_{\text{ff}} < 10$ at small radii. In particular, the combination of low $\min(t_{\text{cool}}/t_{\text{ff}})$ and a power-law entropy slope that does not significantly flatten within the central kiloparsec is a unique feature of NGC 4261 and IC 4296. The other galaxies, in which the central entropy profile is flatter and $\min(t_{\text{cool}}/t_{\text{ff}})$ occurs at a larger radius, could be systems in which AGN feedback has recently elevated the entropy in the central kiloparsec.

3.4.3 Comparison to Simulations

Voit et al. (2015b) showed that the presence of multiphase gas outside the central kiloparsec of an early-type galaxy correlates with the slope of the entropy profile. Galaxies with an entropy slope of $K \propto r^{2/3}$ have multiphase gas present at $r > 1$ kpc, while galaxies with an entropy slope of $K \propto r$ have only single-phase gas beyond $r \sim 1$ kpc. In order to better understand this relationship, we

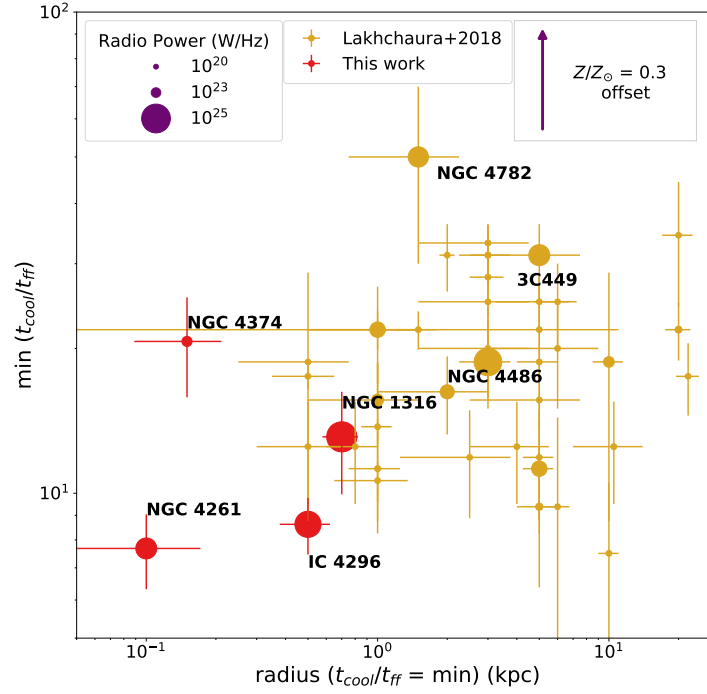


Figure 3.3: **Radio luminosity, $\min(t_{\text{cool}}/t_{\text{ff}})$, and radius at $\min(t_{\text{cool}}/t_{\text{ff}})$** Radius where we measured the minimum value of the $t_{\text{cool}}/t_{\text{ff}}$ profile is plotted against the minimum $t_{\text{cool}}/t_{\text{ff}}$ for the sample of Lakhchaura et al. (2018, gold) offset to solar metallicity (see 3.4.4), with our four galaxies of best data quality (red). The purple arrow represents the offset between the adjusted values and the $0.3 Z_{\odot}$ values from Lakhchaura et al. (2018). The relative size of the points represents their radio power (in W Hz^{-1}) in the 1.4 GHz band. NGC 4261, IC 4296, and NGC 1316 have a small $\min(t_{\text{cool}}/t_{\text{ff}})$ radius, a low $\min(t_{\text{cool}}/t_{\text{ff}})$, and a greater radio power than most galaxies in the sample.

have compared our observed entropy profiles with the profiles of simulated galaxies from Wang et al. (2019). Figure 3.4 shows a comparison between our data for NGC 4261 and IC 4296 and simulated elliptical galaxies with both single- and multiphase gas as well as the entropy profiles for galaxies classified as having extended multiphase gas and no extended multiphase gas from Lakhchaura et al. (2018). The initial conditions for the simulations are chosen to mimic X-ray observations of NGC 5044 (multiphase) and NGC 4472 (single-phase), but the simulations were designed to represent generic single- and multiphase galaxies. The simulations do not resolve the gas profiles at < 1 kpc, meaning that our data have greater effective physical resolution than the simulations. However, we can still make comparisons in the $1 - 10$ kpc range.

We begin by considering whether the simulated galaxies are appropriate comparisons for our

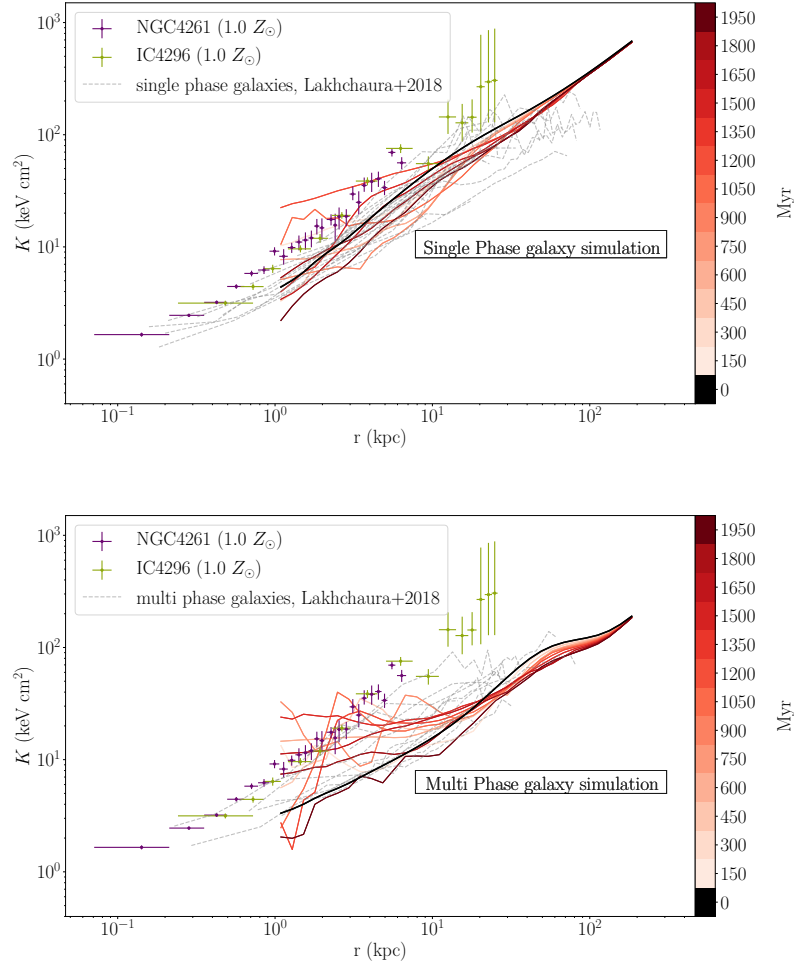


Figure 3.4: Entropy profiles of NGC 4261 and IC 4296 compared to simulations Entropy profiles of NGC 4261 and IC 4296 compared to simulations of somewhat lower mass giant elliptical galaxies with single-phase gas (top) and multiphase gas (bottom) from Wang et al. (2019) along with the single-phase (top) and extended multiphase (bottom) galaxies from Lakhchaura et al. (2018). Simulated profiles are shown at intervals 150 Myr, with earlier snapshots being shown in lighter red. The initial conditions are given by the black line and represent typical entropy profiles for single-phase and multiphase galaxies, based on NGC 4472 and NGC 5044, respectively. Galaxies classified as “N” (no extended multiphase gas) and “E” (extended multiphase gas) from Lakhchaura et al. (2018) are included in gray with errors bars removed for clarity. The simulated galaxies have lower velocity dispersions than the observed galaxies, but the simulations were designed to represent the behavior over time of generic single- and multiphase galaxies, rather than simulating a specific galaxy. The galaxies are referred to as MPG (multi phase galaxy) and SPG (single-phase galaxy) instead of their names throughout Wang et al. (2019). Note that the simulations do not resolve the gas profiles inside 1 kpc, and the flattening of the profiles in the center is likely a numerical effect because the resolution limit of the simulations can result in the simulated AGN affecting a larger region of the galaxy than real jets (Wang et al., 2019). Therefore, we expect the entropy slopes of the single-phase simulations to be the same as our observations, though the normalization can differ. The measured entropy gradients are consistent with those seen in single-phase gas simulations of radio sources in early-type galaxies.

sample. From Makarov et al. (2014), the velocity dispersions of NGC 5044 and NGC 4472 are $224.9 \pm 9.1 \text{ km s}^{-1}$ and $282 \pm 2.9 \text{ km s}^{-1}$, respectively, while the velocity dispersions for IC 4296 and NGC 4261 are $327.4 \pm 5.4 \text{ km s}^{-1}$ and $296.7 \pm 4.3 \text{ km s}^{-1}$, respectively. Voit et al. (2015b) introduced the idea that there may be a correlation between the presence of multiphase gas, the velocity dispersion, and entropy profile slope in early-type galaxies. In contrast, Lakhchaura et al. (2018) found little correlation in entropy profile slope and the presence of multiphase gas when examining a larger sample. However, this is still an area of open study in both theory and observations of early-type galaxies. We do not expect the simulated profiles to match our data exactly because the velocity dispersions of the simulated galaxies are smaller than those of NGC 4261 and IC 4296. However, we can still make useful comparisons between the overall behavior of the simulations and observations that account for how the entropy profile slope correlates with velocity dispersion and the presence or absence of multiphase gas (Voit et al., 2015b).

For the single- and multiphase initial conditions simulated by Wang et al. (2019), the entropy profiles agree with the expectations of the simple physical models shown in Voit et al. (2015b). The multiphase gas simulation has an entropy profile of $K(r) = 3.5 r_{\text{kpc}}^{2/3} \text{ keV cm}^2$, which corresponds to the hypothesized precipitation limit at $t_c/t_{\text{ff}} \approx 10$. The steeper entropy profile characteristic of single-phase galaxies, $K(r) = 5 r_{\text{kpc}} \text{ keV cm}^2$, is consistent with models in which heating by Type Ia supernovae drives an outflow. The implication is that self-regulated AGN feedback can maintain the observed properties of both the single- and multiphase galaxies, consistent with both idealized analytical models (Voit et al., 2015b) and simulations (Wang et al., 2019).

When we compare the single-phase simulation to NGC 4261 and IC 4296, the simulated galaxy does appear to maintain approximately the same entropy slope as our data. Furthermore, the comparison shows that the slopes of the single-phase entropy profiles from Lakhchaura et al. (2018) are different from the slopes of the multiphase entropy profiles, and the multiphase data have similar slopes to the multiphase simulations. The comparison between NGC 4261 and IC 4296, the multiphase galaxy simulations, and the extended multiphase gas data show that our galaxies are better represented by single-phase galaxies. The observed entropy profile slopes between 1 and

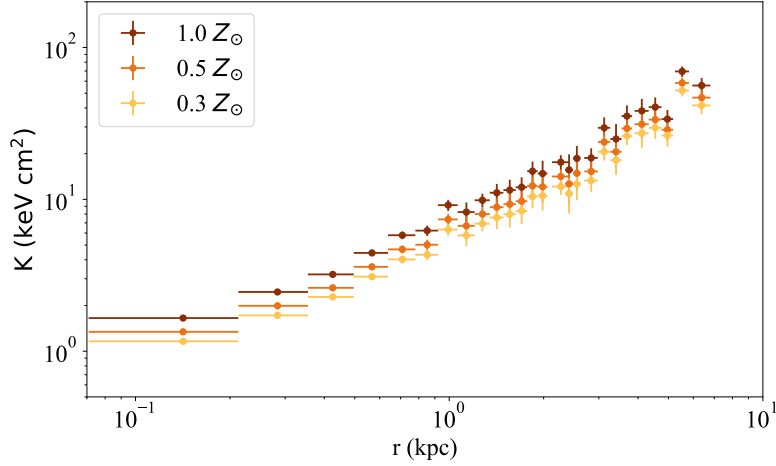


Figure 3.5: **Comparison of the inferred entropy profiles for NGC 4261 for different assumed values of abundance** The points represent $1.0 Z_{\odot}$ (red), $0.5 Z_{\odot}$ (orange), and $0.3 Z_{\odot}$ (yellow). For increasing values of metallicity, the amplitude of the entropy profile increases. Therefore, if a galaxy has a different metallicity than we have assumed, and if we no longer assume that each galaxy has uniform metallicity, the slope of the entropy profile could change as well. However, we would still expect the profile to fall within the metallicity range illustrated in the figure.

10 kpc are consistent with the entropy profile of a single-phase galaxy, which has a steeper slope ($\alpha \sim 1$). For some time steps, the central entropy profile of the simulated single-phase galaxy flattens out, which could represent epochs when the black hole in the simulations is particularly active. However, it could also be a numerical effect because the resolution limit of the simulations can result in the simulated AGN affecting a larger region of the galaxy than real jets (Wang et al., 2019).

3.4.4 Metallicity Analysis

When fitting entropy profiles, we assumed that each galaxy had a constant metallicity of $1.0 Z_{\odot}$ across the profile. The hot gas abundances of early-type galaxies are difficult to obtain from low-resolution X-ray data, so we fixed the gas metallicities while fitting our observations. Here we quantify the sensitivity of our estimates of the X-ray densities and temperatures to the assumed metallicities. In Figure 3.5, we show an example result of the impact of three different metallicity assumptions on the entropy profiles for NGC 4261: 0.3 , 0.5 , and $1.0 Z_{\odot}$. This range of abundances

spans those from various treatments of early-type galaxies in observations and simulations found in the literature (Werner et al., 2012; Li & Bryan, 2014b; Wang et al., 2019). The amplitude of the inferred entropy increases as assumed abundance increases, but the slope of the entropy profile shows little change. Therefore, abundance changes over the full range of expected gas abundances from 0.3 to 1.0 solar would result in a change in amplitude of the inferred entropy profile on the order of 10 – 20% (see Figure 3.5).

In comparing our results with entropy profiles from previous work, we find that different abundance assumptions indeed result in small differences in the inferred entropy profiles. However, when the identical assumptions for abundances are used, the entropy profile results for different authors are the same within the uncertainties. For example, for the giant ellipticals examined by Werner et al. (2012), the assumed abundances in the central kiloparsec were $Z \sim 0.5 Z_{\odot}$ and match our entropy profile results for NGC 4261 for a metallicity of $Z \sim 0.5 Z_{\odot}$.

The relation between abundance and electron density for an *apec* model for a narrow range of gas temperatures (0.5 – 1.2 keV) can be quantified approximately by

$$\log \left[\frac{n_e(Z)}{n_e(Z_{\odot})} \right] = -m \log (Z/Z_{\odot}) \quad (3.4)$$

where Z is metallicity assumed in the determination of n_e and m is the power-law slope. We would expect $m \sim 0.4$ based on the dependence of $\Lambda(T)$ on Z in this temperature range (approximately $\Lambda \propto Z^{0.8}$) and form of the emission integral. To verify this estimate, we found the best-fit n_e and T_X for the four galaxies with the best data quality, sampling a range of assumed metallicities. We determined that the best-fit temperature was insensitive to the metallicity assumption, while density and assumed metallicity were related as in Equation 3.4 with $m = 0.43 \pm 0.04$ (NGC 1316), $m = 0.43 \pm 0.04$ (NGC 4261), $m = 0.39 \pm 0.11$ (NGC 4374), and $m = 0.29 \pm 0.18$ (IC 4296), where uncertainties on m are 1σ . These results are consistent with the expectations from X-ray plasma emission model for $kT \sim 0.6 - 1$ keV. The inferred electron density is therefore inversely related to the assumed abundance.

Furthermore, if the abundance is actually lower in the center than we assume, we have underestimated the central density and overestimated the central entropy. So if a galaxy's gas is less

metal-rich in the center than we have assumed, its central entropy profile could be slightly steeper than shown (e.g. Lakhchaura et al., 2019). However, even if the central metallicity is lower than we assumed, the minima of the $t_{\text{cool}}/t_{\text{ff}}$ profiles are less than $t_{\text{cool}}/t_{\text{ff}} = 20$ for NGC 4261 and IC 4296.

3.5 Conclusions

Our analysis of the entropy profiles for a sample of nearby early-type galaxies with powerful radio sources shows that at least one other galaxy (IC 4296) is like NGC 4261 in having a powerful AGN and $t_{\text{cool}}/t_{\text{ff}} \sim 10$ at < 1 kpc. While the spatial resolution of the X-ray data for IC 4296 is not as good as for NGC 4261, both of their entropy profiles appear to be single power laws, and neither has extended multiphase gas greater than 2 kpc from their nuclei. To be certain of their similarity, we will need additional *Chandra* observations of IC 4296 to match the data quality of NGC 4261.

We produced deprojected temperature and density profiles for the hot gas surrounding seven additional early-type galaxies with powerful radio sources, but these observations lacked sufficient data quality to quantify the slope in the central ~ 0.5 kpc. Unfortunately, these galaxies are likely not good candidates for further study at this time because the additional *Chandra ACIS* observations needed to achieve comparable data quality to NGC 4261 are prohibitively long, given the degradation of *Chandra's* sensitivity to soft X-rays. We found that, in comparing independent analyses of entropy profiles in early-type galaxies, the treatment of abundance affects the amplitude of the entropy profile. Additionally, if the gas is not well mixed, it may have a metallicity gradient, meaning that the slope of the profile could be affected as well. Finally, we compared IC 4296 and NGC 4261 to recent simulations (Wang et al., 2019) and found that they are consistent with a single-phase gas model galaxy. The simulations agree well with our observational results, providing positive evidence for their ability to robustly model the hot ambient medium in early-type galaxies. In this work we were able to show excellent agreement between our observations, the theory of Voit et al. (2015b), and the simulations of Wang et al. (2019).

CHAPTER 4

RELATIONSHIPS BETWEEN CENTRAL VELOCITY DISPERSIONS AND ATMOSPHERES OF EARLY-TYPE GALAXIES

4.1 Abstract

The Voit et al. (2020) black-hole feedback valve analytic model predicts relationships between stellar velocity dispersion and atmospheric structure for massive galaxies. In this work, we test the analytic model using the Chandra archival sample of 49 early-type galaxies from Lakhchaura et al. (2018). We consider the relationships between stellar velocity dispersion and entropy profile slope, multiphase gas extent, and ratio between cooling time and free-fall time simultaneously. We classify sub-samples limited to observations of high data quality and by entropy profile properties to explore the potential relationships between parameters and test the analytic model predictions. We find evidence for agreement with the equilibrium radial profiles from the Voit et al. (2020) model as well as agreement with the analytic model for the sample with low central entropy and limited velocity dispersion.

4.2 Introduction

Early-type galaxies, encompassing both elliptical and lenticular galaxies, are characterized by their elliptical shapes, older stellar populations, and lack of significant active star formation. Star formation in galaxies occurs when there is sufficient molecular gas to form stars and proceeds until the molecular gas supply runs out, either through stars forming more rapidly than the molecular gas can accumulate or the galaxy preventing further accumulation of molecular gas. It follows then, because little star formation is observed in early-type galaxies, that the galaxy must be preventing the molecular gas from accumulating. Molecular gas can accumulate in galaxies via cold streams (e.g. Kereš et al. 2005, 2009; Dekel et al. 2009), cooling flows (e.g. White & Frenk 1991; Fabian 1994; McNamara & Nulsen 2007, 2012; Werner et al. 2019), or stellar mass loss (e.g. Mathews & Brighenti 2003; Leitner & Kravtsov 2011; Voit & Donahue 2011). Therefore, feedback processes

in early-type galaxies must act to prevent each of these sources.

Observationally, the effect of feedback processes on the galactic atmosphere can be captured via observations of the hot X-ray gas. Entropy, in terms of X-ray observables, $K \equiv kT n_e^{2/3}$, where kT is the X-ray temperature and n_e is the electron density, is the preferred quantity for investigating feedback processes in galaxies. Feedback can change the rate at which the gas radiates energy away, affecting the cooling time of the gas. Here, the cooling time (t_{cool}) is defined to be the time needed for a gas at temperature T to radiate an energy $3kT/2$ per particle, and the free fall time from a galactocentric radius r at the local gravitational acceleration g is defined to be $t_{\text{ff}} = (2r/g)^{1/2}$. If the heating due to feedback is gradual compared to the time it takes for the heated gas to expand within the gravitational potential, the temperature of the gas may not change while the gas density lowers, lengthening the cooling time. We turn to entropy to capture gains and losses of energy in the gas. The gas entropy across the radius of a galaxy provides insight into the thermal history of the galactic atmosphere. In general, the galaxy potential well serves as an entropy sorting device, where lower entropy gas sinks to small radii, and higher entropy gas rises to larger radii. The lowest entropy gas is the densest and brightest gas in the galaxy, as observed in the X-ray. Voit et al. (2015b) examined the properties of a sample of 14 massive elliptical galaxies previously studied by Werner et al. (2012, 2014) and showed that the entropy profile slopes of early-type galaxies and the presence of multiphase gas are correlated. In the Werner et al. (2012) sample, inside ~ 2 kpc, the gas entropy levels of the galaxies are similar, but outside ~ 2 kpc, the slopes of the entropy profiles can differ from galaxy to galaxy depending on what thermal processes dominate. Galaxies with extended multiphase gas exhibit entropy profiles with $K \propto r^{2/3}$ from ~ 1 –10 kpc while galaxies with no extended multiphase gas (hereafter referred to as single phase galaxies) exhibit steeper entropy profiles, with $K \propto r$ from ~ 1 –10 kpc. The difference in the entropy profile slopes for galaxies with or without extended multiphase gas could be due to SNIa heating and sweeping gas ejected by the old stellar population out of single phase galaxies into an extended gaseous halo (Voit et al., 2015b). Voit et al. (2015b) also found that the velocity dispersions of the galaxies with extended multiphase gas were $\sigma \leq 255 \text{ km s}^{-1}$ while galaxies with no extended multiphase gas had

velocity dispersions of $\sigma \geq 263 \text{ km s}^{-1}$, indicating that entropy profile slope, velocity dispersion, and multiphase gas extent may be correlated and related to how the black hole interacts with the galactic atmosphere.

Lakhchaura et al. (2018) explored the relationship between entropy profile slope and multiphase gas extent for a larger archival sample (~ 50 galaxies) and did not report evidence for a relationship. However, Lakhchaura et al. (2018) did find evidence that the average behavior of entropy profiles and the ratio of the cooling time and free-fall time of the gas are related. Babyk et al. (2018) explored the relationship between entropy profile slope and velocity dispersion for an archival sample of 40 early-type galaxies (and 110 clusters). They also reported no evidence for a relationship between entropy profile slope and velocity dispersion, but they did find some evidence for a relationship between entropy profile slope and temperature.

Voit et al. (2020) investigated the coupling between supernova sweeping of stellar ejecta, the confining circumgalactic medium (CGM) pressure, and bipolar kinetic feedback fueled by accretion of cooling gas onto the central black hole, forming what they called a black hole feedback valve. They proposed an analytic model, investigating this idea, that predicts a relationship between the velocity dispersion and the entropy profile slope, that determines the effect of feedback on the galactic atmosphere, and whether multiphase gas can form. The model is informed by both numerical simulations and observations and analytically models feedback processes in massive galaxies. The model predicts that the entropy profile slope over the radial range where supernova heating exceeds radiative cooling ($\sim 1\text{--}10 \text{ kpc}$) is determined by the ratio of the specific thermal energy of the ejected stellar gas to the depth of the galactic potential well, as long as the velocity field is subsonic.

The structure of this paper is as follows: Section 4.3 describes our sample selection and data analysis processes, Section 4.4 discusses the connection between observations and theory, and Section 4.5 concludes by discussing how this work adds to the current understanding of precipitation-driven feedback in massive galaxies.

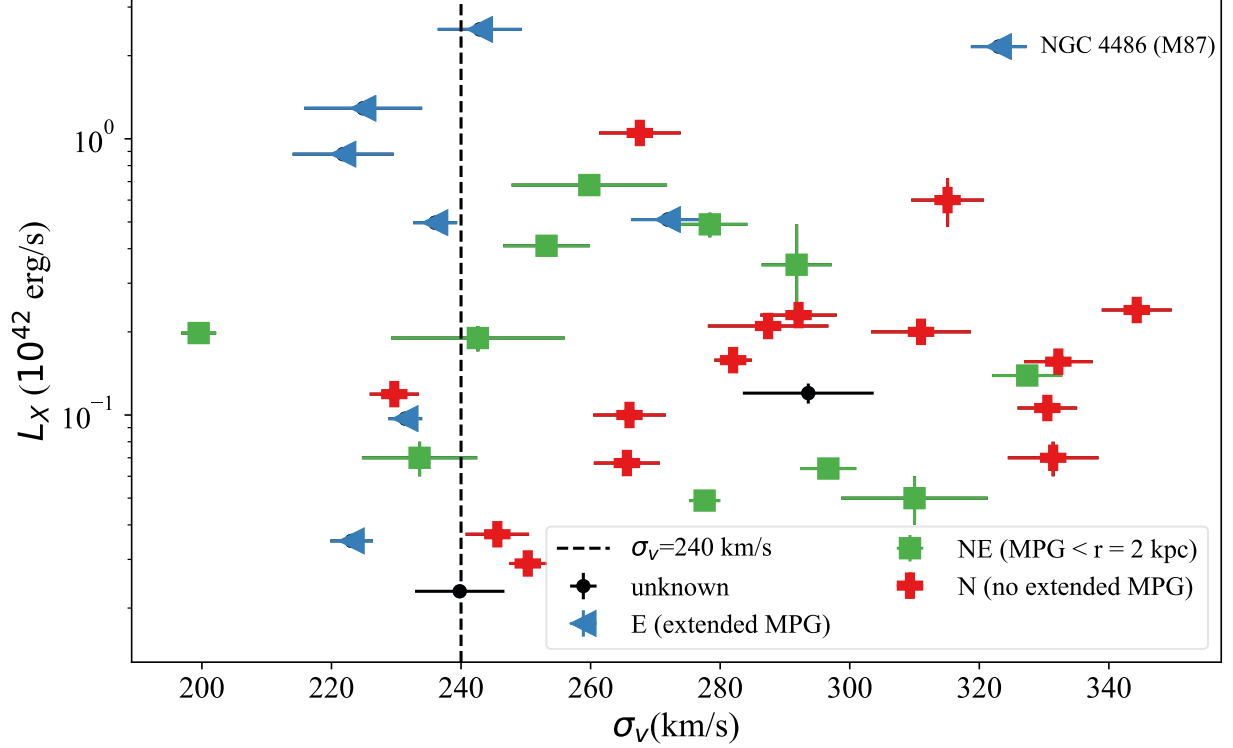


Figure 4.1: **Stellar velocity dispersion vs. X-ray luminosity** Stellar velocity dispersion, σ_v , is plotted with the X-ray luminosity within 10 kpc for the subsample of galaxies from Lakhchaura et al. (2018) with sufficient data to measure an accurate entropy profile slope at 1-10 kpc. The points are also classified by their multiphase gas extent from Lakhchaura et al. (2018). Blue triangles are galaxies with extended multiphase gas, red crosses are galaxies with no extended multiphase gas, green squares are galaxies with multiphase gas contained within 2 kpc, and black dots are galaxies without a gas extent classification. The vertical dashed line indicates the velocity dispersion (240 km s^{-1}) that corresponds to the critical entropy profile slope of $\alpha_K = 2/3$ (Voit et al., 2020).

4.3 Methods

4.3.1 Sample Description

Our primary goal in this work is to determine whether the observed relationship between the velocity dispersion and entropy profile slope is consistent with the analytic model predictions from Voit et al. (2020). Making such a comparison requires sufficient resolution in the observations to measure an entropy profile slope, so we need to use a sample of early-type galaxies with accurate entropy profiles and velocity dispersion measurements to test the analytic model’s predictions.

The main sample explored in this work is the sample of 49 nearby, X-ray and optically bright,

elliptical galaxies with archival *Chandra* data from Lakhchaura et al. (2018). We use the derived radial profile measurements of electron density, n_e , X-ray temperature, kT , entropy, K , and ratio of cooling time to free-fall time, t_c/t_{ff} , as well as their multiphase gas classification scheme¹ and X-ray luminosities. Full details of the galaxy parameters are found in Table 4.1.

Figure 4.1 shows the relationships between X-ray luminosity, stellar velocity dispersion, and multiphase gas characteristics for galaxies with sufficiently resolved entropy profiles (see Section 4.3.3 for details). Voit et al. (2020) showed that $\sigma_v = 240 \text{ km s}^{-1}$ corresponds to $\alpha_K = 2/3$ and represents a critical number for the analytic model. Extended multiphase gas appears to be more common among galaxies with $\sigma_v < 240 \text{ km s}^{-1}$ than among those with $\sigma_v > 240 \text{ km s}^{-1}$ (Voit et al., 2015b). In Figure 4.1, we also see that galaxies with multiphase gas confined to the inner $\sim 2 \text{ kpc}$ are represented across the range of velocity dispersion. The most notable exception to the division between galaxies with extended multiphase gas and those with no extended multiphase gas is M87 (upper right of Figure 4.1) which resides in one of the most massive halos in the sample and has one of the highest X-ray luminosities. The massive halo and high luminosity indicate that the galaxy likely has high external gas pressure, and that its atmospheric characteristics are representative of the entire massive halo rather than only this single galaxy. Apart from M87, the upper envelope of the sample exhibits a decline in X-ray luminosity with increasing velocity dispersion. Voit et al. 2015b, 2020 also predicted that galaxies with higher velocity dispersion will have steeper entropy profile slopes. A steeper entropy profile slope means that the electron density of the gas is lower at 10 kpc (the aperture for measuring L_X) than it would be for shallower entropy profiles slopes.

The sample from Lakhchaura et al. (2018), after initial data quality limits are applied, is well-suited to test the analytic model predictions from Voit et al. (2020). In Section 4.3.4, we will discuss further the ways in which we subdivide the sample in our test of the analytic model.

¹Except for IC 4296 which is correctly identified as NE, rather than E, in Frisbie et al. (2020)

Table 4.1: **Galaxy Parameters** (1) Name; (2) redshift obtained from NED; (3) redshift independent distance; (4) σ_v and error; (5) 0.5–7.0 keV intrinsic X-ray luminosities and their errors estimated from a 10 kpc radius circular region around the X-ray peak (Lakhchaura et al., 2018); (6) $H\alpha$ + $[N\ 2]$ morphology classified as follows: N: no cool gas emission, NE: $H\alpha$ + $[N\ 2]$ extent < 2 kpc, E: $H\alpha$ + $[N\ 2]$ extent $\geq$ 2 kpc and U: galaxies for which the presence/absence of $H\alpha$ + $[N\ 2]$ could not be confirmed (Lakhchaura et al., 2018); (7) α_K and error; (8) minimum ratio of cooling time to free-fall time and error; (9) Fit central entropy and error.

Galaxy	z	D Mpc	σ_v km/s	L_X 10^{42}erg/s	gas extent	α_K	min ($t_{\text{cool}}/t_{\text{ff}}$)	K_0 keV/cm ²
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)
IC1860	0.0229	95.75	259.8 ± 12.0	0.68 ± 0.050	NE	0.66 ± 0.24	18.42 ± 2.40	5.92 ± 0.97
IC4296	0.0124	47.31	327.4 ± 5.4	0.14 ± 0.003	NE	1.23 ± 0.11	11.63 ± 0.78	0.69 ± 0.20
IC4765	0.0150	59.52	278.4 ± 5.8	0.49 ± 0.050	NE	0.88 ± 0.22	11.01 ± 1.20	1.82 ± 0.72
NGC315	0.0164	56.01	293.6 ± 10.1	0.12 ± 0.010	U	0.83 ± 0.24	20.02 ± 4.65	0.91 ± 0.96
NGC410	0.0176	66.00	291.8 ± 5.4	0.35 ± 0.140	NE	0.76 ± 0.35	27.30 ± 5.22	4.93 ± 1.36
NGC499	0.0147	60.74	253.2 ± 6.7	0.41 ± 0.030	NE	0.74 ± 0.24	34.18 ± 3.19	6.90 ± 4.16
NGC507	0.0164	59.83	292.1 ± 5.9	0.23 ± 0.020	N	0.80 ± 0.34	30.15 ± 5.61	7.14 ± 3.38
NGC533	0.0184	61.58	271.9 ± 5.6	0.51 ± 0.030	E	0.88 ± 0.14	12.28 ± 3.75	1.74 ± 0.36
NGC708	0.0162	64.19	221.8 ± 7.8	0.88 ± 0.020	E	0.64 ± 0.06	12.04 ± 0.29	5.38 ± 0.15
NGC741	0.0186	64.39	287.4 ± 9.3	0.21 ± 0.010	N	0.93 ± 0.09	19.16 ± 0.73	2.57 ± 0.39
NGC777	0.0167	58.08	315.1 ± 5.6	0.60 ± 0.120	N	0.59 ± 0.24	24.11 ± 3.01	5.22 ± 1.01
NGC1316	0.0059	19.25	223.1 ± 3.3	0.04 ± 0.002	E	0.72 ± 0.25	32.57 ± 6.72	0.58 ± 0.61
NGC1399	0.0048	17.75	332.2 ± 5.3	0.16 ± 0.004	N	0.94 ± 0.03	26.05 ± 0.40	0.89 ± 0.11
NGC1404	0.0065	19.18	229.7 ± 3.8	0.12 ± 0.001	N	0.80 ± 0.03	20.23 ± 0.51	0.70 ± 0.04
NGC1407	0.0060	23.27	265.6 ± 5.1	0.07 ± 0.003	N	0.83 ± 0.06	41.92 ± 1.67	4.13 ± 0.25
NGC1521	0.0140	50.93	233.6 ± 8.9	0.07 ± 0.010	NE	0.38 ± 0.39	20.90 ± 5.21	0.96 ± 0.83
NGC1600	0.0158	45.77	331.4 ± 7.0	0.07 ± 0.010	N	0.72 ± 0.18	42.60 ± 7.16	5.07 ± 0.39
NGC2300	0.0064	41.45	266.0 ± 5.6	0.10 ± 0.010	N	0.91 ± 0.13	26.09 ± 1.27	4.18 ± 0.43
NGC2305	0.0113	47.88	242.6 ± 13.4	0.19 ± 0.020	NE	0.70 ± 0.25	20.22 ± 3.36	1.54 ± 0.58
NGC3091	0.0122	48.32	311.0 ± 7.7	0.20 ± 0.020	N	0.40 ± 0.10	30.74 ± 4.16	3.48 ± 2.54
NGC3923	0.0058	20.97	245.6 ± 4.9	0.04 ± 0.001	N	0.92 ± 0.12	21.98 ± 1.41	1.34 ± 0.12

Table 4.1 (cont'd)

Galaxy	z	D Mpc	σ_v km/s	L_X 10^{42}erg/s	gas extent	α_K	min ($t_{\text{cool}}/t_{\text{ff}}$)	K_0 keV/cm ²
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)
NGC4073	0.0197	60.08	267.6 ± 6.3	1.05 ± 0.050	N	0.61 ± 0.20	32.22 ± 2.92	8.30 ± 1.18
NGC4125	0.0045	21.41	239.8 ± 6.9	0.02 ± 0.001	U	0.13 ± 0.45	28.22 ± 12.35	1.43 ± 1.40
NGC4261	0.0073	29.58	296.7 ± 4.3	0.06 ± 0.003	NE	1.16 ± 0.06	14.17 ± 1.47	0.52 ± 0.08
NGC4374	0.0033	16.68	277.6 ± 2.4	0.05 ± 0.002	NE	1.18 ± 0.14	25.04 ± 6.58	1.86 ± 0.19
NGC4406	0.0006	16.08	231.4 ± 2.6	0.10 ± 0.004	E	0.54 ± 0.14	26.28 ± 1.60	5.21 ± 3.26
NGC4472	0.0032	15.82	282.0 ± 2.9	0.16 ± 0.001	N	0.96 ± 0.02	26.80 ± 0.23	1.17 ± 0.05
NGC4486	0.0042	16.56	323.0 ± 4.3	2.16 ± 0.004	E	0.61 ± 0.01	22.73 ± 0.27	3.00 ± 0.10
NGC4552	0.0009	15.97	250.3 ± 2.9	0.03 ± 0.001	N	0.95 ± 0.09	11.35 ± 0.63	2.23 ± 0.11
NGC4636	0.0031	15.96	199.5 ± 2.7	0.20 ± 0.002	NE	1.00 ± 0.03	10.79 ± 0.36	1.89 ± 0.08
NGC4649	0.0034	16.55	330.5 ± 4.6	0.11 ± 0.002	N	1.00 ± 0.02	22.63 ± 0.35	1.49 ± 0.02
NGC4696	0.0098	37.48	242.9 ± 6.5	2.49 ± 0.010	E	0.69 ± 0.01	4.73 ± 0.03	2.24 ± 0.07
NGC4782	0.0133	48.63	310.0 ± 11.3	0.05 ± 0.010	NE	0.59 ± 0.26	18.94 ± 9.92	4.30 ± 2.62
NGC5044	0.0090	35.75	224.9 ± 9.1	1.29 ± 0.010	E	0.56 ± 0.03	5.75 ± 0.14	0.08 ± 0.12
NGC5419	0.0139	50.87	344.3 ± 5.4	0.24 ± 0.020	N	1.19 ± 0.28	17.30 ± 2.21	1.38 ± 0.70
NGC5813	0.0064	29.23	236.0 ± 3.4	0.50 ± 0.003	E	0.51 ± 0.02	12.20 ± 0.26	3.44 ± 0.13

4.3.2 Theoretical Model

Voit et al. (2020) presents a basic model for the relationship between the entropy profile slope, α_K , and the stellar velocity dispersion, σ_v . The basic model assumes that the stellar mass distribution can be approximated by a singular isothermal sphere with one-dimensional velocity dispersion, $\sigma_v = v_c/\sqrt{2}$. The outflow driving by Type Ia supernova (SNIa) heating is assumed to be subsonic and therefore close to hydrostatic equilibrium. Combining the contributions to the entropy profile from supernova energy, orbital energy, and gravitational potential energy then gives the following

relation between α_K and σ_v :

$$\alpha_K \approx \frac{5}{3} \left(\frac{\epsilon_*}{v_c^2} - \frac{1}{4} \right)^{-1}, \quad (4.1)$$

where ϵ_* is the mean specific energy of the gas coming from stars. Therefore, the structure of the galaxy’s atmosphere at 1 – 10 kpc depends strongly on the ratio of ϵ_*/v_c^2 . Given this relationship between the entropy profile slope and the stellar velocity dispersion, we will use the velocity dispersions and entropy profile slopes from observations to test the model.

Voit et al. (2020) also presented a more complex form of the basic model by instead assuming that the galaxy’s halo has an NFW density profile and the stellar mass density follows a modified Einasto profile. Numerical integration of the more complex model shows that the basic model overpredicts the entropy profile slope for $\sigma_v > 300 \text{ km s}^{-1}$, and we will compare that modification of the model with data in Section 4.4.1.2.

4.3.3 Entropy Profiles

Equation 4.1 is based on a pressure-bounded, subsonic outflow solution, heated only by SNIa, and predicts a constant radial slope for the entropy profile. Because we want to test that model’s prediction for the relationship between the power-law slope of the entropy profile and stellar velocity dispersion, we limit the range over which we fit the power-law slope to the radial range that is affected as little as possible by other heating processes. If the AGN is as powerful as NGC 4261 and IC4296 (see Frisbie et al. (2020)), it typically deposits its energy rather far from the center ($r > 10 \text{ kpc}$) by drilling through the hot gas via jet. In some systems, some fraction of the AGN energy output might couple to gas closer to the AGN, resulting in flattening or even inversion of the entropy profile near 1 kpc. Therefore, we limit our gas slope measurements to 1-10 kpc to get a “clean” measure of the gas slope where stellar processes are most likely to dominate.

While a few of the galaxies in our sample have entropy profiles that resemble a single power law (Frisbie et al., 2020), most have an excess of entropy over a single power law in the central $\sim \text{kpc}$. Therefore, we have adapted the functional form from Donahue et al. 2005, 2006; Voit &

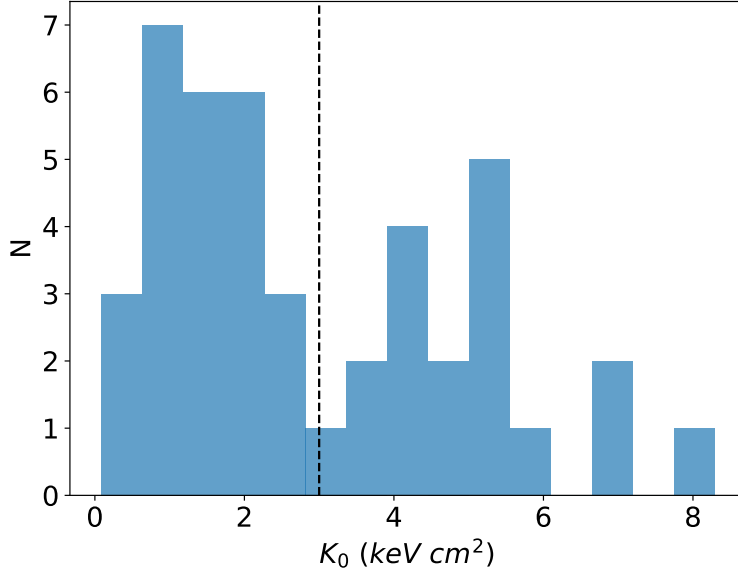


Figure 4.2: **Distribution of K_0 values for the HQ sample** Histogram of K_0 values for the sample where K_0 is fit using Equation 4.2. The black dashed vertical line is at $K_0 = 3 \text{ keV cm}^2$ and represents the limit we applied to remove galaxies with elevated central entropy outside of 1 kpc.

Donahue 2005; Cavagnolo et al. 2009 to the radial range for galaxies instead of clusters,

$$K(r) = K_0 + K_{10} \left(\frac{r}{10 \text{ kpc}} \right)^{\alpha_K}, \quad (4.2)$$

where K_0 is the best fit central entropy, K_{10} is the best fit entropy at a radius of 10 kpc, and α_K is the best fit power law slope. We fit over 1-10 kpc because as discussed in Section 4.3.4, that radial range is where the potential effect of SNIa on the entropy profile is best measured. We calculate the best fit parameters using the python package `emcee`. We establish an initial broad expected range for the parameters in log space with $0 < K_0 < 10^2$, $0 < K_{10} < 10^2$, and $0 < \alpha_K < 2$. Errors were determined from MCMC contours in two dimensions (16 – 84%).

4.3.3.1 Distribution of Central Entropy

Our entropy profile fits suggest that approximately half of the overall sample is clustered near $K_0 \sim 1 - 2 \text{ keV cm}^2$, while the rest have $K_0 \gtrsim 3 \text{ keV cm}^2$ (see Figure 4.2). In the group with low K_0 , the entropy profile at 1–10 kpc is close to a pure power law. This sub-sample is therefore

more suitable for testing the prediction represented in Equation 4.1. The subsample with greater K_0 , on the other hand, clearly deviates from the pure power-law entropy profile predicted by the basic model for an outflow heated by only SNIa. In those galaxies, central heating by the AGN may be producing an entropy floor at small radii, causing a break in the power-law profile. The best-fitting values of α_K in the high- K_0 sample may still be representative of SNIa heating, but the measurements of α_K are not as clean because of greater degeneracy between K_0 and α_K in the fitting procedure.

4.3.4 Sub-Sample Selection

During our analysis, we subdivided the full sample so as to test the prediction in Equation 4.1 as accurately as possible. Our smallest sub-sample with the most restrictive criteria represents the cleanest case to test the model predictions, and the full sample with the least restrictive criteria provides a more general lower limit on the value of α_K associated with a given σ_v . Our analysis requires a statistically significant measurement of the slope parameter α_K , so we only include those galaxies from the sample with sufficient resolution (at least 4 radial bins of any size) from 1–10 kpc. There are 36 galaxies from the sample that fit this criterion, hereafter referred to as the high quality (HQ) sample.

We implement two further limits on our sample for the analysis: central entropy, K_0 , and velocity dispersion, σ_v . Of the 36 galaxy profiles with sufficient data quality, we define a sample of 22 profiles with low central gas entropy (hereafter referred to as the low K_0 sub-sample). We describe the rationale for that selection in Section 4.3.3.1. From that sample, we filter galaxies with velocity dispersions of 210–310 km s⁻¹ (hereafter referred to as the restricted σ_v sample), leaving us with a sample of 16 galaxies. The velocity dispersion limit eliminates more high σ_v galaxies than low σ_v galaxies. We discuss the rationale for that selection in Section 4.4.1.2.

Table 4.2: **Entropy Profile Slope and Velocity Dispersion Relationship Results** Each sample is a subset of the previously listed samples. Descriptions of the sample are in Section 4.3.4. The fit is an ordinary least-squares fit to a linear model ($\alpha_K = A\sigma_{240} + B$, $\sigma_{240} \equiv \frac{\sigma_v}{240 \text{ km s}^{-1}}$) with intrinsic scatter (Akritas & Bereshady, 1996), and errors on the slope are 1σ . Column 1: Sample selection criteria; Column 2: Slope of ordinary least squares fit to the data; Column 3: Reduced chi squared for the fit; Column 3: Intrinsic scatter and error; Column 5: Number of galaxies included in the fit.

Sample	Slope	Reduced χ^2	Intrinsic Scatter	Number
High Quality (HQ)	0.52 ± 0.24	1.03	0.18 ± 0.02	36
Low K_0	0.80 ± 0.33	1.05	0.22 ± 0.03	22
restricted σ_v	1.80 ± 0.51	1.08	0.16 ± 0.02	16

4.4 Discussion

4.4.1 Low Central Entropy, Restricted σ_v , and the Analytical Prediction

4.4.1.1 The Black-hole Feedback Valve Prediction

The Voit et al. (2020) analytic model predicts a relationship for stellar velocity dispersion and entropy profile slope (Equation 4.1) in the radial range where SNIa heating is significant (1–10 kpc). Elevated central entropy, K_0 , beyond 1 kpc, suggests that the central AGN is more strongly coupled to the surrounding medium. Therefore, SNIa heating is not the dominant heating process, and the model is not expected to apply. Because the model applies best to galaxies without elevated central entropy, we will investigate the relationship between velocity dispersion and entropy profile slope for that particular sub-sample.

4.4.1.2 Comparison to the Analytic Prediction

With the criteria explained in Section 4.3, we examine the relationship between the stellar velocity dispersion, σ_v , and the entropy profile slope, α_K . Table 4.2 summarizes the results for our exploration of the relationship between velocity dispersion and entropy profile slope for the three samples. To quantify the potential relationship between α_K and σ_v , we assume that the relationship is approximately linear, $\alpha_K = A\sigma_{240} + B$, with intrinsic scatter and determine the strength of the

relationship by fitting a linear model with intrinsic scatter to the entropy profile slope, α_K , versus the scaled velocity dispersion, σ_{240} , with ordinary least squares (Akritas & Bershadsky, 1996). We use the same process for each sub-sample of galaxies, limited via criteria discussed in 4.3.4.

We first show that, with the requirement for radial resolution and fitting entropy profile slope between 1-10 kpc (see Section 4.3.3 for the entropy profile fitting procedure), some relationship emerges (see Figure 4.3). The slope of the relation is 0.53 ± 0.27 , so while the slope is only about 2σ away from 0, there is some evidence for relationship between σ_v and α_K before any additional limits were placed on the sample. Limiting the radial range of the entropy profile fit and requiring sufficient data resolution over that radial range clearly reduces some of the ambiguity found in previous work (e.g. Babyk et al. 2018; Lakhchaura et al. 2018).

The main sample generally contains massive elliptical galaxies, but there are some that are not necessarily representative of the galaxies the Voit et al. (2020) model sets out to describe. Figure 4.4 shows the fit for the the low K_0 sub-sample determined by the criteria discussed in Section 4.3.3. The slope is 0.80 ± 0.33 , so we see mildly stronger evidence for a relationship when the sample is limited to those galaxies without elevated central entropy.

We determined our final sub-sample, the restricted σ_v sample, shown in Figure 4.4 using the σ_v limiting stated in Section 4.3.4. Galaxies with $\sigma_v < 210 \text{ km s}^{-1}$ may not yet have a well-developed and sufficiently hot circumgalactic medium, so there is no reason to believe that the analytic model would apply. We limit $\sigma_v < 310 \text{ km s}^{-1}$ on the upper end for two reasons. (1) The Voit et al. (2020) model assumes a singular isothermal sphere which simplifies the mass profile (see Section 4.3.2), resulting in the analytic model overpredicting the entropy profile slope for galaxies with high σ_v . (2) Some galaxies in the sample, like M87, are in galaxy groups or clusters, and thus are in a potential well with a stellar velocity dispersion significantly greater than that of the central galaxy, resulting in a shallower entropy profile slope than predicted for an isolated galaxy. Limiting the velocity dispersion in this way limits the sample to galaxies most representative of the restricted σ_v the analytic model describes.

For the restricted σ_v sample, we find a slope of 1.80 ± 0.51 , or $\sim 3\sigma$ away from a flat line.

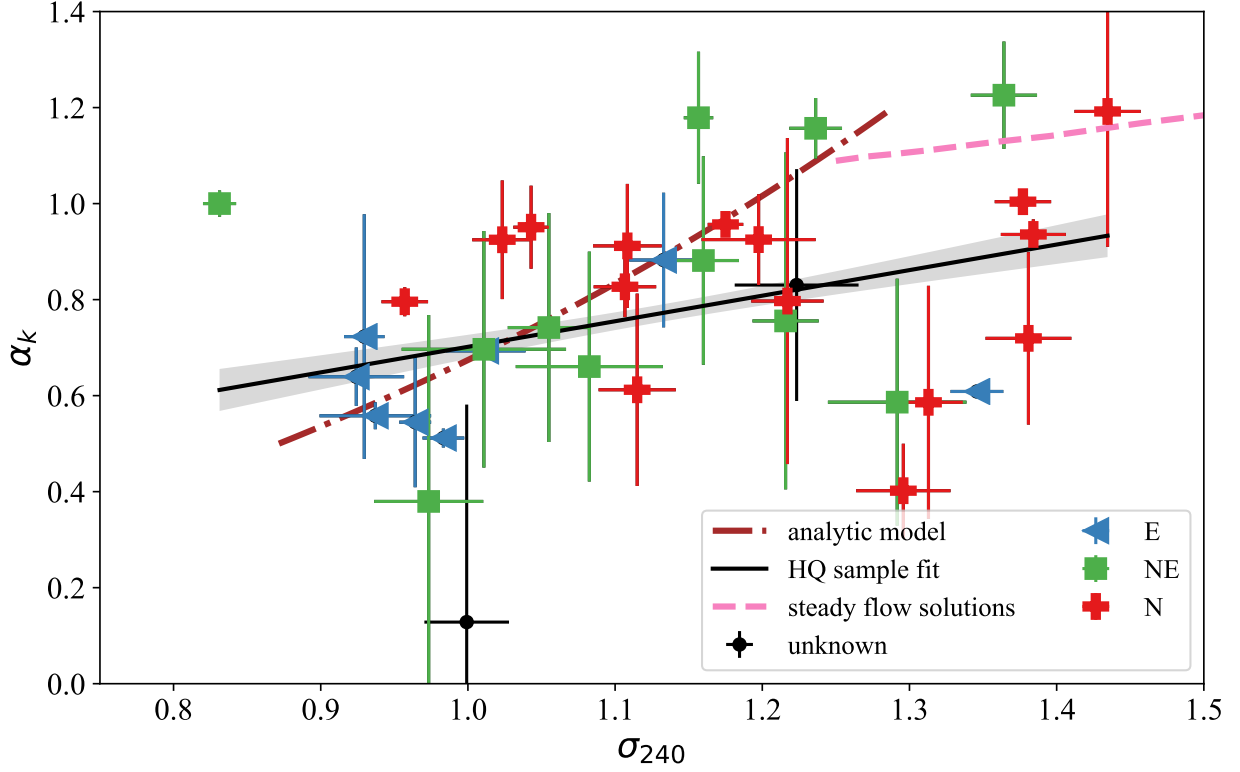


Figure 4.3: **Stellar velocity dispersion vs. entropy profile slope for the HQ sub-sample of galaxies** Scaled stellar velocity dispersion, $\sigma_{240} \equiv \frac{\sigma_v}{240 \text{ km s}^{-1}}$, is plotted with the entropy profile slope, α_K for the sub-sample of galaxies from Lakhchaura et al. (2018) with sufficient data between 1–10 kpc. The points are also classified by their multiphase gas extent from Lakhchaura et al. (2018). Blue triangles are galaxies with extended multiphase gas, red crosses are galaxies with no extended multiphase gas, green squares are galaxies with multiphase gas contained within 2 kpc, and black dots are galaxies without a gas extent classification. The black line is the ordinary least squares fit to the data, and the grey band is the 1σ error. The pink dashed lines represents the steady flow solutions for $\sigma_v \gtrsim 300 \text{ km s}^{-1}$.

The criteria we have applied limit the main sample to those galaxies that are most likely to follow the analytic model, so it is not surprising that the evidence for a relationship is stronger when the sample is limited to the galaxies to which the model is expected to apply: velocity dispersion-limited galaxies with limited direct central coupling ($r < 10 \text{ kpc}$) by the central AGN, as shown by lack of central entropy elevations or inversions in the entropy profile. However, the relationship between entropy profile slope and velocity dispersion may be stronger than indicated by previous works.

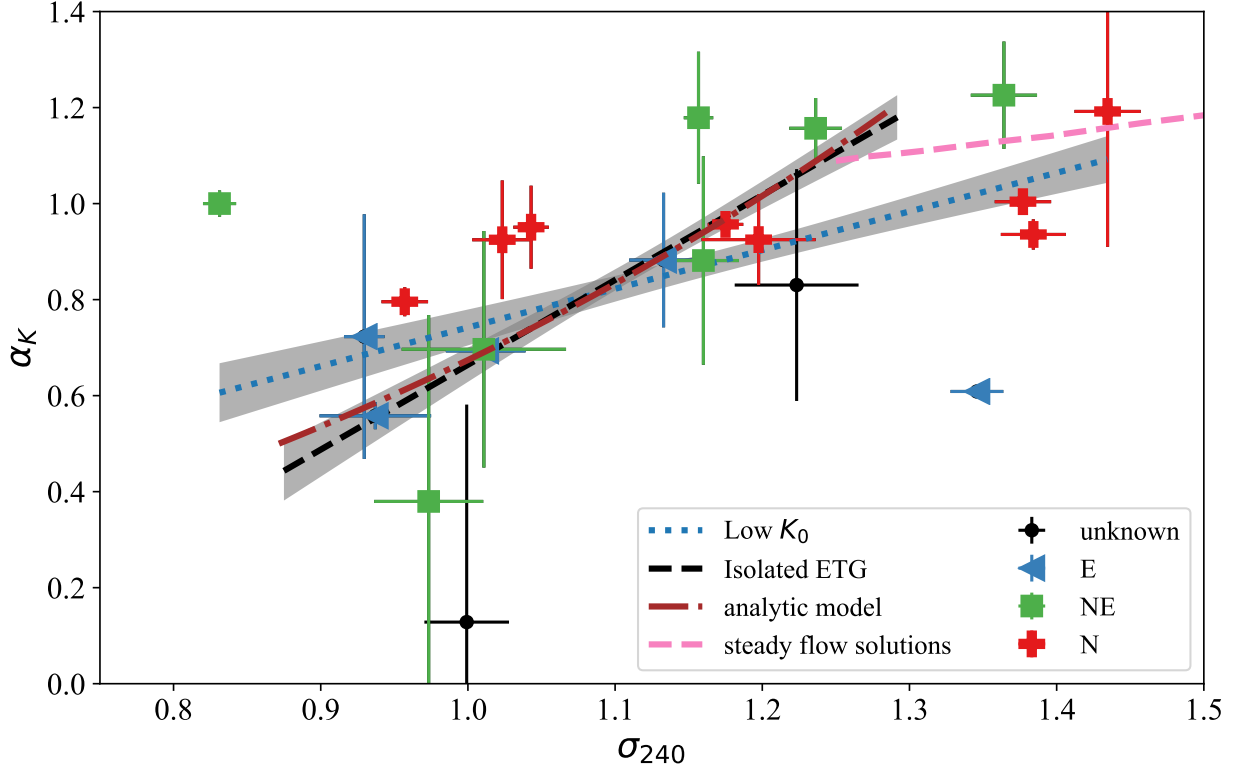


Figure 4.4: **Stellar velocity dispersion vs. entropy profile slope for the low K_0 sub-sample of galaxies** Scaled stellar velocity dispersion, $\sigma_{240} \equiv \frac{\sigma_v}{240 \text{ km s}^{-1}}$, is plotted with the entropy profile slope, α_K , for the subsample of galaxies with $K_0 < 3 \text{ keV cm}^2$, and ordinary least squares fits are given for both the for low K_0 subsample ($K_0 < 3 \text{ keV cm}^2$) and the restricted σ_v subsample ($210 \text{ km s}^{-1} < \sigma_v < 310 \text{ km s}^{-1}$, $K_0 < 3 \text{ keV cm}^2$). The blue dotted line is the fit to the low K_0 sample, and the black dashed line is the fit to the σ_v limited sample. The maroon dash-dotted line is the analytic solution from Voit et al. (2020). The pink dashed lines represents the steady flow solutions for $\sigma_v \gtrsim 300 \text{ km s}^{-1}$. The grey bands are 1σ errors.

4.4.2 Comparison to the Analytic Model and Numerical Integration Results

Figure 4.4 also includes the analytic model plotted with the fit to the K_0 and σ_v limited sample. We find very good agreement over the range of the “best case” fit. Furthermore, as discussed in Section 4.3.2, the analytic model over-predicts the entropy profile slopes for galaxies with higher σ_v , so the numerical integration results may actually better represent the data for $\sigma_v > 310 \text{ km s}^{-1}$ than the analytic model.

The analytic prediction for the relation between σ_v and α_K from Equation 6 of Voit et al. (2020) is a good fit when the data are restricted to the $210 \text{ km s}^{-1} < \sigma_v < 310 \text{ km s}^{-1}$ interval.

However, the 9 points at $\sigma_v > 310 \text{ km s}^{-1}$ all fall below the model prediction and the 1 point at $\sigma_v < 210 \text{ km s}^{-1}$ (NGC 4636) is above it. A closer look at the modeling shows that Equation 6 overpredicts α_K at $\sigma_v < 300 \text{ km s}^{-1}$, compared to the numerical steady flow solutions, shown by the pink dashed line in Figure 4.6.

4.4.3 Best fit Entropy Profile Slope, Multiphase gas extent, and $\min(t_{\text{cool}}/t_{\text{ff}})$

The free-fall time, $t_{\text{ff}} = (2r/g)^{1/2}$, where r is the galactocentric radius and g is the local gravitational acceleration, provides a dynamical timescale to characterize the gravitationally-driven motions of the gas and is based on observations of the stellar light from the galaxy. For our purposes, we will use the equivalent form of $t_{\text{ff}} = r/\sigma_v$. The cooling time is defined as:

$$t_{\text{cool}} \equiv \frac{3}{2} \frac{nkT}{n_e n_H \Lambda(T, Z)}, \quad (4.3)$$

where n is the total number density of particles, n_e is the electron density, n_p is the hydrogen density (where we assume $n_H = n_e/1.2$), and $\Lambda(T, Z)$ is the temperature dependent cooling function for plasma of metallicity Z . The ratio of the cooling time to the free-fall time ($t_{\text{cool}}/t_{\text{ff}}$) indicates if precipitation occurs in the ambient gas. For $t_{\text{cool}}/t_{\text{ff}} \cong 5 - 20$, the galaxy is said to be in the precipitation zone, where multiphase gas is found (Voit et al., 2015b). Voit et al. (2015b) also showed that from 1–10 kpc, galaxies with extended multiphase gas generally track the precipitation zone whereas galaxies without extended multiphase gas generally remain above the precipitation zone. The minimum value of $t_{\text{cool}}/t_{\text{ff}}$ is anti-correlated with the presence of multiphase gas (Voit & Donahue, 2015), so we expect galaxies with greater $\min(t_{\text{cool}}/t_{\text{ff}})$ to have little to no multiphase gas in their centers.

The expectation from theoretical models is that the entropy profile slope should correlate with multiphase gas extent, with galaxies with no extended multiphase gas having entropy profile slopes of $K \sim r$ and galaxies with extended multiphase gas having $K \sim r^{2/3}$. Furthermore, $t_{\text{cool}}/t_{\text{ff}}$ is coupled with the entropy profile because $K = T/n_e^{2/3}$ and $t_{\text{cool}} \propto T/n_e \Lambda(T)$, so entropy and cooling time are approximately proportional to each other. Therefore, profiles with higher $\min(t_{\text{cool}}/t_{\text{ff}})$ have higher entropy gas with longer cooling times.

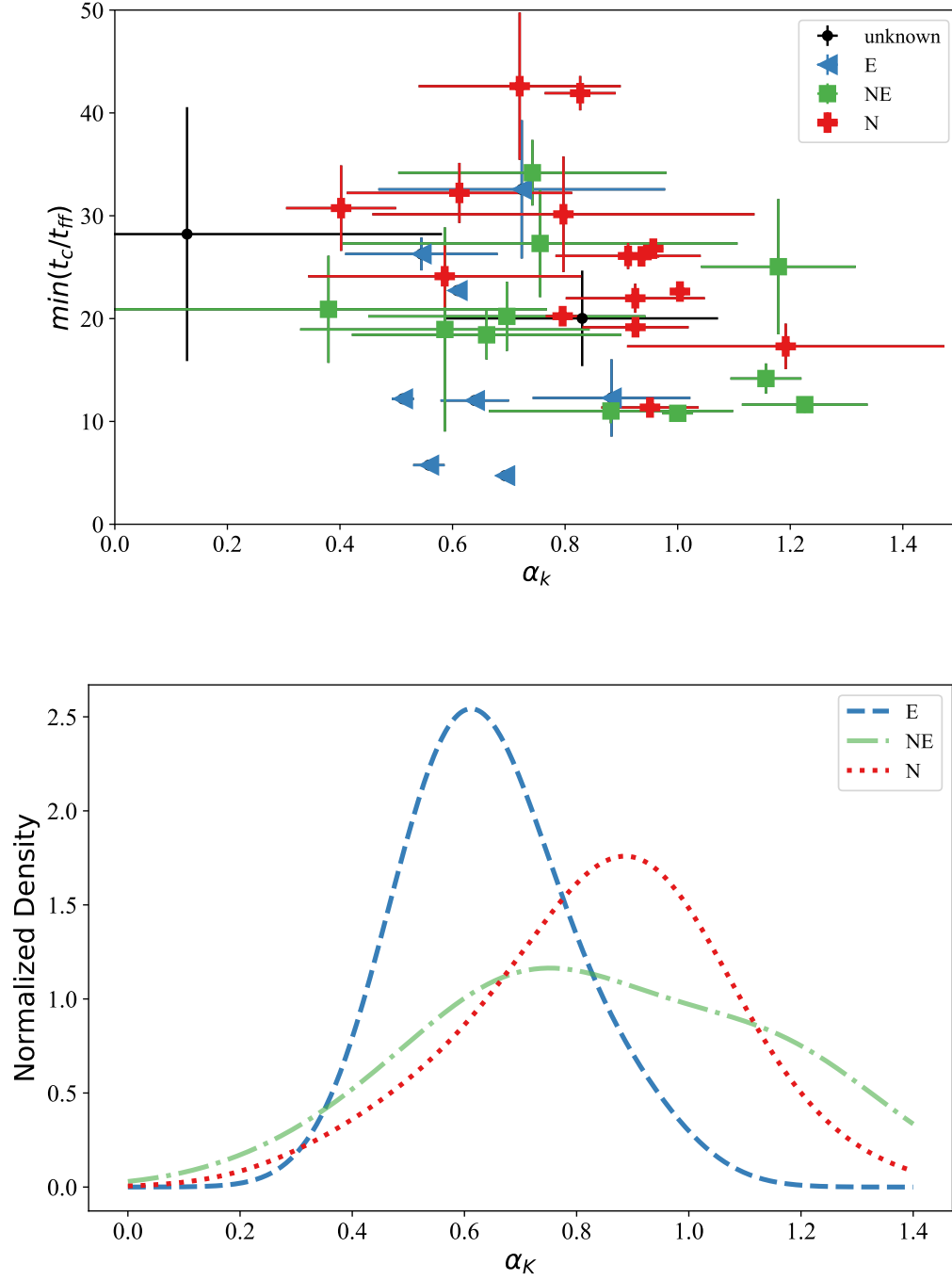


Figure 4.5: α_K vs. $\min(t_{\text{cool}}/t_{\text{ff}})$ for the main sample and α_K by gas extent for the HQ sample. Top: α_K vs. $\min(t_{\text{cool}}/t_{\text{ff}})$ for the main sample. Bottom: Gaussian Kernel Density Estimation of α_K by gas extent for the High Quality sample. Multiphase gas classifications are from Lakhchaura et al. (2018). Blue triangles are galaxies with extended multiphase gas, green squares are galaxies with extended multiphase gas that does not extend past 2 kpc, red crosses are galaxies with no extended multiphase gas, and black dots are galaxies with no multiphase gas extent classification. The colors indicated each category of multiphase gas extent are the same in the bottom plot. Bandwidths used for KDE were E: 0.115, NE: 0.192, and N: 0.152.

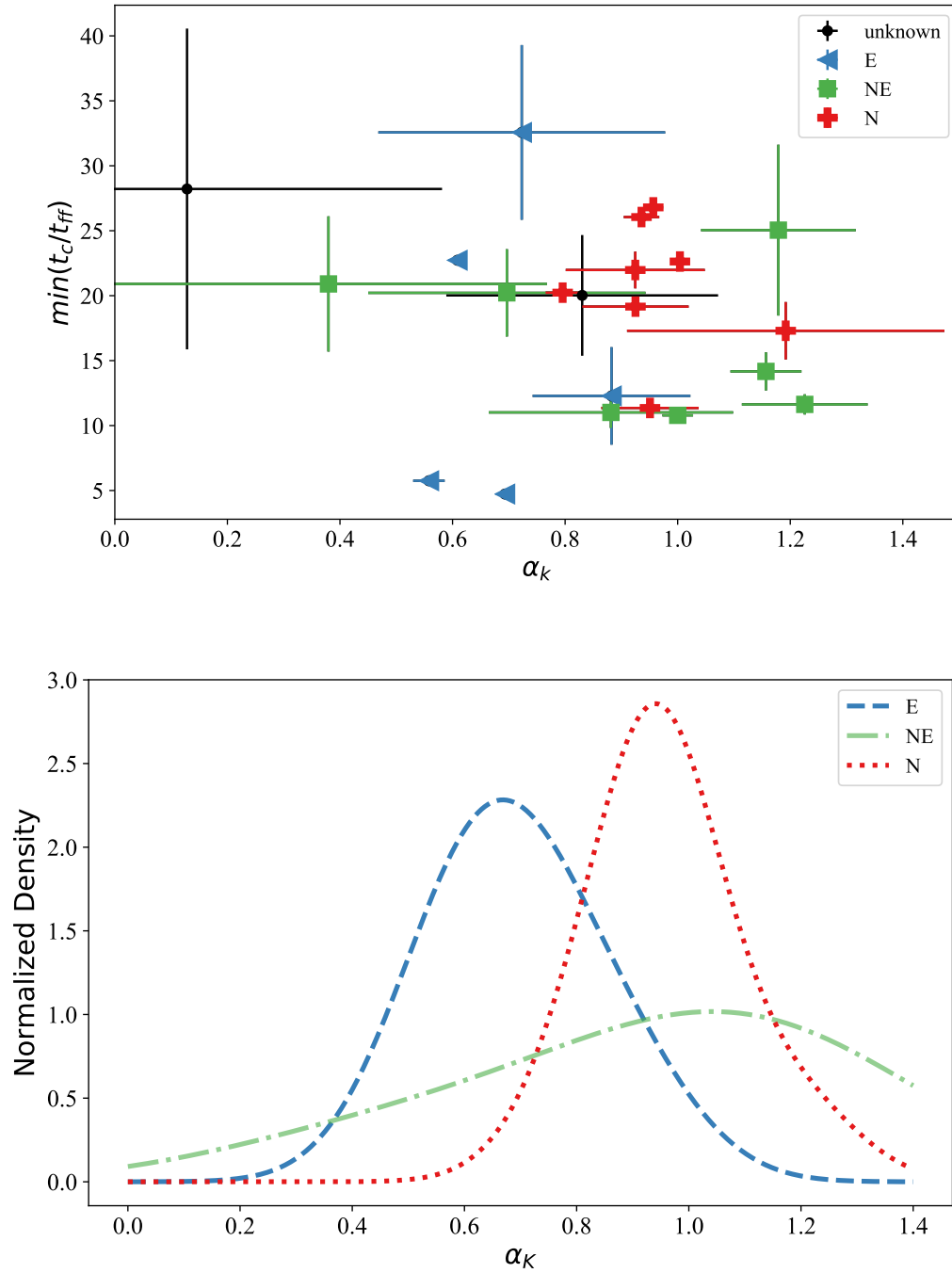


Figure 4.6: α_K vs. $\min(t_{\text{cool}}/t_{\text{ff}})$ and α_K by gas extent for the low K_0 sample Top: α_K vs. $\min(t_{\text{cool}}/t_{\text{ff}})$ Bottom: Gaussian Kernel Density Estimation of α_K by gas extent for the low K_0 sample. Blue triangles are galaxies with extended multiphase gas, green squares are galaxies with extended multiphase gas that does not extended past 2 kpc, red crosses are galaxies with no extended multiphase gas, and black dots are galaxies with no multiphase gas extent classification. The colors indicated each category of multiphase gas extent are the same in the bottom plot. Bandwidths used for KDE were E: 0.132, NE: 0.278, and N: 0.110.

We can expand our understanding of the relationships between galaxy parameters by exploring $\min(t_{\text{cool}}/t_{\text{ff}})$ for both the high quality sample and the sample limited to galaxies with $K_0 < 3 \text{ keV cm}^2$, along with the multiphase gas extent and entropy profile slope, α_K . Figure 4.5 and Figure 4.6 summarize this relationship for the high quality sample and the low K_0 sample, respectively. For the analysis of the relationship between multiphase gas extent and α_K , we employ Kernel Density Estimation to capture the distribution of multiphase gas extent with α_K . Kernel Density Estimation (KDE) convolves the discrete data with a smooth kernel function, in this case, a one-dimensional Gaussian with constant bandwidth. We choose to use KDE because it captures the continuum behavior of the distribution as well as behavior of the distribution that could be masked by choice of bin size in a histogram. The kernel is the function used to take the average of neighboring points, thereby making the distribution continuous instead of discrete. We chose the Gaussian kernel because the data are relatively simple and one-dimensional. The bandwidths were determined for each multiphase gas extent category in each sub-sample by minimizing the mean integrated square error. The bandwidth size is driven by minimizing the mean integrated square error (MISE) by minimizing the integral:

$$MISE(\hat{f}_{\text{kernel}}) = E \int [\hat{f}_{\text{kernel}}(x) - f(x)]^2 dx, \quad (4.4)$$

where $\hat{f}_{\text{kernel}}$ is the chosen kernel function, E represents the expected or mean value, and $f(x)$ is the underlying probability density function. For the KDE shown in the bottom panel of Figures 4.5 and 4.6, we used the KDE tools from the Python package `sci-kit learn` (Pedregosa et al., 2011).

In Figure 4.5, we consider the high quality sample, and we find that galaxies with extended multiphase gas (categorized as “E”) have entropy profile slopes close to $r^{2/3}$, but the galaxies with no extended multiphase gas (categorized as “N”), that we would expect to have entropy profiles slopes around r , extend down to lower entropy profile slopes. However, when we examine the low K_0 sample in Figure 4.6, we find that removing galaxies with elevated central entropy outside 1 kpc effectively removes the galaxies with no multiphase gas and lower entropy profile slopes. Furthermore, we see that removing galaxies with elevated central entropy also removes many of the galaxies with greater $\min(t_{\text{cool}}/t_{\text{ff}})$, likely because the feedback increases the cooling time.

If we consider the high quality sample, we find that the galaxies with high σ_v and no extended multiphase gas extend down to lower entropy profile slopes. Taken together, the two samples show that elevated K_0 outside 1 kpc and higher $\min(t_c/t_{\text{ff}})$ can serve as flags for galaxies that are not representative of SN-heated outflows, but rather have been flattened by feedback, allowing us to better test the analytical model.

4.4.4 Comments on Individual Galaxies

In Figure 4.3, some galaxies stand out as not conforming to the model. Because of the simplicity of the analytic model, we do not necessarily expect it to apply to all galaxies, particularly those with more complex environments. In this section, we discuss those galaxies and what particular characteristics of their environments may explain why they do not conform to the model. NGC 533 conforms to the analytical prediction by entropy profile slope and velocity dispersion, but its multiphase gas extent does not. Other galaxies (M87, NGC 4636, NGC 1521, NGC 1404, and NGC 4125) do not conform to the analytical prediction. Here, we present possible explanations for these notable exceptions to the model.

4.4.4.1 M87

M87 has high σ_v , but an entropy profile near $\alpha_K = 2/3$. However, we do not expect M87 to conform to the model because it is in a potential well with a velocity dispersion significantly greater than the stellar velocity dispersion of the central galaxy, though α_K is consistent with a galaxy near the precipitation limit.

4.4.4.2 NGC 4636

NGC 4636, has $\alpha_K \sim 1$ but is classified as having multiphase gas present inside 2 kpc. The entropy profile is consistent with the precipitation limit from 0.5–8 kpc, but is also consistent with a pure cooling flow from 0–2 kpc (Voit et al., 2020). Voit et al. (2020) also states that at smaller radii, the entropy profile flattens, relative to the the $K \propto r^{2/3}$ precipitation-limited profile, but

reaches $\sim 1 \text{ keV cm}^2$ inside of 100 pc, considerably below the level expected from $\sim 10^{42} \text{ erg s}^{-1}$ of intermittent kinetic feedback power. There are several possible explanations for this low central entropy level: (1) time-averaged kinetic AGN power has been $\sim 10^{41} \text{ erg s}^{-1}$ for the last ~ 100 Myr, (2) the AGN power has been highly collimated, as in NGC 4261, and has penetrated to $\gg 1$ kpc without dissipating much power, (3) AGN power has been too weak to balance cooling for the last ~ 100 Myr. In this last case, a cooling catastrophe is imminent, as suggested by the entropy profile between 0.5 and 2 kpc, and will soon trigger a strong feedback episode.

4.4.4.3 NGC 1521

NGC 1521 does not conform to the model, most likely because of low spatial resolution. The entropy profile only has four radial bins between 1–10 kpc, and only one additional radial bin, interior to 1 kpc. The best fit entropy profile slope has a large uncertainty, and the uncertainty does overlap the analytic prediction. Therefore, improved spatial resolution is necessary to determine if the galaxy conforms to the model.

4.4.4.4 NGC 4125

NGC 4125 does not conform to the analytic model, but like NGC 1521, may not conform due to spatial resolution. However, the shape of the entropy profile and the galaxy parameters are a bit more interesting. The X-ray luminosity (measured inside 10 kpc) is the lowest in the HQ sample ($0.023 \pm 0.001 \times 10^{42} \text{ erg s}^{-1}$). Between 1–10 kpc, the entropy profile is almost flat and based on six radial bins, resulting in a low best fit entropy profile slope (and larger uncertainty). However, interior to 1 kpc, the entropy profile slope is much steeper, and Lakhchaura et al. (2018) find a power law component of the spectrum, indicating the presence of an AGN with luminosity $0.006 \pm 0.001 \times 10^{41} \text{ erg/s}$. Wiklind et al. (1995) detected an upper limit for the molecular gas content, but the measurement is uncertain due to high systematic errors. The combination of the presence of an AGN, $\sigma_v < 240 \text{ km s}^{-1}$, and the flattened entropy profile at larger radii may indicate that this is a galaxy where the steady flow is cooling dominated at larger radii and prone to

developing entropy inversions (Voit et al., 2020).

4.4.4.5 NGC 1404

NGC 1404 is not far from the analytic prediction, but it has low σ_v for its α_K because it has an entropy profile with a sharp increase in slope beyond 7 kpc. It is a satellite of NGC 1399, so the sharp increase in the entropy profile could potentially be a result of ram-pressure stripping by the IGM around NGC 1399.

4.4.4.6 NGC 533

NGC 533 has $\sigma_v = 272 \text{ km s}^{-1}$ and $\alpha_K = 0.9$. The velocity dispersion of the surrounding galaxies is $\sim 464 \text{ km s}^{-1}$, according to Zabludoff & Mulchaey (1998), so we would not expect the basic model to apply.

4.4.5 Predictions for Equilibrium Pressure and Density Profiles

The analytic model of Voit et al. (2020) explores the behavior of the heating/cooling equality based on the Black-hole feedback valve model. Here, we present an observational test of the derived equilibrium profiles with the pressure and density profiles from Lakhchaura et al. (2018). An entropy profile slope of $\alpha_K \approx 2/3$ is a critical slope for the analytic model, meaning that the behavior of the galactic outflows should be fundamentally different above and below $\alpha_K \approx 2/3$. Voit et al. (2020) shows that the ratio of stellar heating to radiative cooling decreases with radius for galaxies with an entropy profile slope below $\alpha_K \approx 2/3$ and rises with radius for galaxies with an entropy profile slope above $\alpha_K \approx 2/3$. Following from Equation 4.1, the velocity dispersion corresponding to this critical entropy profile slope is $\approx 240 \text{ km s}^{-1}$.

When radiative cooling per unit volume equals stellar heating per unit volume, for a given radius, they find an equilibrium pressure profile along which supernova heating equals radiative cooling for a temperature, T (Equation 11 in Voit et al. (2020)):

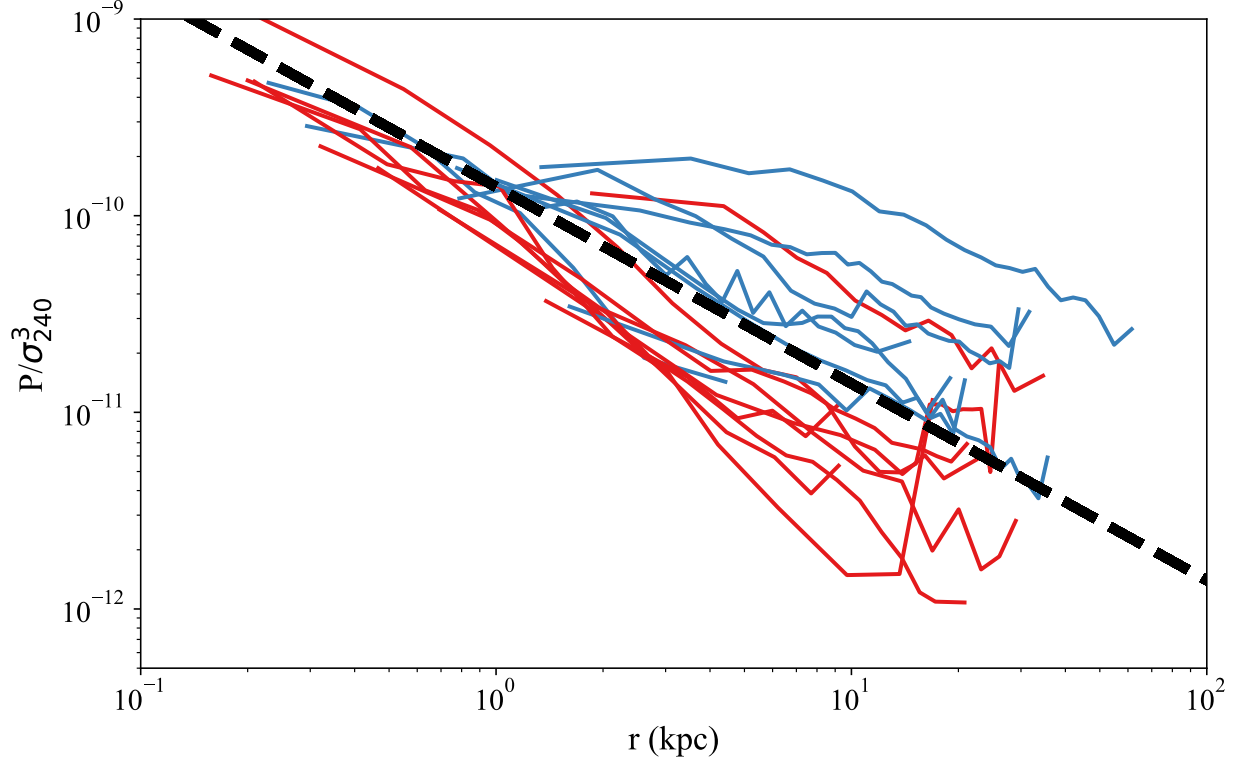


Figure 4.7: **Equilibrium pressure vs. radius for single phase and multiphase galaxies in the HQ sample** The equilibrium pressure at temperature T , for radius r (black dashed line), is plotted with the extended multiphase (E, blue lines) and single-phase (N, red lines) galaxies in the HQ sample. Errors bars are removed from the profiles for clarity. The blue line below the black dashed line is NGC 1316, and the red line above the black dashed line is NGC 4073. See Section 4.4.5 for a discussion of these two galaxies.

$$P_{\text{eq}}(r) \equiv \left[\left(\epsilon_* + \frac{3}{2} \sigma_v^2 \right) \left(\frac{n^2}{n_e n_p} \right) \frac{\rho_*}{t_* \Lambda(T)} \right]^{1/2} kT, \quad (4.5)$$

where n_p is the proton density, ρ_* is the stellar mass density, t_*^{-1} is the specific stellar mass-loss rate, and $\Lambda(T)$ is the radiative cooling function. For the velocity dispersion and temperature corresponding to the critical entropy profile slope ($\sigma_v \approx 240 \text{ km s}^{-1}$, $kT \approx 0.75 \text{ keV}$, $\alpha_K \approx 2/3$), the critical profiles are as follows (Equations 12 and 13 in Voit et al. (2020)):

$$P_{\text{eq}}(r) \approx (1.4 \times 10^{-10} \text{ erg cm}^{-3}) \sigma_{240}^3 r_{\text{kpc}}^{-1} \quad (4.6)$$

$$n_{e,\text{eq}} \approx (0.06 \text{ cm}^{-3}) \sigma_{240} r_{\text{kpc}}^{-1}, \quad (4.7)$$

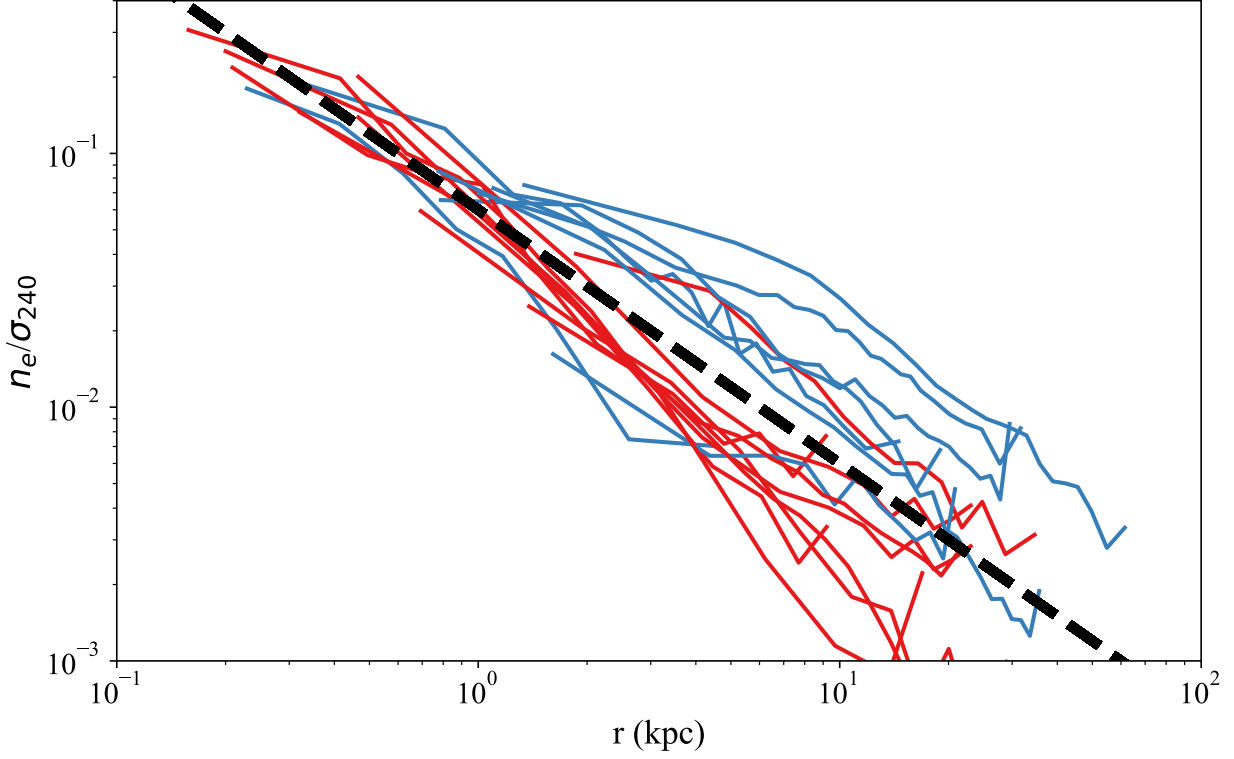


Figure 4.8: **Equilibrium electron density vs. radius for single phase and multiphase galaxies in the HQ sample** The equilibrium electron density at temperature T , for radius r (black dashed line), is plotted with the extended multiphase (E, blue lines) and single-phase (N, red lines) galaxies in the HQ sample. Errors bars are removed from the profiles for clarity. The blue line below the black dashed line is NGC 1316, and the red line above the black dashed line is NGC 4073. See Section 4.4.5 for a discussion of these two galaxies.

where $r_{\text{kpc}} \equiv r/1 \text{ kpc}$, $\sigma_{240} \equiv \sigma_v/240 \text{ km s}^{-1}$, $\rho_* = \sigma_v/2\pi G r^2$, the isothermal stellar mass distribution, and the fiducial values $\mu m_p \epsilon_* \approx 2 \text{ keV}$ and $t_* \approx 200 \text{ Gyr}$, if the weak dependence of $\Lambda(T)$ on σ_v is ignored.

Figures 4.7 and 4.8 show the comparison of the extended multiphase (E) and single-phase (N) galaxies in our sample to the equilibrium pressure and density profiles. The galaxies with multiphase gas confined to the central 2 kpc have been removed for clarity. The model predicts that the equilibrium profiles should divide the profiles of galaxies with extended multiphase gas ($\alpha_K \lesssim 2/3$) from the galaxies with no extended multiphase gas (generally higher $\alpha_K \gtrsim 2/3$). We find that overall, the equilibrium profiles for both P_{eq} and $n_{e,\text{eq}}$ do indeed divide our sample as predicted. However, there are two notable exceptions; one each from the multiphase and

single-phase galaxies. The multiphase galaxy, NGC 1316, does conform to the analytic model within uncertainty, but the entropy profile exhibits an inversion at $r \sim 2.5$ kpc, an entropy profile characteristic of massive elliptical galaxies with $\sigma_v \lesssim 240 \text{ km s}^{-1}$ predicted by Voit et al. (2020). NGC 1316 is also one of the lowest luminosity galaxies in the sample. The single-phase galaxy, NGC 4073, has one of the highest luminosities and one of the highest temperatures in the sample. It is classified as a single-phase galaxy but has $\alpha_K \approx 0.6$ and $\sigma_v \approx 268 \text{ km s}^{-1}$ and does not conform to the analytic model within uncertainty.

4.5 Conclusions

In this work, we were able to show that not only is there evidence for a relationship between the stellar velocity dispersion, σ_v , and the entropy profile slope, α_K , the relationship agrees with the analytic model proposed in Voit et al. (2020). In contrast to previous analyses of this relation, we applied limits to data quality of the archival observations as well as limits on the parameters explored as informed by the data and the analytic model. While the results from the sample limited by both K_0 and σ_v are a more promising comparison to the analytic model, we still see evidence for a relationship between α_K and σ_v for the samples with fewer limits applied. Furthermore, results from the numerical integration of the analytic model suggest that the data may agree with the model for higher σ_v as well. For galaxies in groups with much lower entropy profile slopes than predicted for their velocity dispersion, Voit et al. (2020) proposes that the entropy profile may have a slope of $\alpha_K = 2/3$ that is more representative of a cool-core cluster with extended multiphase gas or galaxies with $200 \text{ km s}^{-1} < \sigma_v < 240 \text{ km s}^{-1}$.

When we set out to characterize the sample by central entropy K_0 , we found that there were two populations of galaxies that we could separate by applying a limit of $K_0 < 3 \text{ keV cm}^2$. Those galaxies with $K_0 > 3 \text{ keV cm}^2$ that were removed from the sample are likely galaxies that have experienced recent feedback, elevating their entropy out to larger radii. When we explored the $\min(t_{\text{cool}}/t_{\text{ff}})$, multiphase gas extent, and entropy profiles slopes of those galaxies as well, we found that galaxies with no extended multiphase gas and lower entropy profile slopes than expected also typically had higher $\min(t_{\text{cool}}/t_{\text{ff}})$, providing further evidence for recent feedback causing lower

α_K . However, we also note that the galaxies with multiphase gas present but inside 2 kpc remain spread across the range of α_K , though more with lower α_K were removed than with higher α_K , indicating that some entropy profiles may be flattened due to feedback, but the effect is less clear. The analytic model requires the galaxies to be in equilibrium, so it is not surprising that galaxies out of equilibrium do not follow it, but model does agree with galaxies close to equilibrium.

Our work shows that while the Voit et al. (2020) analytic model may be relatively simple, it describes the relationship between key galaxy parameters well and can be used to further our understanding of how feedback in massive galaxies works. The comparison of the model to the data supports the notion that SNIa supernova feedback plays an important role in the thermal evolution of massive galaxies. Furthermore, the relationship between entropy profile slope and velocity dispersion is highly dependent on the external gas pressure at larger radii. Current X-ray observations are not able to resolve pressure measurements at large radii, but *Athena* and *LYNX* may be able to. Taking the model predictions and existing observations, one could predict what the gas pressure at large radii and then test that prediction with the next generation of X-ray telescopes.

CHAPTER 5

SUMMARY

5.1 Summary

This dissertation sought to add to our understanding of the thermal properties of galaxy clusters and early-type galaxies. Chapter 2 introduced the ACCEPT 2.0 database, presented my analysis of the entropy profiles for clusters with deprojected radial temperature and density profiles, and provided an example of the science applications of the ACCEPT 2.0 data products including central entropy and morphology measurements. I showed that ACCEPT 2.0 provides robust, uniformly reduced, deprojected entropy profiles and central entropy classifications as well as morphology measurements with great potential for scientific impact. The central entropy measurements I obtained show that ACCEPT 2.0 reproduces the distribution of central entropy, K_0 from ACCEPT and that distribution holds when all of the new clusters with central entropy measurements in ACCEPT 2.0 are introduced. Finally, ACCEPT 2.0 reproduces some of the early morphology work from Cassano et al. (2010) and provides meaningful insights about sample selection for the *XMM* Heritage sample.

In Chapter 3, I explored the properties of early-type galaxies with powerful radio sources. I found that there are other galaxies, like NGC 4261, with powerful radio sources and single power law entropy profiles, namely IC 4296 and potentially NGC 315. Furthermore, if the ratio of the cooling time to the free-fall time is lowest in the central radial bin, it may indicate the presence of a powerful radio source. Finally, when I compared the radial entropy profiles for NGC 4261 and IC 4296, along with the radial entropy profiles from Lakhchaura et al. (2018), to the simulations of Wang et al. (2019), I found good agreement in the general entropy profile slope behavior between observations and simulations for both single phase and multiphase galaxies.

In Chapter 4, I presented an observational test of the black-hole feedback valve model for galactic atmospheres from Voit et al. (2020) using the sample of early-type galaxies from Lakhchaura et al.

(2018). I found that equilibrium pressure and density profiles support the model prediction that galaxies above and below the critical values, $\alpha \sim 2/3$ and $\sigma_v \sim 240 \text{ km s}^{-1}$, behave differently. I also found that, when we select a sub-sample of galaxies from the original sample that the analytic model would be expected to apply to, based on velocity dispersion and central entropy measurements, there is a correlation between the entropy profile slope and velocity dispersion. The slope of the relation between velocity dispersion and entropy profile slope for this sub-sample is at the 3σ level and matches the analytic prediction well. When we broaden our analysis to include all galaxies in the sample with sufficient data resolution to obtain an entropy profile slope, including those that may not be expected to follow the analytic model, we still find some evidence for correlation, although at closer to 2σ significance.

5.2 Future Work

There are many possible avenues to explore from the work completed in this dissertation, both in the realm of galaxies and galaxy clusters. With the release of the ACCEPT 2.0 data products, it will be the largest uniformly reduced, publicly available database of cluster properties and thus can be used to gain insight into the systematics of other X-ray samples. As shown in Chapter 3, X-ray systematics are significant enough to affect conclusions made from data, so the uniform reduction of ACCEPT 2.0 is helpful for exploring the characteristics of large samples of galaxy clusters.

With respect to early-type galaxies, we have obtained additional X-ray data for IC 4296 from both *Chandra* and *XMM* that will hopefully allow us to better probe the inner $\sim \text{kpc}$ of the galaxy. In the more distant future, improved spatial resolution from an X-ray observatory like *LYNX* would allow us to examine the central kpc of the galaxies in our sample that would require prohibitively long observations to achieve the resolution of NGC 4261 with current telescopes. Looking ahead to *Athena* or *LYNX*, we could make predictions about the external gas pressure in early-type galaxies, based on the black-hole feedback valve model for galactic atmospheres, that could be tested with more sensitive X-ray telescopes.

APPENDICES

APPENDIX A

ACCEPT 2.0 PIPELINE DESCRIPTION

A.1 ACCEPT 2.0 Pipeline Details

A.1.1 Sample Selection and Data Processing and Analysis

This Appendix is provided to document choices made within the ACCEPT 2.0 pipeline. This text is heavily drawing on the text that the author team of the ACCEPT 2.0 data release paper, in draft, will be publishing in an Astrophysical Journal Supplement Series.

The ACCEPT 2.0 cluster sample was selected from all Chandra archival observations available as of July 2014. The automated pipeline (integrated with visual inspection of the products, when necessary) was developed to select the sample and perform the data processing and analysis. Using CIAO v4.7, CALDB 4.5, and SHERPA they ran an automated quick spectral analysis of all clusters available in the archive in order to have a rough estimate of the temperature of each cluster. The temperature estimate was used to set a count threshold to decide whether to include a cluster in the sample. To establish a 20% error threshold on the cluster temperature measurement from at least three spatial bins, they determined from simulations that the minimum number of counts required in the 0.5–7 keV band is given by $n_{min,res} = 1377 \cdot kT - 537$, where kT is the cluster temperature in keV. For a 20% error on the temperature in a single spatial bin, the number of counts required in the 0.5–7 keV band is given by $n_{min,glb} = (1377 \cdot kT - 537)/3$. They set $n_{min,glb}$ and $n_{min,res}$ as the minimum counts necessary to include a cluster in our total sample and in the spatially resolved sample, respectively. The total ACCEPT 2.0 sample consists of 606 clusters of which 402 are suitable for a spatially resolved analysis in at least three spatial bins. Of the clusters with spatially resolved analysis, 348 had sufficient counts for deprojection. Chandra data reprocessing was performed in an automated fashion using CIAO task *chandra_repro*, which applies the appropriate ACIS gain maps, the time-dependent ACIS gain correction, and the ACIS charge transfer inefficiency correction. The background light curve during each observation was

used to detect and remove periods of anomalously high background following the recommendations of Markevitch et al. (2003). The automated procedure to remove background flares was followed by a manual visual inspection of the light curves to check for undetected flares or excessive flare cleaning.

A.1.1.1 Initial Pipeline Spatial Analysis

Once clean event files for each ObsID of a cluster are obtained, they used the CIAO tool *merge_obs* to produce fluxed images and exposure maps for each ObsID and for the sum of all ObsIDs. Images and exposure maps are created with a binning factor of 2, corresponding to a pixel size in the image of 0.984 arcsec. At this stage, point sources are detected in the merged image using CIAO tool *wavdetect* in order to create a list of point sources that will be excluded from the spatial and spectral analysis of the cluster and to create an image of the diffuse emission using CIAO tool *dmfilth*. The point source list is also visually inspected using *DS9* in order to prevent false detections or excess exclusion of undetected point sources.

As a default, the X-ray emission peak was used as the center of the radial profiles, however the pipeline allowed manual inspection of the validity of this choice and to use the X-ray emission centroid if necessary. They initially considered concentric annuli with a thickness of 5 arcsec. If the cluster had at least 1500 counts in the 0.5-7 keV band, they grouped the annuli in the radial profile to have at least 300 counts per spatial bin prior to background subtraction. If the cluster had less than 1500 total counts, the minimum number of counts required for each spatial bin was lowered to one-fifth of the total counts in the cluster. The radial profiles extend out until the number of source counts in a given annulus reach 30% of the background counts in that annulus (or when an annulus reaches the chip boundary, in very bright and extended clusters).

The pipeline then performs a fit of the image from a 2-D Lorentz model with a varying power law (also known as a 2-D β model) using *Sherpa*. The function fitted to the image is:

$$f(x, y) = f(r) = A(1 + [r/r_0]^2)^{-\alpha} \quad (\text{A.1})$$

where $r(x, y) = \sqrt{[x_{new}^2(1 - \epsilon)^2 + y_{new}^2]/(1 - \epsilon)}$, $x_{new} = (x - x_o)\cos(\theta) + (y - y_o)\sin(\theta)$, and $y_{new} = (y - y_o)\cos(\theta) - (x - x_o)\sin(\theta)$. The most important parameters of this model are the core radius r_0 , the power-law index α , the ellipticity ϵ , and the angle of ellipticity θ . The last two parameters are used for a morphological analysis of the clusters in relation with other cluster properties (see Section 2.2.2.1). This spatial analysis procedure was repeated for every single cluster in the sample.

A.1.1.2 Global Spectral Extraction and Analysis

To analyze the global properties of each cluster, the pipeline performs a spectral extraction in three different spatial regions: the whole cluster ($r < r_{cluster}$), the cluster core ($r < r_{core}$), and the cluster with the core excised ($r_{core} < r < r_{cluster}$). They used the CIAO tool *specextract* to generate the spectra, appropriate redistribution matrix files (RMFs), and ancillary response files (ARFs). The maximum cluster extension $r_{cluster}$ has been set to $500 h^{-1}$ kpc. If the maximum radius cannot fit in the S2 chip, the cluster was observed only with ACIS-S, or was observed in all four chips of ACIS-I, if at least one of the observations was performed with ACIS-I, the largest radius that can fit in the ACIS field of view is used. The core radius was set by default to $70 h^{-1}$ kpc, but if this radius is larger than $0.3r_{cluster}$, the core radius is set to $r_{core} = 0.3r_{cluster}$. Detected point sources are excised from the spectra of the diffuse emission from the cluster. At this stage background spectra and radial profiles are also built.

A.1.1.2.1 Background Subtraction

The pipeline used the blank-field observations, processed identically to the cluster observations, and reprojected onto the sky using the aspect information from the cluster pointings. The synthetic backgrounds correspond to far longer exposure times (~ 0.5 Msec) than the majority of the cluster observations, giving a good estimate of the background. For clusters observed on ACIS-I, the blank-field background correction is renormalized to the background of the observation, using the ACIS-S2 chip, in a region of the ACIS field of view practically free from cluster emission. For

clusters observed with ACIS-S, the chip used for the renormalization is ACIS-S1. The energy band from 9.5 to 12 keV (mostly dominated by charged particles) is used to perform the normalization. Using the renormalized and reprojected background event files, the pipeline produces radial profiles of background counts and surface brightness for use in the analysis of the cluster surface brightness profiles.

A.1.1.2.2 Spectral Analysis

The spectra of the whole cluster, the cluster core, and the cluster with the core excised were analyzed in an automatic fashion by the pipeline using the Python tools within *Sherpa*. The spectra were analyzed in the 0.5-8 keV spectral band, and the source and background are fitted simultaneously. The source model used is a *mekal* model Kaastra et al. 1996; Liedahl et al. 1995 in which the ratio between the elements is fixed to the solar value as in Anders & Grevesse (1989). They considered line of sight absorption fixed at the Galactic value n_H (Stark et al., 1992), and an additional internal absorption component left free to vary (consistent with zero in the large majority of clusters). The free parameters in the *mekal* model are temperature, kT , metal abundance, Z , and normalization, η . The redshift, z , was fixed at the literature value for the cluster. The background model used was composed of two power-law models, several gaussian emission lines, and an *apec* thermal model at low temperature ($kT = 0.17$ keV, to account for the soft, diffuse X-ray background). The power-law slopes and the quantity, position, and strength of the emission lines depends on the specific ACIS chip used and are adjusted accordingly. The shape of the spectrum is held fixed.

A.1.1.3 Spatially Resolved Spectral Extraction and Analysis

For clusters with a sufficient number of counts and at least three spatial bins, a spatially-resolved spectral analysis was performed. The cluster was divided into concentric annuli that are required to contain at least $n_{min,res}/3$ counts per annulus for the projected spectral analysis and $n_{min,res}$ counts per annulus for the deprojected spectral analysis. In both analyses, the annuli have a minimum thickness of 5 arcsec and extend out until the source counts reach 50% of the background counts.

CIAO tool *specextract* was then used to generate the spectra of each annular region and their relative RMFs and ARFs.

The projected spectral analysis fits each annular region spectrum independently from the others. The X-ray band considered is 0.5-8 keV. Similarly to the global spectral analysis, source and background spectra are fitted simultaneously using the same models described in Section A.1.1.2.2. If at least three spatial annuli exist, the pipeline uses the *deproject* module in *Sherpa* to extract the deprojected temperature and density profiles (see Section 1.3.3 for a discussion of deprojection).

APPENDIX B

ACCEPT 2.0 CENTRAL ENTROPY FITTING RESULTS

Table B.1: **Fit parameters, redshift, and profile center for ACCEPT 2.0 clusters** Column 1: Cluster Name; Column 2: Profile center RA; Column 3: Profile Center DEC; Column 4: redshift; Column 5-6: Best fit K_0 and associated 1σ error; Column 7-8: Best fit K_{100} and associated 1σ error; Column 9-10: Best fit α and associated 1σ error;

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_2717	0.8040	-35.9334	0.0490	29.928	6.217	123.474	9.684	0.933	0.114
NSCS_J000619+105206	1.5845	10.8643	0.1670	58.428	11.089	135.616	17.235	0.908	0.084
ZwCl_0008.8+5215	2.8396	52.5292	0.1040	170.494	40.772	114.906	53.516	0.932	0.198
ABELL_2734_NED01	2.8401	-28.8540	0.0753	25.692	17.051	153.977	24.192	0.578	0.078
MACS_J0011.7-1523	2.9290	-15.3891	0.3780	7.906	6.688	139.323	14.306	0.861	0.076
ABELL_0013	3.4078	-19.5025	0.0943	130.014	31.300	156.989	51.950	0.783	0.219
ABELL_2744	3.5814	-30.3917	0.3080	151.716	90.749	240.949	140.399	0.480	0.240
Cl_0016+16	4.6396	16.4369	0.5410	165.464	19.834	56.168	16.004	1.319	0.126
MACS_J0025.4-1222	6.3737	-12.3761	0.5843	149.815	20.362	63.487	19.633	1.337	0.154
PLCKESZ_G304.84-41.42	7.0238	-75.6301	0.4100	-48.267	43.675	356.033	45.563	0.488	0.083
MCXC_J0035.4-2015	8.8607	-20.2628	0.3640	136.376	16.255	41.337	12.893	1.521	0.149
ABELL_2813	10.8537	-20.6236	0.2924	136.338	22.780	65.152	32.801	1.403	0.287
ZwCl_0040.8+2404	10.9678	24.4054	0.0830	9.048	1.165	113.822	3.727	1.136	0.044
ABELL_0098N	11.6029	20.6221	0.1043	10.598	5.928	179.603	13.006	0.982	0.097
ESO_351-_G_021	13.7493	-35.3209	0.0571	3.731	0.865	104.966	6.880	0.942	0.056
ABELL_0141	16.3933	-24.6330	0.2300	158.783	28.161	61.340	34.694	1.396	0.285
MaxBCG_J016.70077+01.05926	16.7042	1.0538	0.2539	11.751	0.707	68.442	2.757	1.517	0.047
CIZA_J0107.7+5408	16.9381	54.1327	0.1066	305.475	45.983	66.095	52.333	1.107	0.300
IC_1633	17.4809	-45.9308	0.0243	2.846	0.783	845.517	95.777	1.219	0.057
ABELL_2895	19.5469	-26.9653	0.2270	168.687	33.931	90.111	37.843	1.130	0.194
UGC_00842	19.7248	-1.0021	0.0452	-22.632	19.270	215.221	20.699	0.430	0.160

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_0193	21.2819	8.6992	0.0486	134.938	9.008	43.896	16.118	1.559	0.293
ABELL_0209	22.9690	-13.6108	0.2060	86.644	24.724	165.614	30.617	0.812	0.083
Abell_222	24.3940	-12.9930	0.2110	148.849	21.545	48.218	37.967	1.362	0.395
Abell_223	24.4832	-12.8196	0.2070	101.808	21.115	120.946	31.330	1.023	0.160
ABELL_0267	28.1761	1.0126	0.2310	148.399	16.487	54.930	15.020	1.745	0.148
ABELL_0262	28.1929	36.1533	0.0174	3.562	0.138	484.188	26.609	1.459	0.026
NGC_0741_GROUP	29.0876	5.6272	0.0185	1.343	0.211	595.269	91.607	1.369	0.065
NGC_0766	29.6769	8.3471	0.0270	1.208	0.676	294.920	298.944	1.338	0.299
GMBCG_J029.95560-08.83299	29.9560	-8.8336	0.3220	-10.805	5.842	217.850	13.766	0.837	0.049
NGC_0777	30.0620	31.4294	0.0167	4.297	0.287	552.908	227.639	1.838	0.117
ARP_318	32.4095	-10.1459	0.0132	5.359	0.210	157.740	311.625	1.810	0.141
MCXC_J0220.9-3829	35.2358	-38.4809	0.2287	18.431	3.386	110.540	7.654	1.049	0.059
ABELL_3017	36.4712	-41.9179	0.2195	36.916	7.731	109.848	14.716	1.258	0.086
MZ_10451	37.4404	-29.6300	0.0608	11.591	10.746	73.110	20.113	0.569	0.459
SPT-CL_J0232-4421	38.0782	-44.3463	0.2836	16.418	17.740	416.131	384.037	1.254	0.509
ABELL_0370	39.9715	-1.5769	0.3750	256.699	27.693	42.280	25.614	1.614	0.236
MACS_J0242.6-2132	40.6496	-21.5406	0.3140	9.402	1.573	85.748	5.660	1.302	0.059
ABELL_S0295	41.3602	-53.0297	0.3000	165.939	21.589	36.248	26.526	1.560	0.279
ABELL_0376	41.5164	36.9055	0.0484	61.935	16.644	178.494	32.288	0.869	0.157
ABELL_0383	42.0140	-3.5291	0.1871	10.280	1.047	117.585	3.732	1.253	0.041
NGC_1132	43.2163	-1.2740	0.0231	-1.648	0.843	112.294	3.305	0.663	0.035
ABELL_0402	44.4209	-22.1531	0.3224	115.241	24.785	87.153	30.745	1.202	0.192
ABELL_0400	44.4234	6.0243	0.0244	86.201	20.899	104.441	45.919	0.939	0.588
ABELL_0401	44.7372	13.5707	0.0737	154.597	14.171	115.707	17.379	0.940	0.098
MCXC_J0301.6+0155	45.4092	1.9205	0.1695	12.181	1.870	112.084	4.925	1.064	0.045
MCXC_J0303.7-7752	45.9387	-77.8798	0.2742	187.074	34.239	85.939	47.786	1.420	0.316
IC_1880_GROUP	46.6186	-9.7312	0.0338	4.003	1.229	342.980	75.624	1.588	0.154
ABELL_3088	46.7572	-28.6654	0.2534	20.390	9.298	178.149	18.196	0.886	0.074
MACS_J0308.9+2645	47.2330	26.7606	0.3240	144.011	31.449	106.556	32.825	1.111	0.141

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_3094	47.8995	-26.8988	0.0677	31.063	66.818	225.536	83.173	0.399	0.169
ABELL_3120	50.4853	-51.3265	0.0690	17.504	5.900	153.473	12.136	0.982	0.122
UGC_02748	51.9752	2.5616	0.0302	5.953	1.207	329.791	103.667	1.659	0.212
ABELL_3126	52.1524	-55.7180	0.0856	158.222	13.322	50.469	14.489	1.725	0.167
MACS_J0329.6-0211	52.4237	-2.1966	0.4500	4.991	2.268	118.081	5.828	1.095	0.042
ABELL_3128	52.4605	-52.5801	0.0599	12.128	60.726	200.471	72.740	0.341	0.215
MCXC_J0331.1-2100	52.7751	-21.0092	0.1880	11.718	1.574	114.991	5.396	1.252	0.053
3C_089	53.5625	-1.1882	0.1386	30.855	5.532	214.297	11.182	1.347	0.077
ABELL_3140	54.0652	-40.6285	0.0620	91.667	6.967	194.228	16.033	1.808	0.124
Fornax_Cluster	54.6204	-35.4499	0.0046	0.346	0.066	274.257	7.900	0.918	0.011
MCXC_J0338.6+0958	54.6713	9.9669	0.0363	5.111	0.100	106.901	0.925	1.395	0.011
MCXC_J0340.8-4542	55.2241	-45.6765	0.0698	189.419	55.692	68.974	132.705	0.633	0.750
MCXC_J0352.9+1941	58.2437	19.6819	0.1090	7.106	0.380	56.749	2.013	1.494	0.061
MACS_J0358.8-2955	59.7190	-29.9303	0.4250	26.151	14.956	167.420	22.122	0.687	0.062
ABELL_0478	63.3556	10.4653	0.0881	11.580	0.910	216.321	33.585	1.549	0.130
MACS_J0416.1-2403	64.0393	-24.0669	0.4200	76.272	106.522	245.717	116.894	0.421	0.225
MACS_J0417.5-1154	64.3945	-11.9091	0.4400	23.919	8.332	159.396	14.094	0.935	0.061
MCXC_J0425.8-0833	66.4636	-8.5601	0.0397	6.656	1.320	122.571	3.695	0.924	0.042
ABELL_S0463	67.1559	-53.8418	0.0394	117.469	23.573	100.204	44.890	1.076	0.411
MACS_J0429.6-0253	67.4004	-2.8858	0.3990	12.169	3.110	109.873	8.616	1.166	0.071
ABELL_0496	68.4085	-13.2610	0.0329	4.569	0.330	146.262	1.605	0.984	0.014
MCXC_J0437.1+0043	69.2898	0.7322	0.2850	39.939	4.172	88.447	8.153	1.301	0.079
MCXC_J0439.0+0715	69.7530	7.2688	0.2300	57.207	11.956	116.269	17.768	0.946	0.095
MCXC_J0439.0+0520	69.7592	5.3453	0.2080	6.155	1.597	118.345	5.101	1.067	0.044
ABELL_3292	72.4841	-44.6725	0.1723	108.116	10.191	35.724	10.662	1.694	0.168
WEIN_051	72.5271	45.0510	0.0222	84.177	55.153	196.260	94.348	0.420	0.258
MCXC_J0454.1-0300	73.5458	-3.0145	0.5500	193.350	20.561	49.053	20.322	1.387	0.197
ESO_552-G_020	73.7180	-18.1157	0.0314	-6.756	7.226	206.197	10.816	0.609	0.081
ABELL_3322	77.5705	-45.3212	0.2000	104.613	17.637	77.060	21.058	1.196	0.147

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
MCXC_J0510.7-0801	77.6985	-8.0275	0.2195	122.172	51.636	150.700	79.443	0.574	0.180
ABELL_S0520	79.1571	-54.5131	0.2952	320.133	39.838	30.757	30.494	1.656	0.233
ABELL_3343	81.4530	-47.2528	0.1913	163.757	19.794	46.150	19.915	1.602	0.225
MCXC_J0528.2-2942	82.0613	-29.7208	0.1582	88.463	18.810	91.877	37.031	1.361	0.366
RBS_0653	82.2210	-39.4710	0.2839	27.862	4.749	164.154	10.315	0.941	0.043
PLCKESZ_G286.58-31.25	82.8684	-75.1793	0.2100	149.127	62.571	146.327	86.892	0.782	0.226
ABELL_0545	83.1054	-11.5424	0.1540	146.987	21.961	114.049	28.338	0.894	0.125
MCXC_J0532.9-3701	83.2314	-37.0264	0.2747	98.106	24.012	128.772	35.406	1.020	0.151
ESO3060170-A	85.0279	-40.8369	0.0358	2.005	1.515	441.229	108.030	1.170	0.126
MCXC_J0547.0-3904	86.7566	-39.0745	0.2100	10.804	2.398	124.928	8.182	1.179	0.072
ABELL_3364	86.9071	-31.8732	0.1483	205.019	17.037	42.820	17.143	1.635	0.206
ABELL_0548A	87.1596	-25.4779	0.0395	23.658	10.142	170.653	20.112	0.920	0.173
ABELL_0550	88.2147	-21.0536	0.0990	120.247	24.385	102.004	32.615	1.014	0.181
MACS_J0553.4-3342	88.3653	-33.7104	0.4070	106.577	46.574	220.930	82.059	0.694	0.179
ABELL_3378	91.4749	-35.3022	0.1410	8.528	3.136	120.248	6.808	0.910	0.060
SPT-CL_J0615-5746	93.9661	-57.7794	0.9720	77.522	10.404	77.435	14.033	1.314	0.111
CIZA_J0616.3-2156	94.1033	-21.9383	0.1710	234.527	60.544	117.417	92.762	0.854	0.290
ABELL_S0579	94.1339	-39.7968	0.1520	129.966	24.011	68.501	26.967	1.134	0.185
G139.59+24.18	95.4541	74.7014	0.2700	32.801	8.952	150.993	18.544	1.244	0.247
ABELL_3391	96.5896	-53.6958	0.0514	203.151	29.569	122.635	45.810	0.917	0.228
ABELL_3395_SW	96.7019	-54.5498	0.0510	190.118	25.741	99.607	44.497	1.376	0.344
ABELL_3399	99.3105	-48.4719	0.2026	59.435	16.441	164.359	25.195	0.867	0.078
PLCKESZ_G167.65+17.64	99.5154	47.7983	0.1740	209.979	18.217	21.819	7.294	1.886	0.086
ABELL_S0592	99.7025	-53.9740	0.2216	30.594	7.827	141.163	13.223	0.915	0.073
ABELL_3404	101.3709	-54.2286	0.1670	95.055	18.879	119.833	26.794	1.080	0.141
ABELL_0562	103.3397	69.3309	0.1100	122.537	34.177	82.469	77.220	0.761	0.498
ABELL_0576	110.3766	55.7616	0.0389	63.842	7.638	204.151	14.401	1.295	0.200
ABELL_0578	111.2230	66.9854	0.0866	16.576	48.228	216.253	74.410	0.644	0.369
ABELL_0586	113.0848	31.6329	0.1710	116.110	10.286	67.379	10.676	1.443	0.092

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ZwCl_0735.7+7421	115.4344	74.2440	0.2160	16.140	0.906	121.650	2.649	1.155	0.022
FBQS_J074417.4+375317	116.0724	37.8878	1.0686	-18.950	10.380	163.806	19.684	0.802	0.094
MACS_J0744.9+3927	116.2200	39.4568	0.6976	53.443	7.508	86.734	11.754	1.151	0.086
PKS_0745-19	116.8810	-19.2944	0.1028	8.581	0.315	118.197	1.029	1.167	0.012
ABELL_0598	117.8504	17.5147	0.1894	8.320	2.293	144.903	7.307	1.134	0.082
ABELL_0611	120.2368	36.0567	0.2880	63.673	8.437	112.999	14.631	1.241	0.111
SDSS-C4_3062	122.5951	42.2739	0.0638	46.651	12.194	160.145	27.816	1.264	0.264
ABELL_0644	124.3551	-7.5111	0.0704	70.240	6.441	114.365	8.615	1.049	0.047
MCXC_J0819.6+6336	124.8584	63.6240	0.1190	6.464	14.100	160.309	26.200	0.695	0.134
NGC_2563_GROUP	125.1485	21.0679	0.0163	1.362	0.679	930.543	51.355	1.318	0.037
UGCl_120	125.8402	4.3726	0.0293	-2.576	2.074	238.187	25.844	0.979	0.094
ZwCl_0823.2+0425	126.4910	4.2470	0.2248	72.063	16.331	90.869	25.862	1.267	0.215
2MFGC_06756	128.7288	55.5725	0.2411	12.469	1.187	95.625	3.373	1.177	0.034
ABELL_3411	130.4664	-17.4627	0.1687	194.021	24.317	76.514	25.816	1.275	0.179
NSC_J084254+292723	130.7332	29.4575	0.1940	22.863	2.043	100.849	6.189	1.546	0.106
ABELL_0697	130.7398	36.3660	0.2820	229.645	22.639	54.663	16.890	1.620	0.148
ZwCl_0857.9+2107	135.1535	20.8945	0.2300	18.933	2.344	81.093	5.949	1.716	0.129
SDSS_+137.3+11.0+0.18	137.3031	10.9756	0.1800	88.200	16.425	99.194	23.788	1.015	0.165
HCG_037	138.4130	29.9952	0.0223	2.267	1.280	503.351	339.928	1.180	0.198
2MASSi_J0913454+405628	138.4405	40.9416	0.4422	23.534	2.754	96.232	5.776	1.162	0.046
ABELL_0773	139.4690	51.7273	0.2170	179.097	21.364	64.053	17.288	1.388	0.130
Hydra_A	139.5249	-12.0955	0.0549	15.711	0.627	142.639	19.031	1.462	0.111
NSC_J092017+303027	140.1105	30.4938	0.2578	128.487	76.673	206.504	111.689	0.593	0.193
ABELL_0795	141.0241	14.1736	0.1359	27.927	3.477	115.745	6.809	1.085	0.059
UGC_05088_GROUP	143.3570	34.0481	0.0274	3.275	3.277	104.321	8.406	0.788	0.142
WHL_J093820.9+520243	144.5847	52.0482	0.3605	92.459	18.089	60.018	40.973	1.255	0.546
ABELL_0868	146.3589	-8.6568	0.1530	192.942	14.459	14.989	6.529	1.866	0.099
GALEX_J094712.4+762313	146.8029	76.3871	0.3541	25.087	1.568	75.232	4.847	1.915	0.059
ZwCl_0949.6+5207	148.2049	51.8848	0.2140	4.895	8.399	129.424	12.195	1.070	0.097

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
PLCKESZ_G264.41+19.48	150.0067	-30.2770	0.2400	94.347	40.037	176.575	56.786	0.783	0.139
HCG_042	150.0593	-19.6363	0.0133	1.714	0.342	188.993	23.760	1.009	0.060
MCXC_J1000.5+4409	150.1334	44.1444	0.1540	16.985	5.367	111.258	9.943	0.940	0.083
ZwCl_1006.1+1201	152.1978	11.7934	0.2210	122.660	18.258	86.035	23.767	1.187	0.175
MCXC_J1010.5-1239	152.6346	-12.6658	0.3010	22.531	50.166	227.985	73.456	0.561	0.187
ABELL_0963	154.2656	39.0470	0.2060	39.105	8.506	149.621	13.215	0.826	0.057
ABELL_0970	154.3476	-10.6859	0.0587	102.112	18.122	122.251	26.244	1.014	0.145
NGC_3209	155.1602	25.5051	0.0209	3.265	0.846	182.701	297.526	1.217	0.527
MCXC_J1022.0+3830	155.5418	38.5232	0.0491	38.803	9.494	177.395	16.646	1.018	0.255
ABELL_0980	155.6182	50.1061	0.1582	150.124	21.661	62.493	23.780	1.465	0.227
BLOX_J1023.6+0411.1	155.9156	4.1856	0.2906	8.947	0.960	106.958	3.162	1.216	0.026
WHL_J102339.9+490838	155.9162	49.1445	0.1440	128.079	20.582	65.204	22.416	1.253	0.177
ABELL_3444	155.9591	-27.2566	0.2533	20.762	1.805	94.190	4.675	1.340	0.045
ABELL_1033	157.9367	35.0396	0.1259	122.945	16.451	116.344	22.449	1.044	0.126
ABELL_1068	160.1855	39.9529	0.1375	6.283	0.705	109.284	3.042	1.157	0.031
NGC_3402_GROUP	162.6088	-12.8449	0.0154	1.666	0.119	93.942	2.839	0.955	0.020
MCXC_J1053.7+5452	163.3837	54.8792	0.0704	63.300	12.560	60.464	16.230	1.189	0.161
BLOX_J1056.9-0337.3	164.2330	-3.6277	0.8231	126.798	36.945	87.750	78.523	1.248	0.487
MACS_J1108.9+0906	167.2304	9.0991	0.4490	80.249	38.940	111.019	64.854	0.789	0.252
NGC_3551	167.4343	21.7594	0.0320	-21.019	22.510	218.204	37.104	0.491	0.317
1RXS_J111039.6+284316	167.6690	28.7138	0.0220	64.272	9.246	84.537	44.615	1.464	0.397
ABELL_1190	167.9159	40.8402	0.0751	213.612	19.437	25.935	18.426	1.686	0.227
ABELL_1201	168.2270	13.4358	0.1688	63.600	8.173	176.604	19.440	1.575	0.259
ABELL_1204	168.3351	17.5940	0.1706	14.727	1.691	82.196	4.882	1.440	0.138
MACS_J1115.8+0129	168.9669	1.4990	0.3520	22.147	2.422	124.017	9.386	1.667	0.183
HCG_051	170.6101	24.2976	0.0258	-7.020	2.692	130.845	10.858	0.558	0.075
ABELL_1240	170.9098	43.0968	0.1590	196.271	92.437	117.624	160.964	0.672	0.708
MCXC_J1130.0+3637	172.5116	36.6356	0.0600	21.316	2.683	119.264	8.124	1.098	0.124
ABELL_1285	172.5932	-14.5807	0.1061	186.247	29.134	77.426	35.964	0.997	0.214

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_1300	172.9774	-19.9290	0.3072	65.709	26.357	203.526	36.791	0.929	0.083
WHL_J114224.8+583205	175.6006	58.5331	0.3109	459.851	44.958	14.713	25.032	1.732	0.197
ABELL_1361	175.9151	46.3557	0.1171	18.048	3.005	134.880	8.970	1.464	0.144
SDSS-C4-DR3_3018	176.8274	55.7400	0.0510	130.902	37.892	103.086	73.251	1.068	0.657
MACS_J1149.6+2223	177.3994	22.3986	0.5444	267.540	25.887	18.052	9.184	1.859	0.104
ABELL_1413	178.8245	23.4061	0.1427	57.420	5.077	129.718	8.112	1.047	0.044
ABELL_1423	179.3219	33.6104	0.0761	36.915	6.299	359.425	31.006	1.209	0.143
SDSS-C4-DR3_3144	179.9679	55.5349	0.0810	3.097	0.714	393.635	61.006	1.448	0.104
ABELL_1446	180.5156	58.0383	0.1035	157.742	17.020	62.913	31.082	1.462	0.362
MKW_04	181.1134	1.8952	0.0200	2.617	0.437	176.954	7.150	0.863	0.025
MACS_J1206.2-0847	181.5512	-8.8007	0.4400	54.200	11.443	134.851	20.755	1.105	0.102
NGC_4104_GROUP	181.6624	28.1750	0.0283	0.021	0.396	871.354	94.140	1.316	0.044
MCXC_J1215.4-3900	183.8527	-39.0363	0.1190	275.446	63.700	71.648	105.128	1.131	0.466
NGC_4325_GROUP	185.7773	10.6222	0.0252	3.284	0.201	110.659	3.216	1.289	0.026
NSCS_J122648+215157	186.7123	21.8326	0.3700	57.719	16.388	120.674	26.521	0.911	0.164
ABELL_1553	187.6970	10.5540	0.1652	180.396	25.875	57.808	22.365	1.443	0.192
MCXC_J1234.2+0947	188.6014	9.7892	0.2290	194.608	65.368	87.959	115.821	0.842	0.419
MESSIER_089	188.9157	12.5567	0.0011	1.274	0.032	928.754	52.840	1.230	0.016
ABELL_1569	189.1084	16.5383	0.0735	136.449	21.463	44.110	65.747	0.904	0.672
ABELL_1576	189.2428	63.1872	0.2790	103.880	22.109	127.184	31.722	1.026	0.135
NGC_4636	190.7077	2.6877	0.0031	0.801	0.091	228.937	13.528	1.025	0.022
Centaurus_Cluster	192.2053	-41.3110	0.0114	0.873	0.026	324.906	3.843	1.198	0.005
NGC_4759_GROUP	193.2739	-9.2043	0.0147	2.485	0.104	168.149	8.097	1.156	0.025
ABELL_3528B	193.5927	-29.0111	0.0530	15.261	2.683	218.802	8.800	1.157	0.061
NGC_4782-3	193.6504	-12.5607	0.0154	3.197	0.756	309.122	63.049	1.246	0.100
ABELL_1644	194.2986	-17.4091	0.0473	21.222	1.461	616.093	78.787	1.815	0.107
ABELL_3532	194.3430	-30.3626	0.0554	150.770	20.305	83.651	29.757	1.251	0.249
NGC_4839	194.3517	27.4977	0.0246	-26.397	15.220	856.067	103.825	0.777	0.114
WHL_J125933.4+600409	194.8877	60.0706	0.3300	327.114	56.118	36.889	80.081	1.332	0.418

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_1664	195.9276	-24.2449	0.1283	14.131	0.879	107.446	4.539	1.690	0.081
ABELL_1668	195.9445	19.2705	0.0634	6.156	1.317	167.323	7.882	1.033	0.053
NGC4936-offset2	196.0715	-30.5261	0.0103	-0.354	3.076	102.788	11.701	0.593	0.131
MACS_J1311.0-0311	197.7569	-3.1776	0.4940	32.423	3.781	83.827	8.127	1.357	0.085
ABELL_1689	197.8734	-1.3413	0.1832	61.841	4.922	126.295	7.801	1.156	0.048
WHL_J131505.2+514902	198.7700	51.8193	0.2911	159.895	30.277	100.569	32.233	1.161	0.150
NGC_5044	198.8498	-16.3854	0.0093	1.305	0.143	47.729	1.766	0.722	0.022
NGC_5098_GROUP	200.0610	33.1425	0.0366	4.798	0.944	129.716	6.625	1.160	0.060
NGC_5129	201.0416	13.9755	0.0230	-0.457	0.659	212.419	33.879	0.923	0.075
ABELL_1736	201.7060	-27.1634	0.0458	139.631	39.368	75.724	80.890	0.691	0.463
ABELL_3558	201.9869	-31.4955	0.0480	49.877	14.911	194.069	20.760	0.772	0.094
NGC_5171	202.3540	11.7626	0.0229	41.304	17.842	103.650	84.833	0.987	0.735
SSGC_081	202.4490	-31.6065	0.0495	44.512	20.290	165.343	31.988	0.614	0.122
a1750ss	202.5424	-2.1044	0.0914	34.901	69.701	186.190	91.102	0.580	0.361
ABELL_1750C	202.7093	-1.8629	0.0678	159.868	10.761	56.099	9.678	1.886	0.082
ABELL_1750N	202.7956	-1.7282	0.0837	93.060	22.131	117.458	37.180	0.980	0.211
SC_1329-313	202.8650	-31.8213	0.0482	167.330	16.567	61.944	24.466	1.703	0.208
2MASX_J13312961+1107566	202.8737	11.1319	0.0790	4.504	1.406	130.128	37.766	1.462	0.285
ABELL_3560	203.1071	-33.1360	0.0489	40.383	61.464	198.695	85.109	0.367	0.251
ABELL_1758	203.2017	50.5424	0.2790	166.511	49.643	189.649	105.927	1.036	0.431
ABELL_3562	203.4075	-31.6700	0.0490	70.592	12.786	144.224	18.709	0.942	0.101
ABELL_1763	203.8248	40.9988	0.2230	186.619	22.108	54.570	15.648	1.348	0.124
ABELL_1767	204.0327	59.2044	0.0703	136.566	62.655	147.731	126.407	0.647	0.585
ABELL_1775	205.4527	26.3722	0.0717	56.894	3.877	256.554	19.789	1.860	0.099
ABELL_3571	206.8685	-32.8646	0.0391	59.964	13.817	210.881	14.705	0.699	0.151
LCDCS_0829	206.8775	-11.7528	0.4510	3.460	2.299	177.098	6.231	1.074	0.036
ABELL_1795	207.2200	26.5899	0.0625	15.172	0.779	118.037	1.460	1.080	0.018
NSCS_J135021+094042	207.5913	9.6698	0.0900	-6.289	1.541	172.190	5.844	0.831	0.040
MACS_J1359.2-1929	209.7926	-19.4903	0.4470	19.013	3.628	92.645	9.761	1.330	0.106

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_1831	209.8159	27.9756	0.0615	77.354	12.848	100.593	19.688	1.165	0.154
WHL_J135949.5+623047	209.9605	62.5179	0.3216	21.238	6.093	181.636	13.843	1.179	0.131
ABELL_1835	210.2581	2.8787	0.2532	11.299	1.211	105.985	4.227	1.323	0.034
NGC5419-offset2	210.9117	-33.9781	0.0138	-1.973	1.107	423.685	52.658	0.861	0.056
NGC5419-offset1	210.9119	-33.9783	0.0138	0.834	0.936	795.366	128.597	1.053	0.064
ABELL_3581	211.8741	-27.0183	0.0230	8.670	0.667	183.191	51.951	1.380	0.150
A1882a	213.7848	-0.4931	0.1405	61.304	63.352	189.657	97.281	0.529	0.356
WARP_J1415.1+3612	213.7954	36.2022	1.0300	-10.568	40.951	162.648	57.392	0.456	0.217
WHL_J141623.8+444528	214.1164	44.7801	0.3859	18.734	17.583	104.876	37.534	0.900	0.233
GMBCG_J215.94948+24.07846	215.9490	24.0779	0.5431	5.560	3.403	132.581	7.921	1.167	0.105
ABELL_1914	216.5128	37.8244	0.1712	84.494	15.900	139.103	22.462	0.853	0.081
WHL_J142716.1+440730	216.8172	44.1271	0.4981	12.595	3.479	135.389	10.844	1.267	0.078
MACS_J1427.6-2521	216.9177	-25.3541	0.3180	9.474	4.811	133.733	10.335	0.922	0.066
ABELL_1930	218.1579	31.6479	0.1313	-3.707	2.733	182.246	6.827	0.823	0.042
ABELL_1942_AND_CLUMP	219.5912	3.6703	0.2240	67.023	74.345	229.305	101.357	0.466	0.200
WBL_518	220.1651	3.4704	0.0270	89.363	48.075	144.171	94.679	0.405	0.307
NSCS_J144726+082824	221.8610	8.4737	0.1954	11.220	1.701	195.746	8.898	1.556	0.061
ABELL_1991	223.6318	18.6449	0.0587	0.392	0.244	122.804	2.375	0.996	0.017
NSCS_J145715+222009	224.3129	22.3424	0.2578	12.779	1.247	80.992	3.027	1.135	0.027
ABELL_S0780	224.8701	-18.1787	0.2357	20.933	1.311	114.084	10.975	1.817	0.128
ABELL_2009	225.0816	21.3699	0.1532	18.811	2.582	133.365	7.018	1.105	0.044
WHL_J150407.5-024816	226.0309	-2.8044	0.2153	9.077	0.494	82.228	2.096	1.329	0.019
MCXC_J1514.9-1523	228.7606	-15.3892	0.2226	356.147	108.545	137.948	191.290	0.611	0.399
ABELL_2061	230.2923	30.6115	0.0784	224.737	24.178	38.185	19.474	1.434	0.233
MKW_03s	230.4664	7.7089	0.0450	18.044	1.755	117.220	2.724	0.910	0.029
ABELL_2069	231.0474	29.8720	0.1160	297.687	41.573	84.813	46.812	0.963	0.212
MCXC_J1524.2-3154	231.0535	-31.9045	0.1028	5.444	0.428	104.851	2.303	1.195	0.022
MACS_J1532.8+3021	233.2241	30.3496	0.3450	13.365	0.901	82.169	2.761	1.339	0.030
ABELL_2092	233.3234	31.1379	0.0669	79.498	62.076	120.495	105.986	0.690	0.623

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_2107	234.9130	21.7827	0.0411	10.155	3.138	327.958	32.782	0.923	0.086
ABELL_2111	234.9193	34.4244	0.2290	189.915	27.115	88.337	27.655	1.189	0.156
ABELL_2104	235.0339	-3.3042	0.1533	159.142	31.845	127.169	43.602	0.820	0.178
ABELL_2125	235.3090	66.2659	0.2465	171.248	17.883	25.368	19.667	1.588	0.272
ABELL_2124	236.2464	36.1095	0.0656	91.073	24.555	196.375	38.843	0.814	0.173
ABELL_2146	239.0512	66.3516	0.2343	85.695	4.204	66.396	6.028	1.891	0.078
MCXC_J1558.3-1410	239.5910	-14.1661	0.0970	28.469	1.712	115.116	4.612	1.496	0.102
ABELL_2147	240.5709	15.9745	0.0350	163.008	18.556	78.332	36.313	1.411	0.378
ABELL_2151	241.1495	17.7215	0.0366	6.426	3.102	150.802	5.699	0.791	0.056
AWM_4	241.2354	23.9340	0.0318	24.474	5.366	108.464	6.625	0.655	0.094
MACS_J1621.3+3810	245.3533	38.1691	0.4650	-2.413	5.626	168.760	12.893	0.917	0.058
ABELL_2187	246.0584	41.2438	0.1836	91.556	19.021	108.736	26.700	1.194	0.163
ABELL_2204	248.1955	5.5758	0.1522	7.763	0.351	144.831	3.209	1.509	0.028
ABELL_2219	250.0838	46.7119	0.2256	258.490	19.763	49.427	17.930	1.458	0.189
Hercules_A	252.7840	4.9923	0.1550	0.604	2.871	211.799	20.461	1.072	0.127
NGC_6269	254.4921	27.8541	0.0348	1.305	1.932	279.709	47.558	0.892	0.100
ABELL_2256	255.9357	78.6365	0.0581	111.092	33.408	159.118	57.048	0.764	0.272
Ophiuchus_CLUSTER	258.1155	-23.3685	0.0280	-22.281	1.747	259.384	2.600	0.513	0.013
SDSS-C4_3072	260.0414	26.6248	0.1640	19.703	1.477	104.101	4.069	1.308	0.030
MACS_J1720.2+3536	260.0706	35.6066	0.3913	12.312	2.523	131.324	7.489	1.122	0.048
ABELL_2261	260.6136	32.1329	0.2240	41.818	9.759	139.288	14.832	0.899	0.070
ABELL_2294	261.0423	85.8861	0.1694	58.446	48.431	216.653	77.184	0.724	0.174
Abell_2276	263.7693	64.1017	0.1406	24.173	11.035	132.597	18.928	0.818	0.125
ZwCl_1742.1+3306	266.0605	32.9916	0.0757	12.335	1.209	120.830	4.566	1.256	0.064
NSC_J174715+451155	266.8094	45.1960	0.1565	290.692	35.204	11.492	16.426	1.751	0.186
MCXC_J1750.2+3504	267.5691	35.0828	0.1710	-1.787	2.580	163.642	7.589	0.886	0.042
NGC_6482	267.9531	23.0719	0.0131	0.837	0.243	75.095	9.968	0.826	0.057
CIZA_J1804.4+1002	271.1307	10.0568	0.1525	26.268	48.513	220.224	70.702	0.494	0.134
ABELL_2302	274.9918	57.1561	0.1790	190.544	77.382	121.021	139.110	0.764	0.418

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
_HB89_1821+643	275.4876	64.3438	0.2970	-27.371	7.264	182.036	12.223	0.976	0.051
MACS_J1829.0+6913	277.2758	69.2356	0.2030	36.357	4.091	99.984	8.250	1.328	0.094
MCXC_J1852.1+5711	283.0367	57.1952	0.1094	9.598	7.763	168.070	13.197	0.843	0.097
MCXC_J1853.9+6822	283.5094	68.3827	0.0928	76.431	14.615	198.874	33.895	1.402	0.384
PLCKESZ_G337.09-25.97	288.6564	-59.4722	0.2636	78.316	8.850	64.347	13.718	1.663	0.191
MACS_J1931.8-2635	292.9569	-26.5761	0.3520	20.746	1.385	83.339	3.832	1.474	0.049
CIZA_J1938.3+5409	294.5768	54.1597	0.2600	82.492	12.545	63.256	15.557	1.427	0.155
MCXC_J1947.3-7623	296.8121	-76.3958	0.2170	17.503	8.752	179.217	16.098	0.824	0.059
ABELL_3653	298.2623	-52.0358	0.1089	174.821	44.382	172.274	103.031	1.184	0.517
MCXC_J2003.5-2323	300.8624	-23.3827	0.3171	217.888	91.725	134.248	149.277	0.877	0.468
NGC_6868	302.4754	-48.3799	0.0095	0.698	0.554	564.773	228.792	1.219	0.119
MCXC_J2011.3-5725	302.8617	-57.4196	0.2786	39.441	8.622	64.575	15.583	1.194	0.206
MCXC_J2014.8-2430	303.7156	-24.5089	0.1612	5.104	0.800	116.626	3.757	1.208	0.035
SPT-CL_J2023-5535	305.8388	-55.5967	0.2320	186.127	83.939	142.241	160.670	0.758	0.646
ABELL_3695	308.6884	-35.8112	0.0894	307.183	36.618	41.785	35.309	1.562	0.278
SPT-CLJ2043-5035	310.8231	-50.5923	0.7230	15.848	3.138	82.653	8.551	1.242	0.089
MACS_J2046.0-3430	311.5018	-34.5061	0.4230	6.476	2.108	103.510	6.179	1.145	0.061
ABELL_3739	316.0798	-41.3449	0.1651	81.936	38.149	150.954	60.833	0.824	0.202
IC_1365	318.4828	2.5638	0.0493	160.833	19.732	71.379	38.386	1.338	0.371
ABELL_2345	321.7897	-12.1748	0.1765	286.760	114.145	143.127	175.850	0.694	0.381
ABELL_2355	323.8195	1.4177	0.1244	393.293	86.869	75.518	152.817	1.183	0.529
WBL_671	324.2863	0.4459	0.0510	-10.889	12.285	254.081	101.207	0.702	0.395
MACS_J2140.2-2339	325.0632	-23.6613	0.3130	12.928	1.101	91.533	3.478	1.350	0.041
ABELL_3809	326.7462	-43.8985	0.0623	10.277	3.844	132.956	8.319	0.867	0.059
ABELL_2384	328.0882	-19.5478	0.0943	25.359	2.675	139.689	6.597	1.271	0.056
ABELL_2390	328.4034	17.6957	0.2280	14.260	1.929	151.843	5.093	1.070	0.030
CIG_2153.8+3746	328.9678	38.0063	0.2920	55.189	6.284	325.021	34.317	1.708	0.178
ABELL_2409	330.2190	20.9685	0.1479	12.424	34.876	167.993	57.542	0.512	0.194
ABELL_3827	330.4717	-59.9453	0.0984	133.726	15.264	102.414	22.405	1.118	0.186

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_2415	331.4108	-5.5916	0.0581	2.830	0.781	142.155	6.585	1.027	0.046
MCXC_J2211.7-0349	332.9411	-3.8270	0.2700	96.790	15.187	105.826	19.206	1.343	0.111
3C_444	333.6124	-17.0256	0.1530	0.958	1.150	151.868	3.525	1.019	0.039
ABELL_2426	333.6347	-10.3705	0.0978	54.818	12.515	104.600	20.310	1.307	0.201
MACS_J2214-1359	333.7394	-14.0026	0.4830	150.198	23.601	62.246	24.888	1.286	0.195
ABELL_3854	334.4406	-35.7241	0.1492	103.513	18.154	96.946	24.929	1.201	0.151
MCXC_J2218.6-3853	334.6651	-38.9017	0.1379	132.838	15.498	47.640	17.748	1.423	0.219
ABELL_2443	336.5270	17.3654	0.1080	164.262	42.641	113.150	83.701	0.799	0.365
ABELL_2445	336.7321	25.8362	0.1660	69.915	11.294	87.841	15.912	0.992	0.117
ABELL_3880	336.9773	-30.5763	0.0584	0.853	0.950	139.151	4.029	0.890	0.032
MACS_J2229.8-2756	337.4390	-27.9273	0.3240	7.679	1.376	93.191	5.100	1.285	0.054
CGCG_514-050	337.8355	39.3587	0.0171	0.932	2.217	547.804	279.588	0.931	0.163
ABELL_2457	338.9213	1.4862	0.0594	24.049	17.816	181.200	33.329	0.700	0.120
Stephans_Quintet	338.9996	33.9721	0.0215	3.065	2.147	56.500	12.812	0.728	0.205
MACS_J2245.0+2637	341.2694	26.6343	0.3040	38.968	6.768	90.585	13.342	1.271	0.145
ABELL_3911	341.5639	-52.7242	0.0965	338.381	36.851	27.578	34.618	1.607	0.276
ABELL_2485	342.1291	-16.1079	0.2472	74.814	26.486	123.032	41.709	0.956	0.190
ABELL_S1063	342.1846	-44.5301	0.3475	78.785	19.913	127.407	31.163	0.868	0.166
ABELL_3921	342.4910	-64.4284	0.0928	75.119	14.847	162.521	20.746	0.787	0.073
ABELL_2507	344.2191	5.5043	0.1960	198.227	27.121	36.570	18.949	1.743	0.176
ABELL_2537	347.0930	-2.1916	0.2950	78.806	12.795	110.222	21.599	1.101	0.182
MCXC_J2311.5+0338	347.8885	3.6356	0.2998	85.303	14.556	121.671	17.017	1.189	0.071
ABELL_2550	347.8963	-21.7449	0.1226	-0.845	2.976	114.201	5.202	0.714	0.060
ABELL_2556	348.2559	-21.6346	0.0871	12.041	1.363	123.647	4.080	1.084	0.044
ABELL_S1101	348.4948	-42.7263	0.0580	11.526	0.499	77.480	1.226	1.127	0.018
UGC_12491	349.6596	42.9581	0.0174	6.963	0.491	206.196	23.395	1.541	0.084
NGC_7618	349.9501	42.8532	0.0173	-2.042	0.974	54.095	3.680	0.510	0.057
ABELL_2597	351.3324	-12.1243	0.0852	9.501	0.288	99.234	1.035	1.206	0.015
RCS_J2327-0204	351.8647	-2.0775	0.2000	45.275	7.173	212.711	14.883	1.224	0.088

Table B.1 (cont'd)

Cluster Name	RA	DEC	z	K_0	σ_{K_0}	K_{100}	$\sigma_{K_{100}}$	α	σ_α
	deg	deg		(keV cm ²)		(keV cm ²)			
ABELL_2626	354.1269	21.1465	0.0553	13.528	2.257	123.832	4.379	0.861	0.051
ABELL_2631	354.4107	0.2681	0.2730	76.914	106.409	230.751	114.628	0.412	0.218
MCXC_J2344.2-0422	356.0772	-4.3813	0.0786	108.571	13.251	66.387	15.739	1.283	0.157
HCG_097	356.8451	-2.2999	0.0218	-2.965	2.537	90.326	5.550	0.577	0.074
ABELL_2667	357.9142	-26.0841	0.2300	19.098	2.169	93.037	5.623	1.241	0.054
ABELL_2670	358.5368	-10.4252	0.0762	30.012	4.629	125.808	8.013	0.817	0.062

APPENDIX C

ACCEPT 2.0 RADIAL ENTROPY PROFILES

Deprojected radial entropy profiles and fits for the clusters from ACCEPT 2.0 with deprojected profiles. Included on each plot are the best-fit values for K_0 , K_{100} , and α with their 1σ errors as well as the reduced χ^2 for the fit and the radial range of the fit. The RA and DEC of the profile center are given in *hh : mm : ss* in the title of each plot.

Figure C.1 ACCEPT 2.0 Entropy Profiles and Fit Information

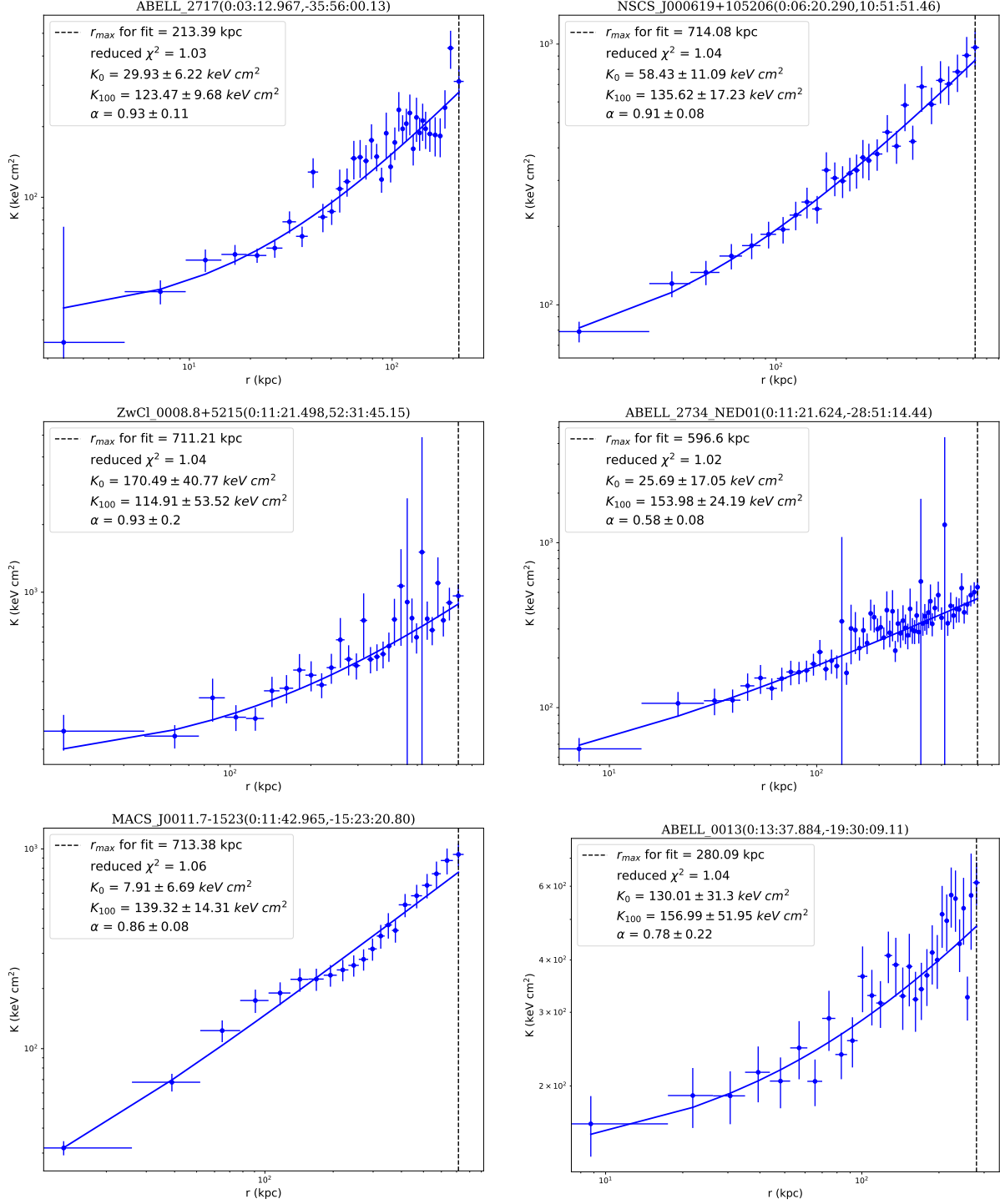


Figure C.1 (cont'd)

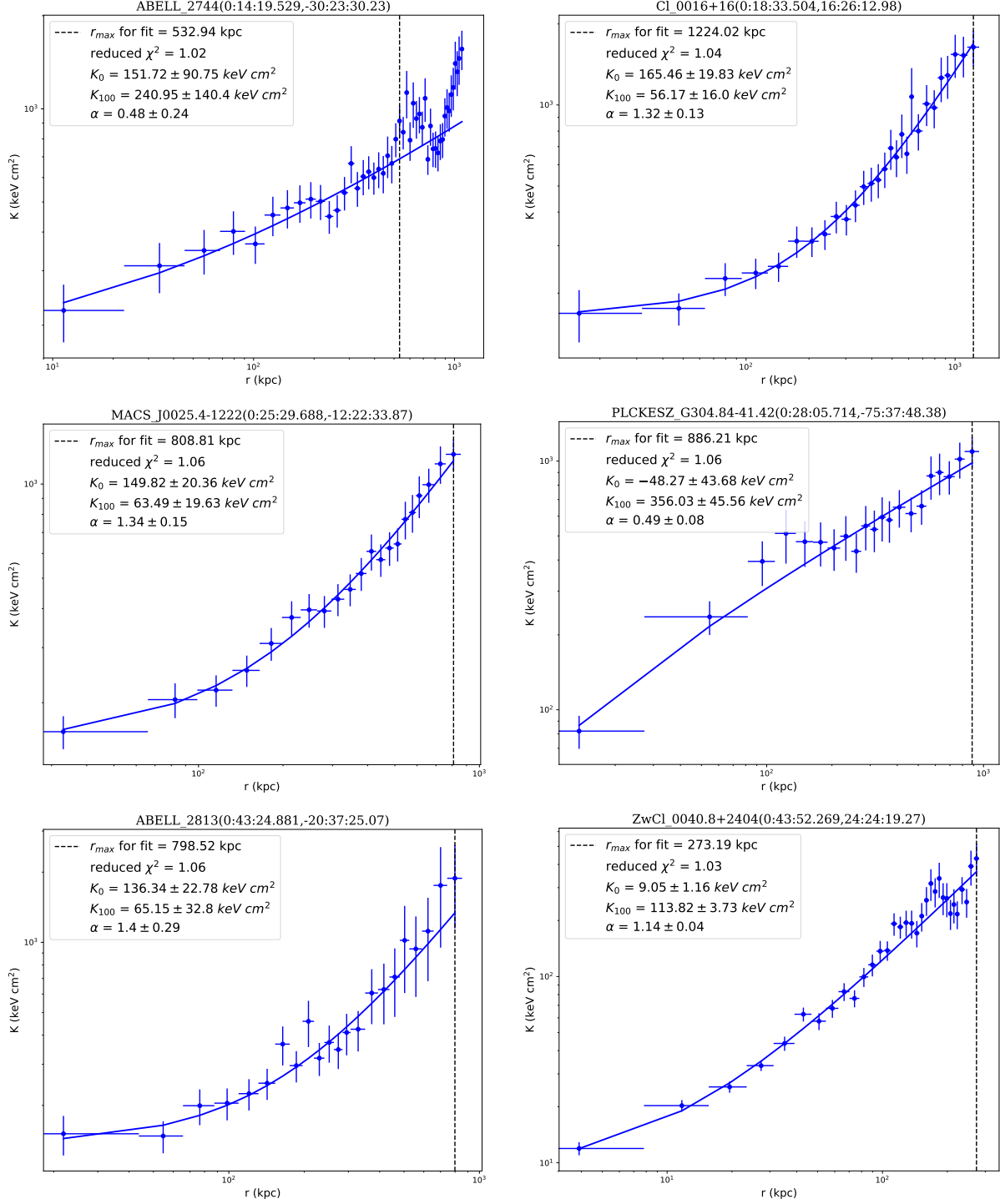


Figure C.1 (cont'd)

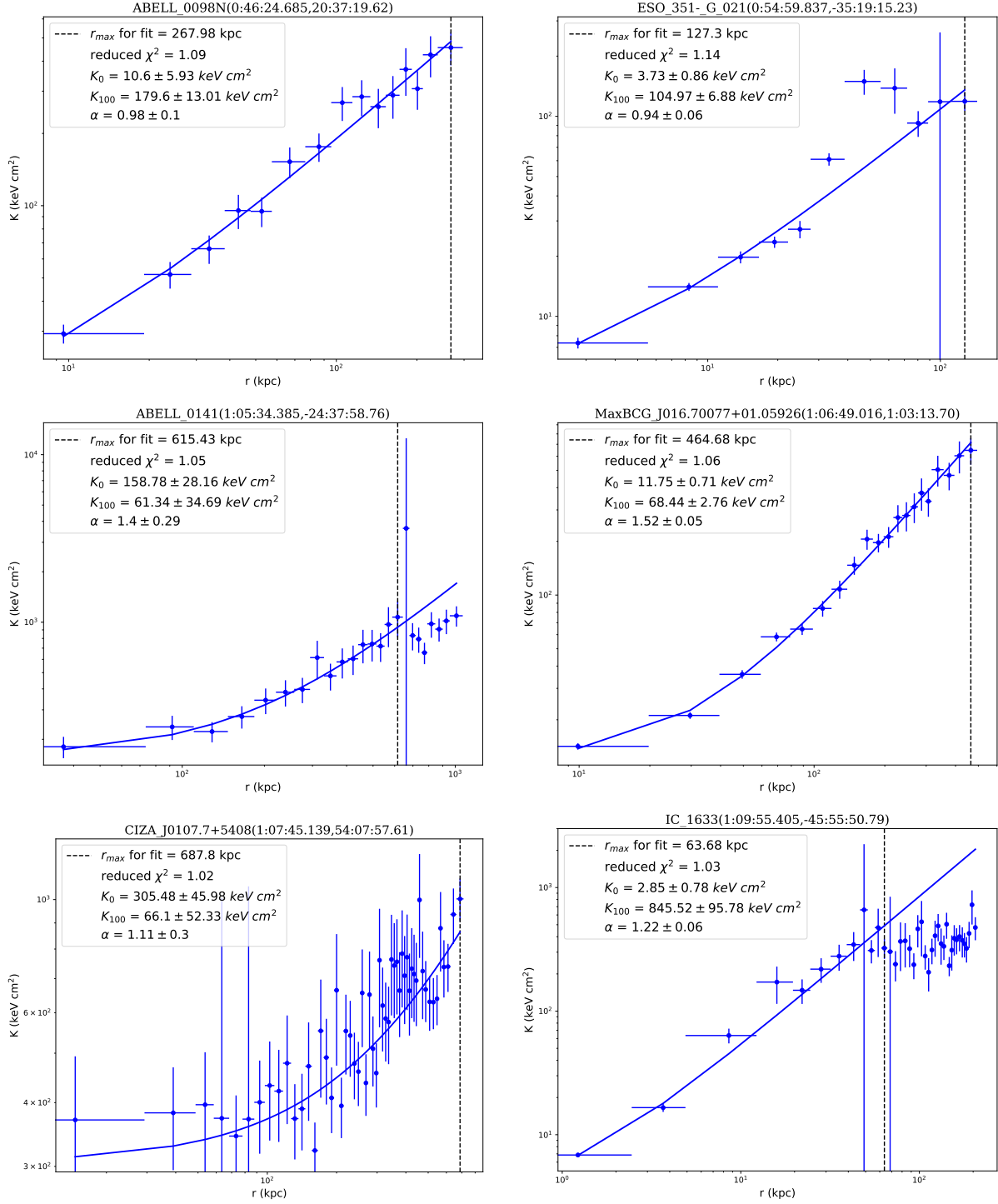


Figure C.1 (cont'd)

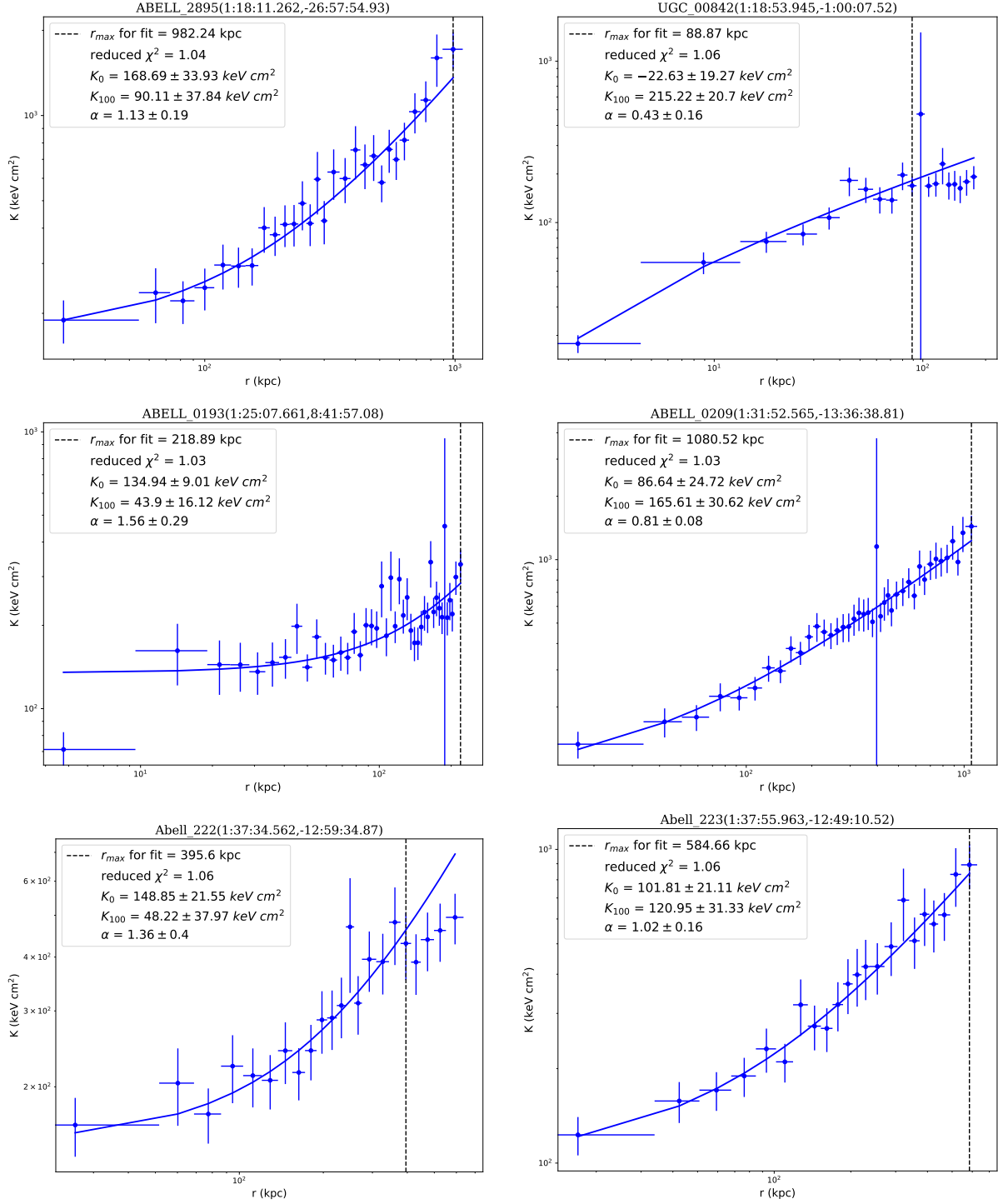


Figure C.1 (cont'd)

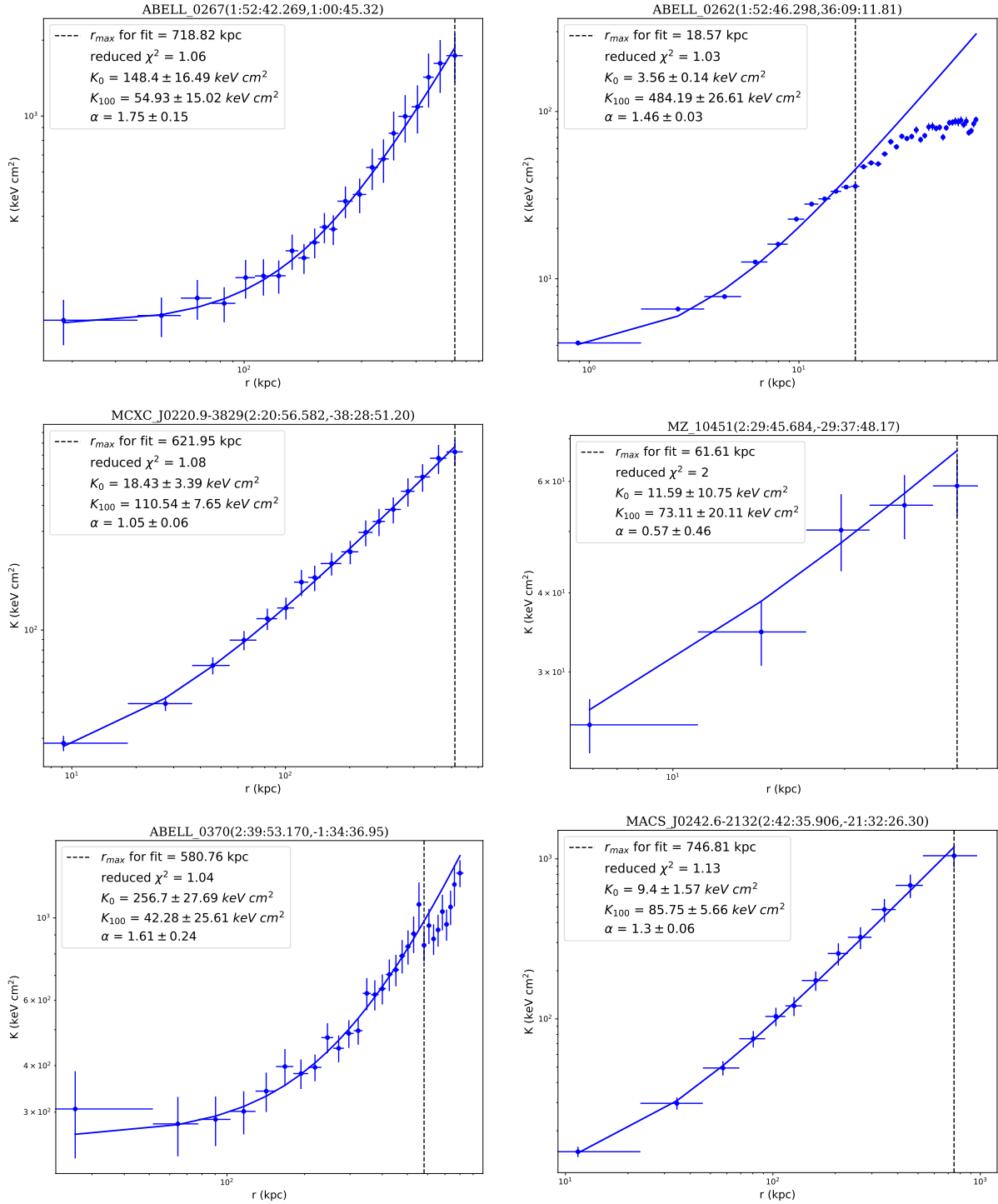


Figure C.1 (cont'd)

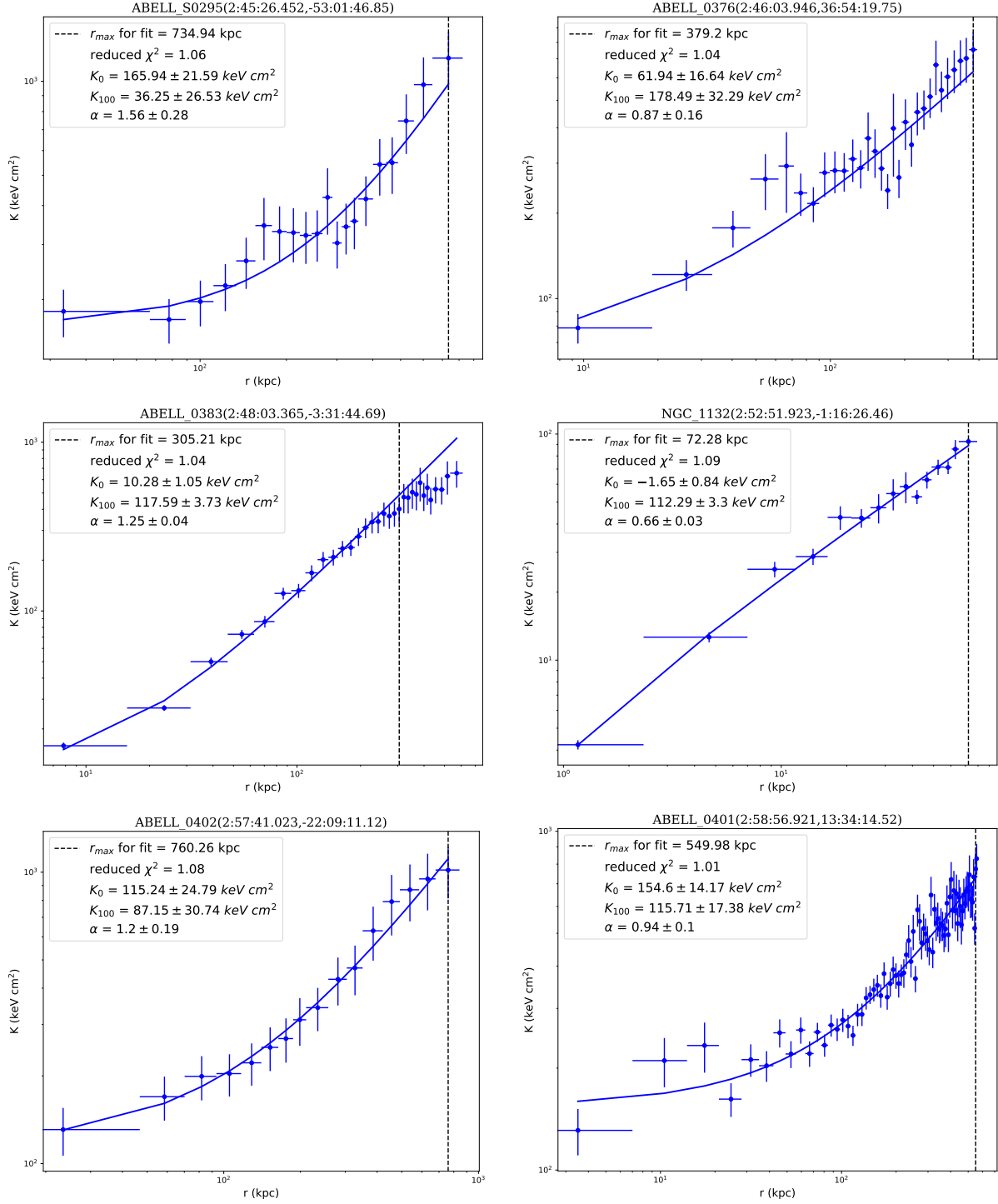


Figure C.1 (cont'd)

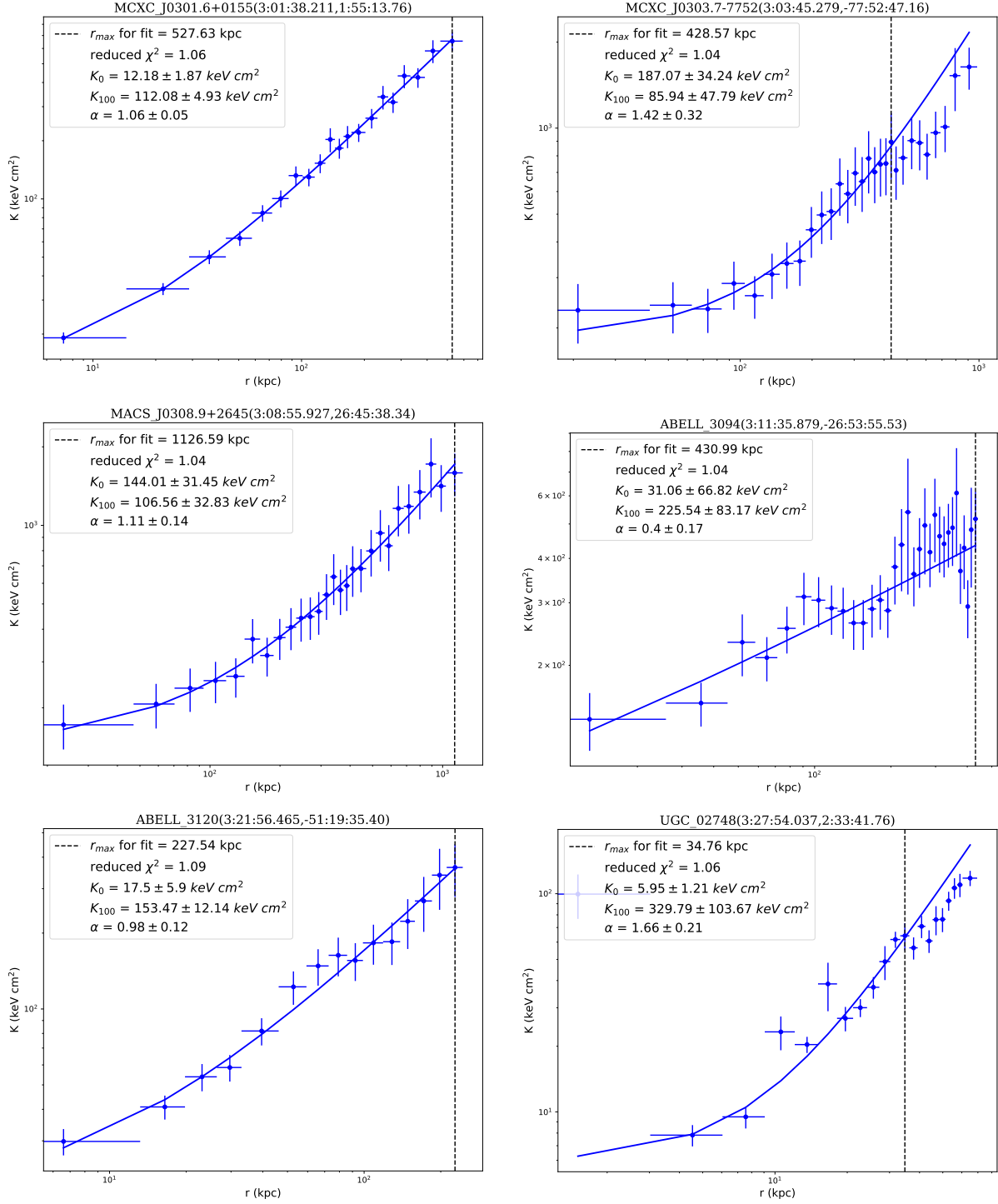


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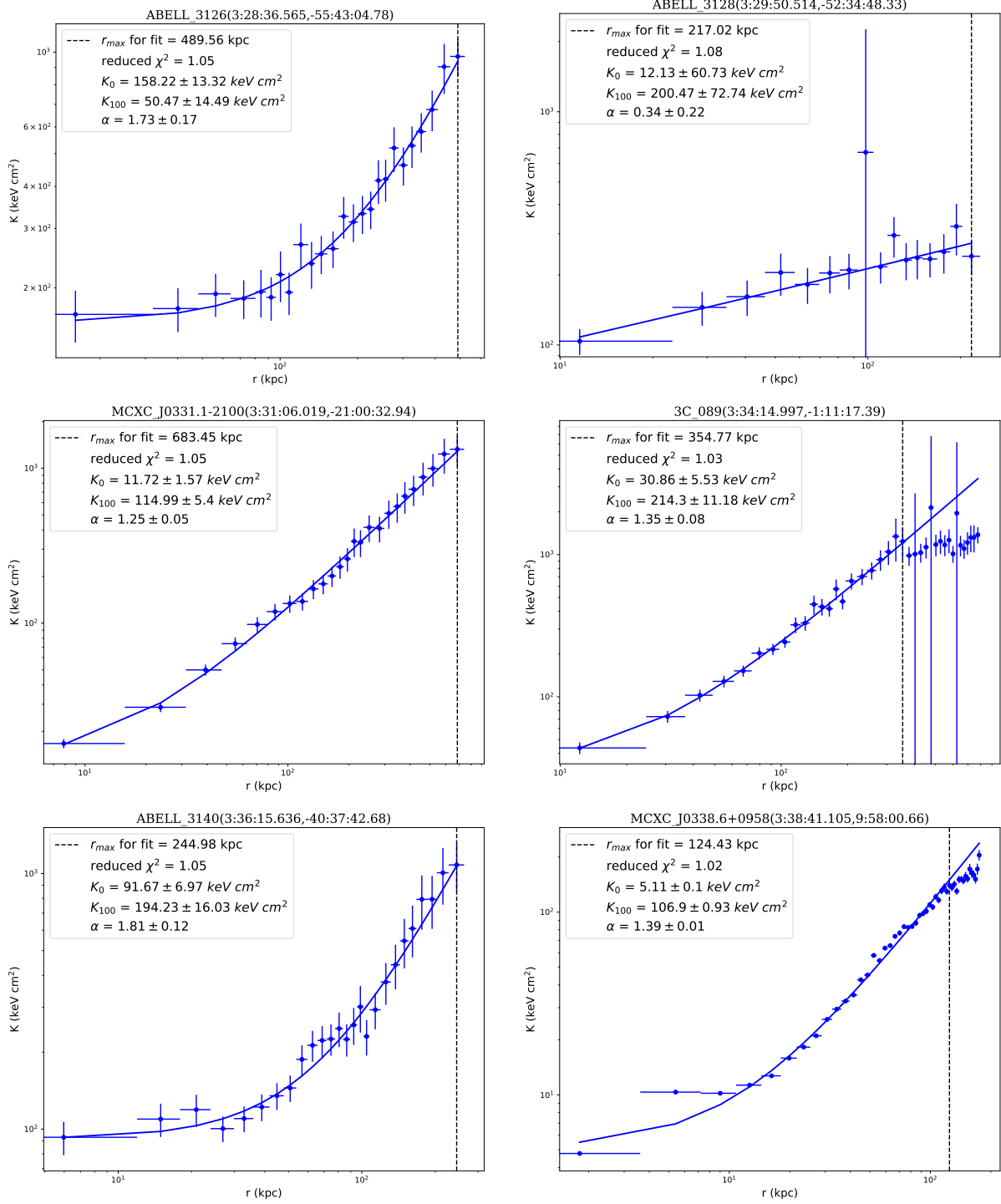


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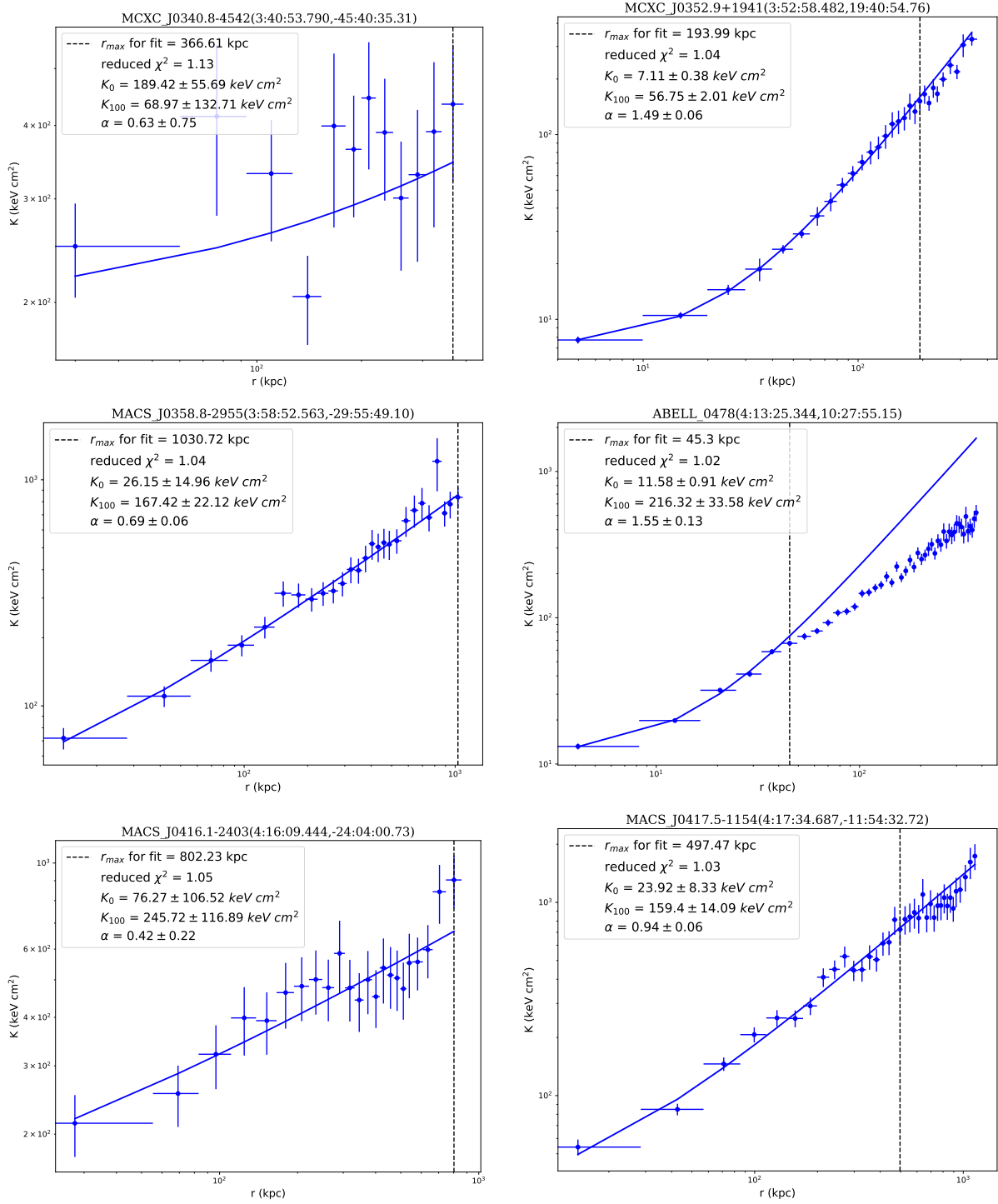


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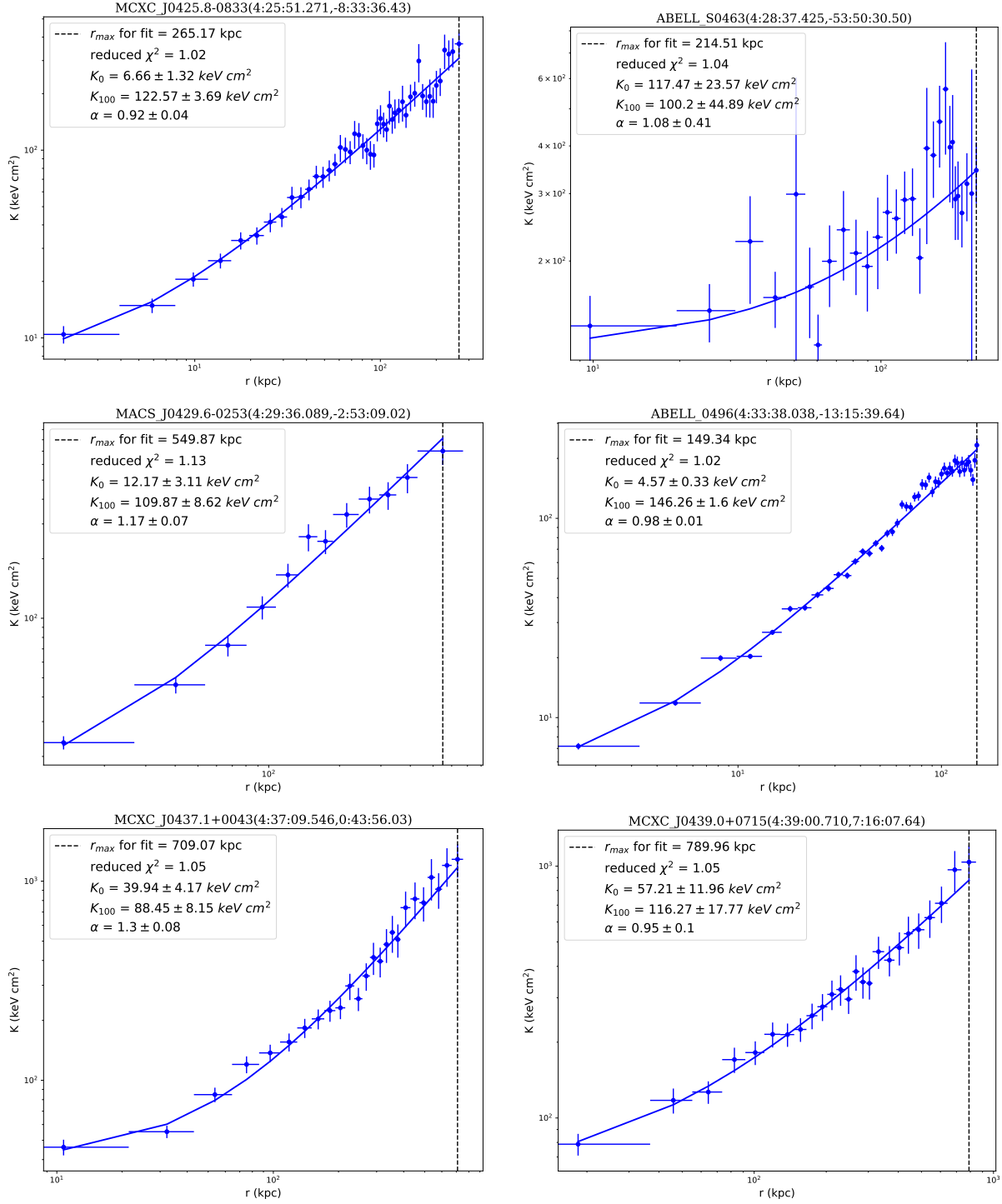


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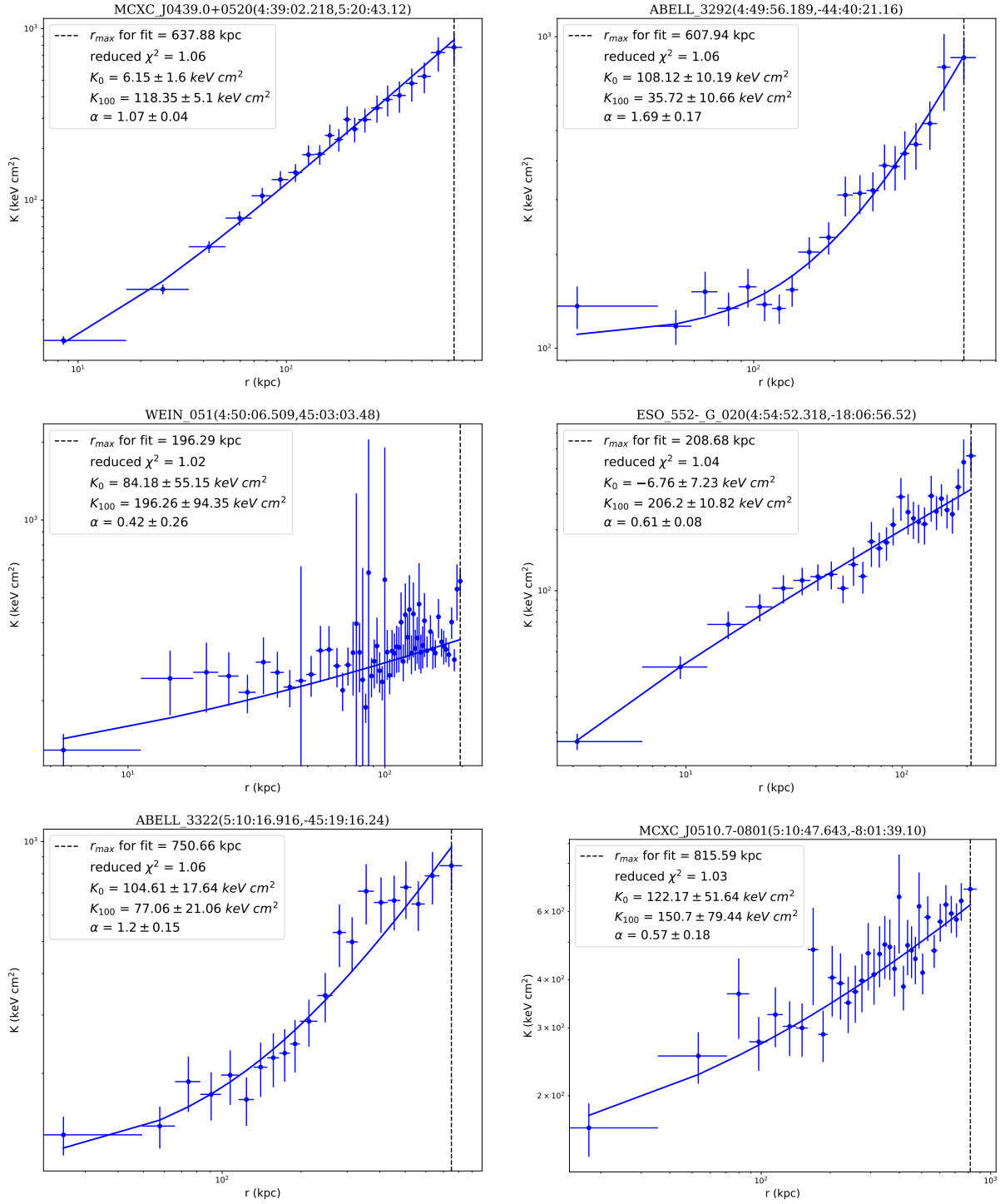


Figure C.1 (cont'd)

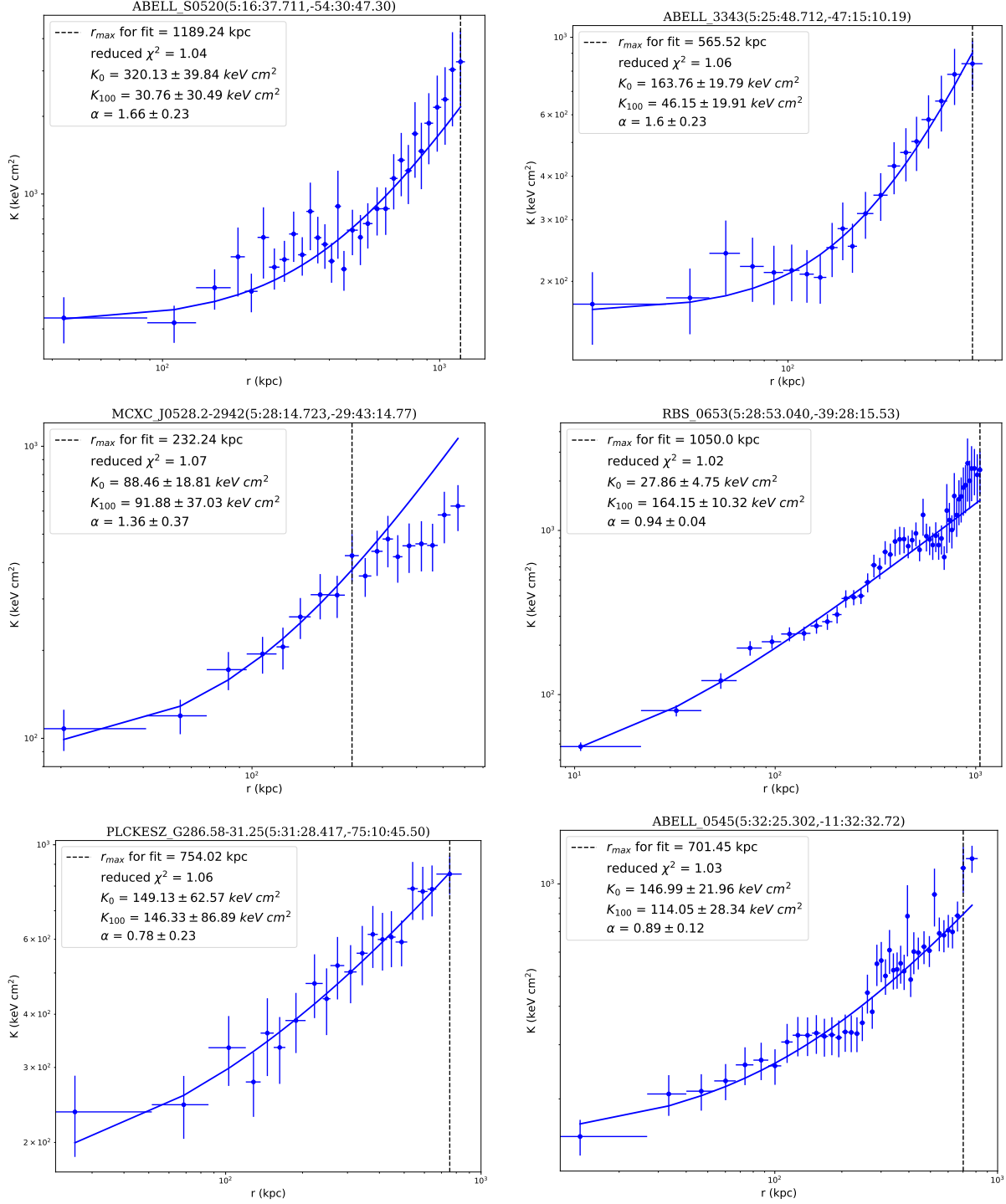


Figure C.1 (cont'd)

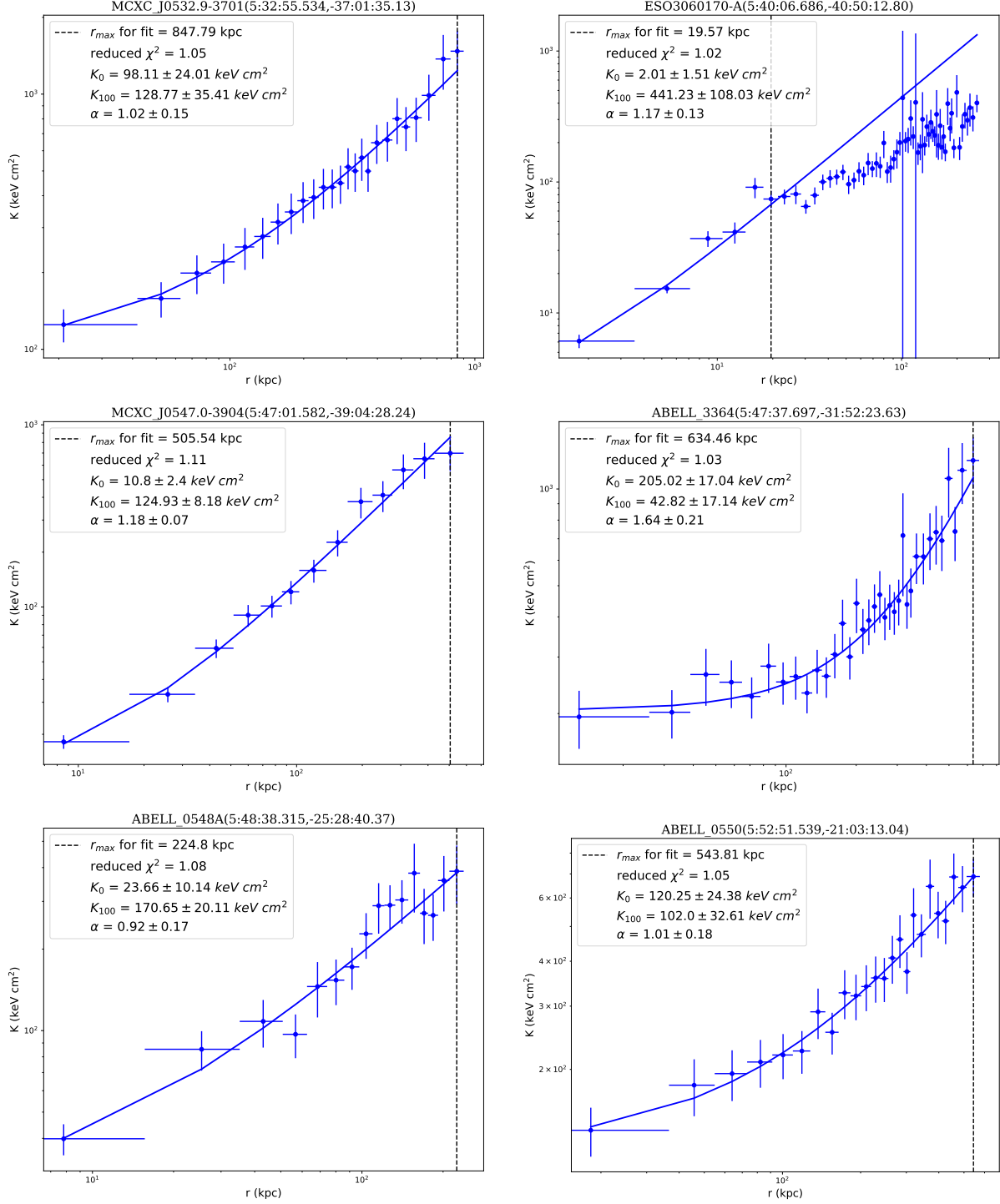


Figure C.1 (cont'd)

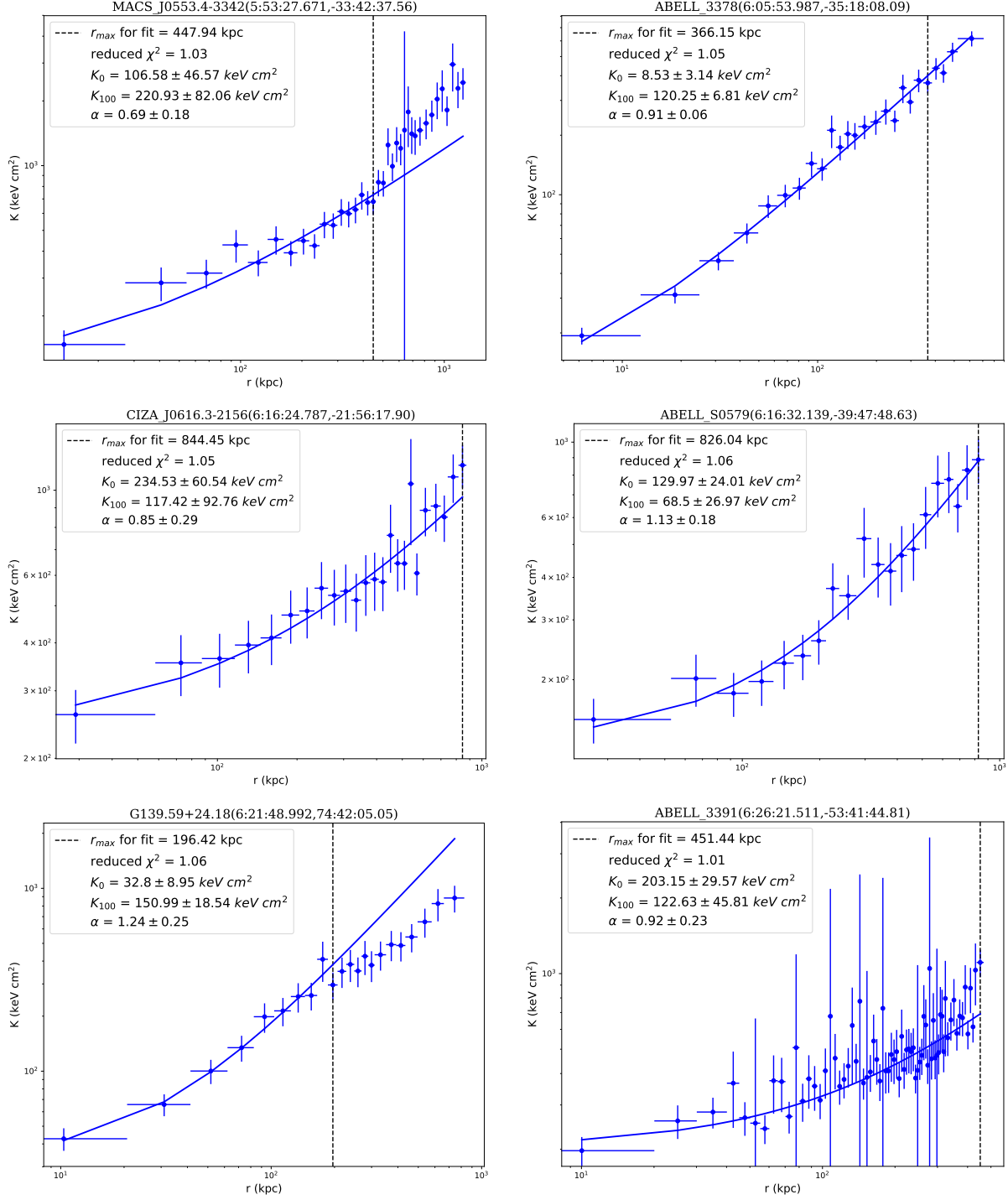


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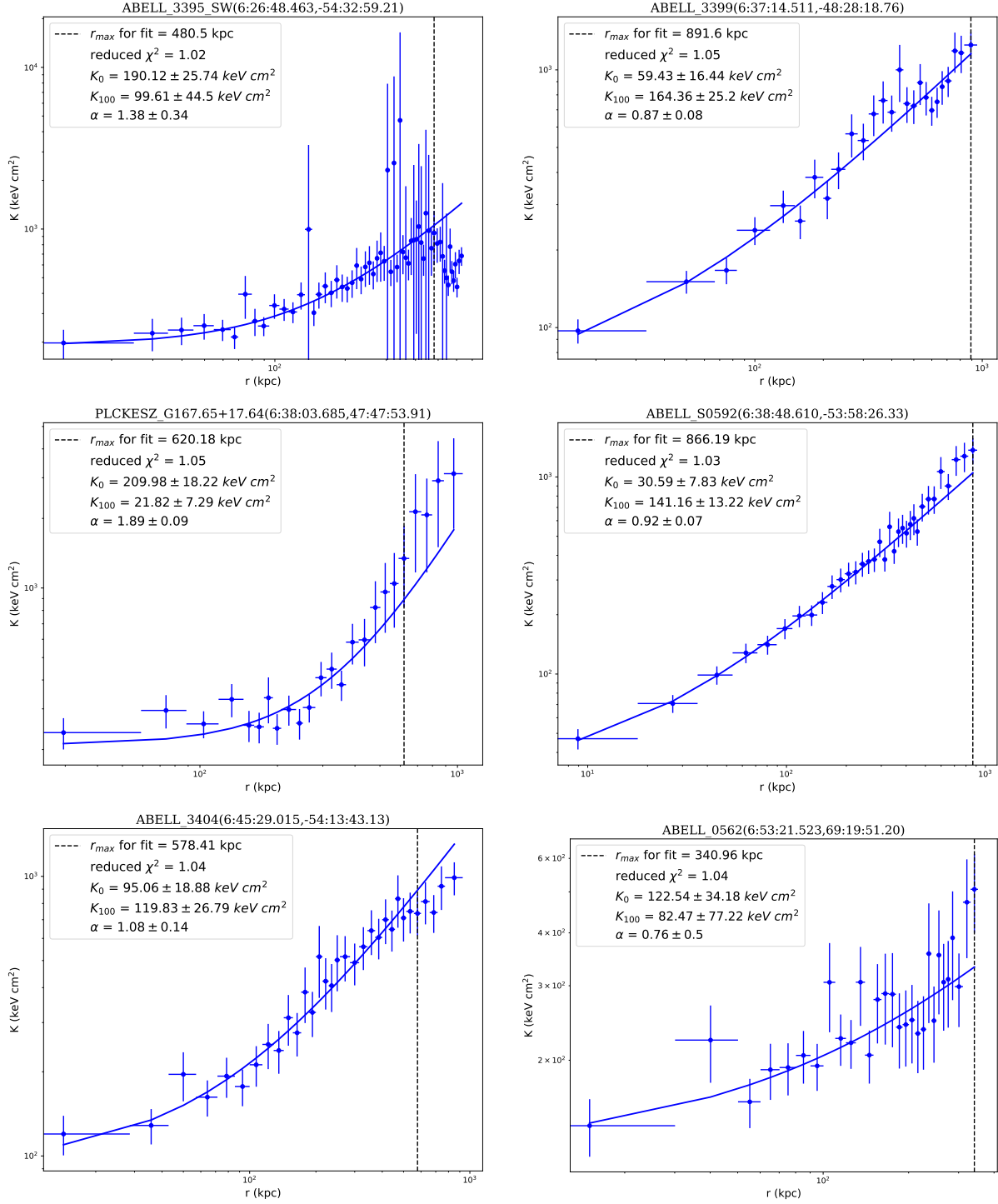


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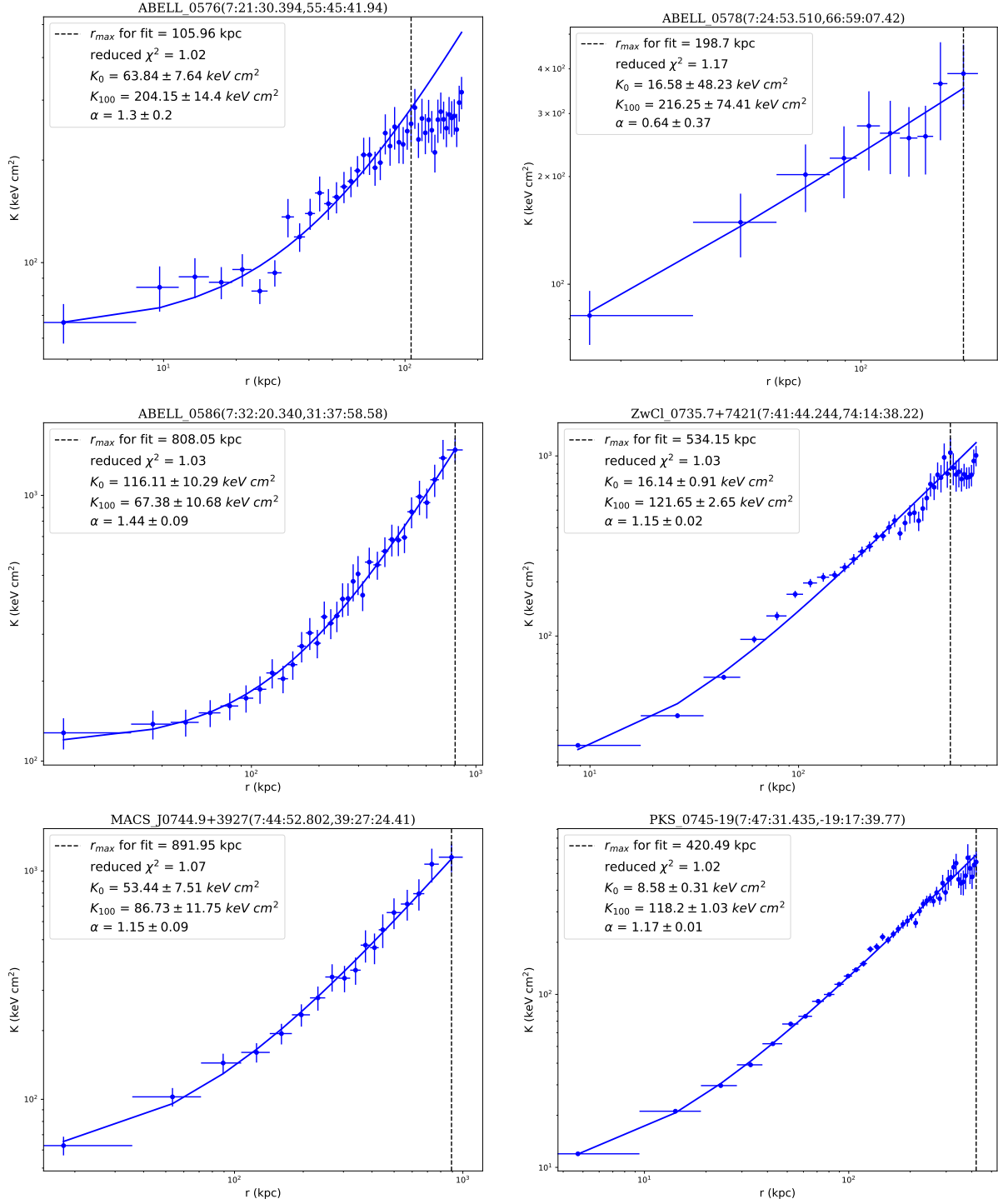


Figure C.1 (cont'd)

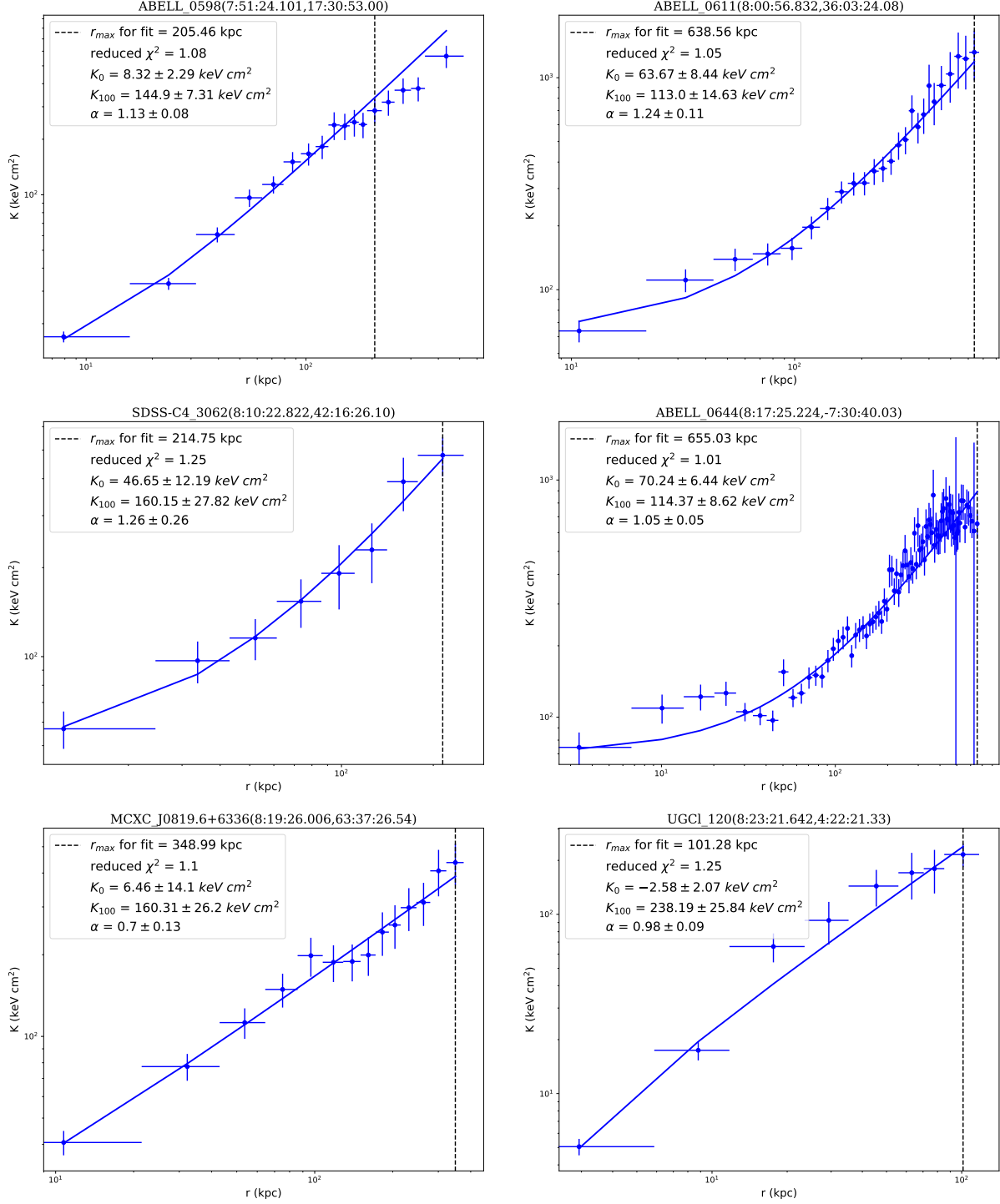


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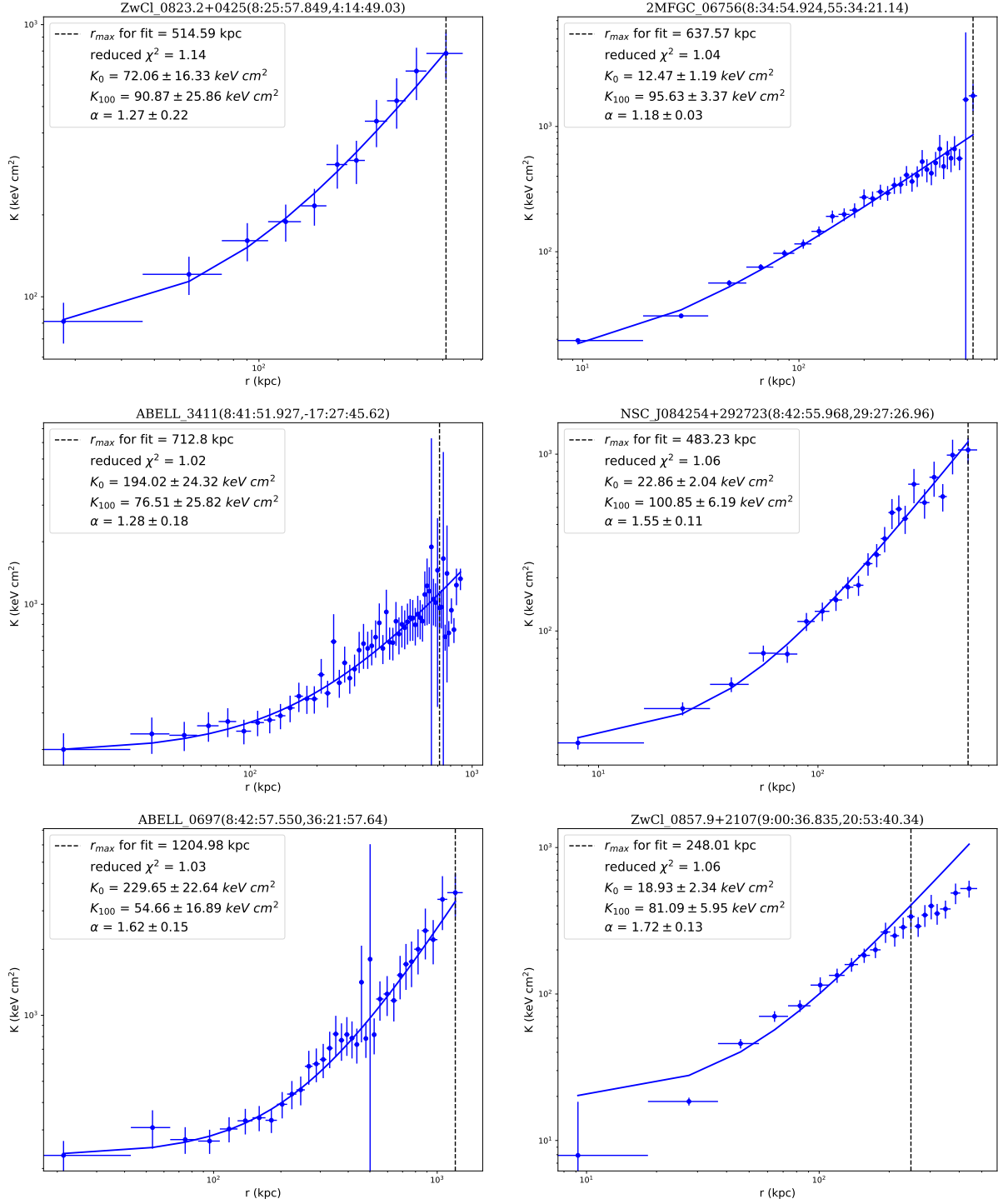


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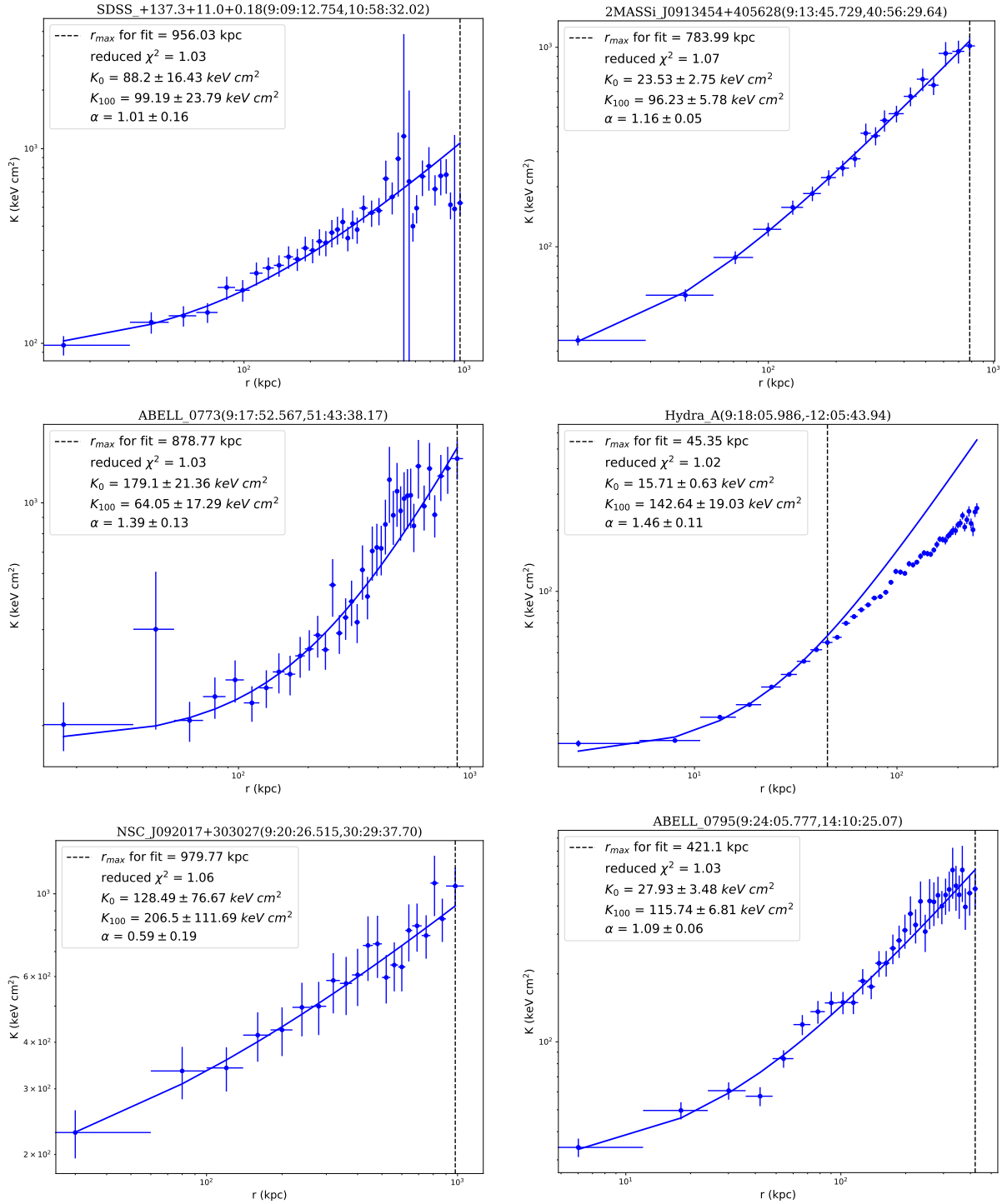


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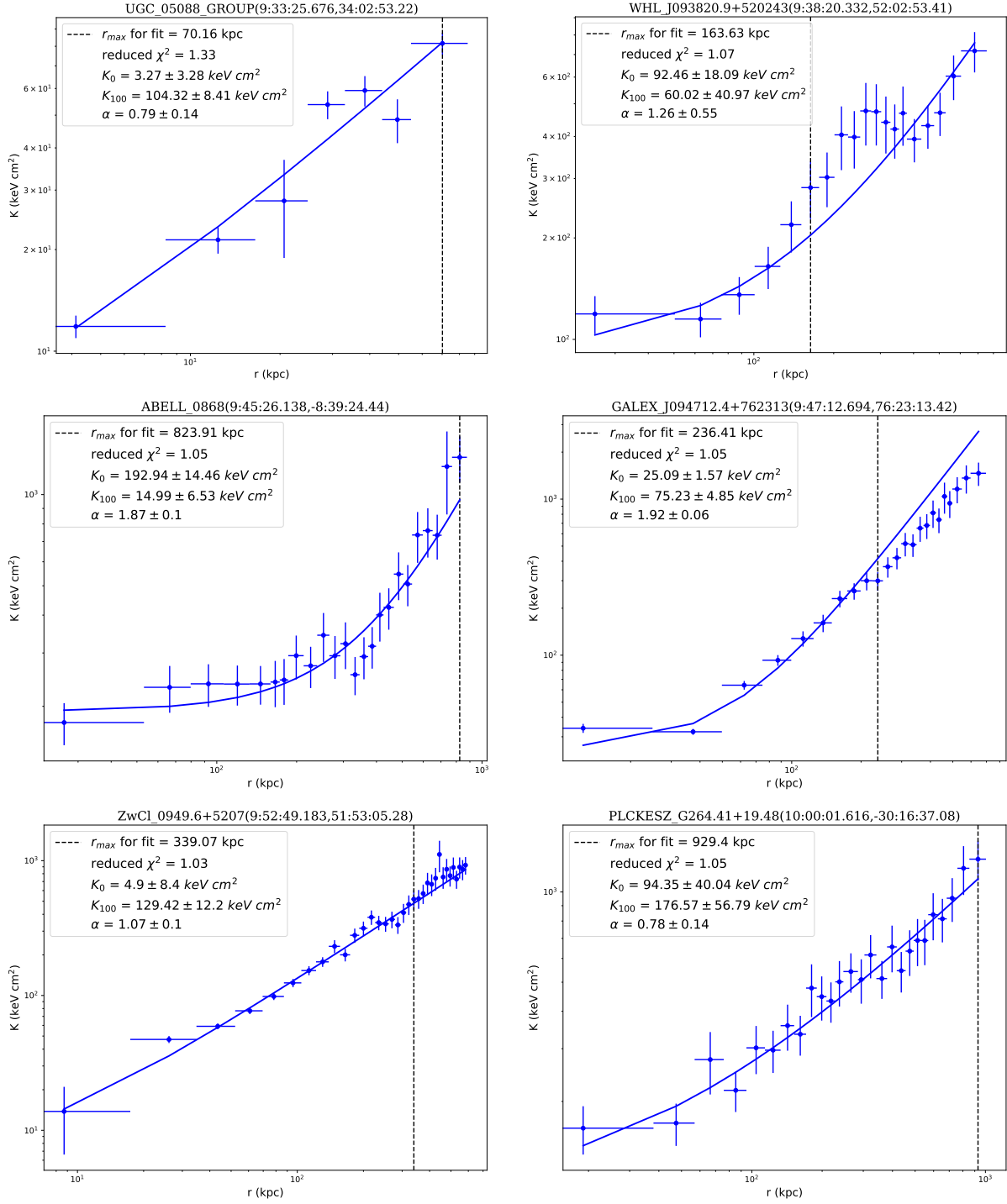


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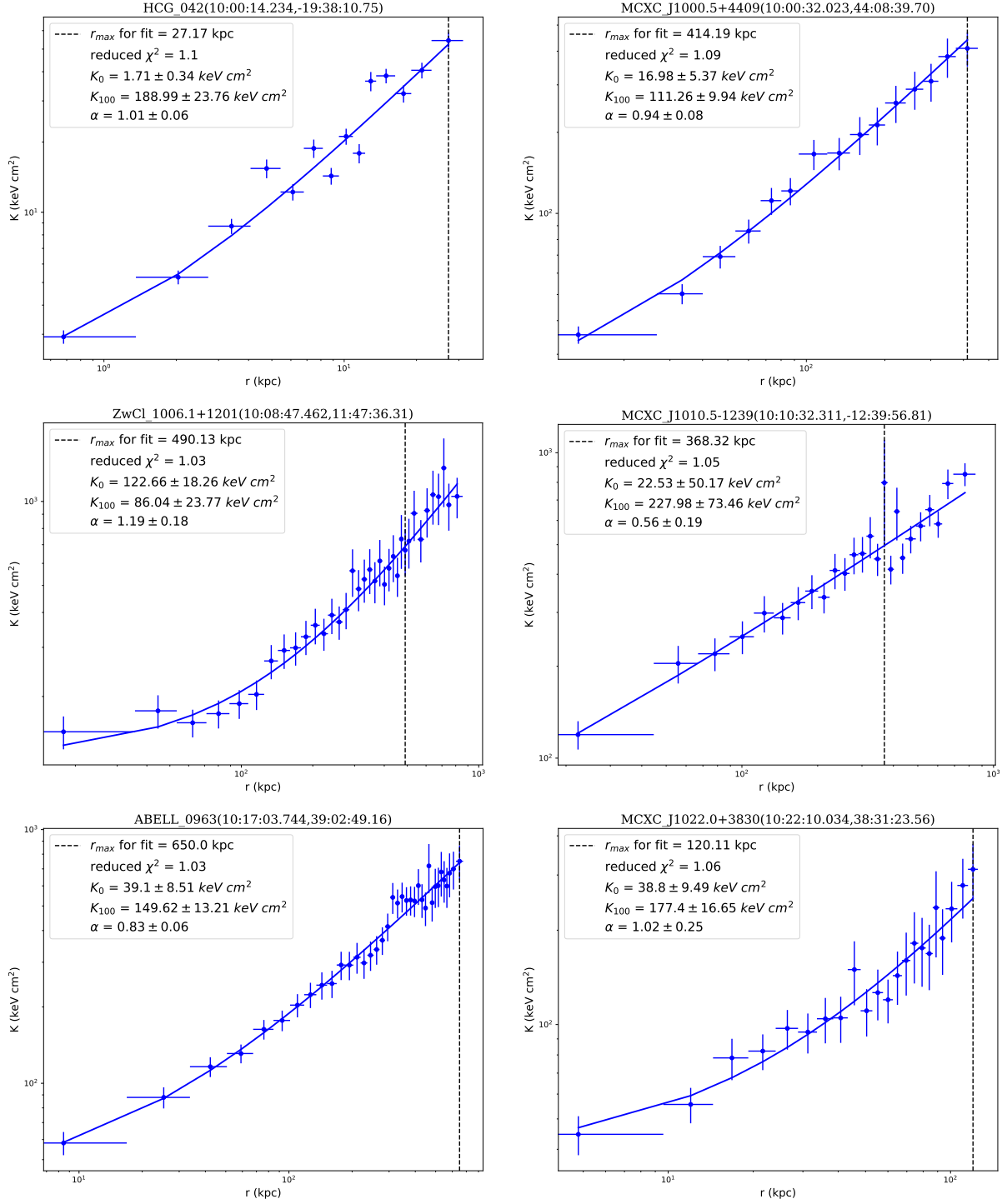


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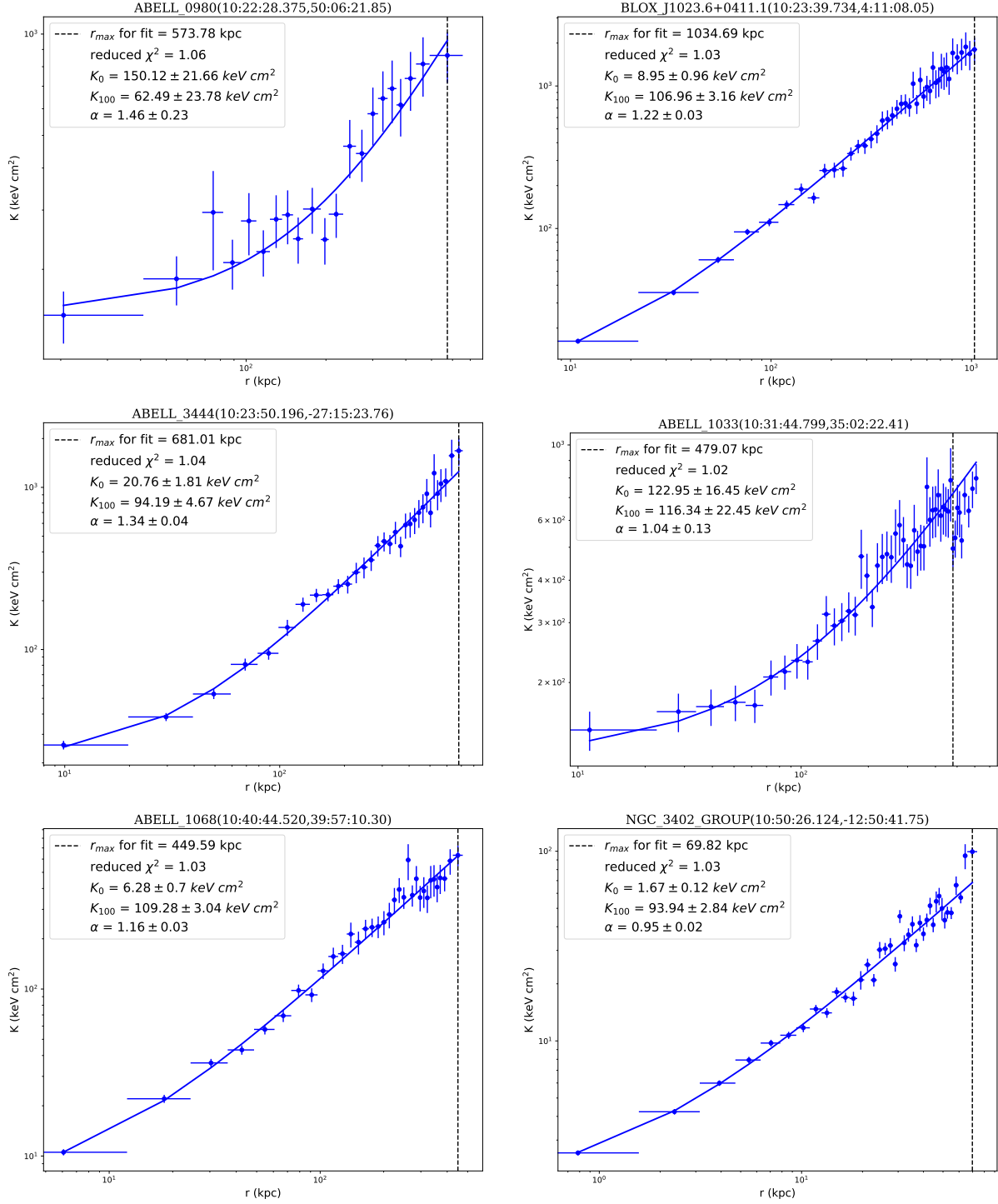


Figure C.1 (cont'd)

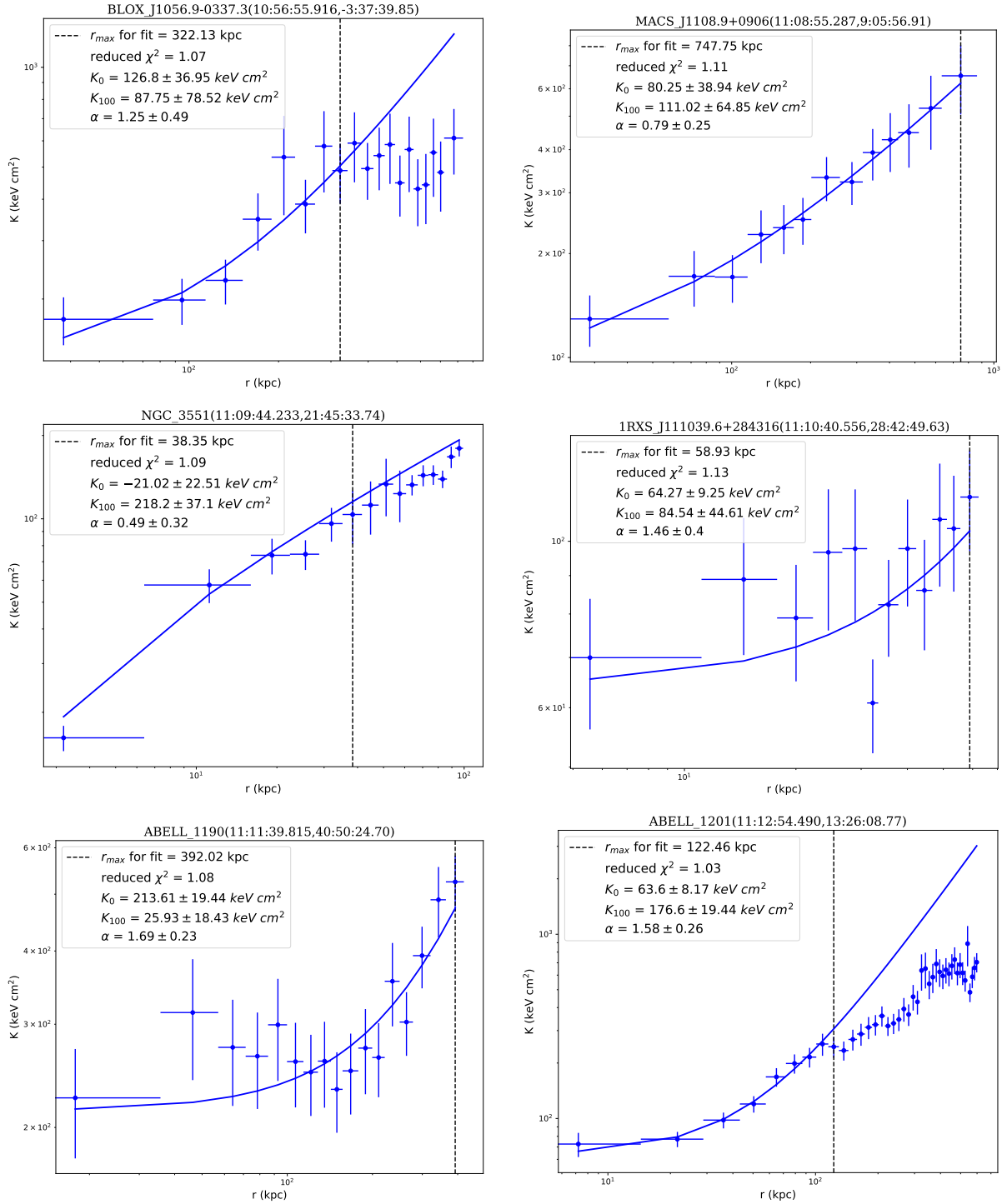


Figure C.1 (cont'd)

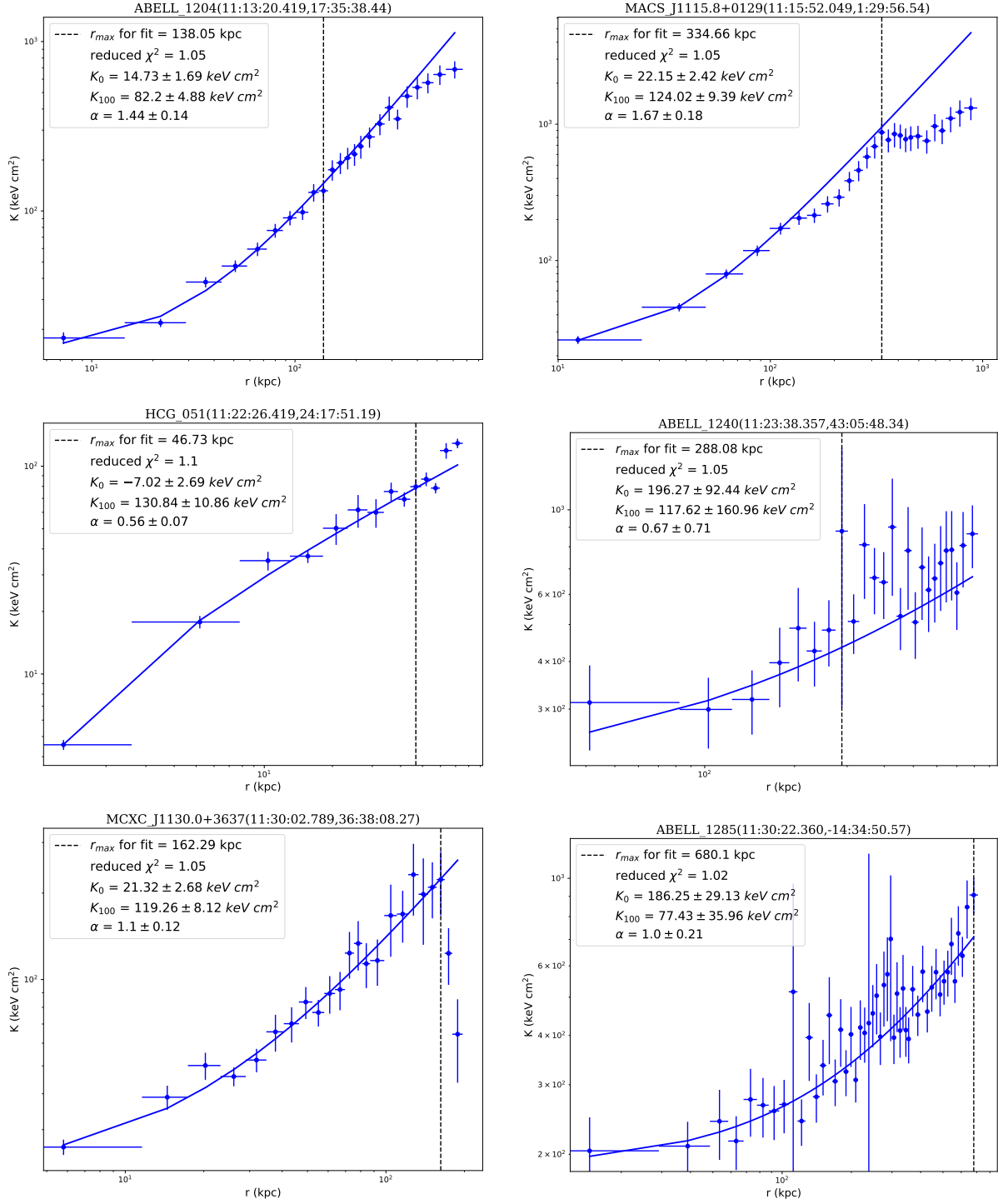


Figure C.1 (cont'd)

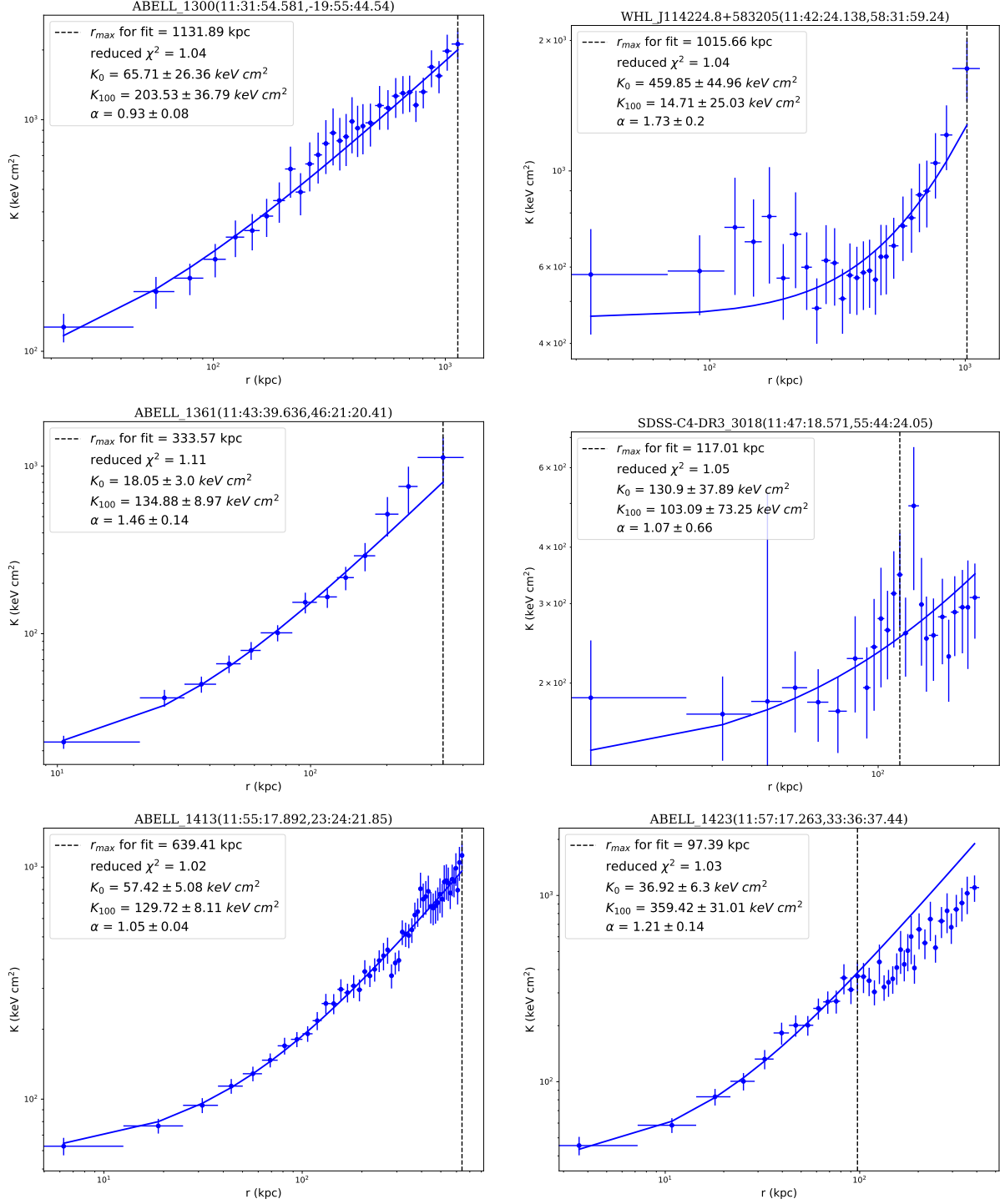


Figure C.1 (cont'd)

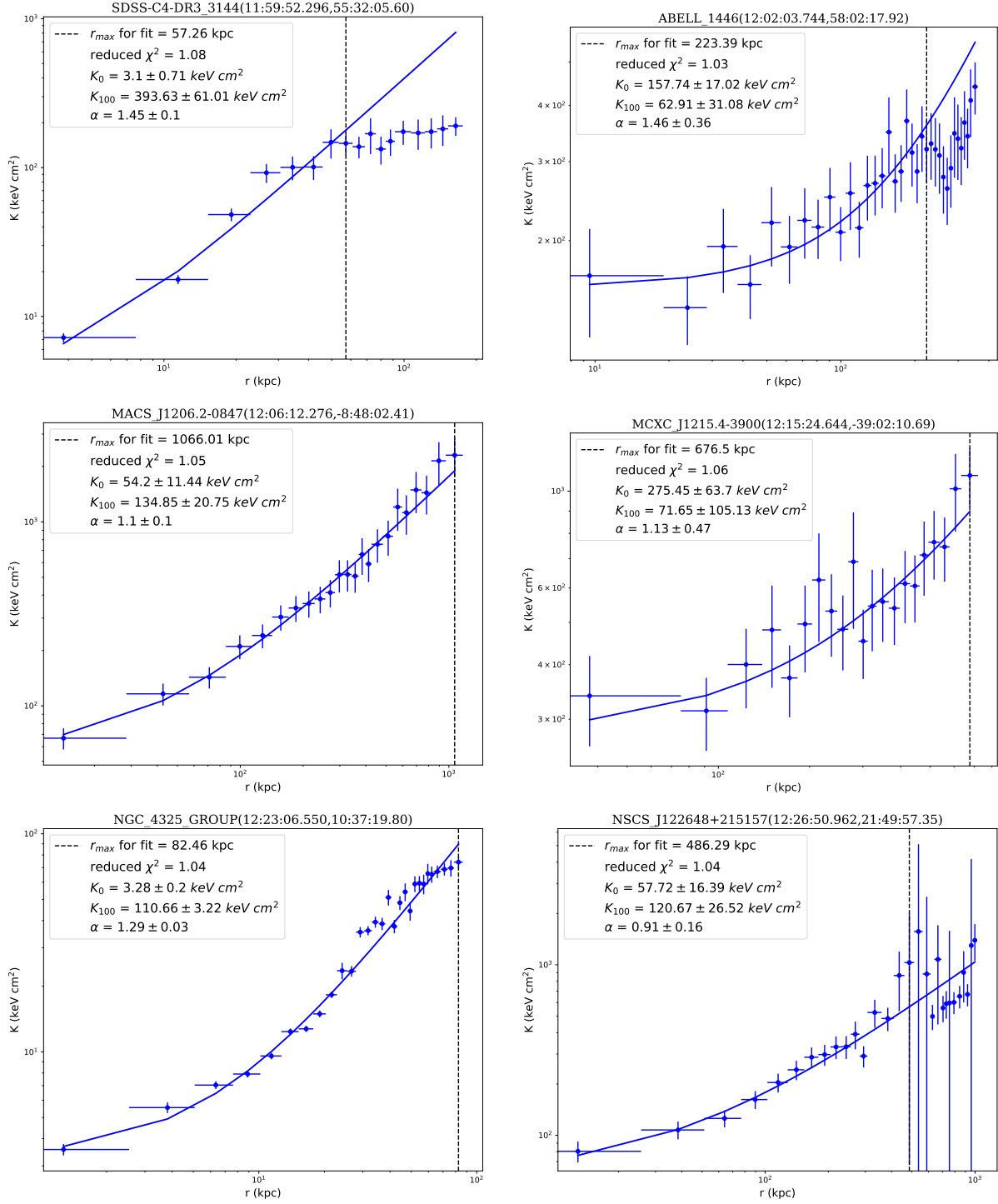


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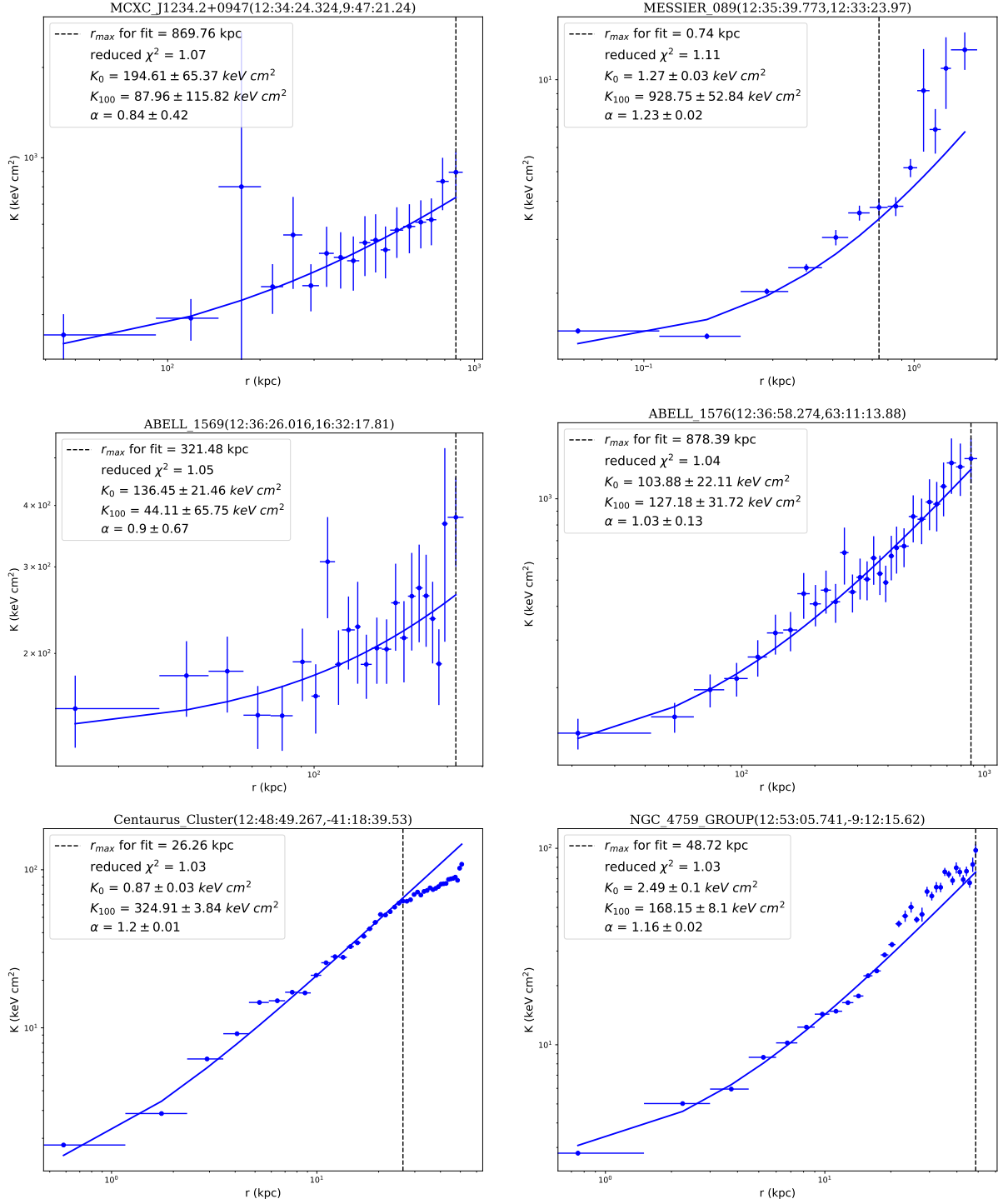


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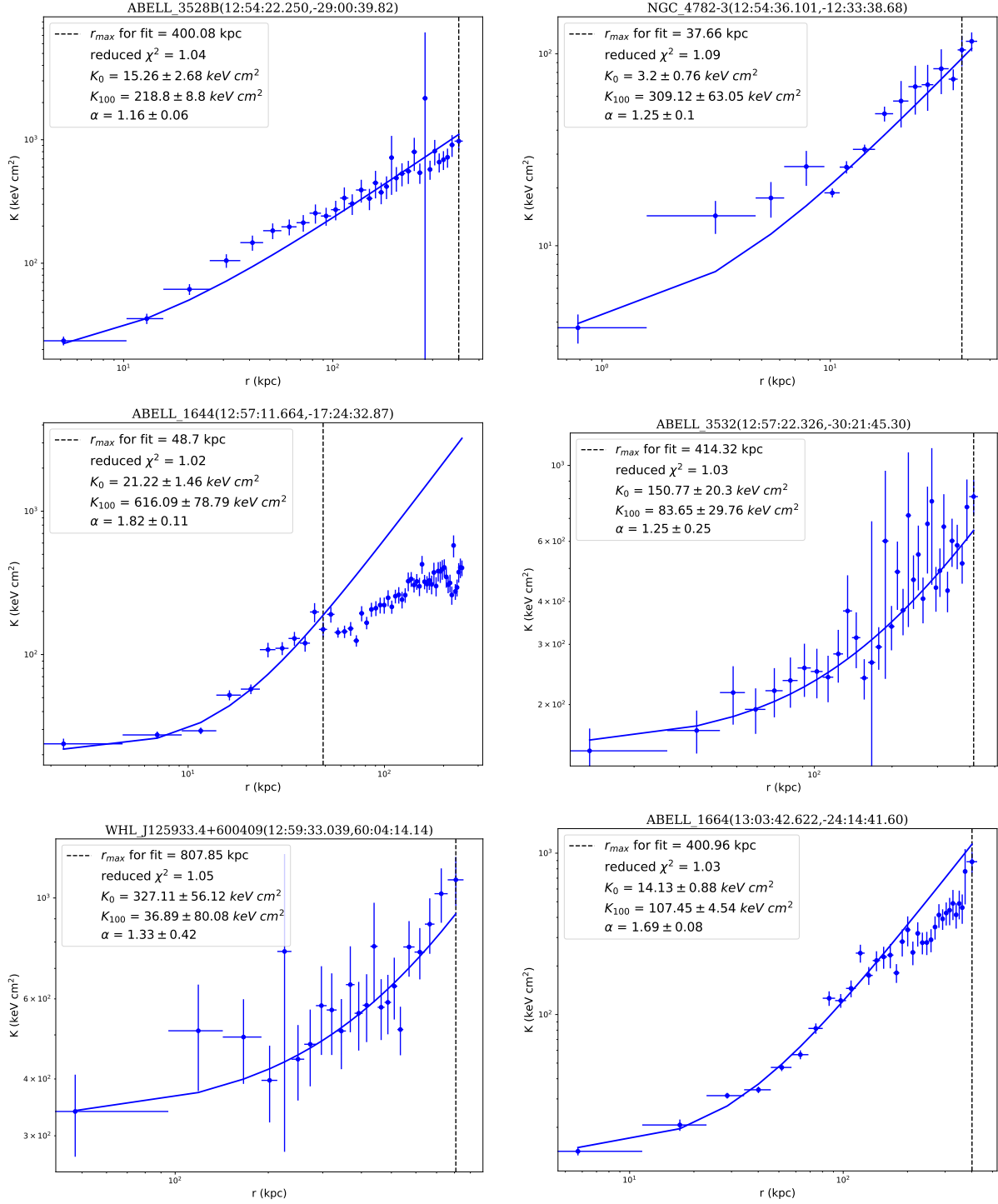


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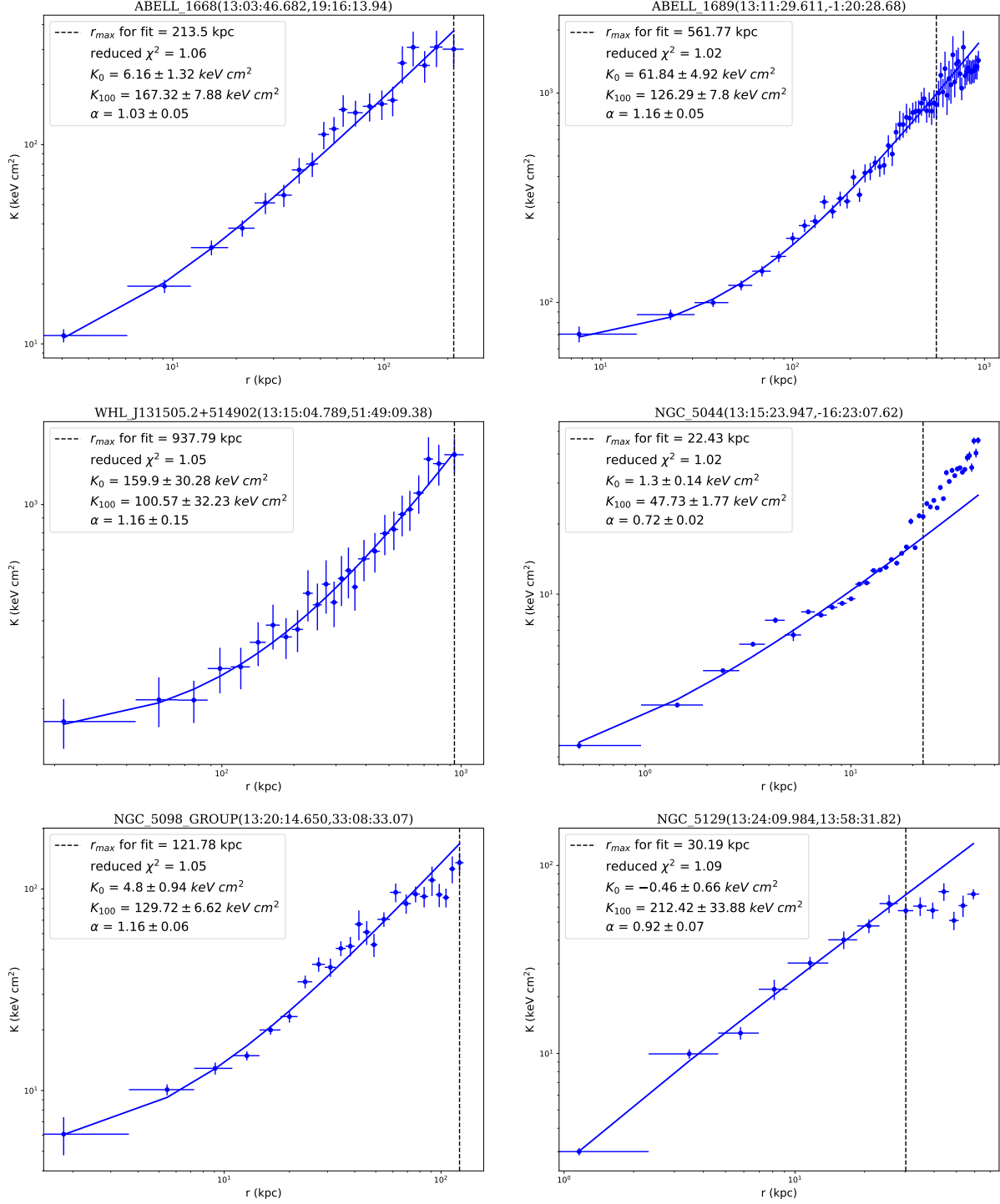


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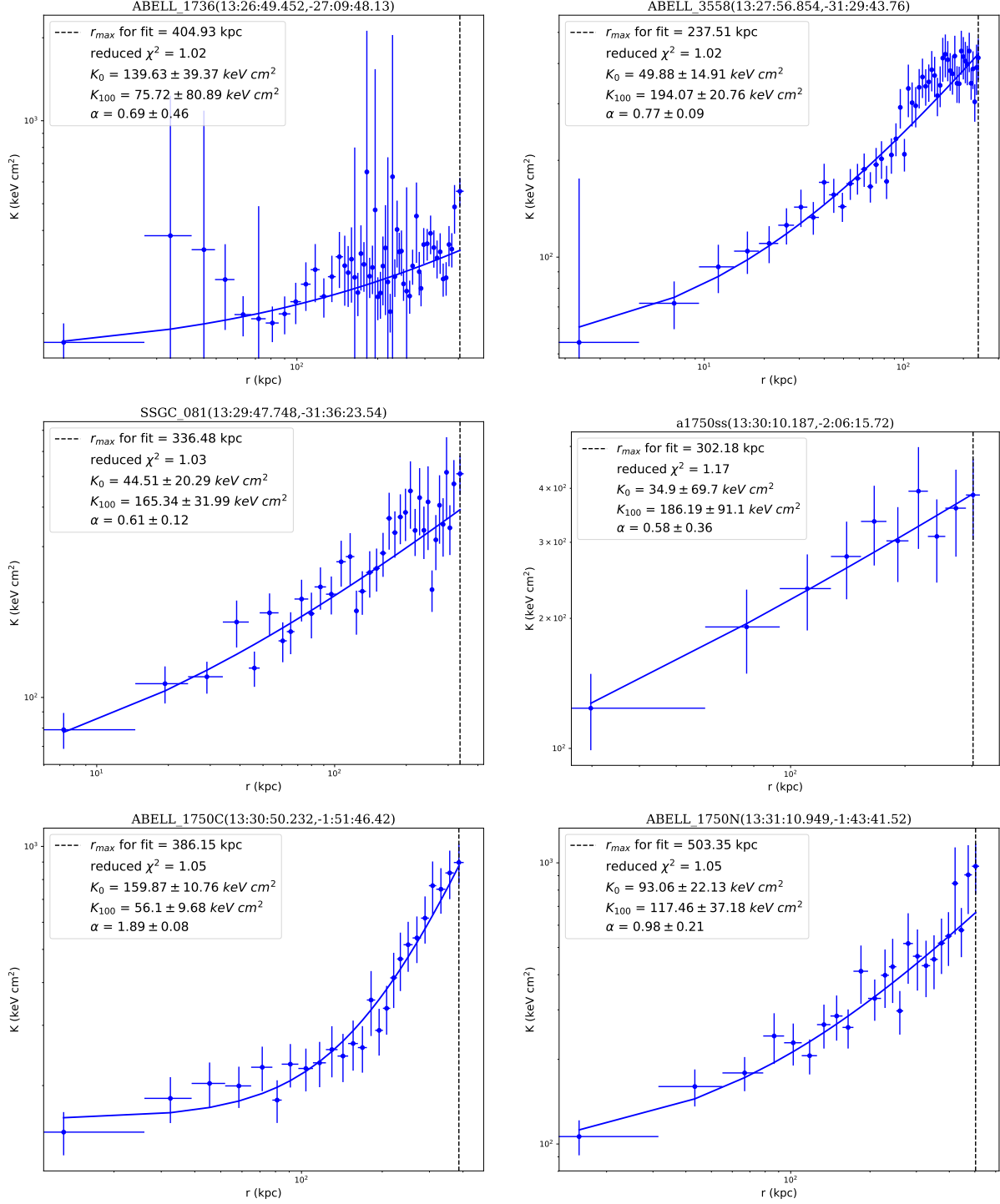


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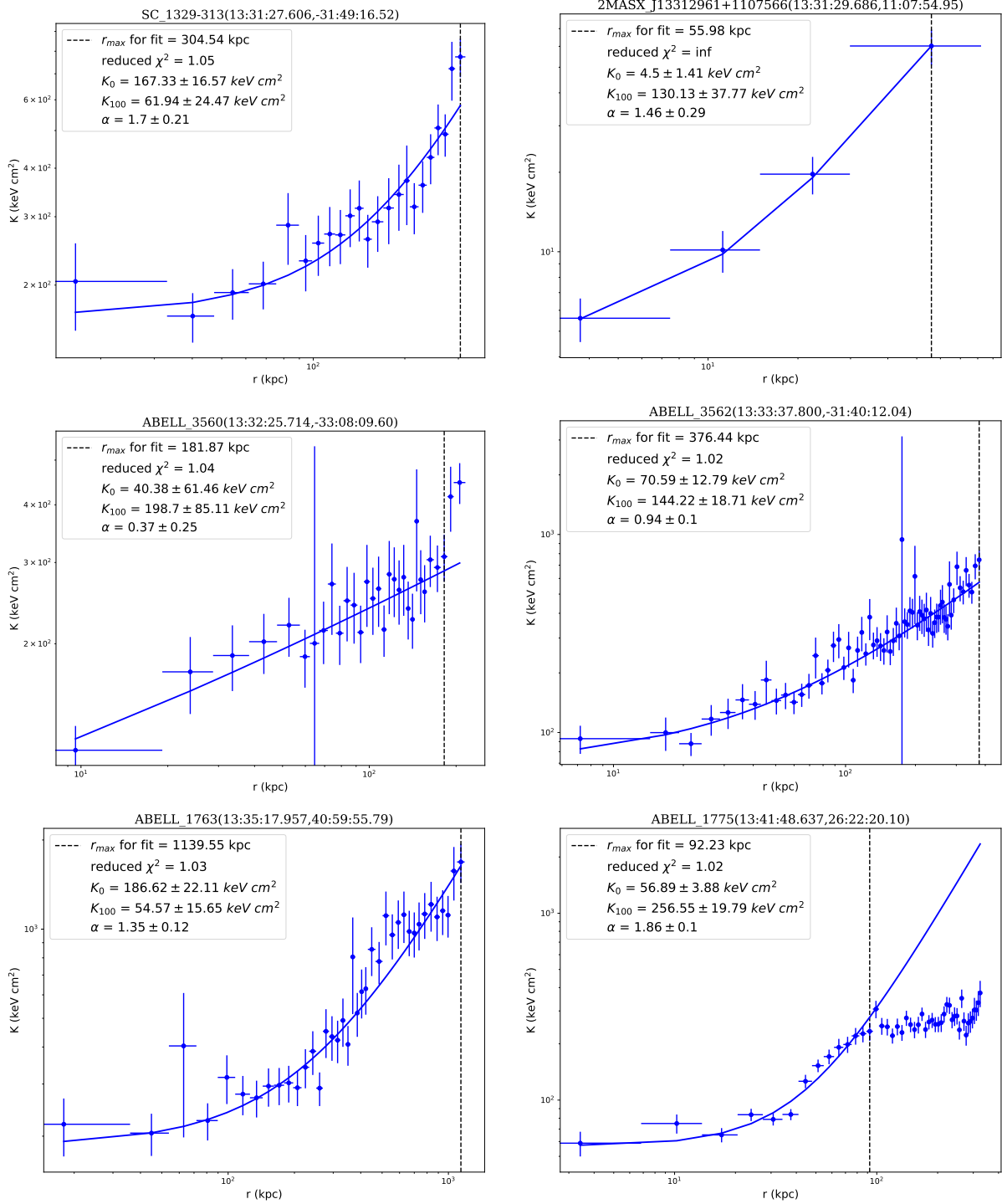


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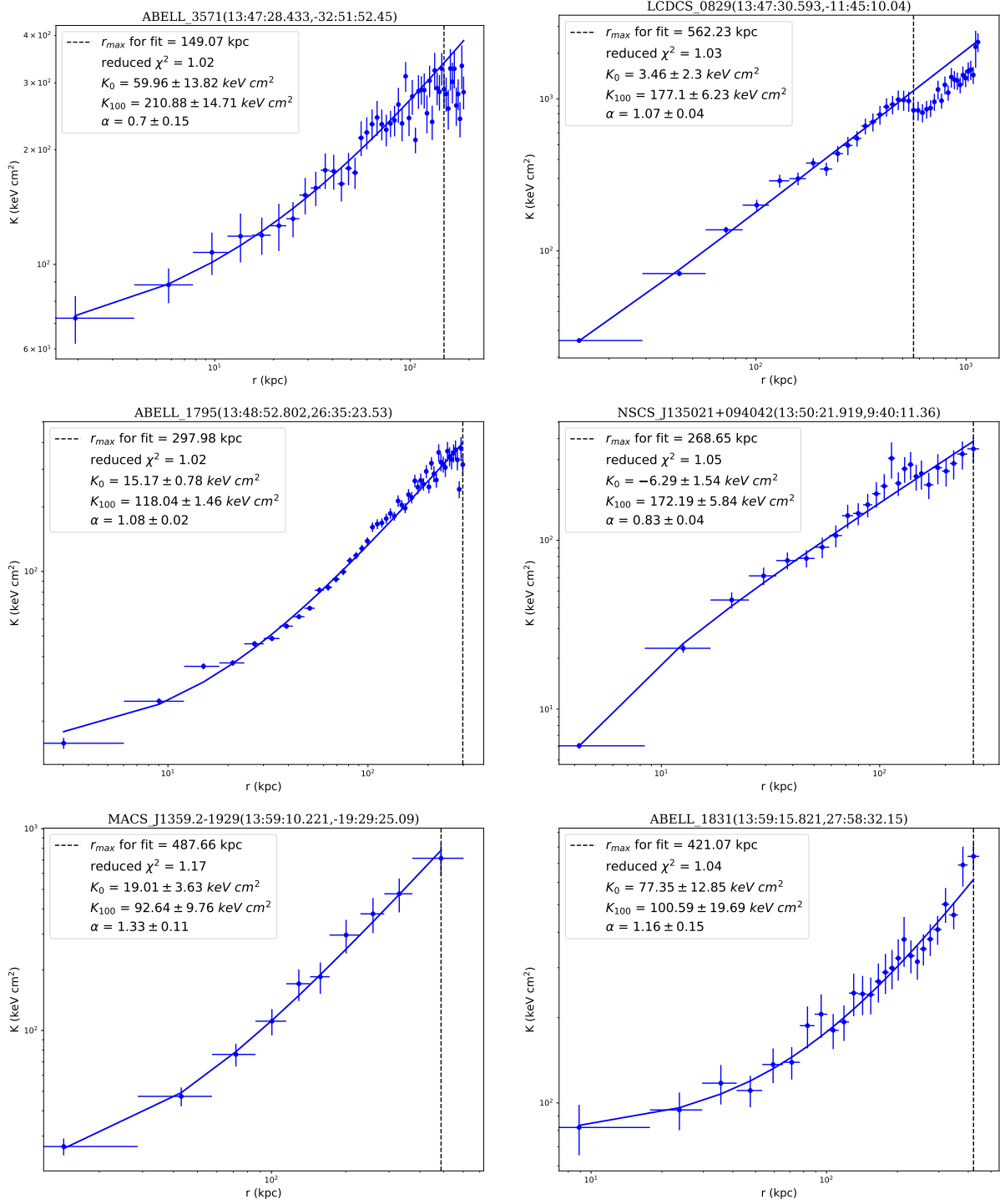


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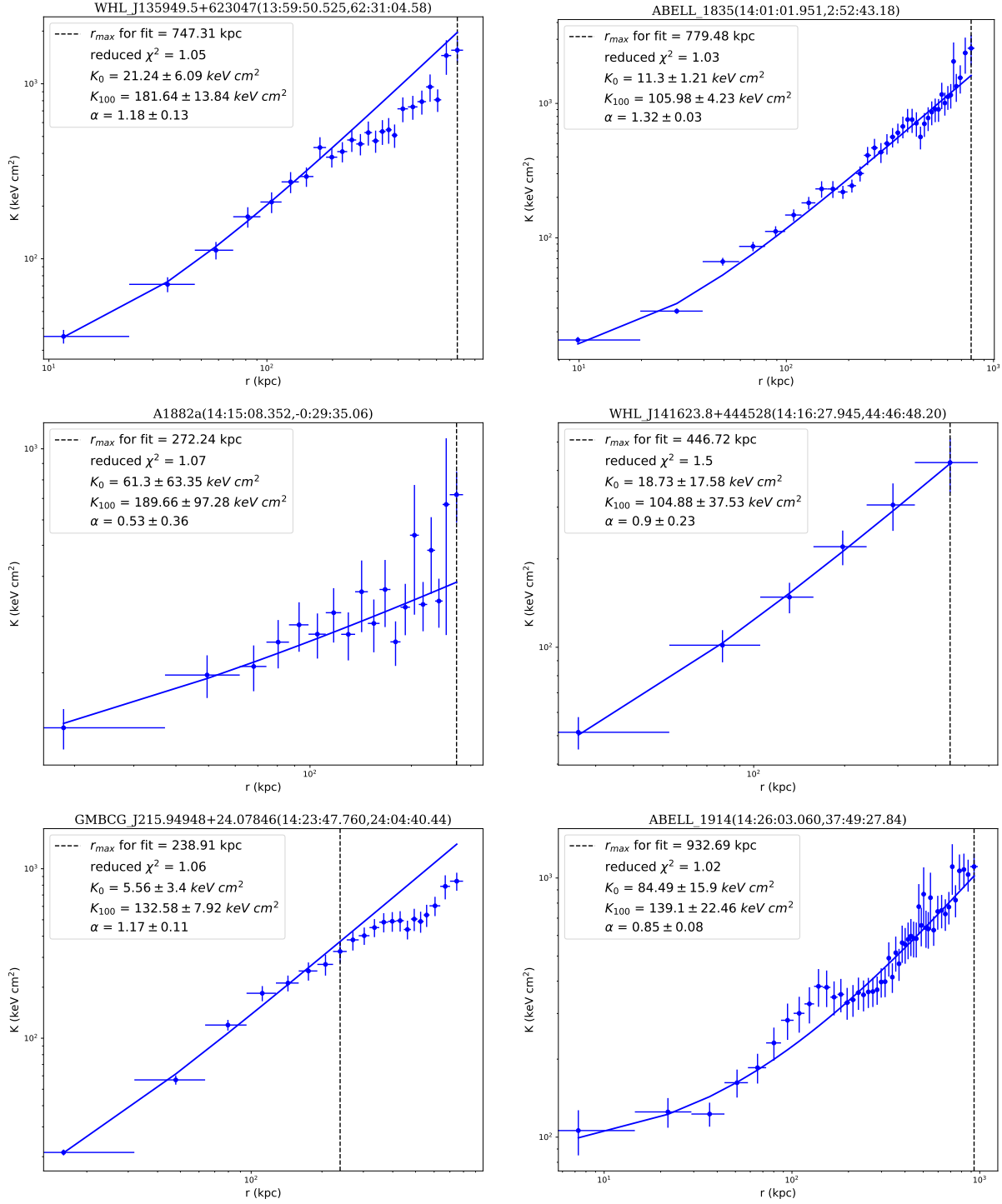


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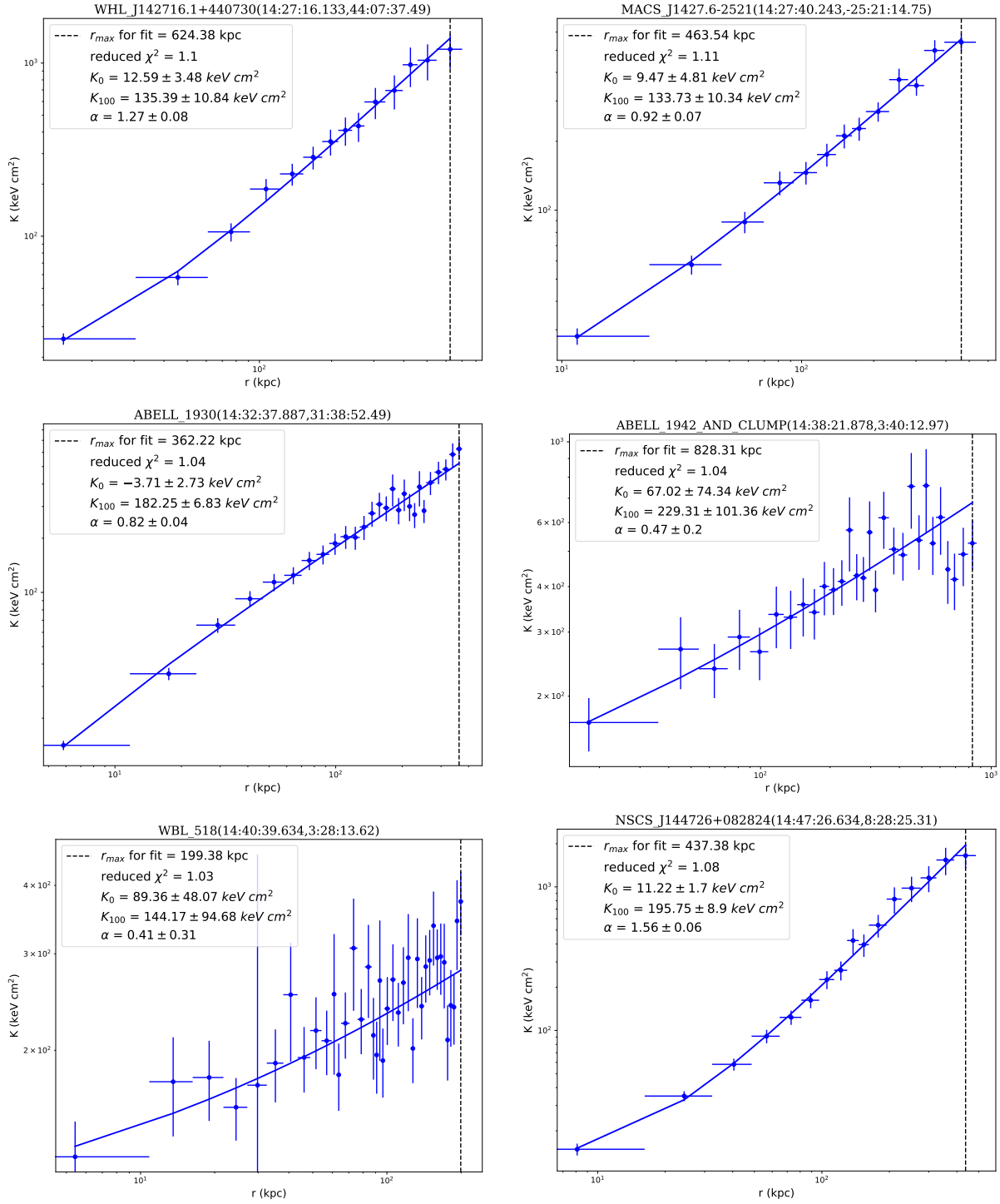


Figure C.1 (cont'd)

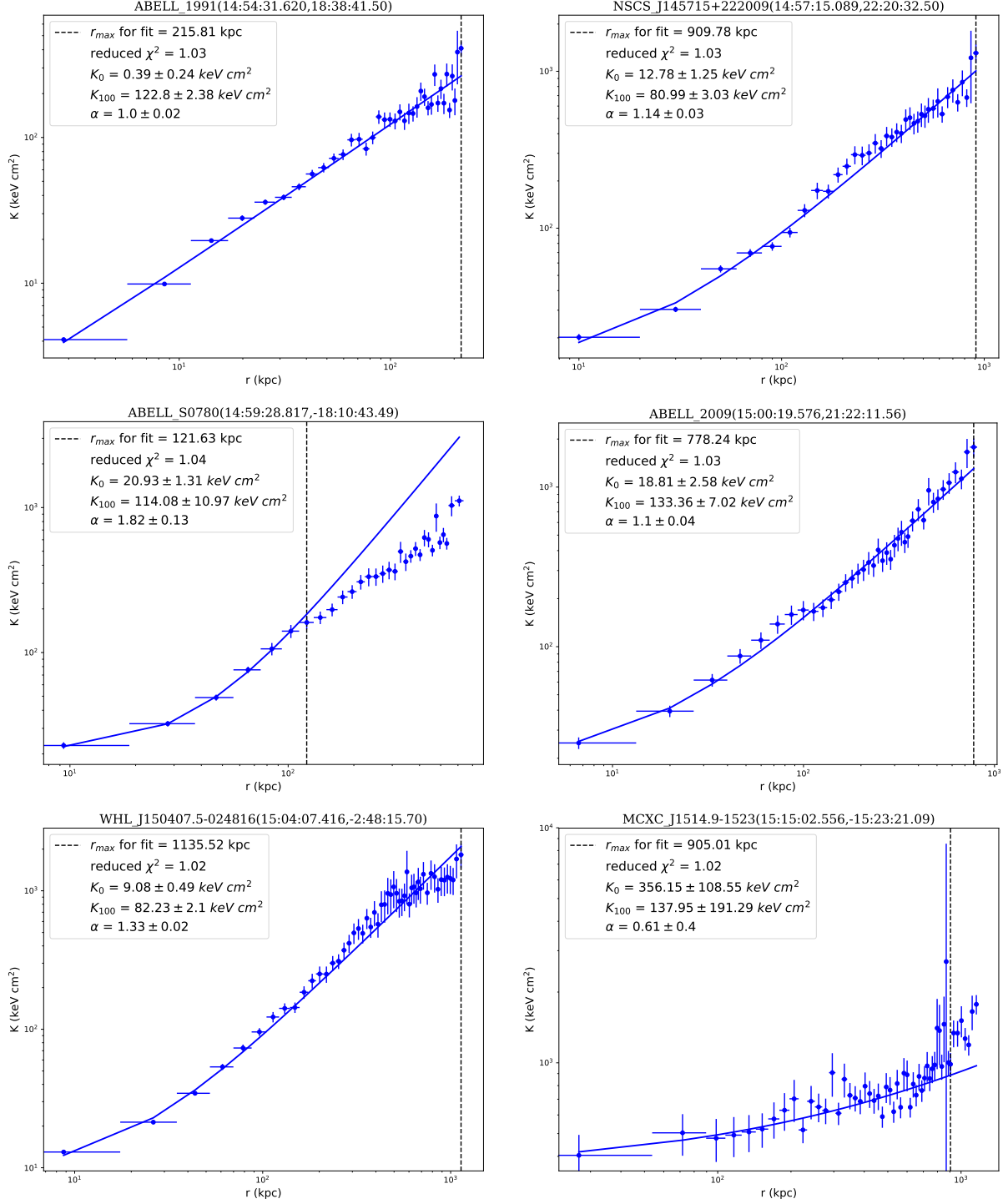


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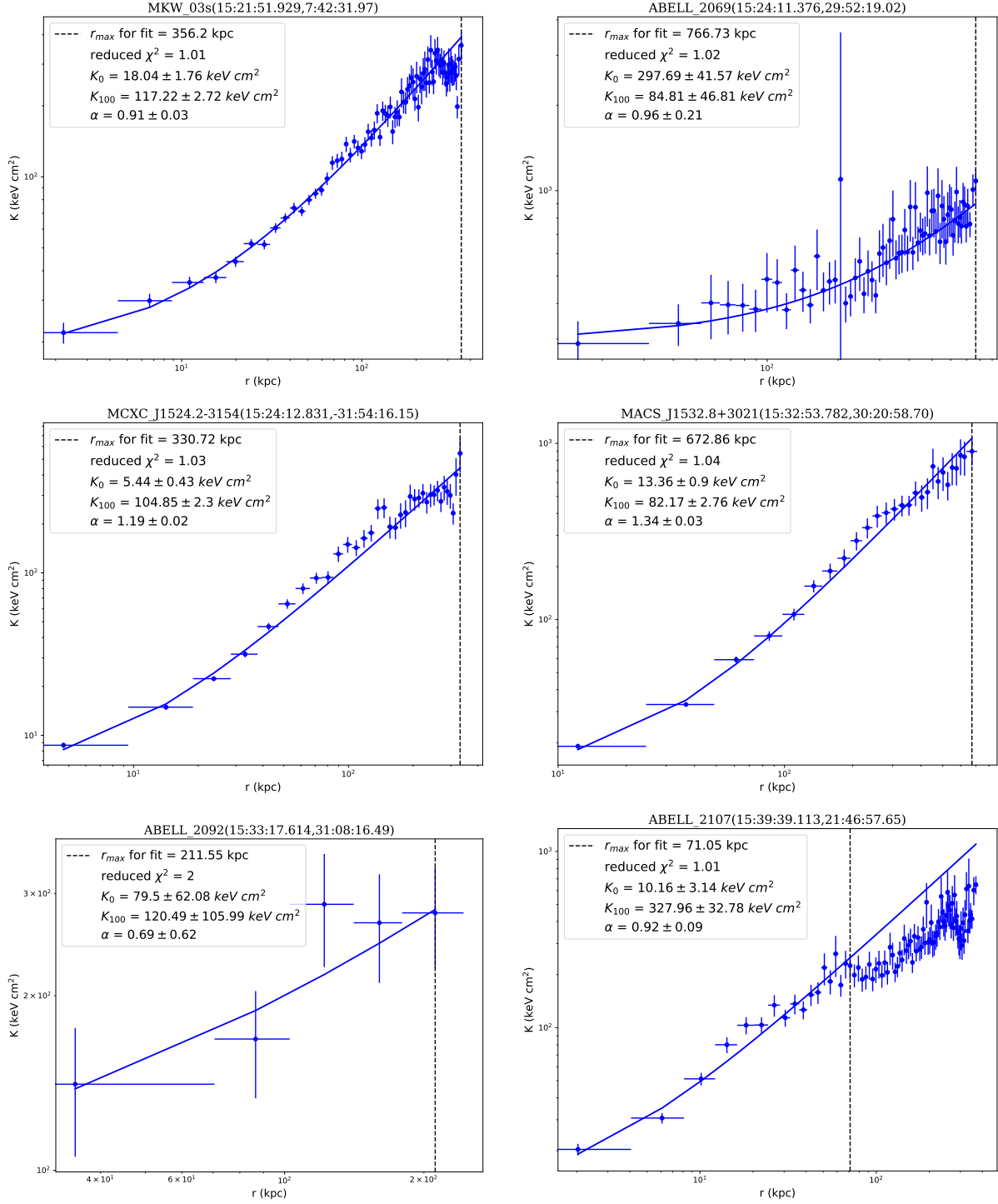


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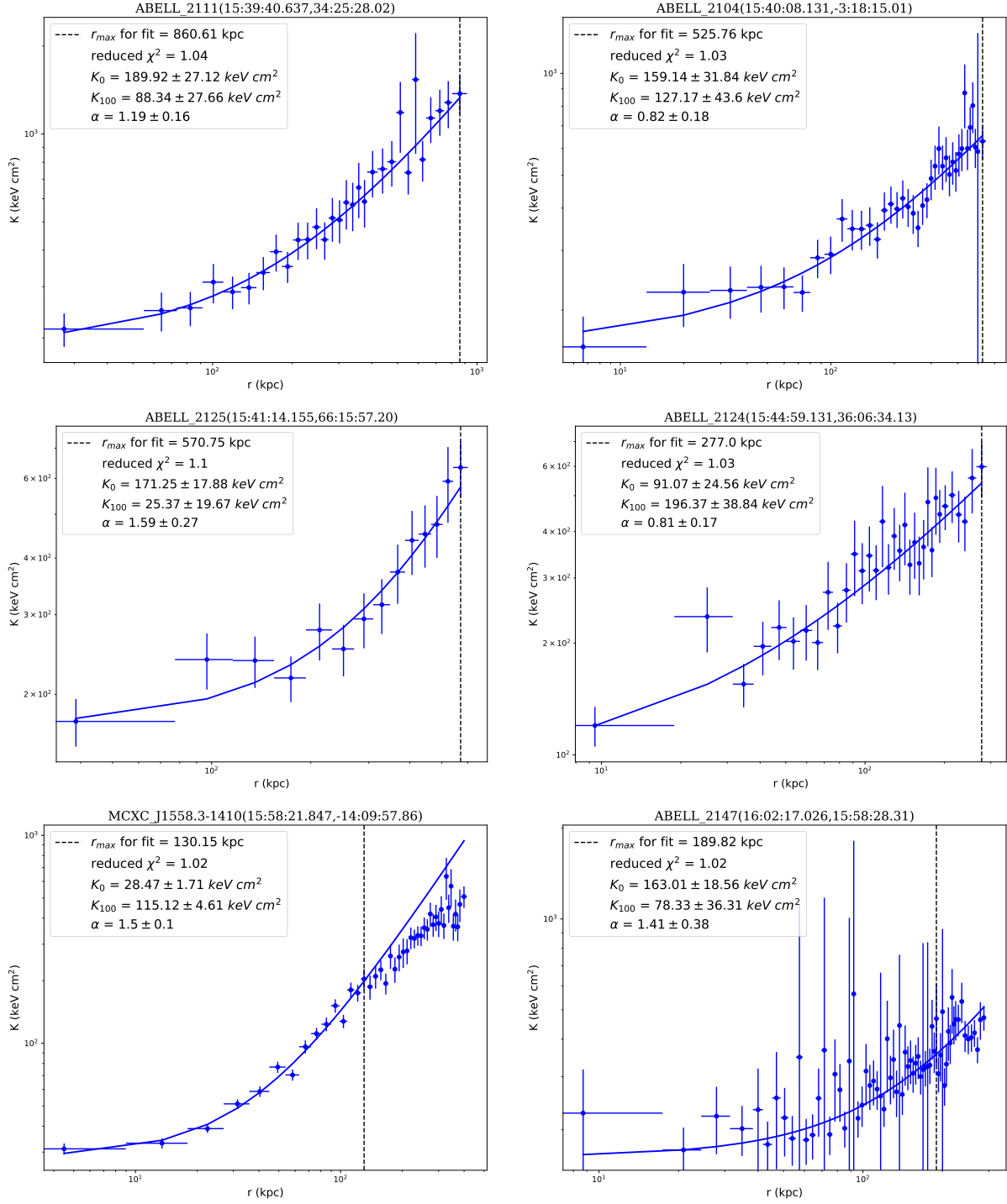


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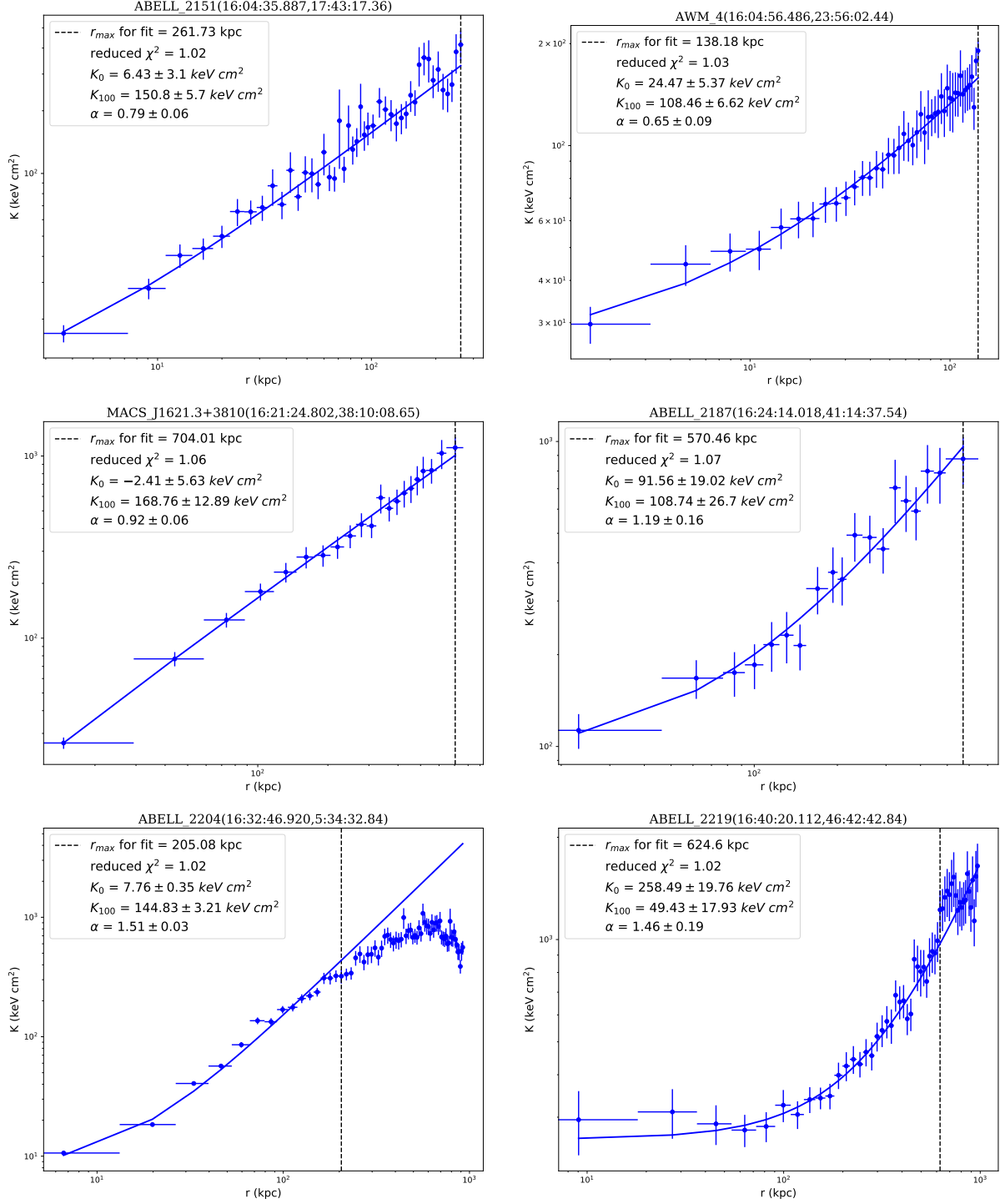


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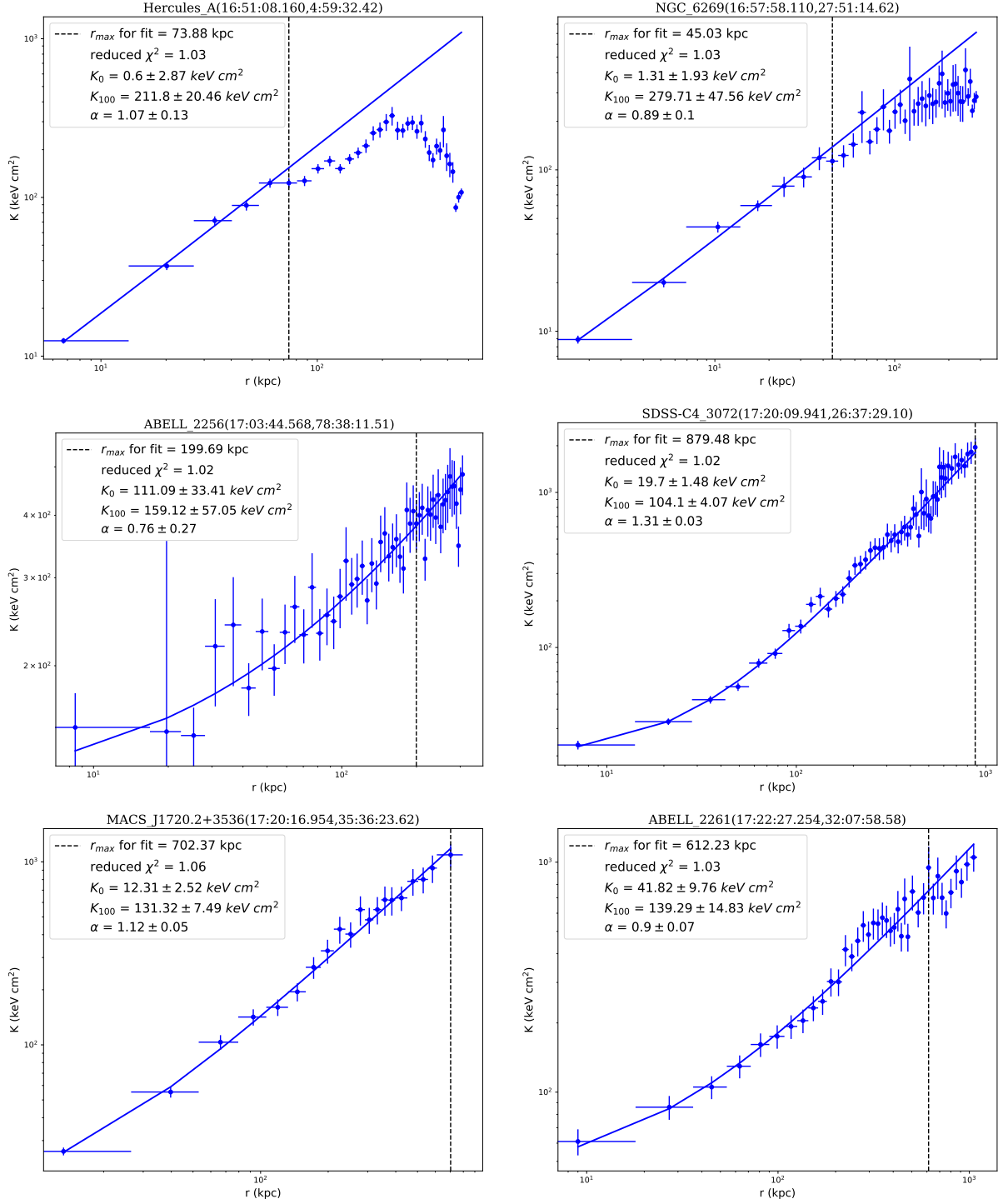


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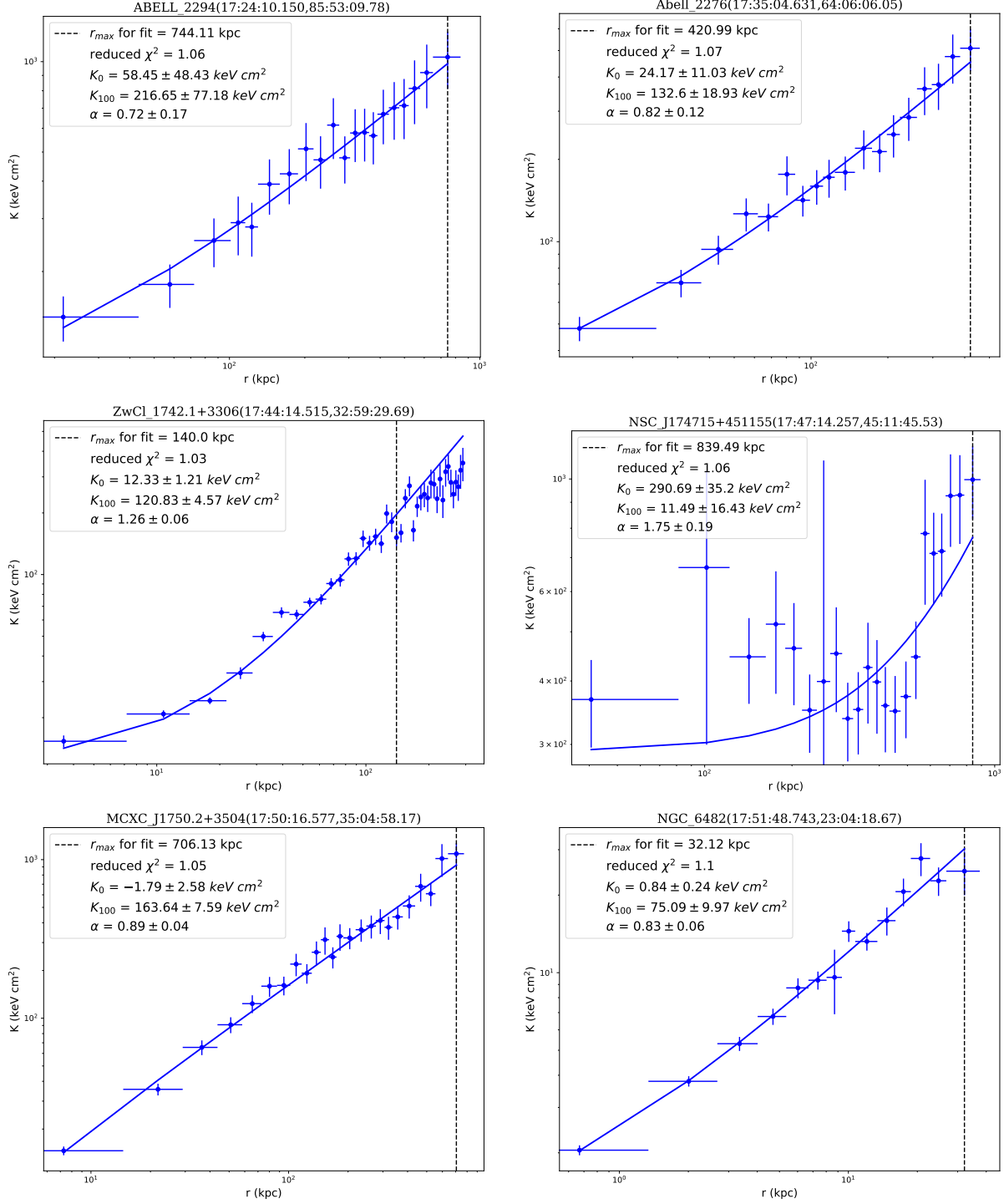


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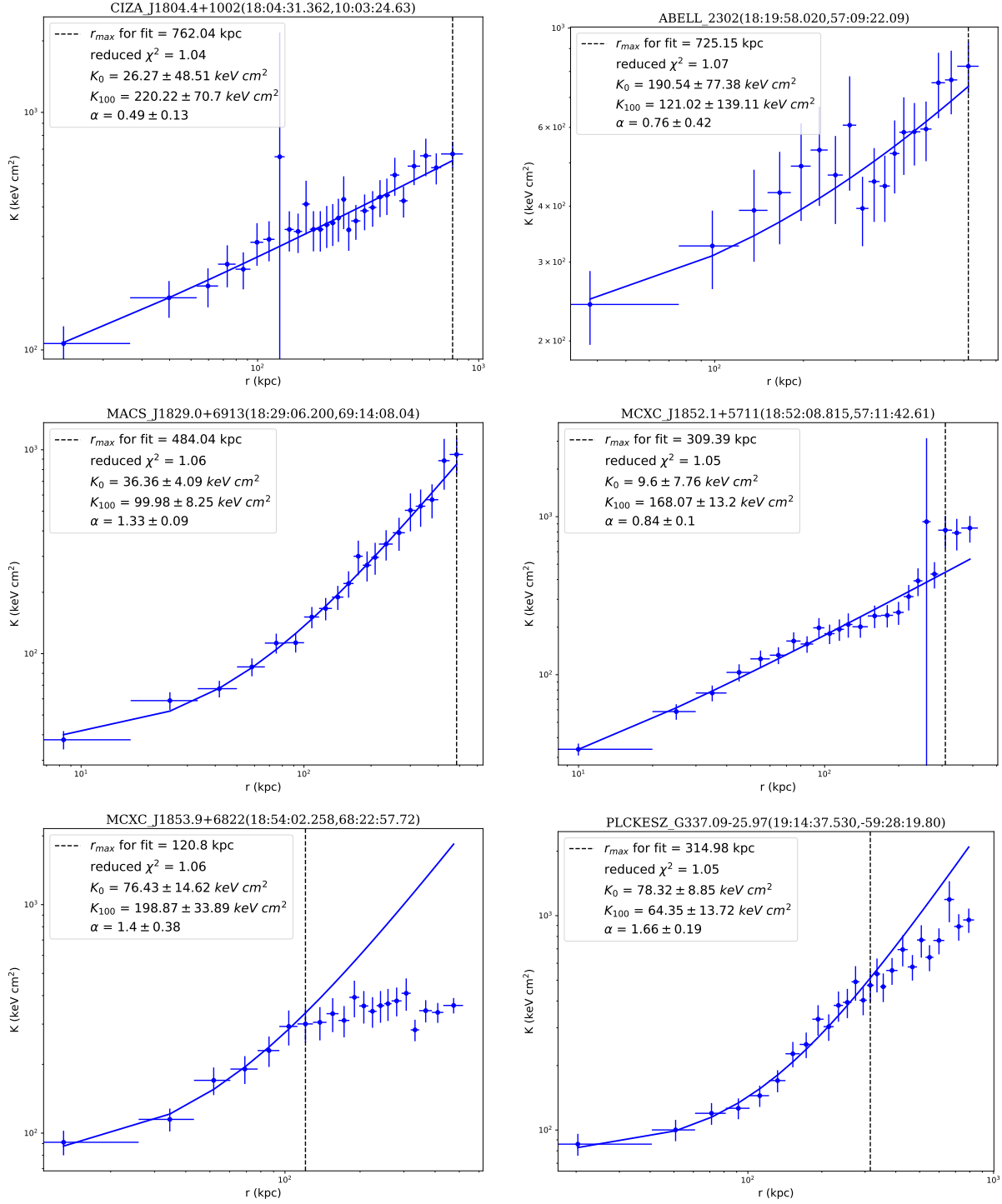


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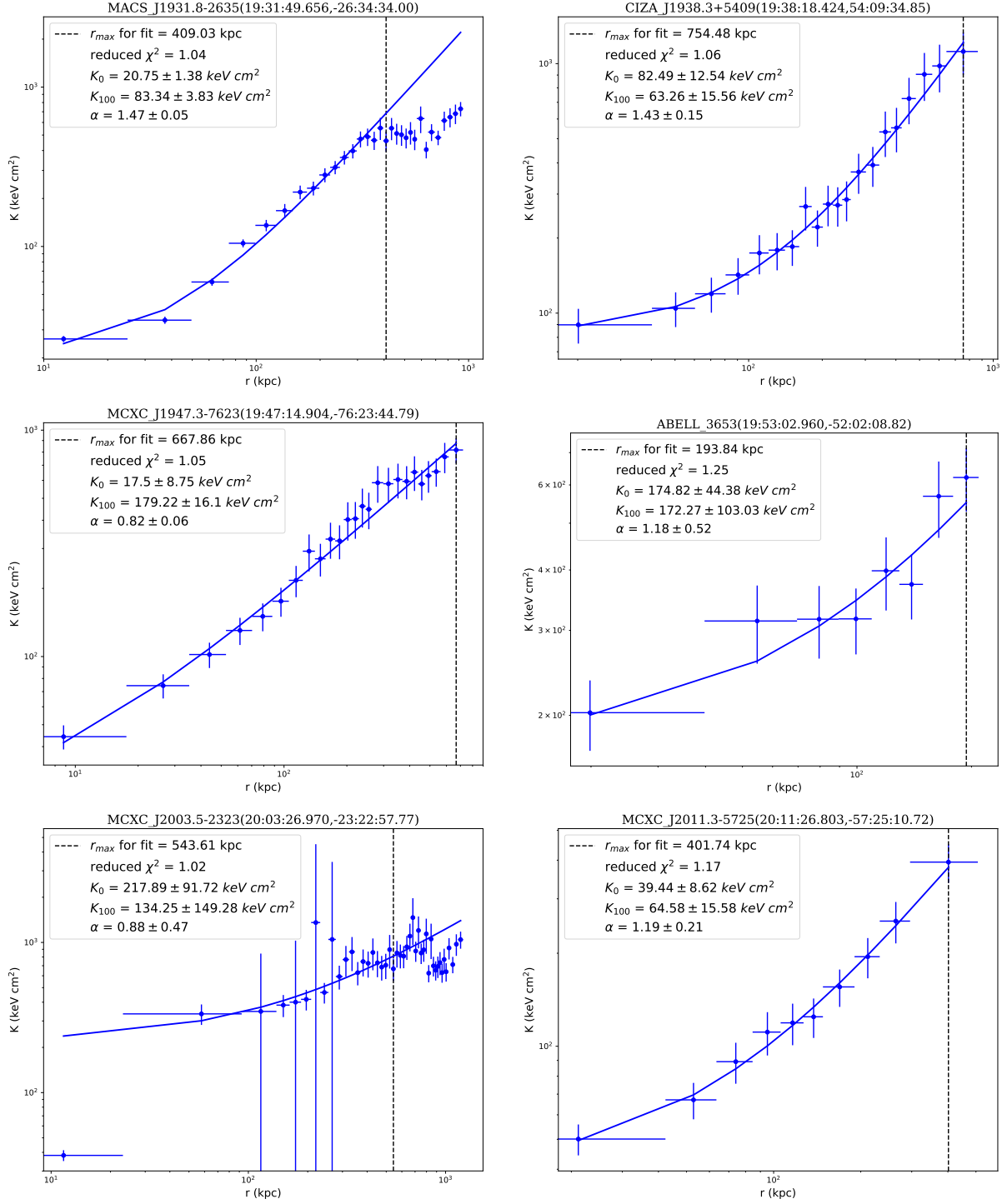


Figure C.1 (cont'd)

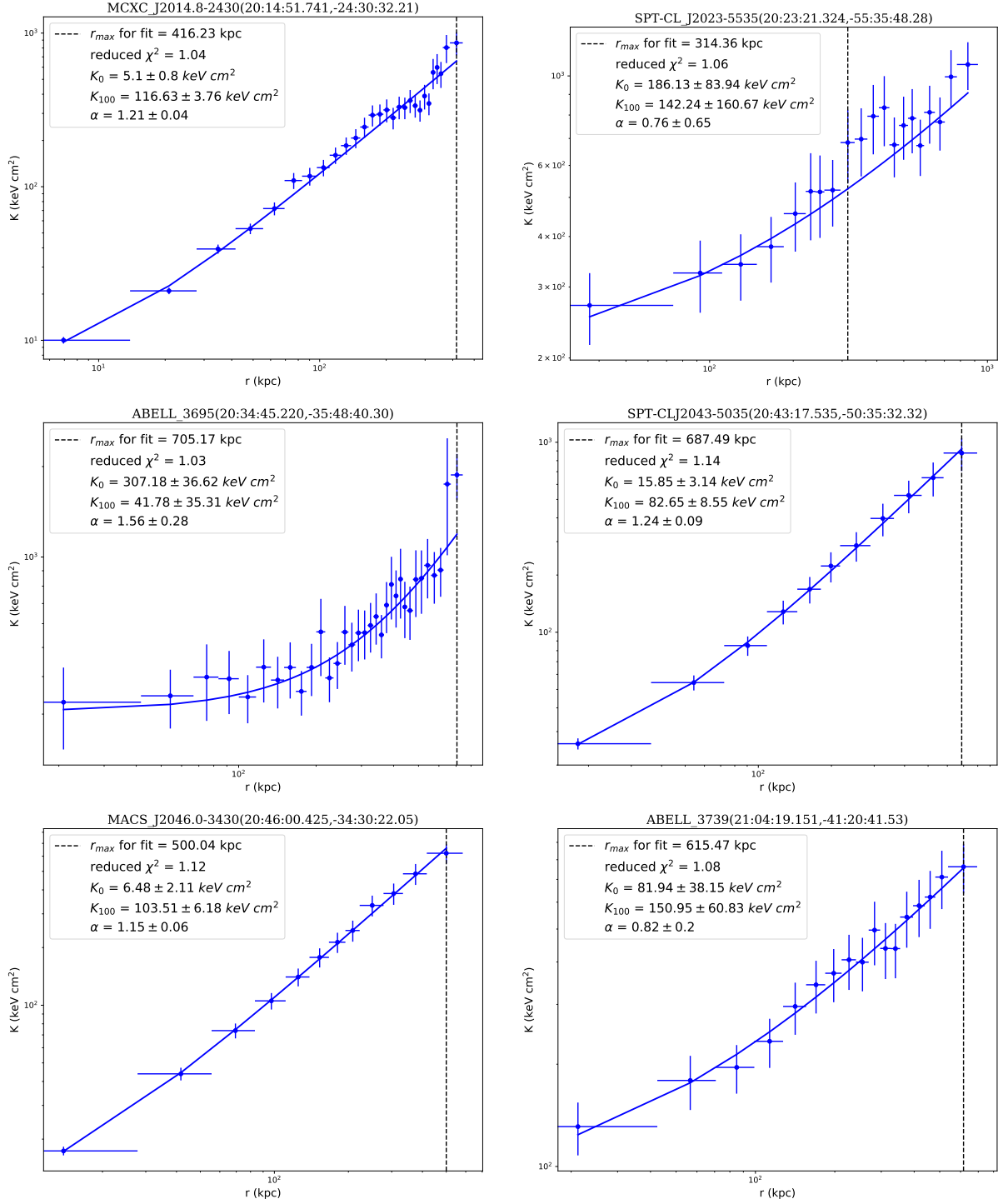


Figure C.1 (cont'd)

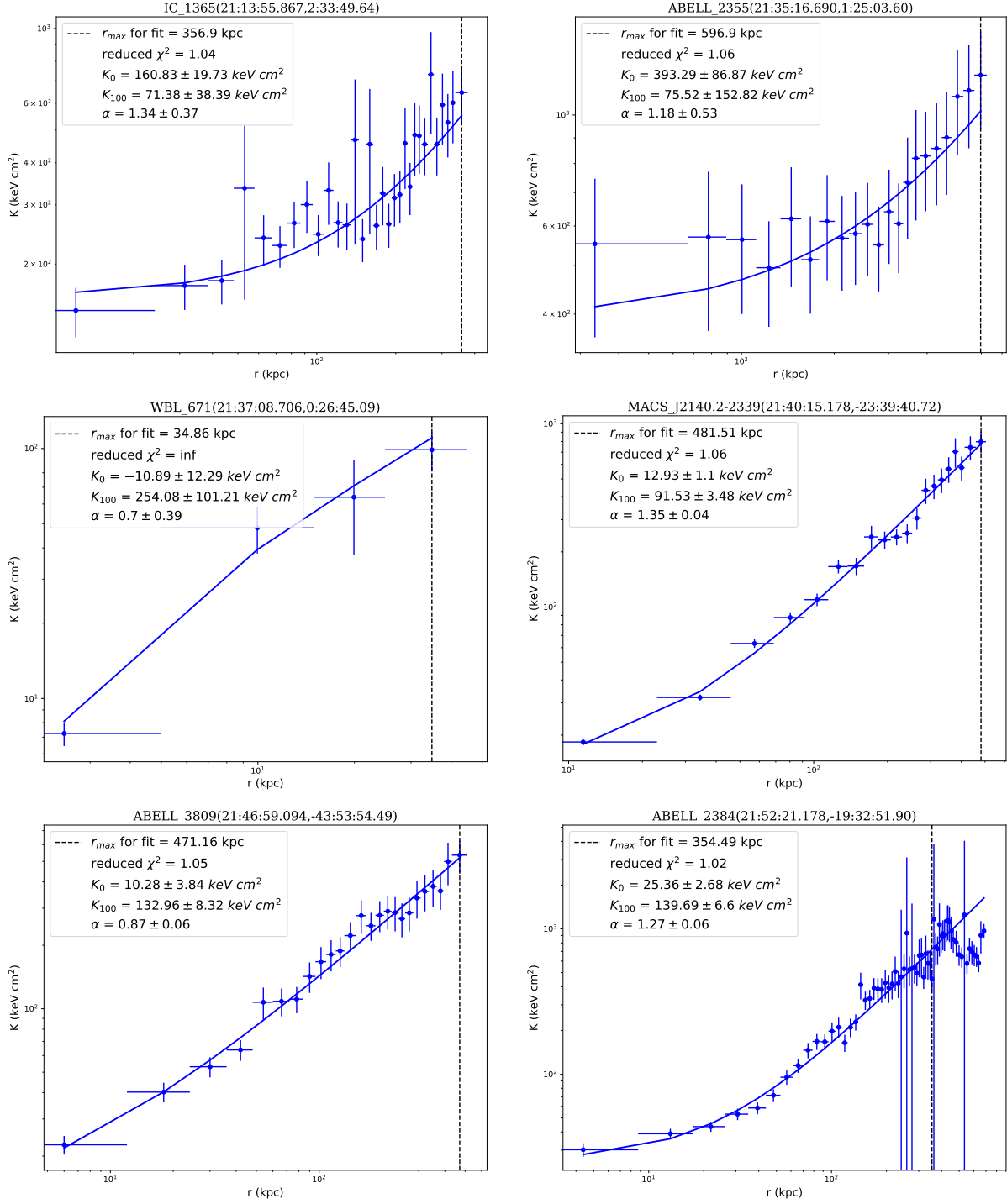


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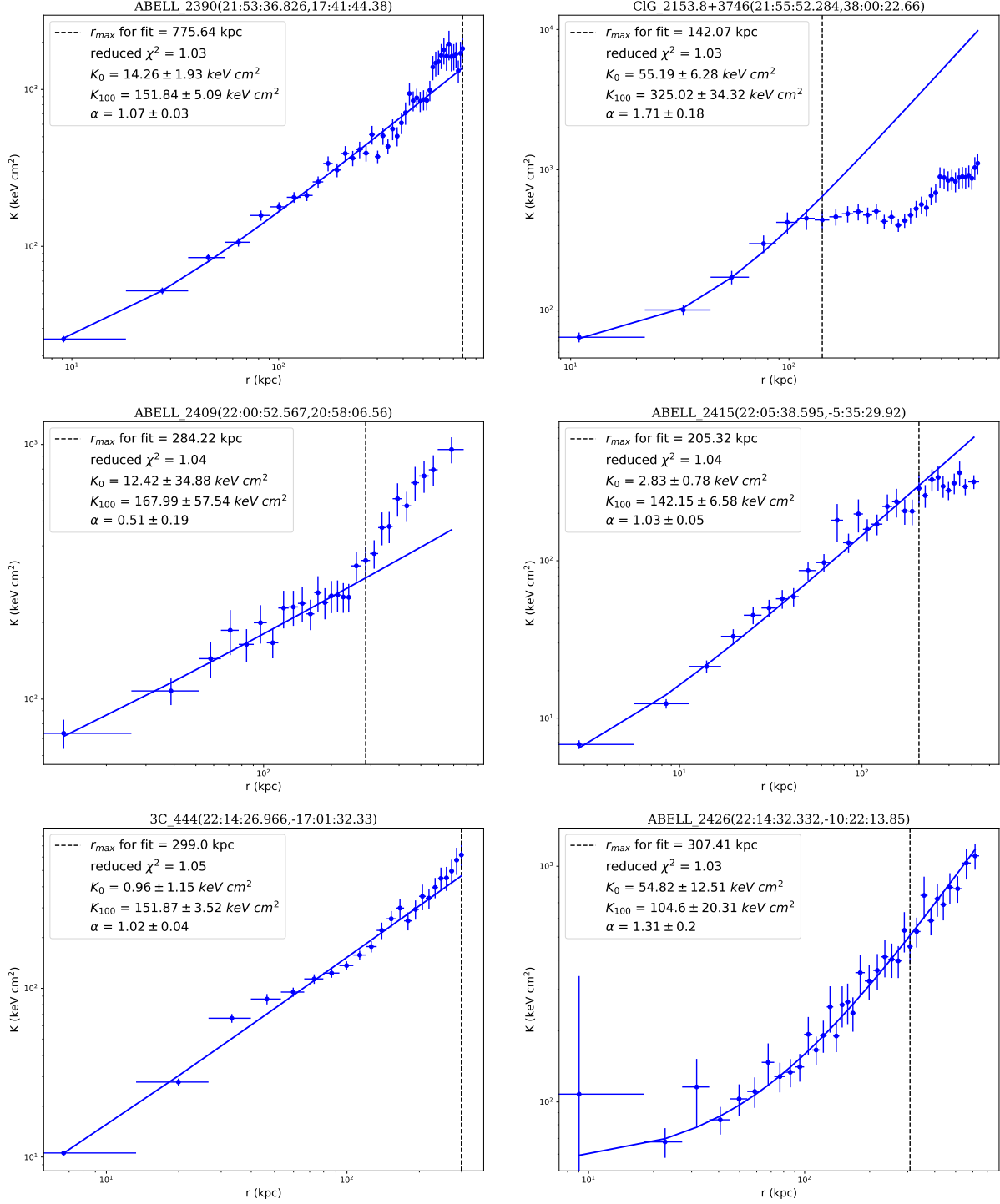


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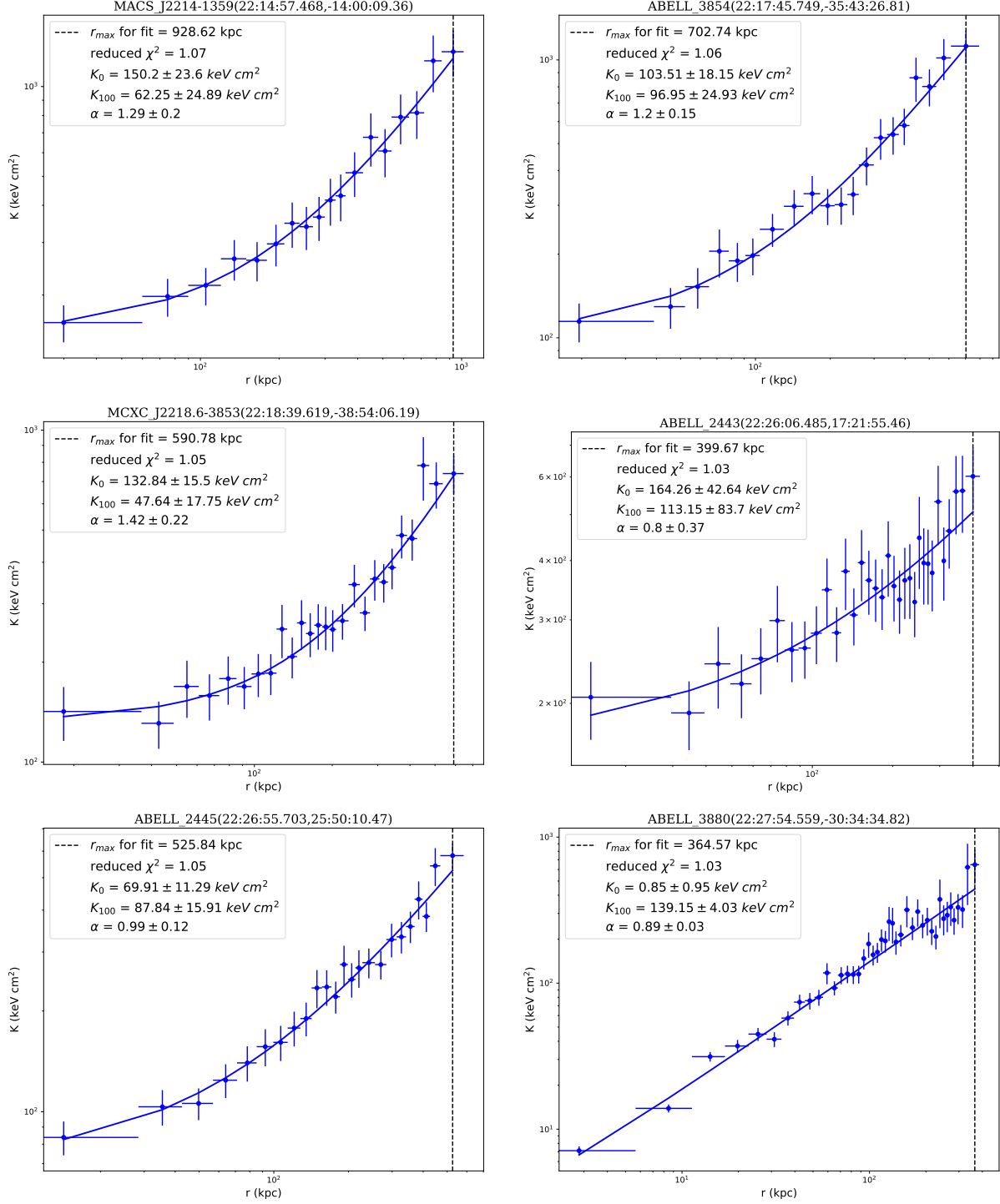


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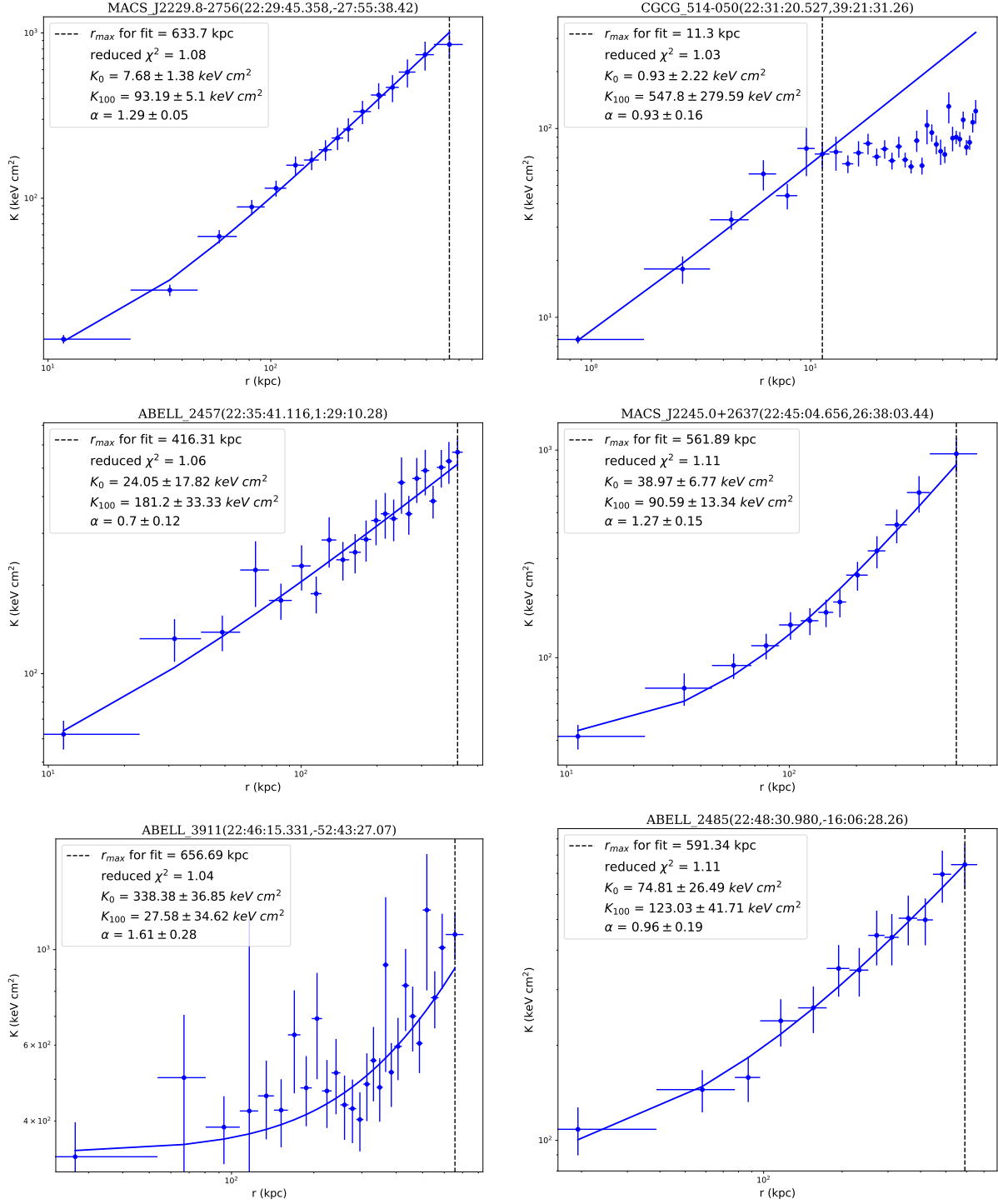


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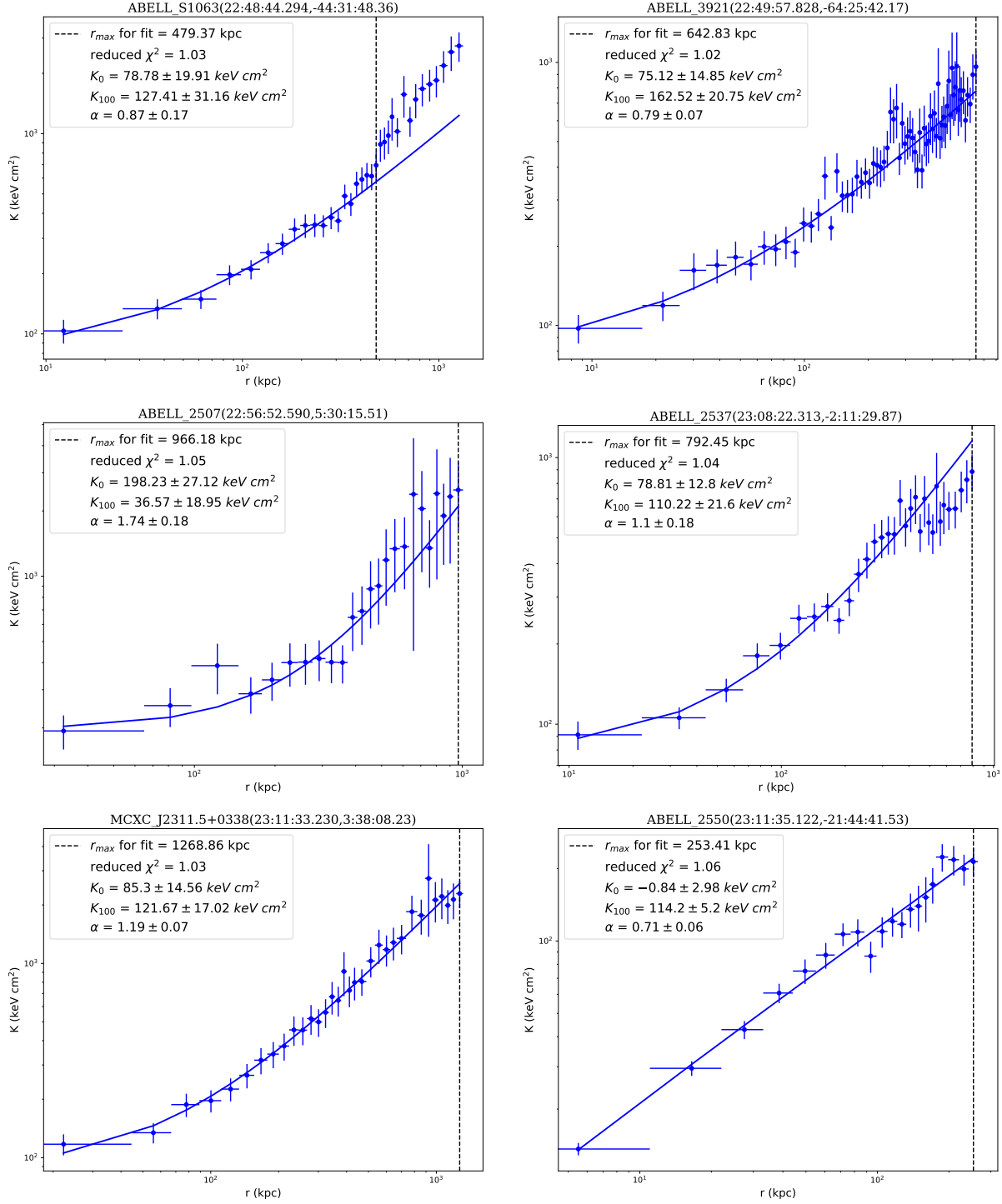
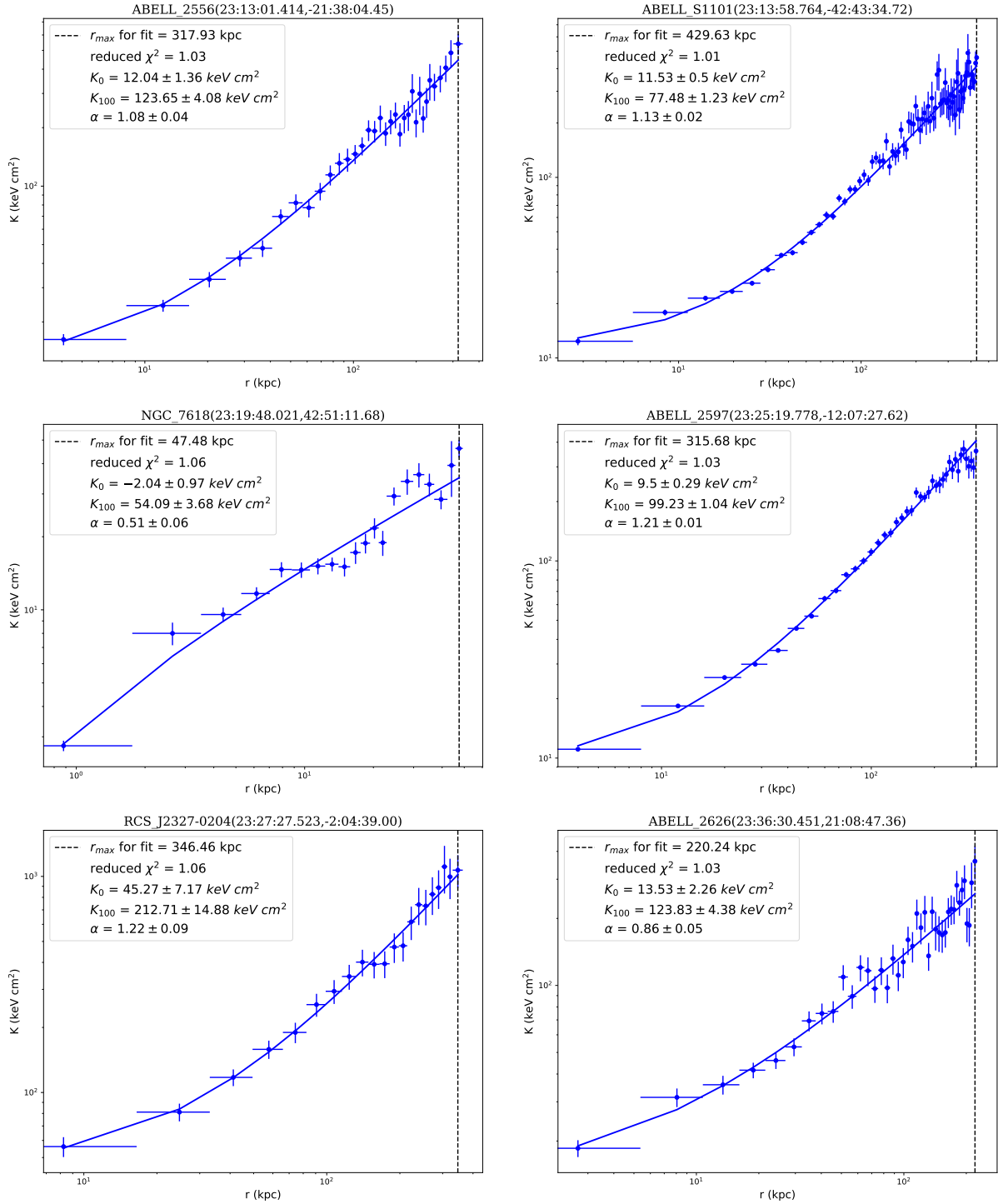


Figure C.1 (cont'd)



APPENDIX D

ACCEPT 2.0 MORPHOLOGICAL PROPERTIES

Table D.1: **Morphological properties for ACCEPT 2.0 clusters** Morphological properties, global temperatures, and global luminosities for ACCEPT 2.0 clusters with deprojected entropy profiles. Column 1: Cluster Name in ACCEPT 2.0; Column 2-3: concentration and error using R_{500} ; Column 4-5: concentration and error with $r = 500$ kpc; Column 6: centroid shift calculated from the data; Column 7: centroid shift from 100 bootstrapped versions of the original data; Column 8: dispersion in the centroid shift from the simulated data; Column 9: power ratio P_3/P_0 calculated from the data; Column 10: power ratio from 100 bootstrapped versions of the original data; Column 11: dispersion in the power ratio from the simulated data; Column 12-13: Best fit global temperatures and errors; Column 14-15: Best fit global luminosity and errors.

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_2717	0.103	0.007	0.219	0.007	0.116	0.116	0.001	21.100	21.200	1.170	5.85	0.39	0.20	0.01
NSCS_J000619+105206	0.163	0.013	0.397	0.025	0.009	0.010	0.001	1.000	1.160	0.521	5.18	0.40	1.85	0.07
ZwCl_0008.8+5215	0.098	0.017	NaN	NaN	0.180	0.179	0.004	0.506	0.692	0.421	4.71	0.30	0.74	0.03
ABELL_2734_NED01	0.146	0.007	0.589	0.034	0.027	0.027	0.001	3.610	3.550	0.448	4.61	0.63	0.07	0.00
MACS_J0011.7-1523	0.192	0.011	0.337	0.016	0.010	0.010	0.001	0.351	0.387	0.202	5.89	0.40	6.23	0.20
ABELL_0013	0.079	0.006	0.199	0.008	0.040	0.039	0.003	6.410	6.650	0.779	4.77	0.26	0.63	0.02
ABELL_2744	0.062	0.004	0.102	0.003	0.044	0.040	0.003	2.200	2.230	0.394	10.48	0.43	14.30	0.18
Cl_0016+16	0.132	0.008	0.151	0.008	0.020	0.018	0.003	0.559	0.735	0.373	9.73	0.74	18.40	0.40
MACS_J0025.4-1222	0.102	0.010	0.197	0.014	0.052	0.047	0.005	1.370	1.550	0.743	7.82	0.60	9.83	0.38
PLCKESZ_G304.84-41.42	0.251	0.018	0.253	0.018	0.047	0.047	0.003	8.750	8.920	1.890	9.80	1.23	8.31	0.30
ABELL_2813	0.160	0.012	0.213	0.014	0.011	0.015	0.007	0.264	0.494	0.402	7.95	0.74	7.71	0.24
ZwCl_0040.8+2404	0.227	0.008	0.383	0.012	0.005	0.005	0.000	0.134	0.161	0.087	4.24	0.23	0.81	0.03

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_0098N	0.511	0.053	NaN	NaN	0.010	0.011	0.002	0.364	0.491	0.387	3.78	0.39	0.32	0.02
ESO_351-G_021	0.307	0.027	NaN	NaN	0.016	0.016	0.002	1.110	1.200	0.790	1.28	0.05	0.03	0.00
ABELL_0141	0.250	0.025	0.242	0.024	0.053	0.042	0.009	16.600	18.000	3.400	6.58	0.77	3.13	0.13
MaxBCG_J016.70077+01.05926	0.252	0.009	0.556	0.021	0.005	0.005	0.000	0.025	0.045	0.037	4.57	0.27	2.12	0.07
CIZA_J0107.7+5408	0.065	0.005	0.127	0.006	0.034	0.025	0.009	1.600	1.660	0.406	7.68	0.37	3.90	0.06
IC_1633	0.284	0.022	NaN	NaN	0.073	0.074	0.002	1.780	1.760	0.465	3.26	0.36	0.08	0.00
ABELL_2895	0.209	0.015	0.209	0.015	0.068	0.065	0.004	2.630	2.500	0.799	8.39	0.86	4.96	0.15
UGC_00842	0.303	0.037	NaN	NaN	0.020	0.021	0.002	2.240	2.160	0.803	1.67	0.08	0.04	0.00
ABELL_0193	0.067	0.006	0.139	0.005	0.023	0.021	0.002	5.390	5.440	0.655	3.64	0.14	0.39	0.01
ABELL_0209	0.140	0.008	0.171	0.009	0.027	0.025	0.002	0.386	0.422	0.279	8.55	0.66	6.57	0.17
Abell_222	0.064	0.012	0.265	0.020	0.080	0.073	0.005	2.640	2.940	0.954	4.14	0.27	1.64	0.07
Abell_223	0.156	0.013	0.444	0.036	0.019	0.019	0.002	1.230	1.410	0.704	5.47	0.50	1.38	0.06
ABELL_0267	0.181	0.013	0.256	0.016	0.031	0.027	0.004	0.202	0.348	0.249	7.57	0.63	4.34	0.11
ABELL_0262	0.066	0.004	0.191	0.002	0.047	0.047	0.000	7.330	7.340	0.212	2.26	0.03	0.02	0.00
NGC_0766	0.038	0.001	0.702	0.250	0.005	0.009	0.003	0.616	3.030	3.050	0.17	0.01	0.04	0.06
MCXC_J0220.9-3829	0.234	0.020	0.492	0.043	0.006	0.007	0.002	0.018	0.171	0.161	4.35	0.37	2.34	0.14
MZ_10451	0.039	0.001	NaN	NaN	0.040	0.037	0.007	0.890	1.960	2.160	0.73	0.04	0.01	0.00
ABELL_0370	0.089	0.006	0.163	0.008	0.048	0.043	0.006	0.580	0.660	0.336	8.75	0.46	7.42	0.12
MACS_J0242.6-2132	0.388	0.025	0.429	0.027	0.002	0.002	0.001	0.285	0.306	0.189	5.75	0.77	4.96	0.31
ABELL_S0295	0.122	0.011	0.189	0.014	0.159	0.158	0.003	0.608	0.788	0.505	7.35	0.71	9.68	0.33

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_0376	0.169	0.019	NaN	NaN	0.038	0.038	0.003	0.651	0.781	0.355	4.53	0.33	0.50	0.01
ABELL_0383	0.216	0.008	0.409	0.012	0.004	0.004	0.001	0.560	0.575	0.125	5.31	0.23	2.02	0.06
NGC_1132	0.223	0.043	0.350	0.029	0.046	0.046	0.003	5.070	5.580	1.790	1.04	0.02	0.01	0.00
ABELL_0402	0.358	0.026	0.313	0.023	0.012	0.012	0.002	1.970	2.040	0.834	8.23	1.24	4.74	0.20
ABELL_0401	0.046	0.002	0.184	0.001	0.015	0.011	0.003	0.883	0.878	0.060	7.57	0.10	4.40	0.02
MCXC_J0301.6+0155	0.222	0.014	0.446	0.026	0.008	0.008	0.001	0.166	0.240	0.171	4.45	0.33	1.67	0.08
MCXC_J0303.7-7752	0.272	0.015	0.193	0.012	0.025	0.025	0.003	0.915	1.160	0.551	9.34	0.89	7.38	0.21
MACS_J0308.9+2645	0.239	0.013	0.215	0.012	0.014	0.013	0.003	0.350	0.464	0.302	9.66	1.05	11.60	0.32
ABELL_3094	0.073	0.013	NaN	NaN	0.053	0.052	0.004	0.153	0.302	0.210	3.11	0.21	0.23	0.01
ABELL_3120	NaN	NaN	NaN	NaN	0.009	0.009	0.002	0.534	0.671	0.446	4.33	0.25	0.10	0.01
UGC_02748	0.386	0.039	NaN	NaN	0.018	0.019	0.002	0.041	0.161	0.141	11.99	1.11	0.01	0.00
ABELL_3126	0.088	0.009	0.360	0.023	0.022	0.020	0.004	0.320	0.362	0.217	4.99	0.31	1.19	0.04
ABELL_3128	0.111	0.015	0.214	0.016	0.168	0.168	0.003	5.490	5.420	1.560	3.14	0.19	0.31	0.01
MCXC_J0331.1-2100	0.334	0.016	0.423	0.020	0.008	0.008	0.001	0.043	0.098	0.077	5.86	0.53	2.63	0.10
3C_089	0.185	0.013	NaN	NaN	0.040	0.039	0.002	0.161	0.269	0.186	4.53	0.33	0.46	0.02
ABELL_3140	NaN	NaN	0.536	0.036	0.010	0.008	0.002	0.466	0.503	0.258	5.46	2.33	0.12	0.02
MCXC_J0338.6+0958	0.060	0.002	0.308	0.002	0.024	0.024	0.000	0.037	0.038	0.008	3.62	0.03	0.48	0.00
MCXC_J0340.8-4542	NaN	NaN	0.014	0.001	0.100	0.099	0.005	0.018	0.304	0.277	2.66	0.30	0.18	0.02
MCXC_J0352.9+1941	0.221	0.008	0.505	0.018	0.004	0.004	0.000	0.034	0.062	0.047	3.14	0.16	0.98	0.04
MACS_J0358.8-2955	0.133	0.008	0.242	0.011	0.069	0.069	0.001	1.740	1.890	0.488	8.75	0.49	14.00	0.37

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_0478	0.135	0.002	0.268	0.002	0.004	0.004	0.000	0.017	0.017	0.008	7.61	0.17	5.86	0.04
MACS_J0416.1-2403	0.158	0.013	0.193	0.015	0.071	0.070	0.005	1.690	1.980	0.938	8.66	0.73	9.94	0.36
MACS_J0417.5-1154	0.219	0.008	0.272	0.009	0.037	0.037	0.001	6.830	7.160	0.714	11.32	0.50	24.50	0.42
MCXC_J0425.8-0833	0.186	0.010	0.473	0.023	0.022	0.022	0.001	5.320	5.430	0.619	3.07	0.16	0.37	0.02
ABELL_S0463	0.439	0.136	NaN	NaN	0.195	0.194	0.003	1.080	1.130	0.421	2.83	0.09	0.09	0.00
MACS_J0429.6-0253	0.383	0.025	0.404	0.027	0.009	0.009	0.001	0.361	0.477	0.262	7.00	0.90	5.72	0.27
ABELL_0496	0.080	0.002	0.218	0.002	0.013	0.013	0.000	0.810	0.805	0.061	4.84	0.13	0.19	0.00
MCXC_J0437.1+0043	0.234	0.013	0.366	0.020	0.009	0.009	0.001	0.229	0.251	0.157	6.83	0.48	5.14	0.14
MCXC_J0439.0+0715	0.181	0.012	0.279	0.015	0.013	0.012	0.002	2.930	3.060	0.648	6.53	0.53	5.46	0.17
MCXC_J0439.0+0520	0.303	0.017	0.401	0.021	0.004	0.004	0.001	0.045	0.090	0.082	5.02	0.48	1.89	0.10
ABELL_3292	0.097	0.012	0.478	0.044	0.022	0.019	0.002	1.090	1.180	0.629	4.20	0.30	1.75	0.08
WEIN_051	0.068	0.005	0.095	0.004	0.050	0.049	0.003	0.596	0.614	0.263	8.10	0.56	0.39	0.01
ESO_552-G_020	0.207	0.018	NaN	NaN	0.031	0.031	0.002	0.039	0.148	0.166	8.23	1.24	0.07	0.00
ABELL_3322	0.217	0.017	0.272	0.019	0.017	0.015	0.002	0.022	0.242	0.221	6.65	0.60	3.87	0.14
MCXC_J0510.7-0801	0.112	0.010	0.166	0.010	0.048	0.047	0.005	0.480	0.661	0.363	7.15	0.43	8.68	0.22
ABELL_S0520	0.099	0.011	0.118	0.012	0.075	0.067	0.008	0.630	1.050	0.777	8.27	0.73	6.96	0.21
ABELL_3343	0.167	0.014	0.301	0.021	0.009	0.008	0.002	0.772	0.926	0.429	6.43	0.53	3.08	0.09
MCXC_J0528.2-2942	0.099	0.014	0.655	0.092	0.027	0.026	0.003	0.262	0.505	0.432	4.49	0.31	1.62	0.09
RBS_0653	0.160	0.008	0.264	0.009	0.071	0.075	0.001	0.583	0.600	0.146	9.15	0.34	8.20	0.12
PLCKESZ_G286.58-31.25	0.104	0.011	0.219	0.018	0.023	0.018	0.006	0.227	0.465	0.423	6.48	0.54	3.85	0.13

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_0545	0.092	0.005	0.151	0.005	0.067	0.067	0.001	5.000	4.960	0.457	7.37	0.31	3.50	0.05
MCXC_J0532.9-3701	0.299	0.016	0.271	0.015	0.011	0.010	0.002	0.112	0.219	0.173	8.64	0.84	7.20	0.21
ESO3060170-A	0.172	0.011	0.417	0.015	0.022	0.022	0.001	0.021	0.049	0.040	2.67	0.10	0.17	0.01
MCXC_J0547.0-3904	0.421	0.033	0.530	0.047	0.010	0.010	0.002	0.243	0.383	0.295	5.19	0.70	0.98	0.06
ABELL_3364	0.130	0.008	0.211	0.010	0.022	0.019	0.005	0.341	0.416	0.237	7.19	0.55	3.19	0.08
ABELL_0548A	0.211	0.037	NaN	NaN	0.031	0.031	0.002	0.571	0.697	0.518	3.07	0.16	0.25	0.01
ABELL_0550	0.080	0.011	0.226	0.014	0.052	0.050	0.004	0.390	0.508	0.304	5.67	0.30	2.15	0.06
MACS_J0553.4-3342	0.126	0.006	0.132	0.006	0.042	0.041	0.002	2.960	3.130	0.656	10.46	0.68	15.00	0.24
ABELL_3378	0.201	0.015	0.377	0.020	0.021	0.021	0.002	0.296	0.366	0.224	4.78	0.27	3.17	0.10
CIZA_J0616.3-2156	0.109	0.012	0.220	0.017	0.034	0.033	0.003	0.100	0.352	0.287	7.24	0.51	3.06	0.07
ABELL_S0579	0.096	0.014	0.337	0.030	0.042	0.040	0.003	1.880	2.420	0.986	4.79	0.35	1.50	0.06
G139.59+24.18	0.281	0.017	0.314	0.018	0.036	0.035	0.002	7.790	7.960	1.520	7.19	0.64	8.23	0.30
ABELL_3391	0.079	0.006	0.318	0.014	0.017	0.017	0.001	1.230	1.300	0.323	8.55	0.66	1.01	0.02
ABELL_3395_SW	0.076	0.008	NaN	NaN	0.029	0.029	0.001	30.100	30.100	1.490	2.26	0.03	0.61	0.02
ABELL_3399	0.156	0.016	0.265	0.020	0.075	0.074	0.004	5.720	5.990	1.510	6.76	0.45	3.38	0.07
PLCKESZ_G167.65+17.64	0.083	0.010	0.145	0.011	0.020	0.020	0.003	0.983	1.160	0.715	6.18	0.41	4.90	0.11
ABELL_S0592	0.203	0.009	0.259	0.011	0.028	0.028	0.001	0.055	0.117	0.087	8.58	0.65	8.34	0.20
ABELL_3404	0.231	0.013	0.247	0.014	0.040	0.039	0.003	1.570	1.850	0.643	7.85	0.64	6.50	0.16
ABELL_0562	0.144	0.017	NaN	NaN	0.025	0.025	0.003	0.133	0.214	0.183	2.94	0.19	0.38	0.02
ABELL_0576	0.072	0.005	0.141	0.004	0.037	0.037	0.002	0.638	0.644	0.203	4.40	0.15	0.21	0.00

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_0578	NaN	NaN	NaN	NaN	0.083	0.083	0.007	0.632	0.740	0.592	2.55	0.23	0.14	0.01
ABELL_0586	0.125	0.008	0.305	0.013	0.005	0.006	0.002	0.733	0.792	0.304	6.24	0.40	3.38	0.09
ZwCl_0735.7+7421	0.114	0.004	0.328	0.004	0.005	0.005	0.000	0.450	0.449	0.070	6.47	0.11	3.97	0.03
MACS_J0744.9+3927	0.285	0.017	0.342	0.020	0.024	0.024	0.002	2.700	2.810	0.775	7.80	0.63	15.00	0.65
PKS_0745-19	0.142	0.001	0.333	0.002	0.009	0.009	0.000	0.438	0.435	0.021	8.55	0.73	5.82	0.03
ABELL_0598	0.294	0.019	0.489	0.034	0.029	0.029	0.001	0.482	0.600	0.274	5.05	0.54	1.69	0.08
ABELL_0611	0.264	0.011	0.324	0.013	0.006	0.006	0.001	0.291	0.308	0.181	8.55	0.73	5.00	0.13
SDSS-C4_3062	NaN	NaN	NaN	NaN	0.026	0.026	0.005	0.007	0.525	0.615	3.76	0.54	0.19	0.04
ABELL_0644	0.065	0.003	0.255	0.004	0.027	0.026	0.001	0.143	0.143	0.037	4.76	0.75	2.86	0.03
MCXC_J0819.6+6336	0.197	0.019	0.462	0.041	0.025	0.025	0.002	0.816	1.060	0.630	3.50	0.32	0.74	0.04
UGCl_120	NaN	NaN	NaN	NaN	0.023	0.023	0.004	1.810	2.410	1.440	1.90	0.12	0.05	0.01
ZwCl_0823.2+0425	0.188	0.022	0.943	0.507	0.006	0.007	0.002	1.140	1.400	0.787	4.66	0.51	1.72	0.11
2MFGC_06756	0.193	0.006	0.439	0.011	0.005	0.005	0.000	0.055	0.062	0.033	5.05	0.23	2.71	0.06
ABELL_3411	0.065	0.008	0.241	0.015	0.146	0.143	0.005	0.105	0.197	0.164	6.25	0.31	2.87	0.05
NSC_J084254+292723	0.300	0.014	0.470	0.022	0.005	0.005	0.001	0.114	0.181	0.115	5.89	0.44	1.51	0.05
ABELL_0697	0.293	0.013	0.155	0.009	0.010	0.009	0.003	0.361	0.534	0.387	11.99	1.11	13.20	0.32
ZwCl_0857.9+2107	0.269	0.010	0.531	0.020	0.002	0.002	0.000	0.093	0.111	0.064	4.39	0.26	2.15	0.08
SDSS_+137.3+11.0+0.18	0.097	0.008	0.298	0.013	0.312	0.310	0.003	141.000	140.000	5.900	5.38	0.31	2.70	0.08
HCG_037	0.007	0.001	0.585	0.214	0.012	0.015	0.005	0.296	8.070	9.790	0.96	0.24	0.00	0.00
2MASSi_J0913454+405628	0.372	0.015	0.492	0.021	0.002	0.002	0.000	0.010	0.038	0.034	6.42	0.59	4.90	0.20

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
	10^7										(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_0773	0.119	0.007	0.195	0.009	0.024	0.020	0.004	0.326	0.400	0.240	8.10	0.56	6.11	0.14
Hydra_A	0.067	0.002	0.272	0.002	0.014	0.014	0.000	4.540	4.540	0.063	4.08	0.03	1.12	0.00
NSC_J092017+303027	0.142	0.017	0.254	0.026	0.598	0.585	0.020	5.610	5.690	1.810	6.35	0.66	3.13	0.13
ABELL_0795	0.179	0.008	0.453	0.018	0.013	0.013	0.001	0.887	0.927	0.275	5.00	0.30	1.85	0.05
UGC_05088_GROUP	0.144	0.034	NaN	NaN	0.047	0.047	0.004	0.562	1.070	0.979	2.33	0.23	0.01	0.00
WHL_J093820.9+520243	0.158	0.015	0.316	0.023	0.098	0.098	0.002	1.360	1.400	0.673	6.42	0.50	5.66	0.19
ABELL_0868	0.060	0.013	0.218	0.016	0.043	0.035	0.006	0.524	0.640	0.403	4.41	0.22	2.58	0.09
GALEX_J094712.4+762313	0.351	0.011	0.424	0.013	0.001	0.002	0.000	0.010	0.020	0.015	8.46	0.57	8.35	0.19
ZwCl_0949.6+5207	0.135	0.005	0.467	0.013	0.029	0.029	0.001	7.640	7.700	0.440	5.64	0.21	2.26	0.05
PLCKESZ_G264.41+19.48	0.204	0.016	0.263	0.020	0.007	0.008	0.001	0.176	0.354	0.347	7.50	0.70	3.76	0.13
HCG_042	0.196	0.020	0.480	0.040	0.017	0.016	0.002	2.810	2.840	0.888	0.82	0.05	0.00	0.00
MCXC_J1000.5+4409	0.206	0.020	0.551	0.052	0.067	0.067	0.003	2.370	2.670	0.948	3.28	0.27	1.03	0.07
ZwCl_1006.1+1201	0.097	0.008	0.301	0.015	0.086	0.079	0.005	0.828	0.914	0.328	5.86	0.33	2.85	0.07
MCXC_J1010.5-1239	0.127	0.012	0.236	0.015	0.098	0.098	0.002	0.335	0.524	0.347	6.48	0.40	4.22	0.09
ABELL_0963	0.130	0.005	0.292	0.008	0.015	0.015	0.001	1.180	1.220	0.237	6.49	0.32	4.63	0.08
MCXC_J1022.0+3830	NaN	NaN	NaN	NaN	0.014	0.013	0.002	4.750	4.580	1.300	2.66	0.24	0.05	0.00
ABELL_0980	0.126	0.011	0.240	0.016	0.019	0.019	0.002	0.216	0.365	0.316	6.54	0.53	3.07	0.09
BLOX_J1023.6+0411.1	0.231	0.005	0.338	0.006	0.010	0.011	0.000	0.121	0.122	0.041	8.75	0.37	9.03	0.14
ABELL_3444	0.247	0.007	0.379	0.009	0.006	0.006	0.001	0.223	0.245	0.090	7.92	0.45	7.84	0.13
ABELL_1033	0.067	0.005	0.362	0.011	0.056	0.055	0.001	7.930	8.040	0.738	5.61	0.18	1.47	0.03

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
	10^7										(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_1068	0.243	0.007	0.409	0.010	0.003	0.003	0.000	0.011	0.019	0.014	4.96	0.25	1.87	0.04
NGC_3402_GROUP	0.122	0.012	0.411	0.013	0.009	0.010	0.001	0.040	0.075	0.056	0.82	0.01	0.01	0.00
BLOX_J1056.9-0337.3	0.375	0.045	0.402	0.048	0.280	0.268	0.007	11.000	11.600	2.890	7.49	1.08	8.11	0.43
MACS_J1108.9+0906	0.164	0.019	0.255	0.027	0.032	0.028	0.004	0.220	0.574	0.567	6.86	0.74	6.09	0.29
NGC_3551	0.324	0.052	NaN	NaN	0.017	0.018	0.004	0.286	0.424	0.361	1.62	0.06	0.02	0.00
1RXS_J111039.6+284316	NaN	NaN	NaN	NaN	0.071	0.069	0.007	4.840	5.060	1.580	1.29	0.09	0.01	0.00
ABELL_1190	0.065	0.015	0.816	0.198	0.017	0.017	0.004	0.690	0.951	0.573	3.60	0.19	0.58	0.03
ABELL_1201	0.111	0.008	0.298	0.010	0.028	0.028	0.001	5.040	4.900	0.622	5.69	0.21	2.54	0.05
ABELL_1204	0.255	0.010	0.422	0.015	0.003	0.003	0.000	0.159	0.171	0.082	4.33	0.28	1.61	0.07
MACS_J1115.8+0129	0.349	0.014	0.345	0.014	0.008	0.008	0.001	0.042	0.081	0.060	9.18	0.73	8.71	0.25
HCG_051	0.312	0.050	NaN	NaN	0.050	0.050	0.004	1.510	1.760	0.860	1.32	0.03	0.01	0.00
ABELL_1240	NaN	NaN	NaN	NaN	0.059	0.060	0.010	10.200	9.950	3.170	4.16	0.42	0.48	0.03
MCXC_J1130.0+3637	0.445	0.048	NaN	NaN	0.016	0.016	0.002	1.760	1.950	0.634	1.74	0.08	0.07	0.00
ABELL_1285	0.061	0.009	0.212	0.010	0.035	0.030	0.005	0.723	0.846	0.290	5.52	0.21	2.25	0.04
ABELL_1300	0.544	0.037	0.204	0.013	0.066	0.064	0.002	2.060	2.160	0.781	11.26	1.17	10.50	0.26
WHL_J114224.8+583205	0.079	0.008	0.096	0.008	0.080	0.080	0.006	1.800	2.160	1.020	8.86	0.68	8.37	0.20
ABELL_1361	0.382	0.023	0.396	0.030	0.005	0.006	0.001	0.252	0.363	0.221	17.00	10.00	0.68	0.16
SDSS-C4-DR3_3018	0.019	0.002	0.020	0.001	0.036	0.037	0.005	1.470	1.540	0.523	2.64	0.17	0.09	0.00
ABELL_1413	0.092	0.003	0.239	0.003	0.007	0.007	0.000	0.557	0.565	0.077	7.41	0.17	4.23	0.04
ABELL_1423	0.172	0.012	0.307	0.016	0.015	0.015	0.002	0.445	0.496	0.190	4.33	0.26	0.32	0.01

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
SDSS-C4-DR3_3144	0.357	0.023	0.911	0.236	0.003	0.004	0.001	0.056	0.252	0.192	1.77	0.08	0.07	0.01
ABELL_1446	0.055	0.008	0.692	0.063	0.034	0.034	0.001	0.566	0.603	0.204	3.53	0.12	0.69	0.02
MACS_J1206.2-0847	0.273	0.014	0.228	0.013	0.028	0.028	0.002	2.640	2.660	0.871	11.37	1.42	19.30	0.56
MCXC_J1215.4-3900	0.073	0.014	0.305	0.033	0.033	0.028	0.007	2.610	2.960	1.320	5.50	0.36	1.60	0.04
NGC_4325_GROUP	0.053	0.010	0.396	0.012	0.010	0.010	0.001	0.355	0.371	0.129	0.98	0.02	0.02	0.00
NSCS_J122648+215157	0.127	0.013	NaN	NaN	0.214	0.212	0.004	9.120	9.760	3.190	4.78	0.47	1.62	0.08
MCXC_J1234.2+0947	0.172	0.038	NaN	NaN	0.101	0.093	0.020	4.680	4.570	2.490	4.55	0.48	1.74	0.11
MESSIER_089	0.040	0.000	0.503	0.019	0.011	0.011	0.001	1.050	1.080	0.264	3.00	10.00	0.00	0.00
ABELL_1569	0.067	0.017	NaN	NaN	0.020	0.021	0.002	1.810	1.770	0.594	6.03	1.00	0.25	0.02
ABELL_1576	0.207	0.013	0.283	0.017	0.016	0.015	0.002	1.530	1.610	0.551	8.00	0.77	3.02	0.11
Centaurus_Cluster	0.109	0.002	0.204	0.001	0.023	0.023	0.000	0.336	0.339	0.017	3.00	10.00	0.00	0.00
NGC_4759_GROUP	0.082	0.005	0.484	0.010	0.019	0.019	0.001	0.003	0.009	0.009	1.38	0.03	0.00	0.00
ABELL_3528B	0.344	0.022	NaN	NaN	0.016	0.016	0.002	1.790	1.880	0.562	7.57	0.63	0.57	0.02
NGC_4782-3	0.450	0.054	NaN	NaN	0.039	0.039	0.002	5.200	5.560	1.710	1.05	0.20	0.00	0.00
ABELL_1644	0.111	0.007	0.157	0.003	0.032	0.032	0.001	5.210	5.190	0.339	4.12	0.87	0.80	0.01
ABELL_3532	0.100	0.009	0.889	0.246	0.045	0.044	0.003	3.160	3.190	0.783	6.08	0.77	0.86	0.03
WHL_J125933.4+600409	0.080	0.014	0.160	0.017	0.094	0.089	0.006	3.430	3.780	1.430	6.89	0.52	4.09	0.15
ABELL_1664	0.166	0.006	0.404	0.011	0.009	0.009	0.001	0.568	0.585	0.126	5.02	0.21	1.91	0.04
ABELL_1668	0.283	0.016	0.367	0.020	0.033	0.033	0.002	1.190	1.310	0.507	3.33	0.29	0.30	0.02
ABELL_1689	0.125	0.002	0.294	0.003	0.004	0.004	0.000	0.030	0.033	0.015	9.95	0.22	8.92	0.11

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
WHL_J131505.2+514902	0.155	0.011	0.196	0.013	0.010	0.011	0.004	0.948	1.060	0.529	8.76	0.77	7.12	0.14
NGC_5044	0.054	0.009	0.206	0.002	0.024	0.024	0.000	0.120	0.122	0.023	1.32	0.03	0.00	0.00
NGC_5098_GROUP	0.187	0.026	NaN	NaN	0.011	0.011	0.001	0.277	0.328	0.213	1.07	0.02	0.02	0.00
NGC_5129	0.321	0.032	NaN	NaN	0.012	0.012	0.003	0.398	0.681	0.496	0.75	0.02	0.01	0.00
ABELL_1736	0.059	0.018	0.195	0.010	0.039	0.040	0.008	6.520	6.660	0.812	2.82	0.67	0.62	0.01
ABELL_3558	0.089	0.004	0.122	0.003	0.036	0.036	0.001	8.860	8.950	0.596	7.44	0.30	0.60	0.01
SSGC_081	0.099	0.012	0.296	0.015	0.034	0.034	0.002	9.040	9.130	1.090	8.32	0.22	0.41	0.01
a1750ss	NaN	NaN	0.013	0.001	0.049	0.046	0.007	0.659	0.894	0.761	2.40	0.20	0.13	0.02
ABELL_1750C	0.117	0.013	NaN	NaN	0.028	0.028	0.002	2.760	2.950	0.707	4.43	0.25	0.44	0.02
ABELL_1750N	0.182	0.027	NaN	NaN	0.031	0.032	0.003	0.006	0.158	0.155	3.59	0.23	0.44	0.02
SC_1329-313	0.068	0.016	NaN	NaN	0.050	0.047	0.007	1.070	1.160	0.601	6.50	0.43	0.26	0.01
2MASX_J13312961+1107566	0.219	0.077	NaN	NaN	0.013	0.016	0.006	14.000	18.400	11.400	0.64	0.04	0.01	0.00
ABELL_3560	0.084	0.010	0.300	0.021	0.029	0.028	0.003	8.370	8.410	1.330	8.75	0.46	0.18	0.00
ABELL_3562	0.074	0.007	0.251	0.009	0.015	0.015	0.002	6.150	6.190	0.567	4.70	0.26	0.81	0.02
ABELL_1763	0.094	0.008	0.158	0.009	0.040	0.040	0.003	0.701	0.741	0.406	7.67	0.59	7.19	0.19
ABELL_1775	0.045	0.007	0.235	0.004	0.063	0.063	0.001	1.450	1.460	0.140	5.47	1.02	0.90	0.01
ABELL_3571	0.069	0.002	0.096	0.001	0.083	0.083	0.001	3.550	3.530	0.172	7.76	0.27	0.55	0.01
LCDCS_0829	0.349	0.005	0.376	0.005	0.017	0.016	0.000	0.434	0.434	0.060	13.75	0.53	23.90	0.29
ABELL_1795	0.078	0.002	0.248	0.002	0.009	0.010	0.000	1.120	1.120	0.047	6.13	0.10	1.88	0.01
NSCS_J135021+094042	0.321	0.015	0.508	0.026	0.005	0.005	0.000	0.065	0.108	0.093	6.24	0.40	0.39	0.02

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
MACS_J1359.2-1929	0.802	0.155	0.685	0.091	0.013	0.013	0.002	0.362	0.490	0.419	6.56	1.20	2.74	0.21
ABELL_1831	0.114	0.018	NaN	NaN	0.012	0.014	0.003	0.655	0.689	0.324	3.54	0.16	0.52	0.02
WHL_J135949.5+623047	0.255	0.013	0.358	0.017	0.011	0.011	0.001	0.319	0.452	0.289	6.86	0.59	4.20	0.13
ABELL_1835	0.312	0.006	0.352	0.007	0.010	0.010	0.000	0.118	0.132	0.043	10.33	0.62	12.80	0.18
A1882a	NaN	NaN	NaN	NaN	0.022	0.021	0.003	0.226	0.468	0.411	3.56	0.23	0.31	0.02
WHL_J141623.8+444528	0.200	0.032	0.623	0.112	0.042	0.040	0.005	83.100	80.800	18.400	3.65	0.28	1.61	0.14
GMBCG_J215.94948+24.07846	0.307	0.010	0.445	0.014	0.002	0.002	0.000	0.049	0.080	0.064	7.26	0.36	6.71	0.13
ABELL_1914	0.115	0.005	0.206	0.006	0.063	0.061	0.003	0.775	0.784	0.190	8.74	0.46	9.73	0.15
WHL_J142716.1+440730	0.683	0.069	0.425	0.029	0.010	0.010	0.001	1.190	1.240	0.519	10.40	1.42	7.62	0.31
MACS_J1427.6-2521	0.256	0.020	0.557	0.049	0.004	0.005	0.001	0.541	0.678	0.368	4.69	0.34	2.07	0.10
ABELL_1930	0.225	0.012	0.615	0.041	0.003	0.004	0.001	0.168	0.202	0.123	4.51	0.26	0.94	0.04
ABELL_1942_AND_CLUMP	0.120	0.014	0.335	0.025	0.027	0.026	0.003	2.500	2.610	0.998	5.39	0.36	1.42	0.05
WBL_518	0.083	0.009	NaN	NaN	0.059	0.051	0.006	0.888	0.911	0.334	8.14	0.36	0.14	0.00
NSCS_J144726+082824	NaN	NaN	0.526	0.027	0.002	0.002	0.001	0.055	0.095	0.072	19.22	7.76	1.71	0.23
ABELL_1991	0.150	0.006	0.338	0.007	0.011	0.011	0.000	0.307	0.312	0.085	2.77	0.06	0.32	0.01
NSCS_J145715+222009	0.176	0.005	0.419	0.009	0.006	0.006	0.000	0.071	0.081	0.031	5.27	0.22	4.00	0.08
ABELL_S0780	0.251	0.009	0.409	0.012	0.015	0.015	0.001	0.321	0.362	0.130	7.33	0.29	5.56	0.08
ABELL_2009	0.227	0.009	0.327	0.011	0.005	0.005	0.001	0.207	0.266	0.108	6.70	0.44	3.58	0.10
WHL_J150407.5-024816	0.265	0.004	0.434	0.008	0.004	0.004	0.000	0.033	0.035	0.013	8.96	0.38	9.96	0.14
MCXC_J1514.9-1523	0.056	0.007	0.087	0.005	0.062	0.052	0.010	0.596	0.816	0.467	8.90	0.48	5.82	0.09

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
MKW_03s	0.064	0.003	0.307	0.003	0.011	0.011	0.000	1.200	1.200	0.076	5.67	0.30	0.63	0.01
ABELL_2069	0.063	0.007	0.152	0.006	0.201	0.198	0.012	6.500	6.520	1.030	4.46	0.58	1.79	0.03
MCXC_J1524.2-3154	0.175	0.005	0.417	0.008	0.003	0.003	0.000	0.032	0.044	0.027	4.47	0.12	1.34	0.02
MACS_J1532.8+3021	0.240	0.005	0.430	0.009	0.002	0.002	0.000	0.012	0.019	0.017	7.49	0.41	6.98	0.12
ABELL_2092	0.218	0.074	NaN	NaN	0.059	0.054	0.010	4.740	5.090	2.790	2.86	0.44	0.12	0.02
ABELL_2107	0.135	0.007	0.649	0.032	0.014	0.014	0.001	0.492	0.508	0.124	3.18	0.76	0.42	0.01
ABELL_2111	0.128	0.011	0.216	0.015	0.028	0.025	0.004	0.286	0.438	0.348	7.70	0.57	3.90	0.08
ABELL_2104	0.065	0.004	0.191	0.005	0.012	0.011	0.001	4.520	4.490	0.414	7.17	0.25	3.79	0.05
ABELL_2125	0.092	0.027	NaN	NaN	0.083	0.080	0.006	3.140	3.920	1.500	3.16	0.21	0.61	0.03
ABELL_2124	0.181	0.010	0.245	0.012	0.021	0.020	0.002	1.880	1.920	0.479	5.09	0.40	0.32	0.01
MCXC_J1558.3-1410	0.104	0.004	0.353	0.006	0.007	0.007	0.000	1.250	1.250	0.143	5.18	0.12	1.98	0.02
ABELL_2147	0.052	0.009	0.153	0.006	0.044	0.043	0.003	2.380	2.410	0.430	7.16	1.08	0.52	0.01
ABELL_2151	0.144	0.013	0.483	0.026	0.011	0.011	0.001	1.990	2.090	0.455	4.48	1.15	0.18	0.02
AWM_4	0.094	0.005	0.345	0.008	0.005	0.005	0.001	0.416	0.436	0.110	2.74	0.09	0.06	0.00
MACS_J1621.3+3810	0.309	0.013	0.447	0.020	0.008	0.008	0.001	0.144	0.187	0.109	7.23	0.58	5.35	0.20
ABELL_2187	0.300	0.024	0.357	0.029	0.043	0.043	0.003	0.358	0.525	0.337	6.71	0.76	2.43	0.10
ABELL_2204	0.283	0.003	0.340	0.003	0.006	0.006	0.000	0.020	0.023	0.011	4.42	1.07	6.43	0.05
ABELL_2219	0.059	0.003	0.140	0.002	0.037	0.036	0.001	0.526	0.515	0.074	11.27	0.20	15.70	0.10
Hercules_A	0.184	0.008	0.392	0.008	0.003	0.004	0.000	0.052	0.061	0.031	4.26	0.10	1.81	0.03
NGC_6269	0.293	0.025	NaN	NaN	0.014	0.015	0.001	3.750	3.850	0.886	2.94	0.19	0.08	0.00

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
ABELL_2256	0.046	0.004	0.073	0.002	0.140	0.136	0.006	4.260	4.220	0.360	3.56	0.23	3.40	0.03
SDSS-C4_3072	0.189	0.005	0.333	0.006	0.004	0.004	0.000	0.129	0.139	0.040	7.11	0.27	4.43	0.08
MACS_J1720.2+3536	0.310	0.014	0.353	0.016	0.016	0.015	0.001	0.070	0.141	0.110	7.87	0.65	7.30	0.22
ABELL_2261	0.223	0.009	0.283	0.010	0.009	0.009	0.001	0.268	0.320	0.137	8.10	0.59	7.49	0.16
ABELL_2294	0.181	0.014	0.216	0.015	0.028	0.026	0.003	0.296	0.486	0.378	7.30	0.74	4.67	0.17
Abell_2276	0.233	0.021	NaN	NaN	0.006	0.007	0.002	1.340	1.640	0.638	2.83	0.28	0.48	0.04
ZwCl_1742.1+3306	0.098	0.004	0.319	0.005	0.010	0.010	0.000	0.567	0.589	0.083	5.00	0.30	1.08	0.02
NSC_J174715+451155	0.094	0.021	0.203	0.023	0.055	0.048	0.008	3.990	4.640	1.670	4.73	0.38	1.51	0.07
MCXC_J1750.2+3504	0.261	0.017	0.436	0.025	0.004	0.004	0.001	0.072	0.141	0.123	4.64	0.36	1.67	0.08
NGC_6482	0.279	0.021	0.494	0.034	0.006	0.006	0.001	0.187	0.285	0.258	0.48	0.04	0.00	0.00
CIZA_J1804.4+1002	0.141	0.011	0.216	0.012	0.079	0.079	0.002	0.082	0.168	0.130	7.10	0.59	5.23	0.15
ABELL_2302	0.090	0.020	NaN	NaN	0.092	0.085	0.008	0.703	0.953	0.638	4.82	0.38	1.32	0.05
MACS_J1829.0+6913	0.188	0.014	NaN	NaN	0.004	0.005	0.001	1.080	1.210	0.492	4.06	0.30	0.83	0.04
MCXC_J1852.1+5711	0.243	0.018	NaN	NaN	0.005	0.005	0.001	1.660	1.880	0.685	4.00	0.28	0.45	0.02
MCXC_J1853.9+6822	0.098	0.015	0.807	0.147	0.046	0.043	0.002	0.101	0.175	0.169	4.13	0.22	1.01	0.02
PLCKESZ_G337.09-25.97	0.183	0.011	0.288	0.015	0.032	0.031	0.002	0.510	0.631	0.345	7.66	0.54	6.85	0.14
MACS_J1931.8-2635	0.270	0.007	0.393	0.009	0.002	0.002	0.000	0.031	0.039	0.028	7.87	0.33	9.42	0.11
CIZA_J1938.3+5409	0.195	0.013	0.271	0.016	0.022	0.021	0.002	3.090	3.030	0.766	7.20	0.73	9.11	0.36
MCXC_J1947.3-7623	0.300	0.017	0.309	0.017	0.027	0.027	0.001	1.320	1.450	0.601	7.34	0.63	5.67	0.18
ABELL_3653	NaN	NaN	NaN	NaN	0.053	0.053	0.008	1.550	2.150	1.190	4.78	0.34	0.52	0.01

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	P_{data}	P_{sim}	δP_{sim}	kT	δkT	L_X	δL_X
	10^7										(keV)		$(10^{44} \text{ erg s}^{-1})$	
MCXC_J2003.5-2323	0.237	0.017	0.233	0.016	0.133	0.133	0.005	0.975	1.150	0.749	9.18	0.80	8.46	0.20
MCXC_J2011.3-5725	0.199	0.019	0.459	0.041	0.006	0.006	0.001	1.420	1.570	0.717	3.63	0.37	2.08	0.15
MCXC_J2014.8-2430	0.332	0.009	0.371	0.009	0.005	0.005	0.000	0.999	1.000	0.153	7.15	0.39	4.26	0.10
SPT-CL_J2023-5535	0.137	0.013	0.145	0.013	0.095	0.092	0.005	2.440	2.620	1.040	8.39	0.76	6.12	0.18
ABELL_3695	0.079	0.012	0.304	0.030	0.057	0.057	0.003	0.263	0.442	0.308	6.46	0.43	1.90	0.04
SPT-CLJ2043-5035	0.432	0.037	0.583	0.060	0.005	0.006	0.001	0.283	0.498	0.452	5.44	0.68	6.60	0.49
MACS_J2046.0-3430	0.343	0.024	0.478	0.036	0.002	0.002	0.000	0.048	0.133	0.122	5.61	0.45	4.33	0.22
ABELL_3739	0.129	0.014	0.274	0.023	0.015	0.014	0.002	0.021	0.272	0.271	6.08	0.53	2.90	0.12
IC_1365	0.083	0.010	NaN	NaN	0.356	0.356	0.003	11.400	11.700	1.370	4.84	0.13	0.56	0.02
ABELL_2355	0.229	0.030	0.441	0.065	0.031	0.034	0.005	0.888	0.966	0.521	7.24	0.69	1.48	0.05
WBL_671	0.037	0.001	NaN	NaN	0.016	0.023	0.008	0.017	4.940	4.990	1.00	0.12	0.01	0.00
MACS_J2140.2-2339	0.300	0.009	0.423	0.011	0.001	0.001	0.000	0.137	0.139	0.059	5.92	0.34	4.10	0.11
ABELL_3809	0.142	0.018	NaN	NaN	0.013	0.013	0.002	0.188	0.291	0.216	2.89	0.12	0.41	0.01
ABELL_2384	0.149	0.008	0.339	0.011	0.056	0.056	0.002	0.955	0.969	0.262	4.77	0.26	1.11	0.02
ABELL_2390	0.165	0.003	0.258	0.003	0.013	0.013	0.000	0.357	0.354	0.066	11.16	0.31	14.50	0.10
CIG_2153.8+3746	0.180	0.007	0.260	0.007	0.069	0.069	0.001	0.458	0.501	0.177	9.49	0.36	11.90	0.16
ABELL_2409	0.111	0.010	0.207	0.012	0.032	0.031	0.002	0.161	0.292	0.214	5.96	0.38	4.31	0.14
ABELL_2415	0.194	0.018	NaN	NaN	0.032	0.031	0.002	0.154	0.271	0.200	2.66	0.12	0.39	0.01
3C_444	0.128	0.005	0.326	0.009	0.016	0.016	0.001	2.080	2.050	0.278	5.40	0.21	0.61	0.02
ABELL_2426	0.163	0.016	0.476	0.039	0.027	0.028	0.002	0.057	0.138	0.105	5.82	0.39	1.66	0.05

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
MACS_J2214-1359	0.224	0.016	0.228	0.016	0.011	0.011	0.003	1.600	2.010	0.942	8.83	0.99	12.10	0.47
ABELL_3854	0.152	0.015	0.329	0.024	0.009	0.009	0.002	0.095	0.213	0.196	5.18	0.39	2.27	0.08
MCXC_J2218.6-3853	0.082	0.009	0.252	0.015	0.018	0.018	0.003	0.066	0.158	0.112	5.60	0.33	3.29	0.09
ABELL_2443	0.085	0.007	0.348	0.019	0.052	0.051	0.003	4.220	4.290	0.681	9.72	1.44	1.70	0.05
ABELL_2445	0.081	0.010	0.488	0.034	0.013	0.013	0.002	1.110	1.140	0.359	4.16	0.22	1.70	0.06
ABELL_3880	0.198	0.011	0.644	0.041	0.023	0.022	0.000	0.028	0.055	0.044	5.31	0.23	0.33	0.02
MACS_J2229.8-2756	0.364	0.018	0.435	0.022	0.001	0.002	0.001	0.171	0.216	0.135	5.88	0.56	4.20	0.20
CGCG_514-050	0.391	0.039	NaN	NaN	0.022	0.022	0.002	4.840	4.750	0.778	1.23	0.04	0.00	0.00
ABELL_2457	0.157	0.025	NaN	NaN	0.015	0.016	0.003	1.080	1.260	0.539	3.81	0.22	0.54	0.02
MACS_J2245.0+2637	0.346	0.025	0.368	0.027	0.008	0.008	0.001	0.414	0.532	0.365	6.71	0.93	4.67	0.25
ABELL_3911	0.063	0.009	0.149	0.013	0.600	0.599	0.004	1.840	1.960	0.644	6.11	0.39	2.18	0.06
ABELL_2485	0.292	0.029	0.526	0.063	0.005	0.007	0.002	0.249	0.515	0.415	6.25	0.71	2.77	0.12
ABELL_S1063	0.181	0.007	0.208	0.007	0.043	0.042	0.002	0.025	0.072	0.066	11.30	0.72	25.00	0.46
ABELL_3921	0.085	0.006	0.236	0.007	0.012	0.012	0.001	0.490	0.543	0.178	6.92	0.41	1.96	0.04
ABELL_2507	0.253	0.040	NaN	NaN	0.315	0.316	0.008	1.180	1.920	1.330	4.31	0.43	0.99	0.06
ABELL_2537	0.142	0.007	0.333	0.013	0.010	0.009	0.001	0.059	0.115	0.093	6.65	0.37	4.84	0.11
MCXC_J2311.5+0338	0.218	0.013	0.241	0.014	0.005	0.006	0.002	0.357	0.469	0.338	9.52	1.00	7.83	0.30
ABELL_2550	0.164	0.015	0.516	0.031	0.004	0.004	0.001	0.058	0.128	0.098	1.93	0.12	0.20	0.01
ABELL_2556	0.174	0.008	0.364	0.011	0.008	0.008	0.000	5.430	5.600	0.554	3.86	0.16	0.79	0.02
ABELL_S1101	0.075	0.002	0.542	0.007	0.005	0.005	0.000	0.025	0.028	0.010	7.57	0.10	0.52	0.01

Table D.1 (cont'd)

Cluster	c_{500}	δc_{500}	c	δc	w_{data}	w_{sim}	δw_{sim}	p_{data}	p_{sim}	δp_{sim}	kT	δkT	L_X	δL_X
								10^7			(keV)		$(10^{44} \text{ erg s}^{-1})$	
NGC_7618	0.205	0.022	0.743	0.089	0.014	0.014	0.001	0.118	0.218	0.163	0.78	0.02	0.01	0.00
ABELL_2597	0.086	0.002	0.379	0.003	0.003	0.003	0.000	0.238	0.242	0.024	4.36	0.06	0.86	0.01
RCS_J2327-0204	NaN	NaN	0.569	0.030	0.005	0.004	0.001	0.011	0.057	0.049	9.55	1.43	0.43	0.02
ABELL_2626	0.126	0.006	0.255	0.007	0.012	0.012	0.001	0.730	0.724	0.174	3.26	0.11	0.43	0.01
ABELL_2631	0.212	0.018	0.185	0.016	0.047	0.038	0.006	1.710	2.020	1.040	8.02	0.96	7.57	0.26
MCXC_J2344.2-0422	0.063	0.010	0.591	0.064	0.034	0.033	0.002	0.325	0.405	0.254	4.46	0.18	1.44	0.04
HCG_097	NaN	NaN	NaN	NaN	0.020	0.020	0.003	2.790	2.960	1.180	0.81	0.04	0.01	0.00
ABELL_2667	0.301	0.012	0.320	0.013	0.011	0.010	0.001	0.672	0.786	0.262	7.60	0.68	8.12	0.22
ABELL_2670	0.079	0.006	0.675	0.042	0.021	0.021	0.001	0.040	0.070	0.060	3.78	0.39	0.73	0.02

APPENDIX E

RADIAL PROFILES OF EARLY-TYPE GALAXIES WITH POWERFUL RADIO SOURCES

Table E.1: **Radial profiles for early-type galaxies with powerful radio sources** Radial profile properties for each galaxy with sufficient counts for temperature deprojection. Errors given for radius represent bin widths, all other errors are 1 sigma. Column 1: galaxy name; Column 2: radial bin center; Column 3: half-width of the radial bin; Column 4: grouping of temperature bins; Columns 5-6: best fit temperatures and their errors; Column 7: electron density bin number; Columns 8-9: best fit densities and their errors; in units of 10^{-2} cm^{-3} for compactness; Columns 10-11: calculated entropies and their errors.

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
NGC 315	1.12	0.56	1	0.70	0.04	1	7.36	0.35	3.99	0.26
NGC 315	1.60	0.24	1	0.70	0.04	2	4.33	0.20	5.69	0.37
NGC 315	2.24	0.32	2	0.77	0.03	3	2.90	0.15	8.17	0.42
NGC 315	3.05	0.40	2	0.77	0.03	4	1.68	0.08	11.77	0.58
NGC 315	4.49	0.72	3	0.75	0.02	5	1.08	0.04	15.35	0.52
NGC 741	2.52	1.26	1	0.82	0.01	6	1.74	0.03	12.23	0.24
NGC 741	5.04	1.26	1	0.82	0.01	7	0.90	0.03	18.98	0.46
NGC 741	8.10	1.53	1	0.82	0.01	8	0.45	0.01	30.07	0.79
NGC 741	11.69	1.80	2	1.24	0.03	9	0.44	0.03	46.21	2.20
NGC 741	16.37	2.34	2	1.24	0.03	10	0.17	0.01	86.63	2.53
NGC 741	26.80	5.22	2	1.24	0.03	11	0.12	0.01	110.53	4.08
NGC 741	37.24	5.22	3	1.74	0.21	12	0.08	0.01	193.51	24.86
NGC 741	46.23	4.50	3	1.74	0.21	13	0.07	0.01	228.92	32.23
NGC 741	54.15	3.96	3	1.74	0.21	14	0.08	0.01	205.20	33.25
NGC 741	61.52	3.69	4	1.64	0.16	15	0.02	0.00	470.97	55.73
NGC 741	68.72	3.60	4	1.64	0.16	16	0.06	0.07	242.57	202.78
NGC 741	75.74	3.51	4	1.64	0.16	17	0.06	0.01	219.04	30.52

Table E.1 (cont'd)

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
NGC 1316	0.24	0.12	1	0.77	0.02	1	15.50	1.14	2.66	0.15
NGC 1316	0.37	0.06	1	0.77	0.02	2	10.30	0.82	3.52	0.21
NGC 1316	0.61	0.12	1	0.77	0.02	3	7.57	0.35	4.30	0.18
NGC 1316	0.85	0.12	2	0.93	0.02	4	4.70	0.23	7.16	0.30
NGC 1316	1.10	0.12	2	0.93	0.02	5	3.54	0.25	8.64	0.46
NGC 1316	1.46	0.18	2	0.93	0.02	6	1.95	0.09	12.87	0.52
NGC 1316	2.07	0.30	3	0.79	0.02	7	1.10	0.05	15.86	0.61
NGC 1316	2.92	0.43	3	0.79	0.02	8	0.57	0.03	24.66	1.01
NGC 1316	3.89	0.49	3	0.79	0.02	9	0.50	0.05	26.87	1.75
NGC 1316	4.99	0.55	4	0.73	0.03	10	0.21	0.01	44.83	2.63
NGC 1316	6.08	0.55	4	0.73	0.03	11	0.46	0.14	26.45	5.46
NGC 1316	6.94	0.43	4	0.73	0.03	12	0.45	0.03	26.82	1.44
NGC 1316	8.15	0.61	5	0.31	0.03	13	0.25	0.01	16.68	1.85
NGC 1316	9.61	0.73	6	0.61	0.07	14	0.26	0.02	32.27	3.85
NGC 4261	0.28	0.14	1	0.72	0.01	1	29.50	0.90	1.62	0.04
NGC 4261	0.43	0.07	1	0.72	0.01	2	19.30	0.40	2.15	0.04
NGC 4261	0.57	0.07	1	0.72	0.01	3	13.80	0.38	2.69	0.06
NGC 4261	0.71	0.07	2	0.74	0.02	4	9.65	0.33	3.54	0.11
NGC 4261	0.85	0.07	2	0.74	0.02	5	6.74	0.31	4.50	0.17
NGC 4261	0.99	0.07	2	0.74	0.02	6	6.21	0.42	4.75	0.24
NGC 4261	1.13	0.07	3	0.76	0.03	7	3.66	0.22	6.87	0.38
NGC 4261	1.28	0.07	3	0.76	0.03	8	4.24	0.57	6.24	0.60
NGC 4261	1.42	0.07	3	0.76	0.03	9	3.28	0.28	7.40	0.50
NGC 4261	1.56	0.07	4	0.73	0.05	10	2.73	0.31	8.07	0.82
NGC 4261	1.70	0.07	4	0.73	0.05	11	2.56	0.35	8.44	0.96
NGC 4261	1.84	0.07	4	0.73	0.05	12	2.38	0.32	8.84	0.99
NGC 4261	1.99	0.07	5	0.78	0.04	13	1.30	0.18	14.14	1.46
NGC 4261	2.27	0.14	5	0.78	0.04	14	2.58	0.44	8.94	1.12
NGC 4261	2.41	0.07	5	0.78	0.04	15	1.54	0.13	12.61	0.95

Table E.1 (cont'd)

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
NGC 4261	2.55	0.07	6	0.72	0.05	16	1.09	0.24	14.58	2.38
NGC 4261	2.84	0.14	6	0.72	0.05	17	1.22	0.21	13.51	1.87
NGC 4261	3.12	0.14	6	0.72	0.05	18	1.14	0.12	14.16	1.42
NGC 4261	3.40	0.14	7	0.85	0.05	19	0.76	0.08	21.79	1.85
NGC 4261	3.69	0.14	7	0.85	0.05	20	0.75	0.13	22.13	2.89
NGC 4261	4.11	0.21	7	0.85	0.05	21	0.53	0.05	27.99	2.42
NGC 4261	4.54	0.21	8	0.94	0.05	22	0.58	0.10	29.05	3.78
NGC 4261	4.96	0.21	8	0.94	0.05	23	0.44	0.06	34.86	3.47
NGC 4261	5.53	0.28	8	0.94	0.05	24	0.48	0.06	32.96	3.32
NGC 4261	6.38	0.43	9	1.29	0.03	25	0.31	0.02	60.35	3.29
NGC 4374	0.14	0.07	1	0.78	0.02	1	16.30	0.63	2.62	0.09
NGC 4374	0.28	0.07	1	0.78	0.02	2	9.75	0.45	3.68	0.14
NGC 4374	0.41	0.07	1	0.78	0.02	3	6.68	0.39	4.74	0.21
NGC 4374	0.55	0.07	1	0.78	0.02	4	4.04	0.30	6.63	0.35
NGC 4374	0.69	0.07	1	0.78	0.02	5	5.24	0.72	5.57	0.53
NGC 4374	0.83	0.07	2	0.61	0.02	6	3.20	0.19	6.09	0.31
NGC 4374	1.03	0.10	2	0.61	0.02	7	1.74	0.14	9.12	0.56
NGC 4374	1.17	0.07	2	0.61	0.02	8	2.61	0.76	6.97	1.36
NGC 4374	1.38	0.10	2	0.61	0.02	9	2.23	0.16	7.74	0.45
NGC 4374	1.59	0.10	2	0.61	0.02	10	1.67	0.12	9.40	0.55
NGC 4374	1.79	0.10	3	0.74	0.02	11	1.33	0.14	13.13	0.97
NGC 4374	2.00	0.10	3	0.74	0.02	12	1.78	0.25	10.84	1.03
NGC 4374	2.21	0.10	3	0.74	0.02	13	0.96	0.06	16.32	0.76
NGC 4374	2.55	0.17	3	0.74	0.02	14	1.02	0.10	15.74	1.07
NGC 4374	2.83	0.14	3	0.74	0.02	15	0.91	0.08	16.92	1.01
NGC 4374	3.17	0.17	4	0.94	0.02	16	0.77	0.06	24.02	1.27
NGC 4374	3.66	0.24	4	0.94	0.02	17	0.62	0.04	27.81	1.20
NGC 4374	4.21	0.28	4	0.94	0.02	18	0.37	0.02	38.95	1.59
NGC 4374	4.97	0.38	4	0.94	0.02	19	0.32	0.02	43.51	2.31

Table E.1 (cont'd)

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
NGC 4374	5.86	0.45	4	0.94	0.02	20	0.48	0.03	33.02	1.47
NGC 4782	7.10	0.62	1	10.89	4.50	1	0.30	0.06	524.03	227.54
NGC 4782	4.81	1.05	1	0.76	0.03	2	0.20	0.03	46.94	4.59
NGC 4782	6.32	0.75	1	0.76	0.03	3	0.28	0.13	38.01	12.19
NGC 4782	7.82	0.75	1	0.76	0.03	4	0.42	0.08	29.27	3.77
NGC 4782	9.03	0.60	2	1.55	0.18	5	0.36	0.07	66.07	11.66
NGC 4782	10.23	0.60	2	1.55	0.18	6	0.52	0.14	51.71	11.05
NGC 4782	11.13	0.45	2	1.55	0.18	7	0.48	0.08	54.52	8.69
NGC 4782	12.04	0.45	2	1.55	0.18	8	0.32	0.04	70.95	10.38
NGC 4782	13.24	0.60	3	0.78	0.03	9	0.25	0.02	42.95	2.72
NGC 4782	14.74	0.75	3	0.78	0.03	10	0.13	0.01	66.32	4.46
NGC 4782	17.45	1.35	3	0.78	0.03	11	0.02	0.00	206.97	17.10
NGC 4782	20.16	1.35	3	0.78	0.03	12	0.19	0.42	50.81	74.84
NGC 4782	22.57	1.20	4	2.13	0.52	13	0.16	0.02	157.71	40.60
NGC 4782	25.28	1.35	4	2.13	0.52	14	0.10	0.02	214.57	56.87
NGC 4782	28.29	1.50	4	2.13	0.52	15	0.15	0.04	159.81	49.71
NGC 4782	30.99	1.35	4	2.13	0.52	16	0.05	0.01	353.88	90.56
NGC 5419	0.28	0.14	1	3.41	0.86	1	22.00	0.57	9.36	2.37
NGC 5419	1.12	0.42	1	3.41	0.86	2	3.48	0.14	31.95	8.11
NGC 5419	3.09	0.98	2	1.13	0.06	3	0.68	0.03	31.53	2.04
NGC 5419	7.59	2.25	2	1.13	0.06	4	0.43	0.04	43.06	3.52
NGC 5419	12.65	2.53	3	2.17	1.98	5	0.25	0.03	116.87	107.10
NGC 5419	17.43	2.39	3	2.17	1.98	6	0.23	0.05	123.06	113.54
NGC 5419	22.50	2.53	4	4.37	1.70	7	0.25	0.09	238.79	110.92
NGC 5419	26.43	1.97	4	4.37	1.70	8	0.34	0.15	194.70	94.62
NGC 7626	0.67	0.33	1	1.01	0.05	1	4.03	0.16	8.62	0.46
NGC 7626	2.00	0.67	2	0.93	0.05	2	1.62	0.06	14.56	0.81
NGC 7626	3.77	0.89	3	0.89	0.05	3	0.66	0.03	25.35	1.65
NGC 7626	6.65	1.44	4	0.86	0.05	4	0.37	0.02	36.08	2.33

Table E.1 (cont'd)

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
NGC 7626	9.98	1.66	5	0.96	0.08	5	0.27	0.02	49.64	4.62
NGC 7626	13.53	1.77	6	0.91	0.08	6	0.18	0.01	62.74	6.20
NGC 7626	17.52	2.00	7	0.87	0.09	7	0.28	0.03	43.40	5.54
IC 1459	0.49	0.24	1	0.94	0.02	1	5.14	0.19	6.78	0.23
IC 1459	0.73	0.12	1	0.94	0.02	2	2.70	0.14	10.41	0.45
IC 1459	0.97	0.12	1	0.94	0.02	3	1.92	0.18	13.08	0.89
IC 1459	1.22	0.12	1	0.94	0.02	4	1.67	0.19	14.34	1.12
IC 1459	1.46	0.12	2	0.38	0.09	5	1.09	0.17	7.81	2.05
IC 1459	1.83	0.18	2	0.38	0.09	6	0.98	0.16	8.38	2.21
IC 1459	2.07	0.12	2	0.38	0.09	7	0.96	0.23	8.53	2.45
IC 1459	2.43	0.18	2	0.38	0.09	8	0.69	0.08	10.59	2.68
IC 1459	2.80	0.18	3	1.53	1.74	9	0.46	0.09	55.32	63.42
IC 1459	3.16	0.18	3	1.53	1.74	10	0.59	0.25	47.01	55.20
IC 1459	3.53	0.18	3	1.53	1.74	11	0.36	0.15	65.09	76.24
IC 1459	3.89	0.18	3	1.53	1.74	12	0.38	0.16	63.09	74.18
IC 1459	4.50	0.30	4	0.59	0.07	13	0.32	0.03	26.94	3.92
IC 1459	4.99	0.24	4	0.59	0.07	14	0.27	0.04	30.16	5.04
IC 1459	5.60	0.30	4	0.59	0.07	15	0.19	0.03	38.20	6.09
IC 1459	6.45	0.43	4	0.59	0.07	16	0.19	0.03	38.17	6.45
IC 1459	7.42	0.49	5	0.74	0.03	17	0.09	0.01	80.01	6.63
IC 1459	8.76	0.67	5	0.74	0.03	18	0.10	0.03	75.13	13.60
IC 1459	10.34	0.79	5	0.74	0.03	19	0.08	0.01	82.70	8.98
IC 1459	12.53	1.10	5	0.74	0.03	20	0.11	0.01	69.92	5.92
IC 4296	0.48	0.24	1	0.75	0.02	1	16.50	0.57	2.48	0.09
IC 4296	0.72	0.12	1	0.75	0.02	2	10.10	0.53	3.44	0.16
IC 4296	0.97	0.12	2	0.78	0.03	3	5.62	0.39	5.29	0.33
IC 4296	1.45	0.24	2	0.78	0.03	4	3.52	0.20	7.23	0.41
IC 4296	1.93	0.24	3	0.84	0.03	5	2.50	0.18	9.82	0.63
IC 4296	2.66	0.36	3	0.84	0.03	6	1.19	0.06	16.08	0.86

Table E.1 (cont'd)

Galaxy	radius	Δr	kT bin	kT	σ_{kT}	n_e bin	n_e	σ_{n_e}	K	σ_K
	(kpc)	(kpc)	ID	(keV)	(keV)	ID	(10^{-2} cm^{-3})	(10^{-2} cm^{-3})	(keV cm^2)	(keV cm^2)
IC 4296	3.87	0.60	4	0.89	0.05	7	0.49	0.03	30.82	1.98
IC 4296	6.28	1.21	4	0.89	0.05	8	0.38	0.03	36.60	2.73
IC 4296	9.42	1.57	5	2.10	1.07	9	0.24	0.03	116.60	59.92
IC 4296	12.56	1.57	5	2.10	1.07	10	0.16	0.03	152.23	79.11
IC 4296	15.46	1.45	6	1.29	0.21	11	0.17	0.03	91.31	18.74
IC 4296	17.88	1.21	6	1.29	0.21	12	0.17	0.05	91.62	22.83

APPENDIX F

RADIAL PROFILES FOR THE GALAXIES IN LAKHCHAURA (2018)

Table F.1: **Radial profiles for the HQ sample** Table of all the radial profiles in the HQ sample from (Lakhchaura et al., 2018). Column 1: Galaxy name; Column 2-3: radius and half-bin widths; Column 4-5: entropy and errors; Column 6-7: Ratio between the cooling time and free-fall time, t_c/t_{ff} , and errors.

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
3C449	2.64	2.64	11.20	0.50	13.46	1.26
3C449	8.79	3.51	61.72	13.87	131.46	48.20
3C449	15.82	3.51	66.22	2.91	105.62	9.78
3C449	22.85	3.51	69.14	2.49	78.87	5.85
3C449	30.21	3.85	72.70	2.79	62.81	5.15
3C449	38.13	4.06	98.98	5.12	81.67	9.70
3C449	46.00	3.80	119.66	9.62	96.02	17.63
3C449	53.26	3.46	110.53	6.61	64.78	9.03
3C449	60.27	3.55	132.54	8.23	116.22	15.84
3C449	67.18	3.37	138.31	9.48	108.28	16.48
3C449	73.76	3.20	118.02	7.78	65.52	9.65
3C449	80.24	3.29	117.18	3.49	56.10	5.22
3C449	86.90	3.37	141.59	10.61	75.67	13.72
3C449	93.73	3.46	96.52	1.91	28.30	2.32
IC1860	2.28	2.28	8.03	0.24	18.42	2.40
IC1860	5.25	0.69	11.85	0.96	19.99	3.30
IC1860	6.62	0.69	16.01	1.50	22.65	4.20
IC1860	7.99	0.69	20.49	2.13	32.70	7.84
IC1860	9.48	0.80	22.30	1.98	32.29	6.80
IC1860	11.31	1.03	25.98	3.07	33.96	7.54
IC1860	13.25	0.91	29.95	3.25	34.22	8.40

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
IC1860	15.07	0.91	39.68	6.48	57.78	21.71
IC1860	16.90	0.91	34.62	3.60	36.49	9.13
IC1860	19.30	1.48	47.60	5.59	62.63	18.04
IC1860	22.15	1.37	44.83	4.31	45.22	11.31
IC1860	25.47	1.94	47.94	3.09	42.35	9.15
IC1860	29.23	1.83	61.94	6.41	61.89	17.78
IC1860	33.12	2.06	72.42	5.83	58.76	15.18
IC1860	41.34	6.17	62.85	1.59	30.89	6.33
IC4296	0.56	0.23	1.42	0.07	11.63	0.78
IC4296	1.86	1.07	4.89	0.17	21.75	1.01
IC4296	8.29	5.36	38.01	1.83	118.54	9.06
IC4296	18.28	4.63	101.73	12.11	211.85	56.02
IC4296	26.74	3.84	88.88	5.21	81.15	10.19
IC4296	36.22	5.64	95.26	7.23	91.68	14.57
IC4296	48.69	6.83	93.74	4.18	53.10	4.23
IC4765	1.21	1.21	3.03	0.23	11.01	1.20
IC4765	2.77	0.35	5.77	0.63	12.49	2.39
IC4765	3.62	0.50	6.91	0.63	13.18	1.84
IC4765	4.68	0.57	8.76	0.75	13.08	1.82
IC4765	5.89	0.64	13.06	1.73	20.35	4.31
IC4765	7.31	0.78	18.17	3.24	26.30	7.31
IC4765	9.09	0.99	18.21	1.65	20.10	2.99
NGC57	1.50	1.50	6.36	0.34	36.48	3.31
NGC57	4.30	1.31	14.23	1.28	42.44	6.59
NGC57	8.42	2.81	33.17	2.49	81.99	12.68
NGC315	0.73	0.20	3.27	0.32	20.02	4.65
NGC315	1.13	0.20	5.31	0.45	33.71	6.73
NGC315	1.59	0.27	7.94	0.74	53.52	11.93
NGC315	2.19	0.33	9.50	0.63	55.42	13.64

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC315	3.12	0.60	14.12	0.91	73.57	21.26
NGC315	4.78	1.06	18.25	0.73	57.08	18.90
NGC410	1.28	1.28	6.57	0.79	42.75	8.62
NGC410	3.52	0.96	9.50	1.11	27.30	5.22
NGC410	7.04	2.56	22.45	1.48	49.70	8.16
NGC499	2.90	2.90	15.11	0.52	45.37	4.27
NGC499	7.82	2.03	33.27	1.99	56.17	7.43
NGC499	11.88	2.03	41.50	2.62	54.75	7.55
NGC499	16.15	2.25	40.00	2.52	35.68	4.35
NGC499	20.79	2.39	46.04	1.53	37.10	3.24
NGC499	25.64	2.46	48.82	1.72	34.18	3.19
NGC499	31.29	3.19	56.55	1.91	39.36	3.64
NGC499	37.96	3.48	74.76	2.88	44.53	4.10
NGC499	44.84	3.40	93.61	4.44	66.19	6.96
NGC499	52.15	3.91	111.35	6.12	107.52	14.76
NGC499	60.05	3.98	115.67	6.23	102.68	15.97
NGC499	67.51	3.48	122.87	6.57	100.94	16.78
NGC499	75.33	4.35	104.18	2.20	58.19	6.79
NGC507	2.78	2.78	13.59	0.42	34.75	3.63
NGC507	7.85	2.28	33.76	3.27	70.77	13.22
NGC507	11.84	1.71	38.59	5.35	63.26	15.56
NGC507	15.05	1.50	35.78	4.36	44.88	9.39
NGC507	17.91	1.36	36.74	3.44	35.32	6.83
NGC507	20.55	1.28	38.55	3.07	30.15	5.61
NGC507	23.19	1.36	42.38	2.21	34.46	5.95
NGC507	25.90	1.36	47.92	4.65	37.88	8.36
NGC507	28.68	1.43	49.56	3.56	34.87	6.78
NGC507	31.61	1.50	60.08	4.98	47.12	10.28
NGC507	34.60	1.50	60.87	3.34	37.26	7.09

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC507	37.60	1.50	73.96	6.31	51.51	12.48
NGC507	40.60	1.50	69.32	4.61	40.99	8.64
NGC507	43.74	1.64	85.57	7.26	59.38	14.44
NGC507	47.16	1.78	110.05	14.22	77.39	24.47
NGC507	50.73	1.78	92.08	8.08	46.80	10.99
NGC507	54.30	1.78	178.55	68.20	188.69	164.01
NGC507	58.01	1.93	103.39	10.78	63.20	17.44
NGC507	61.79	1.86	109.76	11.44	66.12	18.31
NGC507	65.50	1.86	117.09	12.33	67.22	19.24
NGC507	69.28	1.93	105.18	8.21	48.76	11.66
NGC507	73.28	2.07	109.52	6.80	49.37	11.11
NGC507	77.63	2.28	124.67	9.18	60.86	14.80
NGC507	82.34	2.43	88.44	1.90	25.84	4.80
NGC533	0.29	0.29	1.72	0.24	17.79	5.08
NGC533	0.81	0.22	2.52	0.24	12.28	3.75
NGC533	1.25	0.22	3.76	0.54	17.51	6.24
NGC533	1.69	0.22	4.70	0.44	15.41	5.51
NGC533	2.13	0.22	5.92	0.46	17.49	6.44
NGC533	2.57	0.22	6.01	0.69	17.25	6.76
NGC533	3.01	0.22	6.40	0.73	17.54	7.10
NGC533	3.45	0.22	7.17	0.72	15.30	6.05
NGC533	3.89	0.22	6.90	0.56	14.49	5.76
NGC533	4.33	0.22	9.38	1.47	21.93	9.78
NGC533	4.77	0.22	9.07	0.51	15.05	5.97
NGC533	5.29	0.29	12.48	1.31	26.90	11.60
NGC533	5.88	0.29	13.54	1.14	23.51	9.85
NGC533	6.54	0.37	13.97	0.93	25.56	10.54
NGC533	7.27	0.37	16.07	1.33	25.03	10.42
NGC533	8.08	0.44	20.08	1.68	33.04	13.99

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC533	9.03	0.51	21.46	1.50	32.27	13.49
NGC533	10.21	0.66	25.37	2.04	38.78	16.52
NGC533	11.82	0.95	33.13	2.54	49.92	21.23
NGC533	14.54	1.76	33.52	0.56	35.54	14.72
NGC708	1.24	1.24	6.72	0.22	27.72	1.34
NGC708	3.73	1.24	7.37	0.07	12.04	0.29
NGC708	6.22	1.24	11.11	0.08	14.95	0.39
NGC708	8.71	1.24	16.59	0.18	19.69	0.66
NGC708	11.51	1.56	19.86	0.10	20.92	0.60
NGC708	14.94	1.87	24.91	0.19	24.78	0.81
NGC708	18.98	2.18	30.63	0.19	28.69	0.93
NGC708	23.65	2.49	41.75	0.76	37.00	1.56
NGC708	28.76	2.62	47.74	0.97	41.60	1.70
NGC708	34.14	2.76	54.91	0.42	37.51	1.34
NGC708	39.65	2.76	66.31	0.63	46.59	1.79
NGC708	45.32	2.91	66.01	0.52	38.26	1.42
NGC708	50.98	2.76	74.20	0.80	42.97	1.70
NGC708	56.57	2.83	73.05	0.71	35.82	1.38
NGC708	62.31	2.91	85.10	1.09	36.49	1.44
NGC708	68.13	2.91	82.00	1.00	29.58	1.15
NGC708	74.18	3.14	89.29	0.75	33.77	1.30
NGC708	81.15	3.83	70.03	0.47	17.20	0.64
NGC741	1.09	1.09	3.63	0.11	19.16	0.73
NGC741	3.28	1.09	8.73	0.31	25.12	1.19
NGC741	5.62	1.25	16.31	0.35	36.04	1.58
NGC741	8.43	1.56	31.80	2.75	77.95	10.45
NGC741	12.02	2.03	34.75	1.33	48.76	3.41
NGC741	18.47	4.42	77.15	3.94	124.13	12.11
NGC741	27.34	4.45	115.53	15.81	174.54	39.26

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC741	35.63	3.84	173.52	64.03	202.01	92.20
NGC741	42.85	3.38	165.06	25.49	139.44	25.86
NGC741	49.38	3.15	163.43	25.63	123.46	23.88
NGC741	55.60	3.07	195.82	46.47	163.88	63.75
NGC741	61.67	2.99	225.48	50.60	153.69	48.76
NGC741	67.81	3.15	192.79	25.62	108.60	18.33
NGC741	74.18	3.23	242.59	49.36	155.94	39.75
NGC741	80.56	3.15	226.06	25.54	117.48	18.79
NGC741	86.93	3.23	169.45	10.53	74.82	7.02
NGC777	1.73	1.73	6.91	0.25	37.38	2.36
NGC777	4.50	1.04	11.17	1.23	30.25	4.57
NGC777	6.58	1.04	13.70	2.00	29.98	5.81
NGC777	9.07	1.45	19.75	2.02	34.10	5.31
NGC777	12.19	1.66	19.39	1.73	24.11	3.01
NGC777	17.45	3.60	25.56	1.25	26.15	2.17
NGC1132	2.73	2.73	9.68	0.27	33.23	1.47
NGC1132	9.23	3.77	37.08	3.18	80.44	11.76
NGC1132	16.25	3.25	53.59	4.37	64.72	9.54
NGC1132	22.43	2.94	65.36	6.87	80.77	17.85
NGC1132	28.10	2.73	61.24	4.90	50.29	8.02
NGC1132	33.34	2.52	109.38	31.57	146.24	94.86
NGC1132	38.26	2.41	65.53	3.70	37.79	4.66
NGC1132	43.19	2.52	94.64	15.74	70.47	22.89
NGC1132	48.12	2.41	91.03	13.44	56.61	15.87
NGC1132	52.84	2.31	99.82	18.25	59.44	19.51
NGC1132	57.45	2.31	84.29	11.84	40.61	9.67
NGC1132	62.17	2.41	124.78	30.64	84.23	44.69
NGC1132	66.99	2.41	95.26	11.25	41.16	8.90
NGC1132	71.81	2.41	114.65	24.60	55.30	20.21

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC1132	76.63	2.41	151.87	39.02	83.15	45.79
NGC1132	81.46	2.41	123.46	19.23	47.75	14.48
NGC1132	86.28	2.41	128.80	17.19	51.51	14.66
NGC1132	93.30	4.61	160.85	22.27	75.79	22.47
NGC1132	100.33	2.41	117.01	14.13	34.15	7.60
NGC1132	105.25	2.52	159.05	107.10	87.58	83.60
NGC1132	110.49	2.73	157.93	24.34	56.95	19.72
NGC1132	116.36	3.14	110.23	4.26	23.52	2.00
NGC1316	0.23	0.09	2.56	0.28	32.57	6.72
NGC1316	0.41	0.09	3.26	0.26	38.76	5.82
NGC1316	0.62	0.11	4.42	0.33	51.38	7.43
NGC1316	0.87	0.14	5.97	0.36	68.14	11.05
NGC1316	1.17	0.16	7.03	0.37	61.97	10.22
NGC1316	1.65	0.32	11.26	0.67	94.06	17.26
NGC1316	2.61	0.64	20.57	3.01	169.46	43.63
NGC1316	4.42	1.17	17.17	0.53	55.60	12.24
NGC1399	0.47	0.47	2.48	0.03	26.05	0.40
NGC1399	1.86	0.93	8.26	0.06	43.02	0.60
NGC1399	4.17	1.38	22.11	0.44	90.46	2.99
NGC1399	6.84	1.29	36.16	1.13	110.21	6.62
NGC1399	9.32	1.19	39.80	1.23	84.91	5.04
NGC1399	11.65	1.14	44.08	1.57	76.49	5.10
NGC1399	13.91	1.12	53.32	2.32	89.24	7.77
NGC1399	16.09	1.06	50.90	1.59	59.57	3.76
NGC1399	18.21	1.06	60.99	2.52	73.97	6.47
NGC1399	20.37	1.10	59.74	1.63	56.89	3.44
NGC1399	23.04	1.57	54.88	0.68	38.89	1.02
NGC1404	0.18	0.18	1.28	0.01	20.73	0.42
NGC1404	0.55	0.18	2.21	0.02	20.23	0.51

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC1404	0.96	0.23	3.25	0.02	21.72	0.58
NGC1404	1.49	0.30	4.45	0.07	23.16	0.75
NGC1404	2.20	0.41	5.70	0.10	24.59	0.86
NGC1404	3.13	0.53	6.86	0.21	24.82	1.14
NGC1404	4.30	0.64	8.27	0.28	23.63	1.17
NGC1404	5.63	0.69	11.61	0.46	30.10	1.71
NGC1404	7.00	0.69	13.38	0.31	26.85	1.13
NGC1404	8.51	0.82	20.09	0.47	41.34	1.74
NGC1404	10.23	0.89	33.96	1.08	83.59	4.67
NGC1404	11.99	0.87	48.40	2.86	124.43	11.56
NGC1404	13.66	0.80	70.55	4.19	171.97	21.36
NGC1404	15.19	0.73	65.86	2.71	122.90	9.66
NGC1404	16.61	0.69	95.69	3.88	143.38	11.46
NGC1404	17.96	0.66	95.67	3.90	125.45	10.10
NGC1404	19.29	0.66	100.91	4.25	125.25	10.52
NGC1404	20.59	0.64	93.39	3.99	98.56	7.76
NGC1404	21.87	0.64	93.53	3.84	89.24	6.62
NGC1404	23.20	0.69	89.64	1.99	75.46	4.15
NGC1404	24.60	0.71	146.84	11.53	207.38	36.95
NGC1404	25.99	0.69	63.66	0.75	28.79	1.01
NGC1407	0.69	0.53	4.67	0.12	47.66	1.71
NGC1407	1.91	0.69	8.57	0.23	44.98	1.79
NGC1407	3.41	0.80	12.78	0.37	41.92	1.67
NGC1407	5.30	1.08	22.75	0.51	63.10	2.41
NGC1407	7.80	1.42	44.32	1.87	128.12	9.47
NGC1407	10.79	1.58	75.03	5.69	221.54	31.42
NGC1407	13.87	1.50	81.63	6.57	182.53	28.17
NGC1407	16.90	1.53	144.82	31.81	478.53	230.97
NGC1407	20.03	1.61	104.88	8.84	178.53	28.08

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC1407	23.17	1.53	167.67	34.72	410.89	186.47
NGC1407	26.08	1.39	146.99	20.32	260.78	78.22
NGC1407	29.00	1.53	111.38	6.55	118.84	12.30
NGC1521	0.30	0.30	1.79	0.28	20.90	5.21
NGC1521	0.97	0.36	4.57	0.55	23.49	4.05
NGC1521	1.64	0.30	6.36	0.69	23.18	3.77
NGC1521	2.37	0.43	7.38	0.94	25.00	4.33
NGC1521	3.46	0.67	7.76	0.43	18.07	1.47
NGC1550	3.37	3.37	8.18	0.16	19.89	0.66
NGC1550	9.87	3.13	16.35	0.43	18.39	0.84
NGC1550	16.86	3.85	26.07	0.80	24.77	1.45
NGC1550	24.80	4.09	36.88	0.98	29.18	1.63
NGC1550	32.99	4.09	44.67	1.48	30.09	2.15
NGC1550	41.26	4.17	53.74	1.32	27.22	1.27
NGC1550	50.17	4.74	60.31	1.77	26.80	1.59
NGC1550	60.04	5.14	71.01	2.48	30.89	2.27
NGC1550	71.68	6.50	64.94	1.33	19.70	0.73
NGC1600	0.76	0.76	5.88	0.23	44.88	5.97
NGC1600	2.29	0.76	9.77	0.39	42.60	7.16
NGC1600	4.86	1.80	34.77	3.14	177.36	46.25
NGC1600	8.62	1.97	70.62	11.27	265.42	105.16
NGC1600	12.34	1.75	69.26	8.91	140.35	48.72
NGC1600	15.72	1.64	147.03	74.08	581.47	673.76
NGC1600	18.83	1.47	65.86	5.56	70.44	22.06
NGC2300	1.56	1.56	5.28	0.19	26.09	1.27
NGC2300	5.58	2.45	13.94	0.56	30.76	1.73
NGC2300	11.29	3.26	39.79	2.48	76.58	6.88
NGC2300	17.76	3.21	88.62	9.28	184.36	38.41
NGC2300	23.64	2.66	124.68	17.55	256.24	76.15

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC2300	28.46	2.16	91.95	8.23	97.13	16.45
NGC2300	32.65	2.03	118.68	9.60	112.37	17.57
NGC2300	36.58	1.90	121.46	11.19	99.41	18.43
NGC2300	40.47	1.99	127.58	9.58	95.08	13.45
NGC2300	44.53	2.07	129.14	14.20	105.19	25.30
NGC2300	48.38	1.78	85.24	3.61	36.37	2.76
NGC2305	0.30	0.30	2.36	0.21	41.58	5.38
NGC2305	0.84	0.25	2.98	0.34	20.22	3.36
NGC2305	1.48	0.40	4.38	0.41	21.28	2.77
NGC2305	2.47	0.59	7.88	0.92	31.43	4.91
NGC2305	3.86	0.79	13.37	1.24	42.83	6.20
NGC2305	5.98	1.33	11.11	0.75	21.51	2.01
NGC3091	1.38	1.38	8.35	0.29	53.15	3.13
NGC3091	3.46	0.69	12.93	1.31	41.98	5.99
NGC3091	4.78	0.63	16.90	2.68	49.59	11.19
NGC3091	5.99	0.58	15.79	1.49	30.74	4.16
NGC3091	7.43	0.86	22.39	1.22	43.43	5.29
NGC3091	9.16	0.86	17.56	0.39	18.47	0.93
NGC3923	0.20	0.20	1.70	0.05	24.64	1.53
NGC3923	0.58	0.18	2.37	0.09	21.98	1.41
NGC3923	1.03	0.28	3.58	0.14	26.99	2.11
NGC3923	1.75	0.45	5.43	0.33	31.16	3.15
NGC3923	2.93	0.73	9.14	0.65	40.15	4.62
NGC3923	4.45	0.80	14.77	1.03	54.93	7.21
NGC3923	6.08	0.83	17.34	1.23	45.51	6.06
NGC3923	7.70	0.80	30.91	3.27	89.41	18.07
NGC3923	9.20	0.70	24.90	2.01	39.50	5.27
NGC4073	1.86	1.86	10.88	0.16	32.22	2.92
NGC4073	4.37	0.64	16.45	0.96	36.79	4.46

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4073	5.66	0.64	20.85	1.88	48.00	8.11
NGC4073	7.02	0.72	25.59	4.44	60.18	17.78
NGC4073	8.53	0.79	29.43	3.38	64.60	13.36
NGC4073	10.25	0.93	36.49	4.56	83.06	19.56
NGC4073	12.18	1.00	45.92	5.86	99.94	23.59
NGC4073	14.19	1.00	51.76	7.29	105.51	28.85
NGC4073	16.48	1.29	58.44	6.06	93.81	19.94
NGC4073	19.13	1.36	65.44	6.40	96.61	19.43
NGC4073	21.78	1.29	88.26	16.27	150.39	54.05
NGC4073	24.79	1.72	75.50	10.83	86.32	21.58
NGC4073	28.73	2.22	101.42	7.59	136.11	26.40
NGC4073	34.61	3.65	90.14	3.79	77.10	11.08
NGC4125	0.23	0.08	2.79	0.52	28.22	12.35
NGC4125	0.43	0.13	4.08	0.63	47.29	14.58
NGC4125	0.69	0.13	6.45	0.54	62.15	14.88
NGC4125	0.94	0.13	8.88	0.88	90.95	24.93
NGC4125	1.23	0.15	10.90	1.24	103.67	32.99
NGC4125	1.53	0.15	8.84	0.96	68.73	22.40
NGC4125	1.86	0.18	8.87	0.73	51.68	15.83
NGC4125	2.27	0.23	11.35	0.96	65.98	21.90
NGC4125	2.81	0.31	10.30	0.59	36.36	12.17
NGC4261	0.28	0.14	1.10	0.02	14.17	1.47
NGC4261	0.67	0.25	2.65	0.04	23.73	3.06
NGC4261	1.44	0.53	5.37	0.14	35.27	5.04
NGC4261	3.09	1.12	14.20	0.30	67.99	10.05
NGC4261	5.94	1.72	42.49	1.49	169.87	27.13
NGC4261	12.02	4.36	107.04	12.62	423.98	104.52
NGC4261	18.55	2.18	155.67	21.64	256.21	86.33
NGC4261	24.88	4.15	151.31	11.96	143.39	31.33

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4261	31.77	2.74	113.86	3.91	66.51	11.23
NGC4374	0.41	0.24	2.29	0.10	26.32	7.29
NGC4374	0.99	0.34	3.69	0.20	25.04	6.58
NGC4374	1.70	0.38	5.73	0.29	27.13	6.77
NGC4374	2.57	0.49	8.74	0.41	34.74	8.40
NGC4374	3.78	0.71	15.22	0.53	54.98	13.26
NGC4374	5.60	1.11	27.25	1.42	109.44	28.74
NGC4374	7.95	1.25	49.78	6.53	220.69	73.07
NGC4374	10.29	1.09	52.64	6.24	159.97	54.33
NGC4374	12.45	1.07	88.41	15.26	234.34	103.59
NGC4374	14.55	1.03	100.51	20.91	228.15	119.14
NGC4374	16.52	0.95	139.13	47.12	314.58	260.58
NGC4374	18.37	0.89	80.61	5.49	73.56	27.31
NGC4406	1.61	1.61	10.87	0.43	52.64	2.85
NGC4406	4.37	1.15	26.62	2.02	75.24	8.69
NGC4406	6.41	0.88	22.78	1.43	40.44	3.63
NGC4406	8.09	0.81	23.33	1.41	32.98	2.84
NGC4406	9.68	0.79	31.29	1.60	46.93	5.00
NGC4406	11.24	0.77	26.31	0.87	26.28	1.60
NGC4406	12.85	0.84	36.03	2.79	37.36	4.16
NGC4406	14.56	0.86	41.99	3.04	42.70	4.82
NGC4406	16.22	0.81	48.19	2.20	48.56	5.09
NGC4406	17.80	0.77	45.89	1.69	38.67	3.30
NGC4406	19.33	0.77	53.74	2.78	49.64	5.88
NGC4406	20.85	0.75	35.32	0.48	17.54	0.67
NGC4472	0.32	0.32	2.38	0.02	35.65	0.31
NGC4472	0.96	0.32	4.34	0.03	26.80	0.23
NGC4472	1.74	0.45	7.56	0.05	36.00	0.31
NGC4472	2.79	0.60	13.36	0.10	57.55	0.60

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4472	4.02	0.62	18.41	0.14	68.27	1.13
NGC4472	5.24	0.60	23.27	0.27	65.79	1.50
NGC4472	6.47	0.62	28.48	0.32	70.12	1.29
NGC4472	7.71	0.62	30.79	0.32	65.14	1.16
NGC4472	8.98	0.64	35.70	0.44	71.43	1.47
NGC4472	10.28	0.66	36.53	0.49	63.97	1.58
NGC4472	11.62	0.68	41.73	0.42	68.92	1.30
NGC4472	13.00	0.70	42.78	0.64	66.14	1.79
NGC4472	14.41	0.72	47.72	0.59	68.29	1.77
NGC4472	15.86	0.74	52.30	0.85	71.82	2.26
NGC4472	17.39	0.79	56.39	1.07	72.96	2.33
NGC4472	19.09	0.91	65.54	1.31	82.28	2.76
NGC4472	21.13	1.13	57.87	0.41	53.03	0.72
NGC4486	0.79	0.71	5.85	0.02	35.11	0.34
NGC4486	1.93	0.43	8.85	0.03	25.09	0.21
NGC4486	2.78	0.41	9.55	0.07	22.73	0.27
NGC4486	3.64	0.45	11.58	0.06	24.81	0.28
NGC4486	4.61	0.51	15.79	0.04	34.31	0.27
NGC4486	5.65	0.53	17.34	0.07	34.25	0.36
NGC4486	6.76	0.57	20.26	0.08	42.41	0.38
NGC4486	7.88	0.55	22.27	0.11	43.45	0.43
NGC4486	8.96	0.53	23.80	0.11	41.89	0.42
NGC4486	10.01	0.51	25.16	0.09	41.47	0.40
NGC4486	11.03	0.51	29.61	0.10	38.78	0.40
NGC4486	12.08	0.53	32.75	0.12	43.18	0.47
NGC4486	13.14	0.53	34.68	0.13	43.94	0.49
NGC4486	14.22	0.55	35.46	0.17	45.43	0.52
NGC4486	15.33	0.55	35.05	0.14	40.02	0.41
NGC4486	16.45	0.57	37.68	0.17	42.82	0.49

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4486	17.61	0.59	40.95	0.21	45.24	0.50
NGC4486	18.81	0.61	42.03	0.21	43.85	0.48
NGC4486	20.09	0.67	45.76	0.20	44.67	0.47
NGC4486	21.47	0.71	49.27	0.21	48.17	0.52
NGC4486	22.93	0.75	52.68	0.36	50.12	0.60
NGC4486	24.49	0.81	57.37	0.16	54.30	0.57
NGC4486	26.12	0.83	56.33	0.14	47.66	0.48
NGC4486	27.80	0.85	74.52	0.30	67.70	0.76
NGC4486	29.51	0.87	47.13	0.16	22.21	0.22
NGC4552	0.21	0.13	2.21	0.07	11.35	0.63
NGC4552	0.49	0.15	3.14	0.07	29.71	1.33
NGC4552	0.76	0.11	3.58	0.08	30.87	1.31
NGC4552	1.02	0.15	3.79	0.04	24.97	0.61
NGC4552	1.44	0.27	5.34	0.13	42.30	1.54
NGC4552	2.12	0.42	7.32	0.33	50.44	3.61
NGC4552	3.04	0.49	8.65	0.45	37.89	2.51
NGC4552	4.21	0.68	12.27	0.92	47.33	4.18
NGC4552	6.22	1.33	25.35	1.72	93.05	9.31
NGC4552	9.71	2.16	42.84	3.29	116.47	15.27
NGC4552	13.66	1.78	66.67	9.14	108.97	18.49
NGC4552	16.92	1.48	110.80	4.91	55.62	4.20
NGC4636	0.38	0.38	1.84	0.06	11.53	0.61
NGC4636	1.01	0.25	2.88	0.05	10.79	0.36
NGC4636	1.56	0.30	4.03	0.07	15.23	0.51
NGC4636	2.21	0.34	5.61	0.15	22.73	0.99
NGC4636	2.87	0.32	6.83	0.09	22.67	0.89
NGC4636	3.50	0.30	7.37	0.09	20.35	0.80
NGC4636	4.09	0.29	9.15	0.15	20.36	0.83
NGC4636	4.68	0.30	10.52	0.07	22.00	0.85

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4636	5.31	0.32	11.60	0.10	22.47	0.90
NGC4636	6.01	0.38	13.68	0.18	25.04	1.04
NGC4636	6.85	0.46	16.02	0.22	29.59	1.25
NGC4636	7.86	0.55	20.94	0.23	42.09	1.80
NGC4636	9.00	0.59	23.86	0.26	45.93	2.00
NGC4636	10.22	0.63	30.68	0.57	58.10	2.73
NGC4636	11.51	0.67	30.65	0.54	48.90	2.46
NGC4636	12.90	0.72	35.06	0.63	55.77	2.76
NGC4636	14.41	0.78	40.63	0.80	63.86	3.28
NGC4636	16.00	0.82	41.87	0.65	56.66	2.79
NGC4636	17.70	0.88	43.13	0.63	54.50	2.53
NGC4636	19.54	0.97	51.92	0.42	64.54	2.93
NGC4636	21.67	1.16	44.10	0.24	38.01	1.66
NGC4649	0.16	0.16	1.94	0.01	41.35	0.60
NGC4649	0.41	0.10	2.15	0.02	22.63	0.35
NGC4649	0.63	0.12	3.27	0.02	34.06	0.59
NGC4649	0.89	0.14	3.78	0.02	28.17	0.47
NGC4649	1.18	0.16	4.80	0.05	31.04	0.61
NGC4649	1.56	0.22	6.30	0.07	37.58	0.77
NGC4649	2.05	0.28	8.44	0.05	44.00	0.81
NGC4649	2.69	0.36	11.88	0.07	60.06	1.20
NGC4649	3.46	0.41	14.62	0.08	61.41	1.22
NGC4649	4.36	0.49	18.30	0.13	67.92	1.43
NGC4649	5.41	0.55	22.61	0.28	75.44	1.78
NGC4649	6.55	0.59	27.24	0.42	80.43	2.09
NGC4649	7.78	0.63	32.78	0.43	82.11	2.02
NGC4649	9.10	0.69	37.71	0.51	85.97	2.15
NGC4649	10.56	0.77	44.71	0.64	96.94	2.53
NGC4649	12.16	0.83	55.05	1.09	124.87	4.00

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4649	13.86	0.87	72.12	1.94	177.98	7.48
NGC4649	15.56	0.83	87.41	4.10	232.40	15.59
NGC4649	17.22	0.83	109.43	3.45	295.42	20.67
NGC4649	20.81	2.76	110.16	1.42	216.18	5.71
NGC4696	1.34	1.34	4.60	0.01	4.73	0.03
NGC4696	3.53	0.85	9.31	0.03	9.98	0.06
NGC4696	5.14	0.76	10.25	0.03	10.61	0.06
NGC4696	6.71	0.80	13.92	0.05	14.35	0.07
NGC4696	8.32	0.80	15.46	0.05	15.63	0.09
NGC4696	10.06	0.94	19.34	0.03	20.27	0.08
NGC4696	11.98	0.98	22.60	0.05	24.63	0.10
NGC4696	14.04	1.07	28.72	0.20	30.41	0.40
NGC4696	16.23	1.12	35.10	0.16	37.86	0.23
NGC4696	18.51	1.16	40.86	0.20	44.91	0.29
NGC4696	20.83	1.16	43.92	0.22	46.48	0.31
NGC4696	23.16	1.16	47.73	0.39	47.27	0.44
NGC4696	25.48	1.16	49.99	0.25	48.21	0.33
NGC4696	27.85	1.21	52.32	0.51	45.66	0.49
NGC4696	30.31	1.25	54.33	0.53	44.14	0.47
NGC4696	32.90	1.34	61.77	0.50	46.99	0.42
NGC4696	35.76	1.52	76.92	0.45	63.02	0.48
NGC4696	38.85	1.56	85.48	0.52	70.96	0.57
NGC4696	42.07	1.65	90.99	0.44	68.23	0.47
NGC4696	45.55	1.83	92.99	0.44	63.64	0.41
NGC4696	49.71	2.32	111.19	0.44	78.21	0.47
NGC4696	54.85	2.82	138.93	0.59	110.93	0.77
NGC4696	61.47	3.80	122.94	0.44	70.47	0.35
NGC4778	1.56	1.56	5.64	0.09	19.99	3.95
NGC4778	4.24	1.13	8.68	0.17	15.02	2.54

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4778	6.65	1.27	13.36	0.37	21.79	3.58
NGC4778	9.26	1.34	17.32	0.36	22.67	3.72
NGC4778	12.09	1.49	21.82	0.51	27.70	4.70
NGC4778	15.27	1.70	26.51	0.43	26.28	4.60
NGC4778	19.38	2.40	41.99	1.05	51.95	9.70
NGC4778	24.61	2.83	60.55	2.66	69.48	14.28
NGC4778	30.34	2.90	71.67	3.45	73.01	15.86
NGC4778	36.28	3.04	92.59	3.51	89.51	19.96
NGC4778	42.36	3.04	112.56	5.16	108.76	25.71
NGC4778	48.44	3.04	125.81	7.29	114.13	28.61
NGC4778	54.52	3.04	165.52	16.63	174.48	54.92
NGC4778	60.46	2.90	172.03	17.33	162.30	51.99
NGC4778	66.19	2.83	155.83	10.82	111.77	31.07
NGC4778	71.99	2.97	158.50	9.20	100.89	26.95
NGC4778	78.71	3.75	192.48	12.97	142.79	40.02
NGC4778	86.91	4.46	376.99	119.33	547.61	411.43
NGC4778	96.74	5.37	183.69	6.19	90.45	23.03
NGC4782	0.75	0.41	5.50	1.48	18.94	9.92
NGC4782	2.15	0.99	15.51	4.78	60.42	35.43
NGC4782	4.06	0.93	35.57	7.78	60.68	26.70
NGC4782	5.80	0.81	41.77	9.18	56.43	24.84
NGC4782	7.31	0.70	31.66	3.80	26.26	8.03
NGC4782	10.49	2.48	72.69	12.90	113.93	47.16
NGC4782	15.91	2.95	78.48	11.64	70.41	21.24
NGC4936	2.99	2.99	12.45	1.21	29.89	12.12
NGC4936	8.53	2.54	33.59	2.55	56.79	23.32
NGC4936	13.39	2.32	47.46	5.02	63.84	29.86
NGC4936	17.69	1.98	48.95	4.20	55.72	28.13
NGC4936	21.50	1.83	48.71	5.28	43.94	23.41

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC4936	25.17	1.83	42.63	2.06	24.12	12.98
NGC5044	1.10	0.93	3.17	0.05	5.75	0.14
NGC5044	2.54	0.51	5.86	0.10	9.39	0.23
NGC5044	3.51	0.47	7.56	0.14	11.54	0.33
NGC5044	4.40	0.42	8.18	0.13	11.29	0.25
NGC5044	5.20	0.38	8.57	0.20	10.83	0.30
NGC5044	5.97	0.38	8.81	0.20	10.65	0.29
NGC5044	6.73	0.38	8.59	0.14	9.26	0.22
NGC5044	7.49	0.38	9.47	0.17	10.03	0.23
NGC5044	8.25	0.38	10.04	0.14	9.71	0.19
NGC5044	9.01	0.38	10.99	0.13	10.10	0.19
NGC5044	9.77	0.38	11.89	0.15	11.14	0.21
NGC5044	10.53	0.38	12.30	0.07	10.66	0.16
NGC5044	11.34	0.42	13.16	0.07	11.39	0.17
NGC5044	12.18	0.42	14.32	0.08	12.60	0.20
NGC5044	13.03	0.42	14.64	0.08	12.19	0.19
NGC5044	13.88	0.42	15.35	0.08	12.54	0.20
NGC5044	14.76	0.47	16.14	0.09	12.79	0.20
NGC5044	15.69	0.47	16.39	0.08	12.23	0.19
NGC5044	16.75	0.59	17.95	0.14	13.70	0.23
NGC5044	18.20	0.87	19.47	0.15	14.40	0.22
NGC5044	20.11	1.04	22.32	0.14	16.03	0.24
NGC5044	22.27	1.13	24.20	0.13	16.24	0.25
NGC5044	24.70	1.30	25.72	0.11	16.01	0.24
NGC5044	27.73	1.73	34.60	0.16	23.80	0.37
NGC5044	31.63	2.17	30.19	0.06	13.64	0.20
NGC5129	2.28	2.28	6.31	0.16	22.58	0.85
NGC5129	9.53	4.97	29.13	1.00	55.41	2.89
NGC5129	19.47	4.97	59.89	4.74	76.66	8.84

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC5129	28.80	4.35	67.05	3.83	61.04	6.20
NGC5129	36.98	3.83	77.18	5.34	56.21	7.59
NGC5129	44.33	3.52	96.76	12.18	75.95	21.18
NGC5129	51.17	3.31	79.98	4.89	39.75	5.11
NGC5129	57.69	3.21	93.98	7.86	47.58	8.65
NGC5129	64.01	3.11	77.60	6.54	36.60	6.91
NGC5129	70.23	3.11	91.52	8.53	46.06	9.63
NGC5129	76.24	2.90	59.99	1.54	15.71	0.88
NGC5419	0.74	0.25	2.72	0.22	17.30	2.21
NGC5419	1.52	0.54	5.59	0.25	27.59	2.04
NGC5419	6.25	4.19	34.25	3.47	137.55	22.11
NGC5419	14.62	4.19	112.67	42.19	269.90	127.33
NGC5419	22.21	3.40	130.20	29.23	167.59	51.70
NGC5419	28.76	3.15	227.11	83.42	396.10	286.55
NGC5419	34.46	2.55	120.04	21.95	107.06	34.07
NGC5419	39.56	2.55	138.21	20.53	119.27	25.03
NGC5419	44.47	2.37	124.60	19.58	87.61	18.42
NGC5419	49.14	2.31	162.76	44.73	137.41	75.48
NGC5419	53.15	1.70	115.08	14.66	47.57	13.62
NGC5419	56.85	2.00	98.25	3.32	30.10	2.28
NGC5813	1.00	0.59	4.30	0.04	15.90	0.33
NGC5813	2.04	0.45	4.61	0.06	12.20	0.26
NGC5813	2.93	0.45	5.78	0.10	15.01	0.39
NGC5813	3.86	0.48	7.04	0.08	17.24	0.41
NGC5813	4.79	0.45	8.38	0.11	19.33	0.55
NGC5813	5.65	0.41	7.53	0.23	14.27	0.60
NGC5813	6.48	0.41	9.59	0.19	17.10	0.56
NGC5813	7.27	0.38	10.12	0.18	15.61	0.49
NGC5813	8.03	0.38	11.53	0.13	15.48	0.41

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC5813	8.79	0.38	11.70	0.06	13.87	0.32
NGC5813	9.58	0.41	13.02	0.12	15.15	0.38
NGC5813	10.44	0.45	14.10	0.14	15.28	0.38
NGC5813	11.44	0.55	14.84	0.13	15.42	0.38
NGC5813	12.68	0.69	16.95	0.12	17.99	0.44
NGC5813	14.23	0.86	19.05	0.07	19.77	0.49
NGC5813	16.26	1.17	24.84	0.10	29.16	0.74
NGC5813	18.95	1.52	20.56	0.06	14.12	0.35
NGC5846	0.78	0.78	3.58	0.07	24.27	2.61
NGC5846	2.23	0.68	5.39	0.15	16.49	2.27
NGC5846	3.66	0.74	7.70	0.23	19.06	2.81
NGC5846	5.15	0.74	8.82	0.38	17.71	2.75
NGC5846	6.67	0.78	12.08	0.65	21.72	3.56
NGC5846	8.19	0.74	14.13	0.85	22.07	3.67
NGC5846	9.64	0.71	16.14	0.91	22.33	3.74
NGC5846	11.07	0.71	19.05	0.69	24.36	3.98
NGC5846	12.49	0.71	22.39	0.98	26.67	4.44
NGC5846	13.91	0.71	22.30	1.07	25.24	4.30
NGC5846	15.31	0.68	23.86	0.96	23.23	3.86
NGC5846	16.66	0.68	26.16	1.64	26.31	4.62
NGC5846	18.06	0.71	30.41	1.06	26.03	4.29
NGC5846	19.54	0.78	37.67	2.23	37.80	6.63
NGC5846	21.13	0.81	39.17	1.79	37.19	6.26
NGC5846	22.81	0.87	50.80	2.97	49.44	8.79
NGC5846	24.56	0.87	63.26	4.35	65.45	11.97
NGC5846	26.34	0.91	68.29	5.89	74.94	18.25
NGC5846	28.15	0.91	76.42	4.18	78.23	13.87
NGC5846	29.93	0.87	81.23	8.33	85.34	18.46
NGC5846	31.71	0.91	76.60	4.90	72.61	14.39

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC5846	33.59	0.97	84.31	5.29	82.49	16.21
NGC5846	35.63	1.07	68.99	1.81	44.99	7.43
NGC6407	3.17	3.17	8.76	0.88	21.33	3.89
NGC6407	8.59	2.25	19.73	3.10	26.60	7.09
NGC6407	12.39	1.55	31.33	9.25	35.81	16.87
NGC6407	17.58	3.64	37.16	4.60	33.46	9.05
NGC6861	0.68	0.47	3.48	0.14	30.62	1.68
NGC6861	2.80	1.65	16.57	1.01	97.89	7.85
NGC6861	6.46	2.01	46.80	6.57	192.75	43.17
NGC6861	10.26	1.79	60.70	6.22	146.87	29.10
NGC6861	13.42	1.36	91.42	23.97	239.61	138.52
NGC6861	16.04	1.26	64.62	6.61	82.39	16.37
NGC6861	18.41	1.11	77.31	18.72	135.92	73.87
NGC6861	20.53	1.00	56.31	5.50	55.90	12.11
NGC6861	22.47	0.93	94.94	37.13	156.62	137.71
NGC6861	24.26	0.86	56.17	5.45	42.49	9.16
NGC6861	25.95	0.83	62.13	7.54	47.84	12.96
NGC6861	27.60	0.83	46.73	2.12	22.82	2.21
NGC6868	1.68	1.45	9.34	0.47	50.39	3.28
NGC6868	5.09	1.96	31.71	2.34	110.13	11.29
NGC6868	9.79	2.74	30.68	1.98	57.11	4.87
NGC6868	15.29	2.75	41.39	1.52	53.53	3.43
NGC6868	20.35	2.31	51.77	2.44	57.78	5.55
NGC6868	24.63	1.97	61.39	4.81	64.21	9.93
NGC6868	28.49	1.89	50.57	2.44	33.77	3.04
NGC6868	33.84	3.47	47.78	1.69	27.55	1.70
NGC7619	1.57	1.57	7.77	0.16	46.50	1.40
NGC7619	6.27	3.13	24.52	0.68	70.33	3.06
NGC7619	13.08	3.68	45.58	1.53	82.17	4.96

Table F.1 (cont'd)

Galaxy	radius	Δr	K	σ_K	t_c/t_{ff}	$\delta t_c/t_{ff}$
	(kpc)	(kpc)	(keV cm ²)			
NGC7619	20.19	3.43	69.45	4.73	93.46	10.42
NGC7619	27.00	3.37	85.75	7.93	100.42	15.96
NGC7619	33.44	3.07	126.62	13.10	136.13	30.50
NGC7619	39.35	2.83	129.04	12.44	109.57	22.74
NGC7619	45.14	2.95	170.22	24.55	164.14	52.27
NGC7619	50.74	2.65	137.25	13.70	89.14	19.56
NGC7619	55.98	2.59	132.82	10.56	70.76	12.21
NGC7619	61.22	2.65	176.45	26.14	116.21	38.35
NGC7619	66.83	2.95	113.74	4.09	38.42	2.76
NGC7796	1.82	1.82	10.73	0.48	69.46	4.38
NGC7796	5.70	2.06	31.72	3.97	115.70	17.54
NGC7796	11.76	4.00	47.47	2.82	89.39	7.78
NGC7796	22.21	6.45	76.14	9.61	121.91	20.60
NGC7796	33.85	5.19	76.09	11.90	70.50	13.04

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